Fundamental constraints of quantum thermodynamics in the Markovian regime

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TEAM-NET

Outline

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- 2. Statement of the problem
- 3. Main technical tool
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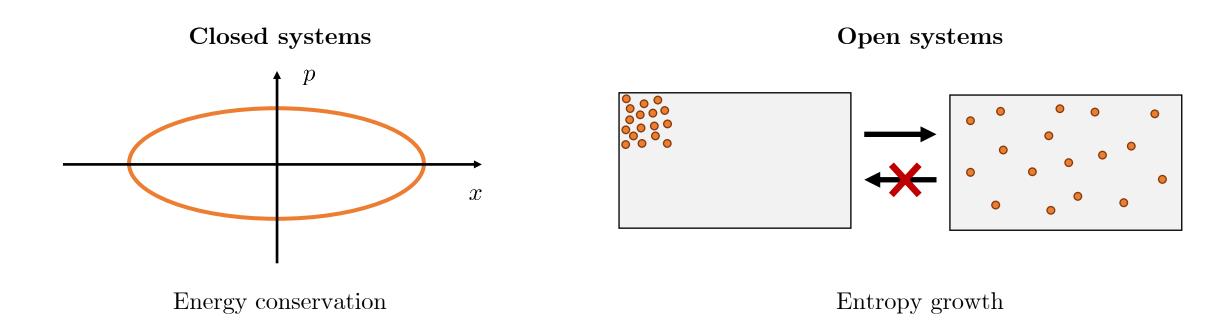
In collaboration with:



 ${\it Matteo \ Lostaglio} \\ {\it QuTech, \ Delft \ University \ of \ Technology} \\$

Motivation

What can we say about the dynamics without solving equations of motion?

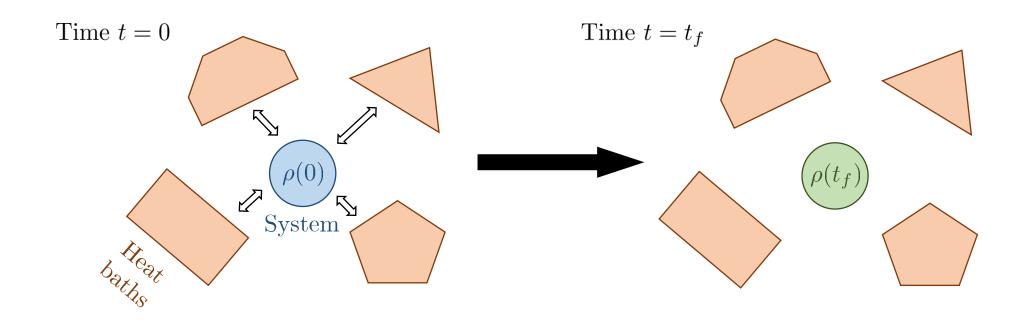


Quantum thermodynamics:

Using minimal assumptions of the quantum theory, find constraints on the evolution of a quantum system interacting with thermal baths.

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Statement of the problem



Original question: Given $\rho(0)$ and \Leftrightarrow denoting arbitrary energy-conserving unitary, what can $\rho(t_f)$ be?

M. Horodecki, J. Oppenheim Nature Commun. 4, 2059 (2013)

Our question: Given $\rho(0)$ and \Leftrightarrow denoting Markovian energy-conserving interaction, what can $\rho(t_f)$ be?

Formal statement of the problem

General Markovian open quantum dynamics:

$$\frac{d\rho(t)}{dt} = -i[H, \rho(t)] + \mathcal{L}_t(\rho(t)), \quad \text{where} \quad H = \sum_{i=1}^d E_i |E_i\rangle\langle E_i| \quad \text{and} \quad \mathcal{L}_t(\rho) = \sum_{i=1}^{d^2 - 1} r_i(t) \left(L_i \rho L_i^{\dagger} - \frac{1}{2} \{ L_i^{\dagger} L_i, \rho \} \right)$$

Markovian thermal process (MTP) defined by additional two properties:

- Stationary thermal state: $\forall t: \mathcal{L}_t \gamma = 0$, where $\gamma = e^{-\beta H}/\text{Tr}\left(e^{-\beta H}\right)$
- Covariance: $\forall t, \rho : \quad \mathcal{L}_t([H, \rho]) = [H, \mathcal{L}_t(\rho)]$

Solving for possible final states:

$$\rho(0) \stackrel{\text{MTP}}{\longmapsto} \rho(t_f)$$

Formal statement of the problem

Covariance encodes total energy conservation at each infinitesimal time and enforces populations and coherences,

$$p_i(t) = \langle E_i | \rho(t) | E_i \rangle, \qquad c_{ij}(t) = \langle E_i | \rho(t) | E_j \rangle$$

to evolve independently.

We can thus analyse a simplified problem:

Solving for possible final energy populations:

$$\boldsymbol{p}(0) \stackrel{ ext{MTP}}{\longmapsto} \boldsymbol{p}(t_f)$$

Main technical tool

Equation of motion restricted to population reads:

$$\frac{d\boldsymbol{p}(t)}{dt} = L_t \boldsymbol{p}(t), \quad \text{with} \quad L_t \boldsymbol{\gamma} = 0 \quad \left(\text{simplified version of } \frac{d\rho(t)}{dt} = -i[H, \rho(t)] + \mathcal{L}_t(\rho(t)) \right)$$

Possible final populations are Tp(0) with T being a stochastic matrix satisfying:

- Stationary thermal state: $T\gamma = \gamma$
- Embeddability: $T = S(t_f)$ with $dS(t)/dt = L_t S(t)$ and S(0) = 1

Solution to the first constraint is known as **thermomajorisation**, i.e.,

$$\exists T \text{ such that } T\gamma = \gamma \text{ and } Tp(0) = p(t_f) \iff p(0) \succ_{\gamma} p(t_f)$$

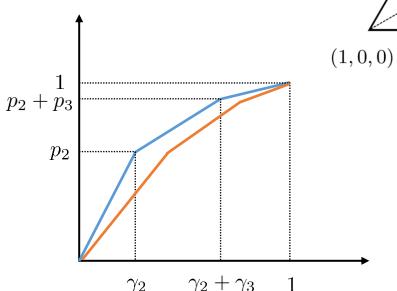
Main technical tool

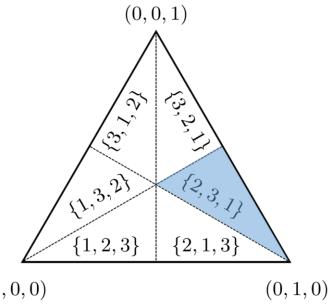
How to verify thermomajorisation condition $p \succ_{\gamma} q$?

- 1. Find the γ -ordering of \boldsymbol{p} , i.e. arrange p_i/γ_i in a non-increasing order \downarrow .
- 2. Create the Lorenz curve of p, i.e. piecewise linear curve with elbow points given by:

$$x_j = \sum_{i=1}^j \gamma_i^{\downarrow}, \qquad y_j = \sum_{i=1}^j p_i^{\downarrow}$$

3. Do the same for q and check whether the Lorenz curve of q lies below that of p.





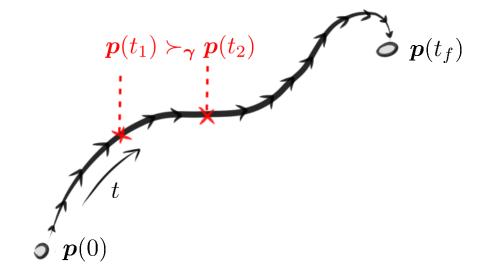
Main technical tool

We introduce **continuous thermomajorisation**:

$$\boldsymbol{p}(0) \gg_{\boldsymbol{\gamma}} \boldsymbol{p}(t_f)$$

iff there exists a thermomajorising trajectory p(t) such that:

$$\forall t_1, t_2 \in [0, t_f]: \quad t_1 \leq t_2 \Rightarrow \boldsymbol{p}(t_1) \succ_{\boldsymbol{\gamma}} \boldsymbol{p}(t_2)$$



Why?

Becuase it yields a complete description of population dynamics:

$$\boldsymbol{p}(0) \stackrel{\text{MTP}}{\longmapsto} \boldsymbol{p}(t_f)$$
 if and only if $\boldsymbol{p}(0) \gg_{\boldsymbol{\gamma}} \boldsymbol{p}(t_f)$

Results

Exhaustive H-type theorem:

A dynamical evolution p(t) of populations can be generated by a Markovian thermal process if and only if:

$$\forall a \in [0,1]: \quad \frac{d\Sigma_a(t)}{dt} \ge 0, \quad \text{where} \quad \Sigma_a := -\sum_{i=1}^d \left| p_i(t) - a \frac{\gamma_i}{\gamma_d} \right|.$$

Universality of elementary thermalisations:

 $p(0) \stackrel{\text{MTP}}{\longmapsto} p(t_f)$ is possible if and only if there exists a sequence of elementary thermalisations such that:

$$\boldsymbol{p}(t_f) = T^{i_f, j_f}(\lambda_f) \dots T^{i_1, j_1}(\lambda_1) \boldsymbol{p}(0), \quad \text{where} \quad T^{i, j}(\lambda) = \begin{bmatrix} (1 - \lambda) + \frac{\lambda \gamma_i}{\gamma_i + \gamma_j} & \lambda \frac{\gamma_i}{\gamma_i + \gamma_j} \\ \lambda \frac{\gamma_j}{\gamma_i + \gamma_j} & (1 - \lambda) + \frac{\lambda \gamma_i}{\gamma_i + \gamma_j} \end{bmatrix} \oplus \mathbf{1}_{\backslash (i, j)}$$

Results

Algorithmic verification of $\mathbf{p}(0) \stackrel{\text{MTP}}{\longmapsto} \mathbf{p}(t_f)$:

- Only finite set of conditions needs to be verified.
- If the path exists, the algorithm returns the Lindbladian realising it.
- One can also find the full set of states achievable via MTP from given.

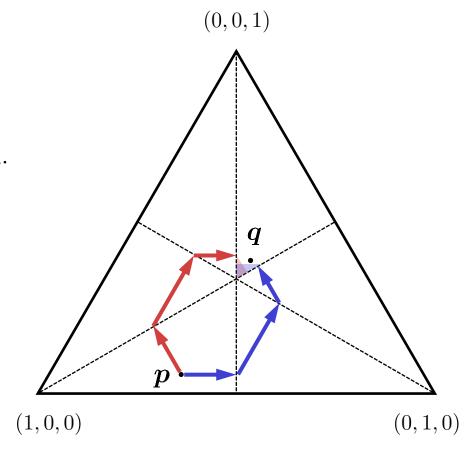
 github.com/KorzekwaKamil/continuous thermomajorisation

How is this done?

Idea 1: when initial and final state have the same ordering, it's easy.

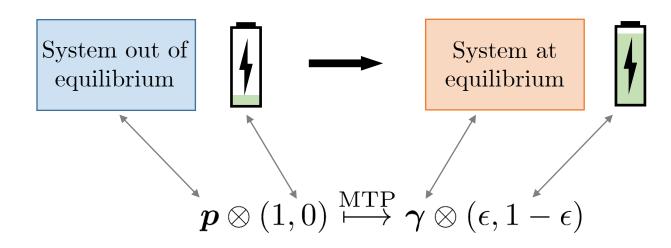
Idea 2: when changing orderings, there is a unique optimal way to do it.

Idea 3: there are finite number of paths connecting different orderings.



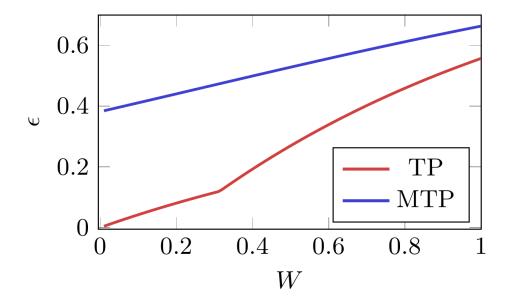
Applications

Role of memory in work extraction:



versus

$$\boldsymbol{p} \otimes (1,0) \stackrel{\mathrm{TP}}{\longmapsto} \boldsymbol{\gamma} \otimes (\epsilon, 1 - \epsilon)$$



System spectrum $\{0, 1\}$ Battery spectrum $\{0, W\}$ System initially thermal with $\beta_S = 2$ Bath with $\beta_E = 1$

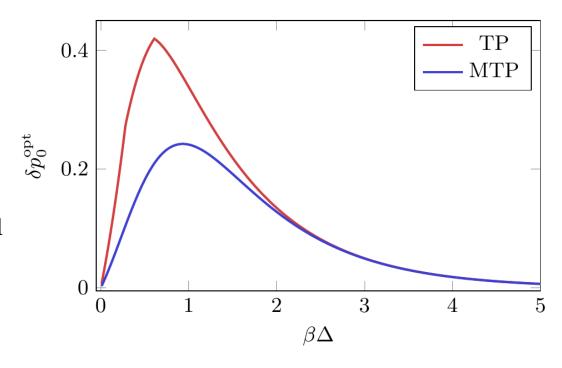
Applications

Role of memory in cooling:

One step of heat-bath algorithmic cooling protocol:

- Take a thermal system and unitarily invert its populations.
- Interact it with the bath and try to maximise ground state population.

Again, we can compare optimal MTP with optimal TP protocols.



System spectrum $\{0, \Delta, 2\Delta, 3\Delta\}$ System initially in equilibrium with bath at β

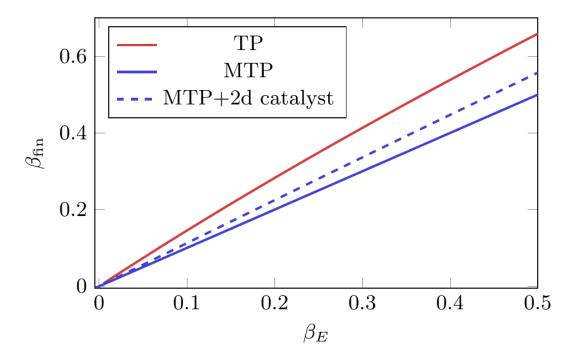
Applications

Catalysts and memory in thermodynamic protocols

Catalyst c is a system that is returned unchanged at the end of the process:

$$oldsymbol{p}\otimesoldsymbol{c}\overset{ ext{MTP}}{\longmapsto}oldsymbol{q}\otimesoldsymbol{c}$$

Thermal catalysts can be used as a memory that enhances or unlocks otherwise impossible tasks to be performed, with catalyst's dimension quantifying the amount of memory.



System and catalyst spectrum $\{0, 1\}$ System initially thermal with $\beta_S = \beta_E/2$ Bath and catalyst thermal with β_E

Outlook

1. Apply to study non-Markovian boosts to relevant processes.

Non-Markovianity boosts the efficiency of bio-molecular switches arXiv:2103.14534 (2021)

- 2. Understand the asymptotic behaviour of continuous thermomajorisation.
- 3. Extend the formalism to treat states with coherence.
- 4. Optimise the runtime of the algorithmic verification procedure.

More soon on arXiv:

Continuous thermomajorisation and a complete set of laws for Markovian thermal processes

M. Lostaglio, K. Korzekwa arXiv:2110.xxxxx (2021)

Thank you!