

$$E(r=R) = \frac{Q}{4\pi\varepsilon_0 R} = g \cdot 10^8 \cdot 10^{-9} \cdot 10^2 = g \frac{10^1}{m}$$

$$U_{ext} = \frac{1}{4\pi\varepsilon_0} \frac{1}{R} = g \cdot 10^8 \cdot 10^{-9} \cdot 10^2 = 0.9 \text{ kJ}$$

$$E(r) = \frac{U(r)}{r}$$

$$W = 10^{-8} \frac{As}{m^2}$$

$$\oint_E dA = Q$$

$$2\pi r l E_{in} = 0$$

$$E_{in} = 0$$

$$2\pi r l E_{out} = 2\pi r l E_{in}$$

$$E_{out} = \frac{W R}{C_0} \frac{1}{r}$$

A graph showing the electric field E as a function of distance r . The vertical axis is labeled E and the horizontal axis is labeled r . A curve starts at a positive value for small r , decreases rapidly, and then levels off towards zero as r increases. Two points on the curve are labeled V_{∞} and V_R .

$E = \frac{Q}{4\pi\epsilon_0 r^2}$

A cylindrical Gaussian pillbox of radius R and height l is shown around a point charge Q . The electric field E is constant across the top face of the pillbox. The flux through the pillbox is given by $\Phi_E = E l A$, where $A = \pi R^2$.

$E l A = \frac{Q}{\epsilon_0} \Rightarrow E = \frac{Q}{4\pi\epsilon_0 R^2}$

The total flux Φ_E through a closed Gaussian surface enclosing the charge Q is given by $\Phi_E = \frac{Q}{\epsilon_0}$.

$E = \frac{Q}{4\pi\epsilon_0 R^2}$

The potential V at a distance r from a point charge Q is given by $V = \frac{Q}{4\pi\epsilon_0 r}$.

$V = \frac{Q}{4\pi\epsilon_0 r} = \frac{Q}{4\pi\epsilon_0 R^2} \cdot r = \frac{1}{4\pi\epsilon_0 R^2} \cdot Q r = \frac{1}{4\pi\epsilon_0 R^2} \cdot \frac{Q}{2\pi R l} \cdot 2\pi R l r = \frac{Q}{2\pi R l} \cdot \frac{r}{R} = \frac{Q}{2\pi R l} \cdot \frac{V_R}{V_\infty}$

$$\oint E(\Sigma) d\Sigma = 0$$

$$E = g(E_\text{in})$$

$$W = \int_A^B F(s) ds$$

$$\int_A^B E dr = -\frac{N}{g}$$

$$\int_A^B (-E(\Sigma)) dr = U_B - U_A$$

$$\frac{F}{g} = B \begin{pmatrix} V \\ 0 \end{pmatrix} \begin{pmatrix} V \\ 0 \end{pmatrix}^\top \quad \text{Ed} = V$$

$$E(\Sigma)$$

$$\underline{\text{Method}}$$

$$E_\text{in} \cdot 4\pi r^2 R = Q = 0$$

$$E_\text{in} = 0$$

$$E_\text{ext} = 4r^2 \pi = Q$$

$$E_\text{ext} = \frac{Q}{4\pi \epsilon_0 r^2}$$

$$U(\Sigma) = - \int_{\infty}^R \frac{Q}{4\pi \epsilon_0} \frac{1}{r^2} dr = \frac{Q}{4\pi \epsilon_0} \left[-\frac{1}{r} \right]_{\infty}^R$$

$$= \frac{Q}{4\pi \epsilon_0} \left[\frac{1}{r} \right]_{\infty}^R = \frac{Q}{4\pi \epsilon_0 R}$$

$$C = \frac{Q}{U} = 4\pi \epsilon_0 R = 0.11 \cdot 10^{-9} = \frac{1}{3} \mu\text{F}$$

$$F = \begin{bmatrix} A & 0 \\ 0 & B \end{bmatrix}$$

The diagram illustrates a hollow spherical shell (bulb) with an internal charge Q . The outer surface has a uniform positive charge density σ . Electric field lines originate from the positive charges on the outer surface and terminate at the negative charges inside the shell. Equipotential surfaces are shown as concentric spheres centered at the shell's center. The potential is zero both inside and outside the shell.

$$\begin{aligned} C_{(R)} &\stackrel{\nu}{\downarrow} \\ Q_{(R)} &\Rightarrow \alpha^* \\ R &\Rightarrow R^* \quad u = \frac{\alpha^*(r)}{C(R^*)} \end{aligned}$$

