Finite-size effects in quantum thermodynamics

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TEAM-NET

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3. Asymptotic reversibility
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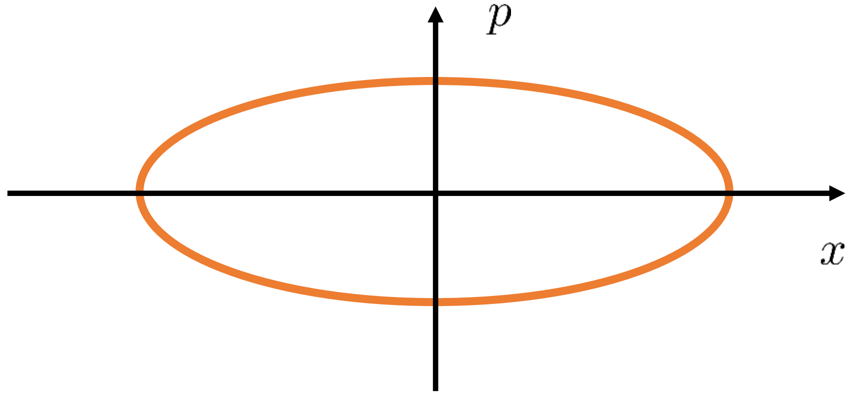
Motivation

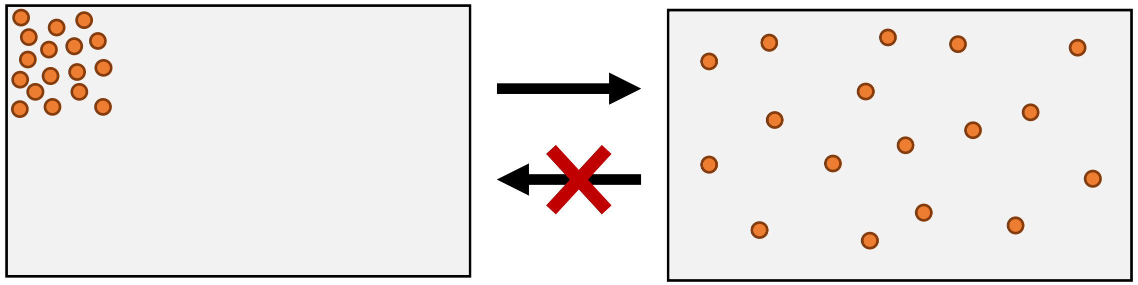
What can we say about the dynamics without solving equations of motion?

Closed systems



Energy conservation

Open systems



Entropy growth

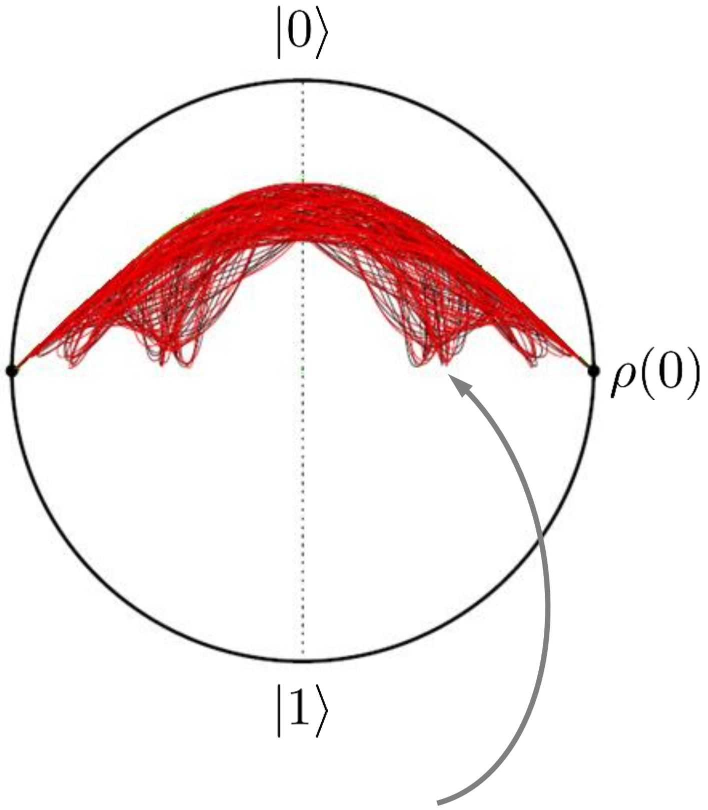
Quantum thermodynamics:

Using minimal assumptions of the quantum theory, find constraints on the evolution of a quantum system interacting with thermal baths

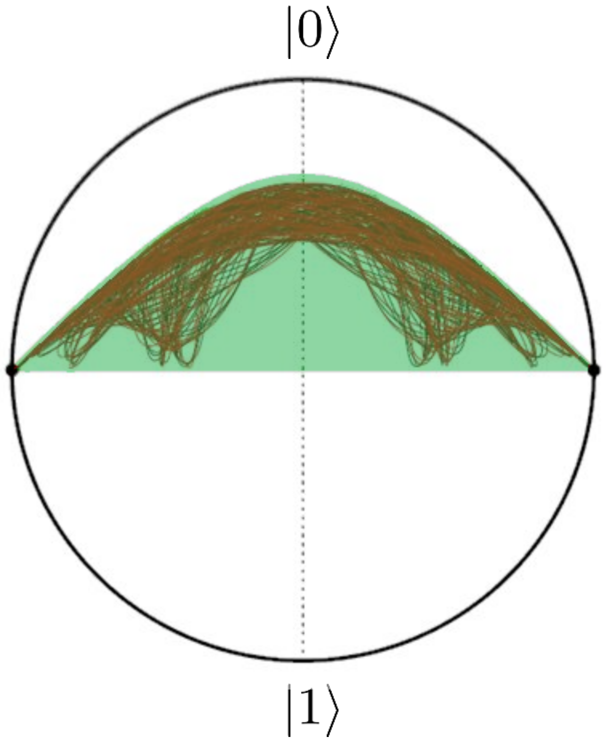
Motivation

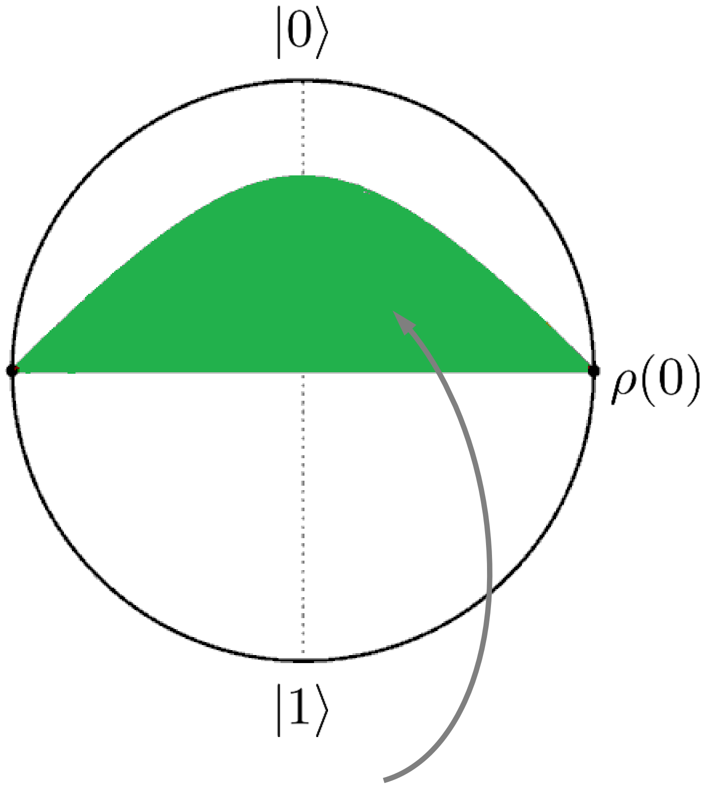
Open dynamics approach:

Resource-theoretic approach:



Exact time evolution
for a given model





Allowed final states
compatible with the laws
of thermodynamics

Thermodynamic setting

Thermodynamic transformations modelled by **thermal operations***:

$$\mathcal{E}^{T}(\cdot) = \mathrm{Tr}_B\left(U\left(\cdot \otimes \gamma_B\right)U^{\dagger}\right) \quad \text{with} \quad [U, H + H_B] = 0$$

Thermal bath γ_B — Energy-conserving interaction U — System ρ → System σ

Hamiltonian: H_B | Hamiltonian: H | Hamiltonian: H

Gibbs state γ of the system at temperature T: $\quad \gamma = e^{-\frac{H}{T}}/\mathcal{Z}, \quad \mathcal{Z} = \mathrm{Tr}\left(e^{-\frac{H}{T}}\right)$

Note: all results with units such that $k_B = 1$.

*M. Horodecki, J. Oppenheim Nature Commun. **4**, 2059 (2013)

Thermodynamic setting

Setting: Initial state ρ, target state σ, background temperature T

Single-shot interconversion: Does there exist $\mathcal{E}^T$ such that $\mathcal{E}^T(\rho) = \sigma$?

Many-copies interconversion: Does there exist $\mathcal{E}^T$ such that $\mathcal{E}^T\left(\rho^{\otimes n}\right) \approx_\epsilon \sigma^{\otimes R_n n}$?

Optimal rate R_n for error ϵ?

Note: $\sigma \approx_\epsilon \tilde{\sigma}$ means $1 - F\left(\sigma, \tilde{\sigma}\right) \leq \epsilon$ with fidelity F

Restrictions:

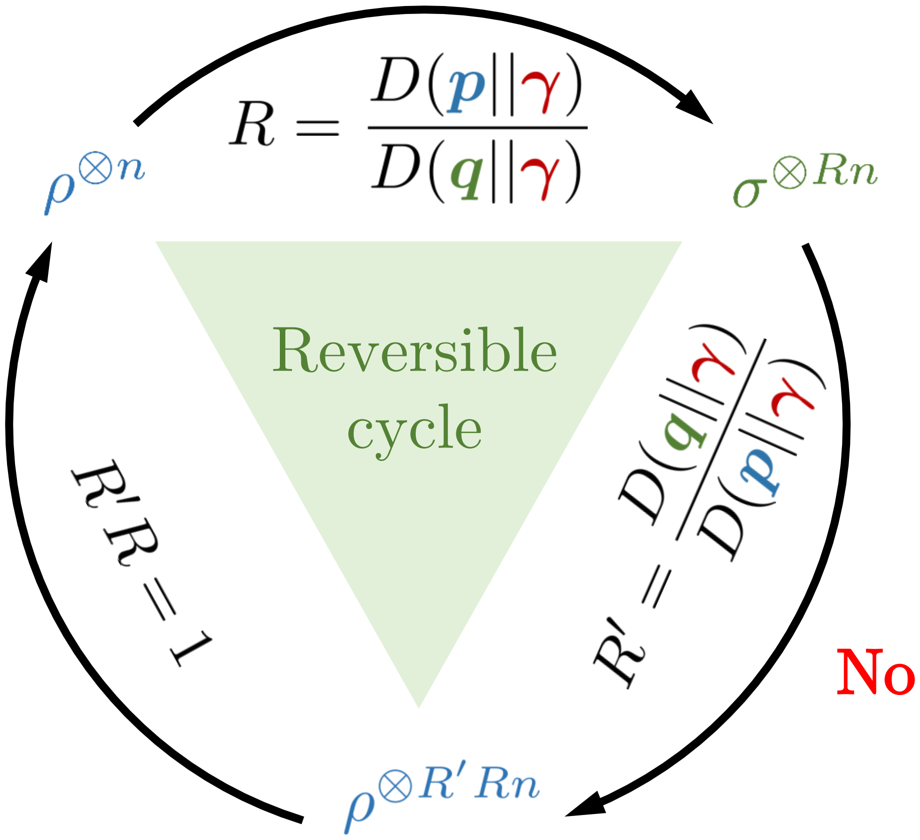
Focus on many copies (large but finite n) and *energy-incoherent* states:

$$[\rho, H] = [\sigma, H] = 0 \quad \Longrightarrow \quad \text{states represented by: } \boldsymbol{p} = \mathrm{eig}(\rho),\ \boldsymbol{q} = \mathrm{eig}(\sigma).$$

$$[\gamma, H] = 0 \quad \Longrightarrow \quad \text{thermal state represented by: } \boldsymbol{\gamma} = \mathrm{eig}(\gamma)$$

Asymptotic reversibility

Asymptotic rate for $n \to \infty$*: $R_\infty(\boldsymbol{p} \to \boldsymbol{q}) = \dfrac{D(\boldsymbol{p}\|\boldsymbol{\gamma})}{D(\boldsymbol{q}\|\boldsymbol{\gamma})}$



Relative entropy: $D(\boldsymbol{p}\|\boldsymbol{\gamma}) := \sum_{i=1}^{d} p_i \log \dfrac{p_i}{\gamma_i}$

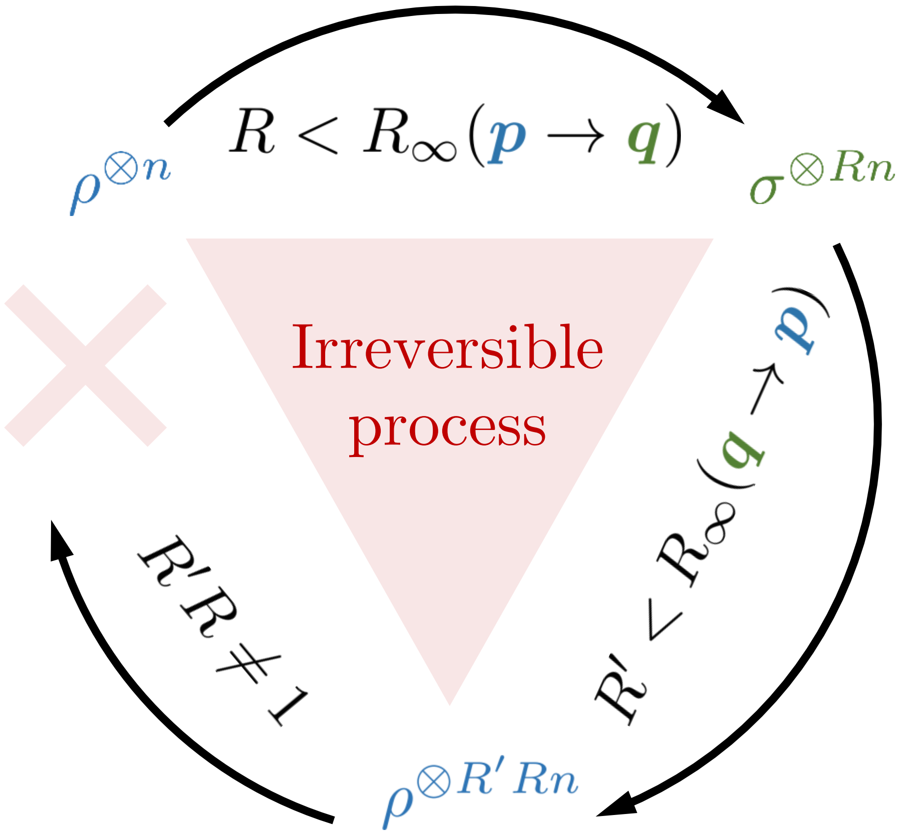
Physical interpretation: $\dfrac{1}{T}\Big[\underbrace{\langle E\rangle_{\boldsymbol{p}} - T H(\boldsymbol{p})}_{\substack{\text{Free energy} \\ F = U - TS}} - \underbrace{(-T\log\mathcal{Z})}_{\text{Free energy of } \gamma}\Big]$

No dissipation of free energy in the thermodynamic limit!

*F. Brandão *et al.*,
Phys. Rev. Lett. 111, 250404 (2013)

Finite-size irreversibility

Rate for large but finite n: $R_n = R_\infty - f(\boldsymbol{p}, \boldsymbol{q}, \boldsymbol{\gamma}, n, \epsilon)$



Relevant quantity quantifying irreversibility:

Relative entropy variance:

$$V(\boldsymbol{p}\|\boldsymbol{\gamma}) := \sum_{i=1}^{d} p_i \left(\log \frac{p_i}{\gamma_i} - D(\boldsymbol{p}\|\boldsymbol{\gamma})\right)^2$$

Physical interpretation:

$$V(\boldsymbol{\gamma}'\|\boldsymbol{\gamma}) = \underbrace{\frac{\partial \langle E \rangle_{\gamma'}}{\partial T'}}_{\text{Specific heat capacity}} \cdot \underbrace{\left(1 - \frac{T'}{T}\right)^2}_{\text{Carnot factor}}$$

Finite-size irreversibility

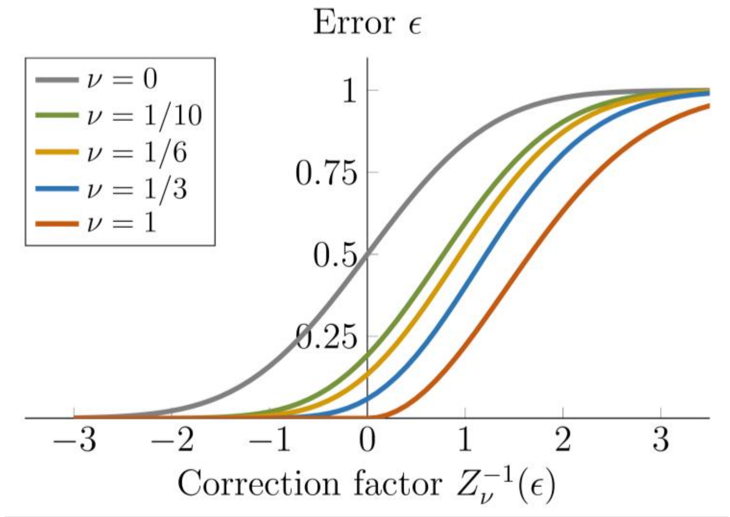
Optimal conversion rate R_n with constant error ϵ:

$$R_n(\epsilon) \simeq R_\infty + \sqrt{\frac{V(\mathbf{p}\|\boldsymbol{\gamma})}{D(\mathbf{q}\|\boldsymbol{\gamma})^2}} \frac{Z_\nu^{-1}(\epsilon)}{\sqrt{n}}$$

Irreversibility parameter:

$$\nu = \frac{V(\mathbf{q}\|\boldsymbol{\gamma})/D(\mathbf{q}\|\boldsymbol{\gamma})}{V(\mathbf{p}\|\boldsymbol{\gamma})/D(\mathbf{p}\|\boldsymbol{\gamma})}$$

Rayleigh-normal distribution Z_ν*:



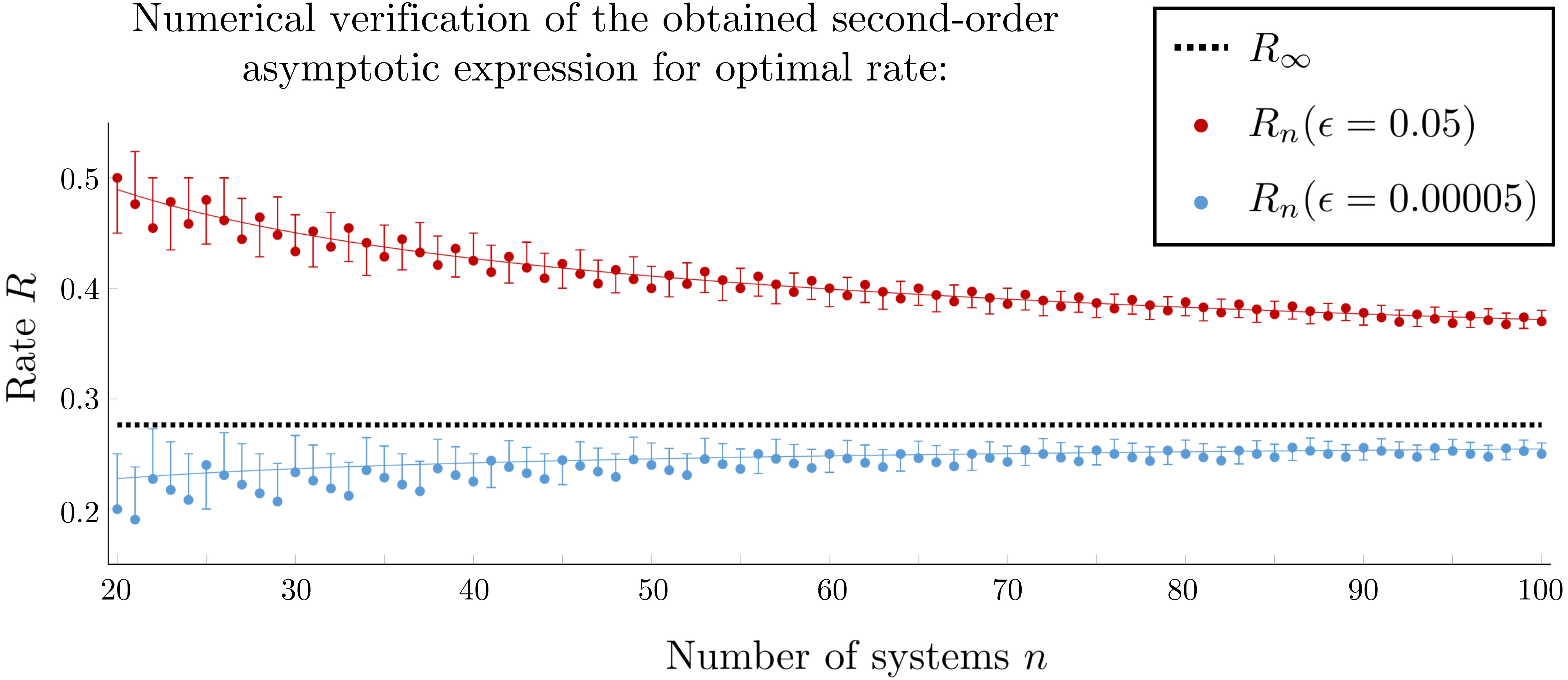
Z_0 - standard normal distribution Φ

Z_1 - Rayleigh distribution ($Z_1(x) = 0$ for $x \leq 0$)

*W. Kumagai *et al.*, IEEE Trans. Inf. Theory **63**, 1829–1857 (2017)

Finite-size irreversibility

Numerical verification of the obtained second-order asymptotic expression for optimal rate:



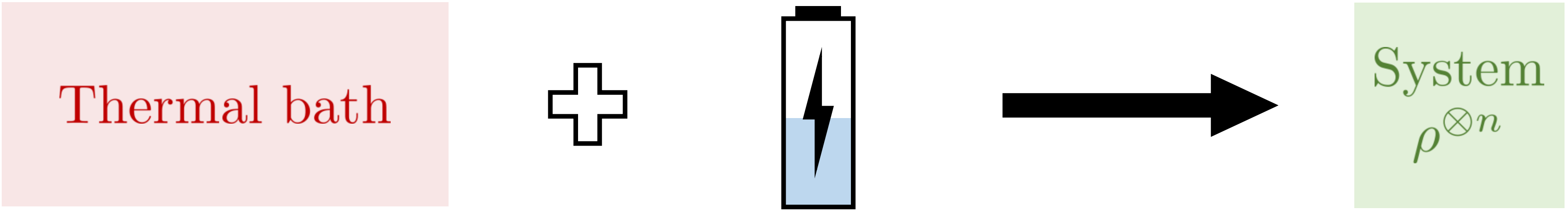
Finite-size irreversibility

Effects of finite-size irreversibility on work distillation and dilution processes:

Work distillation process:

Distillable work: $n \cdot W_D$

Work dilution process:

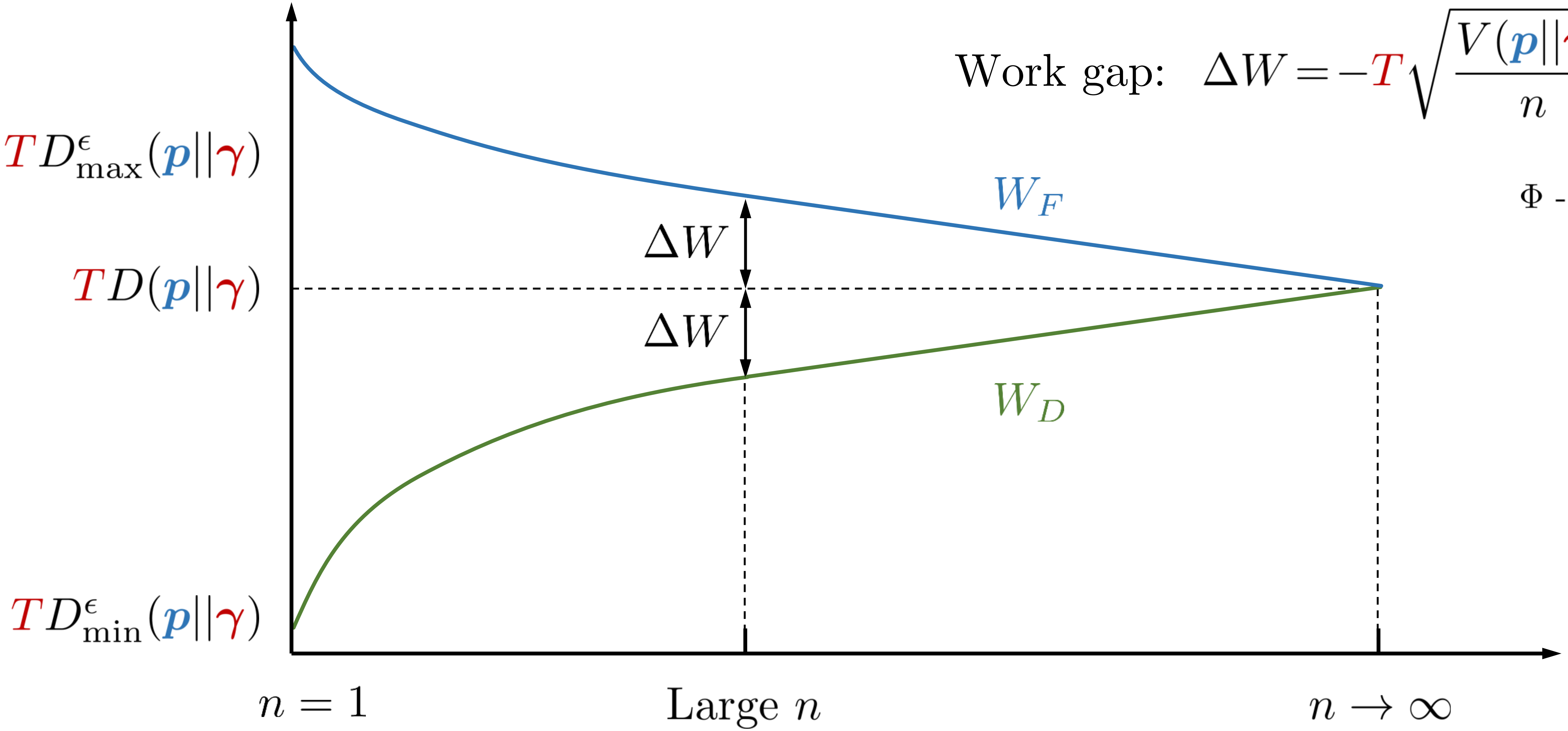


Work of formation: $n \cdot W_F$

Finite-size irreversibility

Work gap: $\Delta W = -T\sqrt{\frac{V(\boldsymbol{p}||\boldsymbol{\gamma})}{n}}\Phi^{-1}(\epsilon)$

Φ - Normal cdf



Dissipation of free energy beyond the thermodynamic limit!

Resource resonance

Optimal conversion rate R_n with vanishing error $\epsilon = e^{-n^\alpha}$ and $\alpha \in (0,1)$:

$$R_n(\epsilon) \simeq R_\infty - \sqrt{\frac{V(\boldsymbol{p}\|\gamma)}{D(\boldsymbol{q}\|\gamma)^2}}\,\frac{\left|\sqrt{1/\nu}-1\right|}{\sqrt{n^{1-\alpha}}}$$

When $\nu = 1$ correction term disappears for every error ϵ

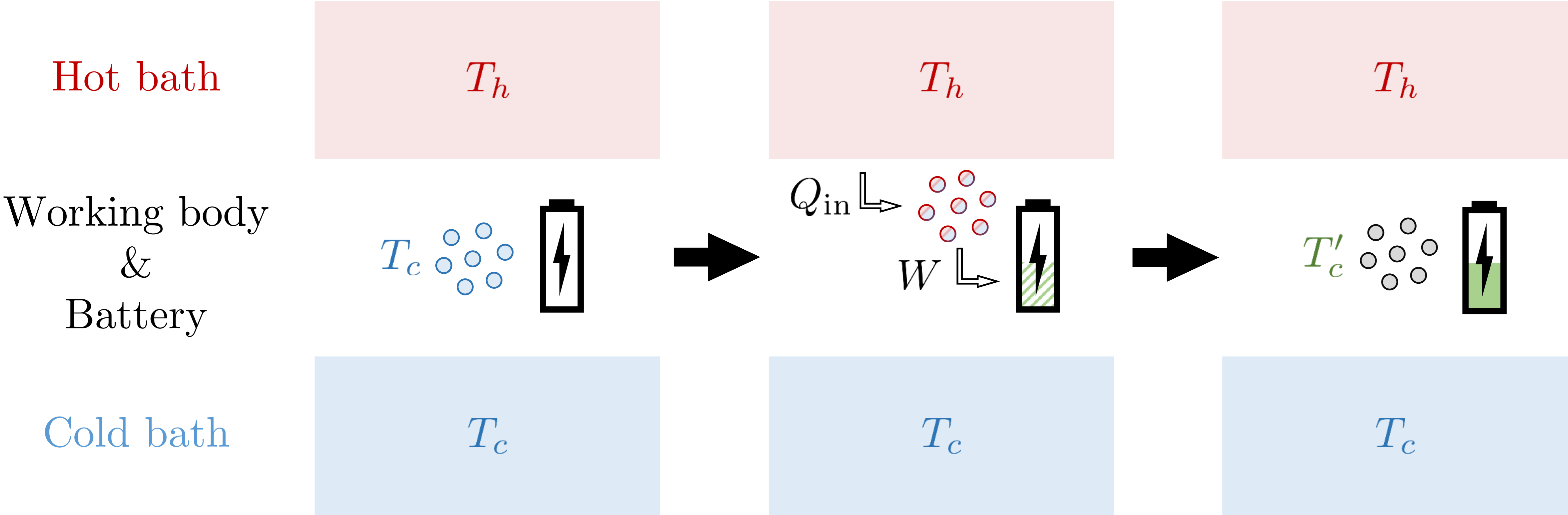
$$\Downarrow$$

No free energy dissipation!
(at least up to second order asymptotics)

Recall that $\nu = 1$ means that the relative fluctuations of free energy are the same for the initial state ρ and target state σ

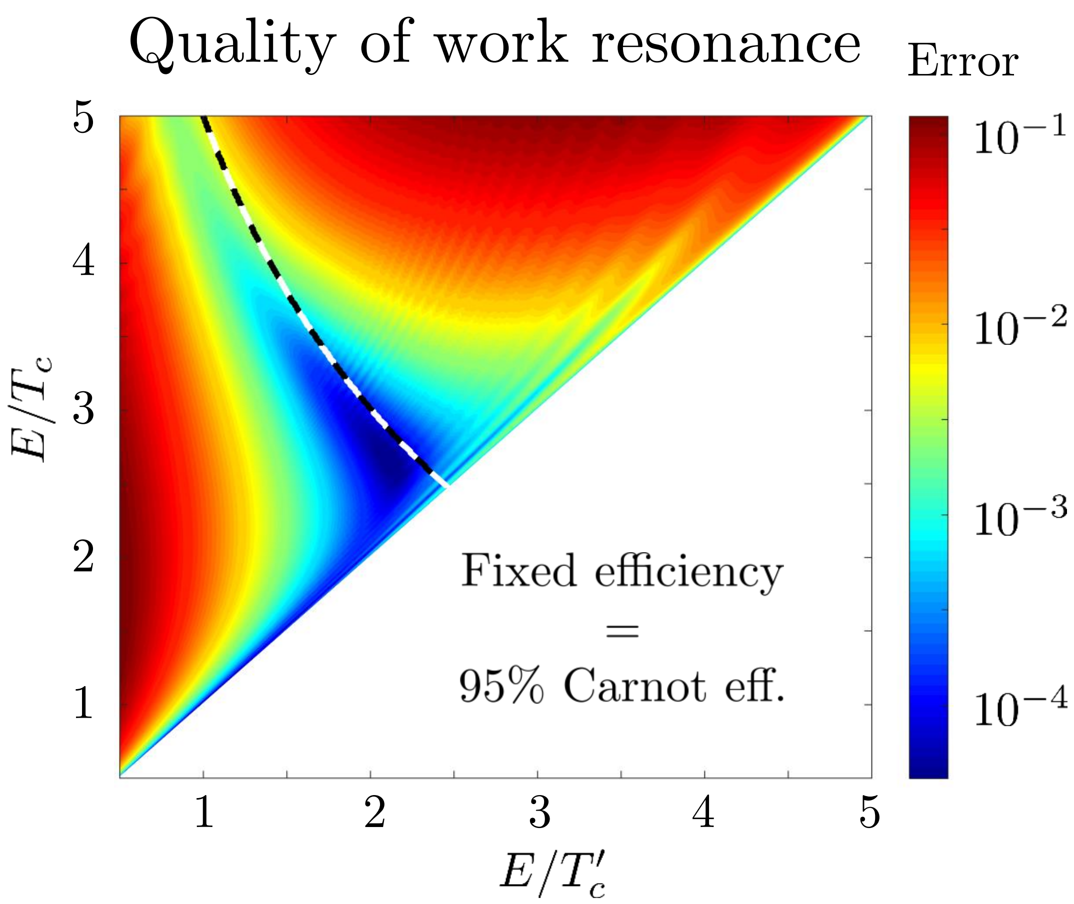
Resource resonance

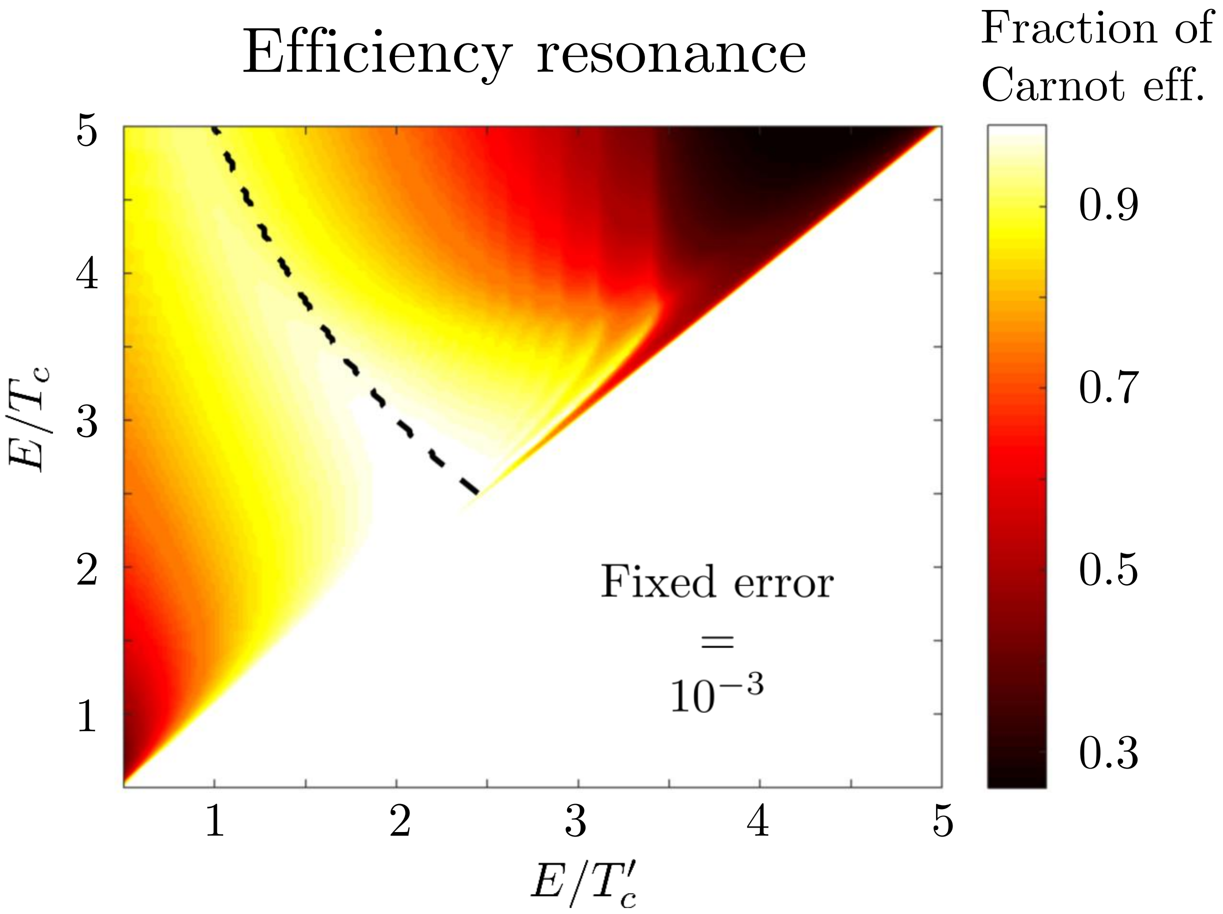
Resonance example: Heat engine with a finite-size working body:



Resource resonance

Working body: $n = 200$ qubits, energy gap E
Background (hot) bath: $T_h = 10E$

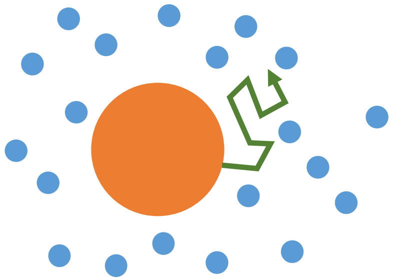




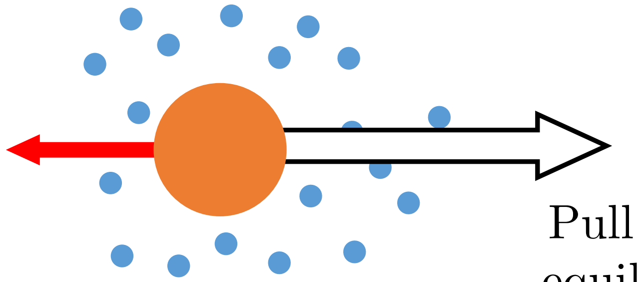
Fluctuation-dissipation relations

Einstein-Smoluchowski relation for a Brownian particle:

In thermal equilibrium



Position fluctuations $\sim$ Dissipative friction



Pull out of equilibrium

Thermodynamic distillation process

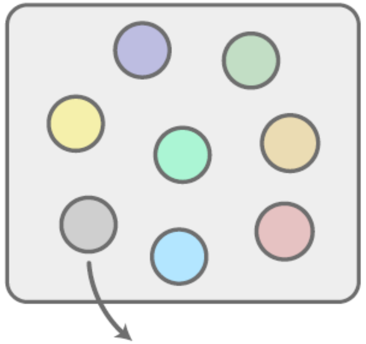
Non-zero free energy:

$$F^N := \frac{1}{\beta}\sum_{n=1}^{N} D(\rho_n^N \| \gamma_n^N)$$

Non-zero free energy fluctuations:

$$\sigma^2(F^N) := \frac{1}{\beta^2}\sum_{n=1}^{N} V(\rho_n^N \| \gamma_n^N)$$

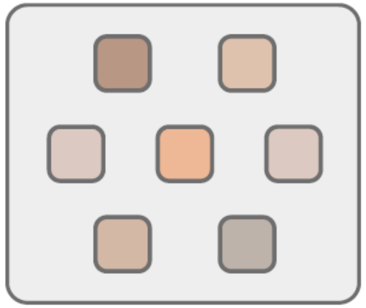
(ρ^N, H^N)



(ρ_n^N, H_n^N)

$\xrightarrow{\mathcal{E}^T}$

$(\tilde{\rho}^N, \tilde{H}^N)$



$\tilde{\rho}^N = |\tilde{E}_k^N\rangle\langle\tilde{E}_k^N|$

Non-zero free energy, but vanishing free energy fluctuations

Free energy fluctuations $\overset{?}{\sim}$ Free energy dissipated in the process

Fluctuation-dissipation relations

Optimal error in thermodynamic distillation process:

$$\lim_{N\to\infty} \epsilon_N = \lim_{N\to\infty} \Phi\left(-\frac{\Delta F^N}{\sigma(F^N)}\right)$$

ΔF^N - Free energy difference between initial and **target** state

Minimal amount of free energy dissipated in the optimal distillation process:

$$\boxed{F^N_{\text{diss}} \simeq a(\epsilon_N)\sigma(F^N)}$$

F^N_{diss} - Free energy difference between initial and **final** state

$$a(\epsilon) = -\Phi^{-1}(\epsilon)(1-\epsilon) + \exp(-[\Phi^{-1}(\epsilon)]^2/2)/\sqrt{2\pi}$$

Three regimes:

$$\lim_{N\to\infty} \frac{\Delta F^N}{\sqrt{N}} = \begin{cases} \infty, & \Longrightarrow \quad \epsilon = 0, \quad F^N_{\text{diss}} = \Delta F^N \\ -\infty, & \Longrightarrow \quad \epsilon = 1, \quad F^N_{\text{diss}} = 0 \\ \alpha \in \mathbb{R} & \end{cases}$$

Also holds for initial pure states with coherence!

Outlook

- Extend finite-size analysis to other resource-theories (asymmetry, contextuality).
- Design experimental protocols employing the resonance phenomenon.
- Generalise the formalism to include quantum states with coherence.
- Look for resonance phenomena in other quantum information processing tasks.
- Extend resource-theoretic fluctuation-dissipation theorem to continuous variable systems

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Thank you!