

# Quantum coherence, time-translation symmetry and thermodynamics

Kamil Korzekwa, Matteo Lostaglio, David Jennings, Terry Rudolph

Department of Physics, Imperial College London, London SW7 2AZ, United Kingdom

Imperial College  
London



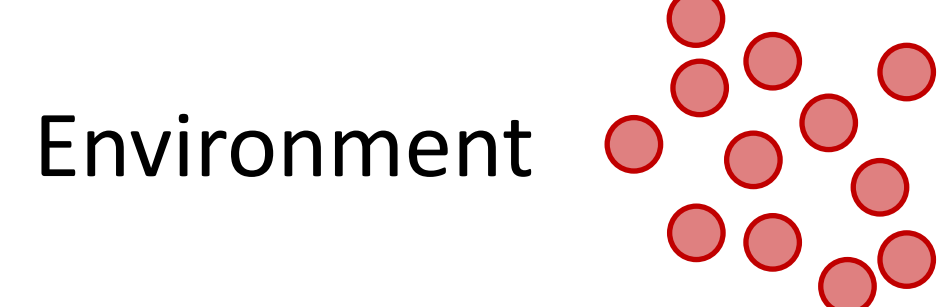
EPSRC

## 1. Thermodynamic setting



System

Arbitrary state:  $\rho_S$   
Hamiltonian:  $H_S$



Environment

Thermal State:  $\gamma_E \propto e^{-\beta H_E}$   
Hamiltonian:  $H_E$

Joint energy-conserving unitary evolution

$$U(\text{System} \otimes \text{Environment})U^\dagger \quad [U, H_S + H_E] = 0$$

Hence the evolution of the system is described by *thermal operations*:

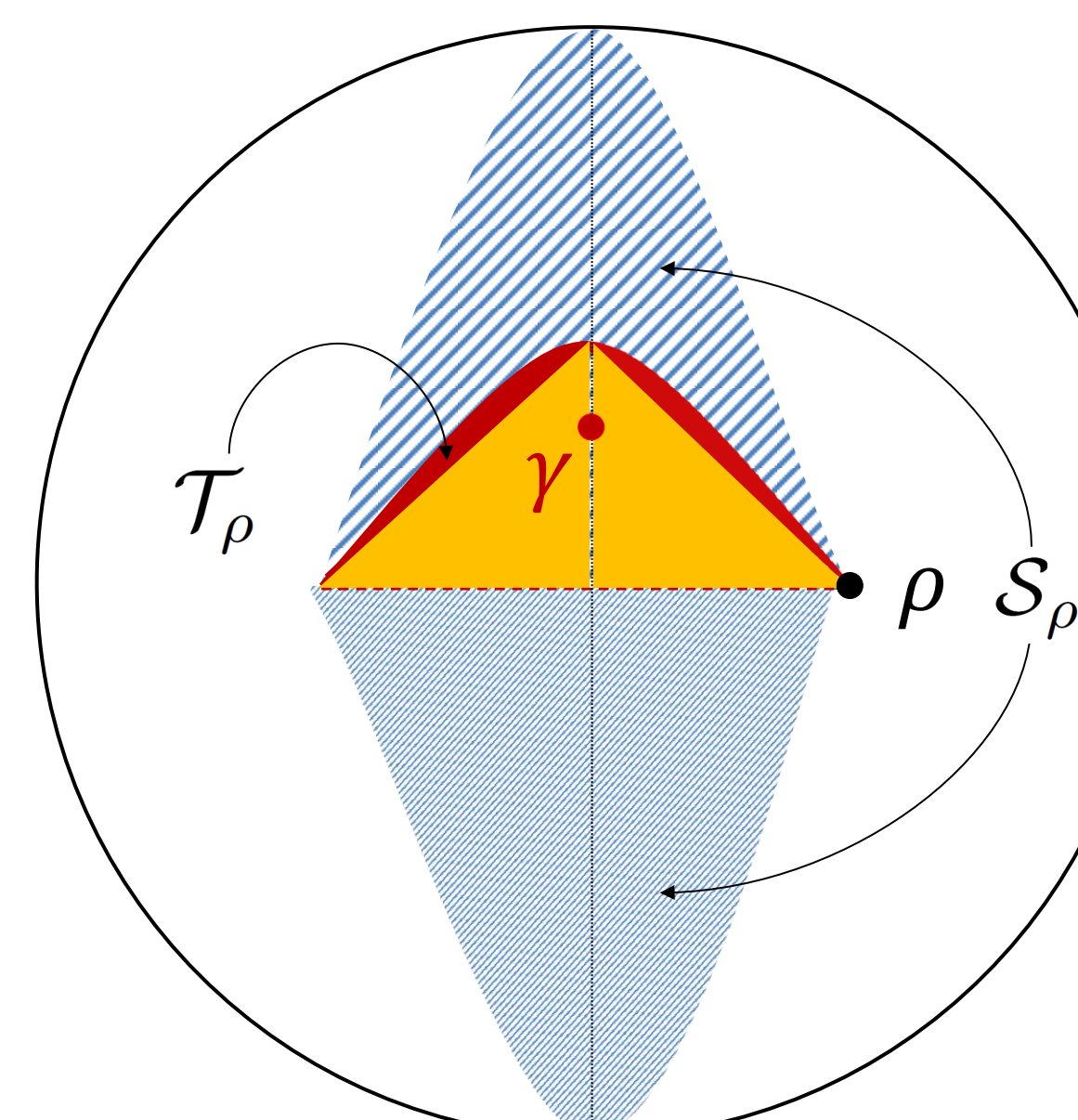
$$\mathcal{E}_T(\rho_S) = \text{Tr}_E(U(\rho_S \otimes \gamma_E)U^\dagger),$$

that form a subset of time-translation symmetric operations<sup>1</sup>:

$$\mathcal{E}_T(e^{-iH_S t} \rho_S \rho_S e^{iH_S t}) = e^{-iH_S t} \mathcal{E}_T(\rho_S) e^{iH_S t}$$

## 2. Thermal transformations of states with coherence – elementary scenario

$$H_S = |1\rangle\langle 1|$$



Coherence is **actively** contributing to enlarge the set of thermodynamically accessible states.

Work is **not** the universal resource of thermodynamics.

Coherence contribution to free energy is **locked** - no trivial extension to quantum Szilard engine.

$\mathcal{T}_\rho$ : Set of states accessible from  $\rho$  via thermal operations (orange region if coherence is passive)

$\mathcal{S}_\rho$ : Set of states accessible from  $\rho$  via thermal operations and the access to infinite amount of work

## 3. How to deal with coherences? Modes of coherence

### Definition:

Assuming non-degenerate Hamiltonian

$$H_S = \sum_n \hbar \omega_n |n\rangle\langle n| \quad \rho_S = \sum_{n,m} \rho_{nm} |n\rangle\langle m|$$

The free evolution of the system is given by

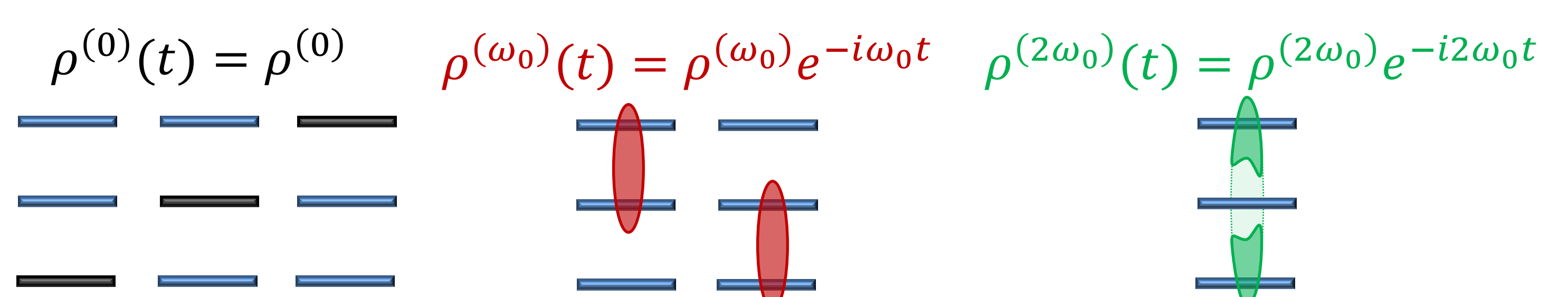
$$\rho_S(t) = e^{-iH_S t} \rho_S e^{iH_S t} = \sum_{n,m} \rho_{nm} |n\rangle\langle m| e^{-i\hbar(\omega_n - \omega_m)t}$$

We can decompose any state into *modes of coherence*<sup>2</sup> - 1-dimensional irreps of the U(1) time-translation group action.

$$\rho = \sum_\omega \rho^{(\omega)} \quad \rho^{(\omega)} := \sum_{\omega_n - \omega_m = \omega} \rho_{nm} |n\rangle\langle m|;$$

### Example:

$$H_S = \sum_{n=0}^2 n \hbar \omega_0 |n\rangle\langle n| \quad \rho = \begin{pmatrix} p_0 & c_{01} & c_{02} \\ c_{10} & p_1 & c_{12} \\ c_{20} & c_{21} & p_2 \end{pmatrix}$$



### Decomposing thermal operations using modes:

Because of time-translation symmetry each mode in the initial state is **independently mapped** by a thermal operation to the corresponding mode of the final state:

Intensity of each mode has to **decrease**:

$$\sigma = \mathcal{E}_T(\rho) \\ \sigma^{(\omega)} = \mathcal{E}_T(\rho^{(\omega)}) \\ \|\sigma^{(\omega)}\| \leq \|\rho^{(\omega)}\|$$

## 4. Bounds on coherence transformation under thermal operations

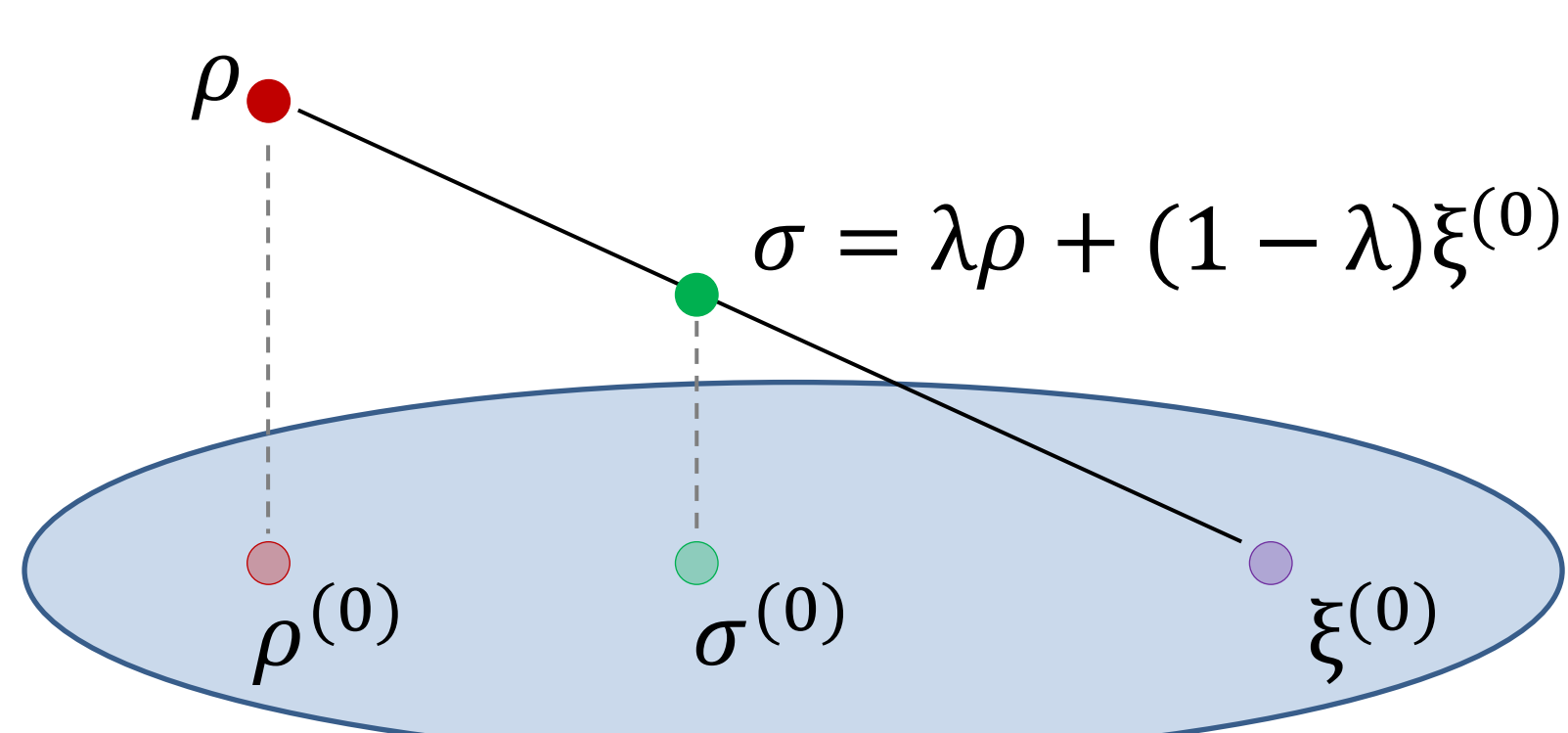
Upper bound for final coherence (based on transition probabilities for diagonal elements):

$$|\rho'_{nm}| \leq \sum_{c,d} |\rho_{cd}| \sqrt{p_{n|c} p_{m|d}}$$

$$p_{n|c} = \langle n | \mathcal{E}_T(|c\rangle\langle c|) | n \rangle$$

Guaranteed lower bound for final coherence (based on thermomajorization condition for incoherent states):

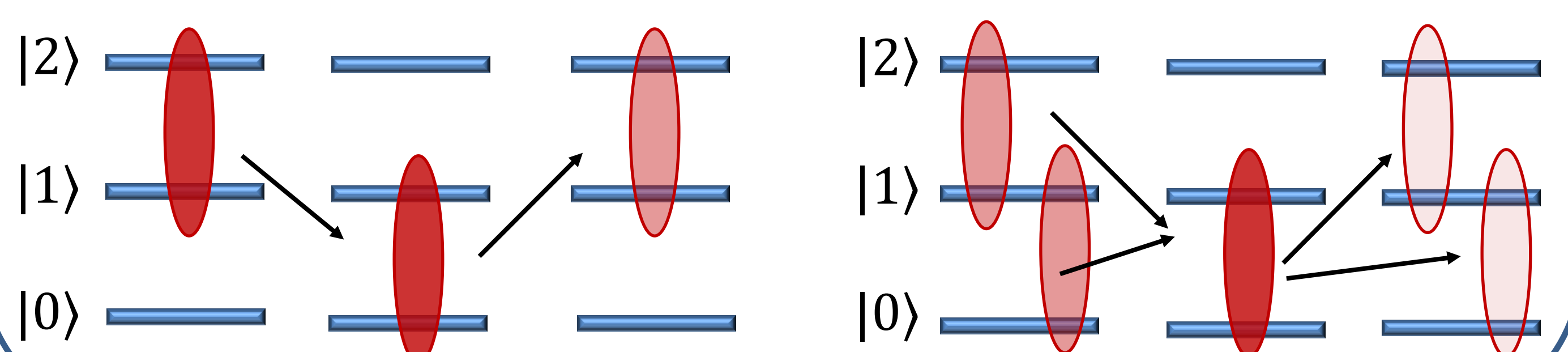
$$|\rho'_{nm}| \geq \lambda^* |\rho'_{nm}|$$



## 5. Irreversibility of coherence transfer

Using the fact that transition probabilities must preserve thermal state one arrives at:

$$|\rho'_{nm}| \leq \sum_{\substack{c,d \\ \omega_c - \omega_d = \omega_n - \omega_m \\ \omega_c > \omega_n}} |\rho_{cd}| + \sum_{\substack{c,d \\ \omega_c - \omega_d = \omega_n - \omega_m \\ \omega_c \leq \omega_n}} |\rho_{cd}| e^{-\beta \hbar(\omega_n - \omega_c)}$$



## 6. Outlook

Recent results on *catalytic coherence*<sup>3</sup> show that coherence, unlike other quantum resources, does not have to degrade while being used to lift time-translation symmetry. This, however, requires investing work. The question that remains open is: can external coherence (reference) be used catalytically to extract work locked in the coherence of the system?

### References:

1. M. Lostaglio, D. Jennings, and T. Rudolph, arXiv:1405.2188 (2014).
2. I. Marvian and R. W. Spekkens, Phys. Rev. A **90**, 062110 (2014).
3. J. Åberg, Phys. Rev. Lett. **113**, 150402 (2014).

### For more details check:

M. Lostaglio, K. Korzekwa, D. Jennings, and T. Rudolph, arXiv:1410.4572 (2014).