Relativistic Limitations of Quantum Mechanics

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Table of Contents

- Introduction
- 2 QM & Relativity
 - Is Quantum Mechanics relativistic?
 - Can we *make* it relativistic?
 - The Klein-Gordon Equation
 - The Dirac Equation
- 3 Where to from here?

Motivation & History

- Relativity and modern Quantum Mechanics were developed more or less simultaneously
- ② Historic attempts were made to integrate them, both fell short in several regards
- We will explore the ways physicists of yonderyear attempted to integrate them
- We will ponder the question; can the two be fully unified?

A quick aside about notation

I will be making use of Einstein Summation Notation:

$$\underbrace{\sum_{i=1}^{3} a_i b^i = a_i b^i}_{\text{Latin indices}} \qquad \underbrace{\sum_{\mu=0}^{3} a_\mu b^\mu = a_\mu b^\mu}_{\text{Greek indices}}$$

My metric signature will be (+--). Operators will typically be boldface and have carets, e.g., $\hat{\mathbf{p}}$ for momentum.

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Is Quantum Mechanics relativistic?

For a theory to be relativistic, it has to be *Lorentz invariant*, i.e., invariant under the transformations

$$x' = \gamma(x - vt)$$
 $t' = \gamma \left(t - \frac{v}{c^2}x\right)$

Is the T.D.S.E:

$$-\frac{\hbar^2}{2m}\frac{\partial^2}{\partial x^2}\psi(x,t) + V(x)\psi(x,t) = i\hbar\frac{\partial}{\partial t}\psi(x,t)$$

Lorentz Invariant?

The relativistic Hamiltonian is

$$H=\sqrt{p^2c^2+m^2c^4}$$

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$$H=\sqrt{p^2c^2+m^2c^4}$$

Can we work with this? **Dirac couldn't.**

You would have $\sqrt{\nabla}$ which would involve expansion in a power series, which is not simple.

Let's work with the *square* of that Hamiltonian, in operator form:

$$\hat{\mathbf{H}}^2 = \hat{\mathbf{p}}^2 c^2 + m^2 c^4$$

$$\implies ((-i\hbar \nabla)^2 c^2 + m^2 c^4) \psi = \left(i\hbar \frac{\partial}{\partial t}\right)^2 \psi$$

The Klein-Gordon Equation

$$(\Box + \mu^2)\psi = 0$$

Think of the D'Alembertian as the four-dimensional analogue of the Laplacian.

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In natural units, where $\hbar=c=1,\,\mu=m$ and hence the Klein-Gordon equation becomes

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In natural units, where $\hbar = c = 1$, $\mu = m$ and hence the Klein-Gordon equation becomes

$$(\Box + m^2)\psi = 0$$

Or for a massless particle

$$\Box \psi = 0$$



The Dirac Equation

$$(i\gamma^{\mu}\partial_{\mu} - m)\psi = 0$$

with

$$\gamma^1 = \begin{pmatrix} \mathbb{I}_2 & \mathbf{0}_2 \\ \mathbf{0}_2 & -\mathbb{I}_2 \end{pmatrix} \qquad \gamma^\mu = \begin{pmatrix} \mathbf{0}_2 & \sigma^\mu \\ -\sigma^\mu & \mathbf{0}_2 \end{pmatrix} \qquad \gamma^5 = \begin{pmatrix} \mathbf{0}_2 & \mathbb{I}_2 \\ \mathbb{I}_2 & \mathbf{0}_2 \end{pmatrix}$$

and

$$\partial_{\mu} = \begin{pmatrix} \frac{\partial}{\partial t} & \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \end{pmatrix}$$

This is a relativistic Schrödinger equation!



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Mankind has uncovered two extremely efficient theories: one that describes our universe's structure (Einstein's gravity: the theory of general relativity), and one that describes everything our universe contains (quantum field theory), and these two theories won't talk to each other.

- Christophe Galfard

Conformal theory of everything

F. F. Faria.* Centro de Cièncias da Natureza, Universidade Estadual do Piauí, 64002-150 Teresina, Pl. Brazil

Abstract

By using conformal symmetry we unify the standard model of perticle physics with gravity in a consistent quantum field theory which describes all the fundamental particles and forces of nature.

Conformal theory of everything

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IS MADE OF THIS WINGSTAM STRINGS.

OLKEY, VANCT WINGSTAM STRINGS.

TOUTHO.

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STRING THEORY GUMMARIZED:

I TUT HAD AN AND-SOME DEAD.
SHOPPIGE AUTHORS AND DISERBY
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OKAY, MART WOULD
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