

Entropy and thermodynamic potentials — 11.03, 2019

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1 Overview

In this lecture we continued talking about entropy. Then we started discussing thermodynamic potentials.

2 Reminding

Let's remind us first law of thermodynamics:

$$dU = \partial Q + \partial W - \partial Q - p dV \quad (1)$$

Where for magnets we have

$$dU = \partial Q - M dH \quad (2)$$

where M - magnetisation.

We also defined Boltzmann entropy as

$$S_B = k_B \ln \Omega \quad (3)$$

3 Boltzmann entropy

Let's assume we have constant temperature $T = \text{const.}$ Then if we divide our volume V into small parts V_0 , we have $\Omega^{(N)} = (\frac{V}{V_0})^N$ and then entropy depends on V and equals

$$S_B(V) = k_B \cdot N \ln \frac{V}{V_0} \quad (4)$$

However, if we remove some volume dV from our box, we get

$$\begin{aligned} S_B(V + dV) &= k_B \cdot N \ln \frac{V + dV}{V_0} \\ \partial S_B &= S_B(V + dV) - S_B(V) = k_B N \left\{ \ln \frac{V + dV}{V_0} - \ln \frac{V}{V_0} \right\} = k_B N \ln \left(\frac{V + dV}{V} \right) \quad (5) \\ \partial S_B &= k_B N \ln \left(1 + \frac{dV}{V} \right) \approx k_B N \frac{dV}{V}, \text{ since } \ln(1 + x) \approx x \text{ if } |x| \ll 1, \frac{dV}{V} \ll 1 \end{aligned}$$

3.1 Thermodynamic definition of change of entropy

Let's recall

$$\begin{aligned}
 pV &= nRT \quad \rightarrow \quad \frac{k_B N}{V} = \frac{p}{T} \\
 \partial S_B &= \frac{pdV}{T}, \quad T = \text{const} \quad \rightarrow \quad U_{\text{ideal gas}} = U(T) \text{ or } U(V) \text{ or } U(p) \quad \rightarrow \quad U_{\text{ideal gas}} = \text{const} \quad (6) \\
 dU &= 0 \quad \rightarrow \quad 0 = \partial Q - pdV \quad \rightarrow \quad \partial S_B = \frac{(\partial Q)_{\text{rev}}}{T}
 \end{aligned}$$

So we have our definition of thermodynamic change of entropy:

$$\partial S = \frac{(\partial Q)_{\text{rev}}}{T} \quad (7)$$

3.2 Examples

- On the graphic we can see a process with constant temperature. We have $pV = nRT = \text{const}$. Then

$$W = \int dV = \int_{V_1}^{V_2} p dV = \int_{V_1}^{V_2} \frac{nRT}{V} dV = nRT(\ln V_1 - \ln V_2) = nRT \ln \left(\frac{V_1}{V_2} \right) \quad (8)$$

So we have

$$Q = W = nRT \ln \left(\frac{V_1}{V_2} \right) \quad (9)$$

And then

$$\Delta S = \int \frac{(\partial Q)_{\text{rev}}}{T} = \frac{1}{T} \int (\partial Q)_{\text{rev}} = \frac{Q}{T} \quad (10)$$

so finally

$$\Delta S = nR \ln \left(\frac{V_1}{V_2} \right) \quad (11)$$

then if $V_2 > V_1 \rightarrow \Delta S > 0$.

- Another process we have shown on graphic 2. In this process $\Delta S = kR \ln \left(\frac{V_1}{V_2} \right)$ is the same for each way, so it is not dependent on path. Then S is a function of state.
So **S is thermodynamic potential**.

4 Thermodynamic potentials

As we already know

$$dU = TdS - pdV \quad \rightarrow \quad U = U(V, S)$$

where V, S are "neutral" variables. So U is another thermodynamic potential.

Definition 1. Free energy is defined as

$$F = U - TS$$

Definition 2. *Gibbs free energy is defined as*

$$G = F + pV$$

All of them are functions of state, so they are thermodynamic potentials. They are "equivalent" in a way - which means that if you have one of them, you can derive others.

4.1 Example 1

$$\begin{aligned}
 U &= U(v, S) \text{ - given, } F = ? \\
 dU &= TdS - p dV, \quad \text{let } V = \text{const} \rightarrow dV = 0 \\
 T &= \frac{dU}{dS} = \left(\frac{dU}{dS} \right)_V, \quad V = \text{const} \\
 p &= - \left(\frac{dU}{dV} \right)_S, \quad \text{so from that we have} \\
 p &= p(V, S), \quad \text{and } T = T(V, S). \quad \text{Then} \\
 F &= U - TS = U - \left(\frac{dU}{dS} \right)_V \cdot S.
 \end{aligned} \tag{12}$$

So finally we have

$$F = U - \left(\frac{dU}{dS} \right)_V \cdot S \tag{13}$$

where U is given.

4.2 Example 2

Free energy $F = F(?)$

$$\begin{aligned}
 F &= U - TS \rightarrow dF = dU - d(TS) = T dS - p dV - (T dS + dT \cdot S) \\
 dF &= -S dT - p dV \rightarrow F = F(T, V)
 \end{aligned} \tag{14}$$

So finally we have

$$F = F(T, V)$$

How to find U ?

$$\begin{aligned}
 U &= F + TS, \quad S = - \left(\frac{\partial F}{\partial T} \right)_V, \quad p = - \left(\frac{\partial F}{\partial V} \right)_T \\
 U &= F - T \cdot \left(\frac{\partial F}{\partial T} \right)_V = -T^2 \left(\frac{\partial}{\partial T} \frac{F}{T} \right)_V
 \end{aligned} \tag{15}$$