

CONE ANGLES AND LCTS

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ABSTRACT. We add details to the proof of [1, Theorem 8].

Let (X, ω) be a compact Kähler manifold of dimension n . Let $f : X \rightarrow \Delta^*$ be a holomorphic submersion over the punctured unit disk $\Delta^* = \Delta \setminus \{0\} \subset \mathbb{C}$. Let $K_{X/\Delta^*} := K_X \otimes f^* K_{\Delta^*}^{-1}$ denote the relative canonical bundle. We will assume $K_{X/\Delta^*} \simeq \mathcal{O}_X$, i.e., the fibers $X_s := f^{-1}(s)$ are Calabi–Yau for all $s \in \Delta^*$. A trivializing section $\Omega \in H^0(X, K_{X/\Delta^*})$ yields a family of trivializing sections $\Omega_s := \Omega|_{X_s}$ of K_{X_s} . In particular, we have a smooth family of volume forms $\mu_s := (\sqrt{-1})^{(n-1)^2} \Omega_s \wedge \overline{\Omega}_s$.

The following theorem, due to Yashan Zhang [1], shows that

$$f_* \mu = \int_{X_s} \mu_s$$

is of conical-type, modulo some logarithmic poles. That is, if we let ω_{cone} be the conical Kähler metric

$$\omega_{\text{cone}} = \sqrt{-1} \frac{ds \wedge d\bar{s}}{|s|^{2(1-\beta)}},$$

for some $\beta \in (0, 1] \cap \mathbb{Q}$, we show that

$$f_* \mu \sim (-\log |s|)^{N-1} \omega_{\text{cone}},$$

for some $N \in \mathbb{N}$.

Remark. In the course of the proof, we will see that the cone angle β of the metric ω_{cone} is given by the log canonical threshold of the central fiber X_0 .

Theorem. Let X be a compact Kähler manifold of dimension n with μ as above. Let Δ denote the unit disk in \mathbb{C} with coordinate s . Let $f : X \rightarrow \Delta$ be a family of Calabi–Yau manifolds with singular fiber $X_0 = f^{-1}(0)$.¹ There exists $\beta \in \mathbb{Q} \cap (0, 1]$, $N \in \mathbb{N}$, and $C \geq 1$ such that

$$C^{-1}(-\log |s|)^{N-1} \omega_{\text{cone}} \leq f_* \mu \leq C(-\log |s|)^{N-1} \omega_{\text{cone}}.$$

¹That is, $f : X \rightarrow \Delta$ is a proper surjective holomorphic map of relative dimension $n - 1$ with connected fibers whose generic fiber is Calabi–Yau.

In particular, since the canonical ω_{can} on the base solves the Monge–Ampère equation

$$\omega_{\text{can}}^\kappa = (\omega_0 + \sqrt{-1}\partial\bar{\partial}\varphi)^\kappa = \frac{f_*\mu}{V_0\binom{n}{\kappa}} e^\varphi \sim f_*\mu,$$

this determines the asymptotics of the canonical metric when the base is a curve.

Proof. Let $\pi : \tilde{X} \rightarrow X$ be a resolution of X_0 such that $\tilde{X}_0 = \pi^{-1}(X_0)$ has simple normal crossings, i.e.,

$$\tilde{X}_0 = \sum_j a_j E_j,$$

where E_j are irreducible divisors with defining sections σ_j . By adjunction, we know that

$$K_{\tilde{X}} = \pi^* K_X + \sum_j k_j E_j.$$

In particular, we can write

$$\pi^*\mu = \prod_j |\sigma_j|^{2k_j} \tilde{\mu},$$

where $\tilde{\mu}$ is a smooth nondegenerate volume form on \tilde{X} . We want to compute the asymptotics of $f_*\mu$ near the central fiber. It suffices to calculate the asymptotics of

$$\int_{\tilde{X}_s} \pi^*\mu,$$

as s approaches 0, where $\tilde{X}_s := \pi^{-1}(X_s)$.

Fix a point $x_0 \in \tilde{X}_0$ and let $J(x_0) := \{j : x_0 \in E_j\}$. In a chart $U_0 \ni x_0$ with coordinates (z_1, \dots, z_n) , the map $\tilde{f} := f \circ \pi : \tilde{X} \rightarrow \Delta$ can be expressed as

$$f(z_1, \dots, z_n) = \prod_{j \in J(x_0)} z_j^{a_j} =: s.$$

Shrinking U_0 if necessary, write $\tilde{\mu} = h(z) dz \wedge d\bar{z}$, where $dz = dz_1 \wedge \dots \wedge dz_n$, and h is a smooth non-vanishing function on U_0 . For simplicity, we will take $h \equiv 1$. Letting $V_s := \tilde{X}_s \cap U_0$, we see that

$$\int_{V_s} \pi^*\mu = \int_{V_s} \prod_{j \in J(x_0)} |z_j|^{2k_j} dz \wedge d\bar{z}.$$

An elementary calculation gives

$$f^* \frac{ds}{s} = \sum_{j \in J(x_0)} \frac{a_j}{z_j} dz_j.$$

Fix $j_0 \in J(x_0)$ to be the index which minimizes the ratio $\frac{k_j+1}{a_j}$, and set

$$\xi_0 := \frac{(-1)^{j_0}}{s} \frac{z_{j_0}}{a_{j_0}} dz^{\widehat{j_0}},$$

where $dz^{\hat{j}_0}$ is a short-hand for $dz_1 \wedge \cdots \wedge dz_{j_0-1} \wedge dz_{j_0+1} \wedge \cdots \wedge dz_n$. Then

$$\int_{V_s} \prod_{j \in J(x_0)} |z_j|^{2k_j} dz \wedge d\bar{z} = \int_{V_s} \prod_{j \in J(x_0)} |z_j|^{2k_j} f^* ds \wedge \xi_0 \wedge \overline{f^* ds} \wedge \overline{\xi_0},$$

and since $f^* ds \wedge \overline{f^* ds}$ is constant along the fibers, we simply need to calculate

$$\int_{V_s} \prod_{j \in J(x_0)} |z_j|^{2k_j} \xi_0 \wedge \overline{\xi_0}.$$

For $j \in J(x_0)$, set $z_j = e^{\rho_j/a_j} e^{\sqrt{-1}\vartheta_j}$, and for $k \notin J(x_0)$, set $z_k = r_k e^{\sqrt{-1}\vartheta_k}$. For $j \in J(x_0)$,

$$dz_j = \frac{1}{a_j} e^{\sqrt{-1}\vartheta_j} e^{\rho_j/a_j} d\rho_j + \sqrt{-1} e^{\rho_j/a_j} e^{\sqrt{-1}\vartheta_j} d\vartheta_j,$$

and for $k \notin J(x_0)$,

$$dz_k = e^{\sqrt{-1}\vartheta_k} dr_k + \sqrt{-1} r_k e^{\sqrt{-1}\vartheta_k} d\vartheta_k.$$

Recall that the region of integration is the disk $\Delta = \{s \in \mathbb{C} : |s| < 1\}$. With $s = \prod_{j \in J(x_0)} z_j^{a_j}$, it follows that $|z_j| \leq 1$ for $j \in J(x_0)$. Thus, in these new coordinates,

$$\rho_j \leq 0, \quad \sum_{j \in J(x_0)} \rho_j = \log |s|,$$

$$0 \leq r_k \leq 1, \quad \sum_{j \in J(x_0)} a_j \vartheta_j = \arg(s).$$

Now

$$\begin{aligned} & dz^{\hat{j}_0} \\ &= \bigwedge_{j \in J(x_0)^*} dz_j \bigwedge_{k \notin J(x_0)} dz_k \\ &= \left(\bigwedge_{j \in J(x_0)^*} \frac{1}{a_j} e^{\sqrt{-1}\vartheta_j} e^{\rho_j/a_j} d\rho_j + \sqrt{-1} e^{\rho_j/a_j} e^{\sqrt{-1}\vartheta_j} d\vartheta_j \right) \wedge \left(\bigwedge_{k \notin J(x_0)} e^{\sqrt{-1}\vartheta_k} dr_k + \sqrt{-1} r_k e^{\sqrt{-1}\vartheta_k} d\vartheta_k \right) \\ &= \bigwedge_{j \in J(x_0)^*, k \notin J(x_0)} \left(\frac{1}{a_j} e^{\sqrt{-1}\vartheta_j} e^{\rho_j/a_j} e^{\sqrt{-1}\vartheta_k} d\rho_j \wedge dr_k + \sqrt{-1} \frac{1}{a_j} e^{\sqrt{-1}\vartheta_j} e^{\rho_j/a_j} e^{\sqrt{-1}\vartheta_k} d\rho_j \wedge d\vartheta_k \right) \\ &\quad + \bigwedge_{j \in J(x_0)^*, k \notin J(x_0)} \left(\sqrt{-1} e^{\rho_j/a_j} e^{\sqrt{-1}\vartheta_j} e^{\sqrt{-1}\vartheta_k} d\vartheta_j \wedge dr_k - e^{\rho_j/a_j} e^{\sqrt{-1}\vartheta_j} d\vartheta_j \wedge d\vartheta_k \right) \\ &=: \bigwedge_{j \in J(x_0)^*, k \notin J(x_0)} (A_{j,k} d\rho_j \wedge dr_k + B_{j,k} d\rho_j \wedge d\vartheta_k + C_{j,k} d\vartheta_j \wedge dr_k + D_{j,k} d\vartheta_j \wedge d\vartheta_k). \end{aligned}$$

Then

$$\begin{aligned}
dz^{\hat{j}_0} \wedge d\bar{z}^{\hat{j}_0} &= \bigwedge_{j \in J(x_0)^*, k \notin J(x_0)} (A_{j,k} d\rho_j \wedge dr_k + B_{j,k} d\rho_j \wedge d\vartheta_k + C_{j,k} d\vartheta_j \wedge dr_k + D_{j,k} d\vartheta_j \wedge d\vartheta_k) \\
&\wedge \bigwedge_{j \in J(x_0)^*, k \notin J(x_0)} (\bar{A}_{j,k} d\rho_j \wedge dr_k + \bar{B}_{j,k} d\rho_j \wedge d\vartheta_k + \bar{C}_{j,k} d\vartheta_j \wedge dr_k + \bar{D}_{j,k} d\vartheta_j \wedge d\vartheta_k) \\
&= \bigwedge_{j \in J(x_0)^*, k \notin J(x_0)} A_{j,k} \bar{D}_{j,k} d\rho_j \wedge dr_k \wedge d\vartheta_j \wedge d\vartheta_k + D_{j,k} \bar{A}_{j,k} d\vartheta_j \wedge d\vartheta_k \wedge d\rho_j \wedge dr_k \\
&= \bigwedge_{j \in J(x_0)^*, k \notin J(x_0)} \frac{1}{a_j} e^{2\rho_j/a_j} e^{2\sqrt{-1}\vartheta_j} d\rho_j \wedge dr_k \wedge d\vartheta_j \wedge d\vartheta_k,
\end{aligned}$$

and subsequently,

$$\begin{aligned}
\xi_0 \wedge \bar{\xi}_0 &= \frac{1}{|s|^2} \frac{|z_{j_0}|^2}{a_{j_0}^2} \frac{1}{a_j} e^{2\rho_j/a_j} e^{2\sqrt{-1}\vartheta_j} d\rho_j \wedge dr_k \wedge d\vartheta_j \wedge d\vartheta_k \\
&= \frac{1}{|s|^2} \prod_{j \in J(x_0)} \frac{1}{a_j^2} e^{2\rho_j/a_j} d\rho_j \prod_{\ell \in J(x_0)^*} e^{2\sqrt{-1}\vartheta_\ell} d\vartheta_\ell \bigwedge_{k \notin J(x_0)} dr_k.
\end{aligned}$$

The main integral is then given by

$$\begin{aligned}
&\int_{V_s} \prod_{j \in J(x_0)} |z_j|^{2k_j} \xi_0 \wedge \bar{\xi}_0 \\
&= \frac{1}{|s|^2} \int_{V_s} \prod_{j \in J(x_0)} \frac{1}{a_j^2} e^{2k_j \frac{\rho_j}{a_j} + 2\rho_j/a_j} d\rho_j \bigwedge_{\ell \in J(x_0)^*} e^{2\sqrt{-1}\vartheta_\ell} d\vartheta_\ell \bigwedge_{k \notin J(x_0)} dr_k \\
&= \frac{1}{|s|^2} \int_{V_s} \prod_{j \in J(x_0)} \frac{1}{a_j^2} e^{2\left(\frac{k_j+1}{a_j}\right)\rho_j} d\rho_j \bigwedge_{\ell \in J(x_0)^*} e^{2\sqrt{-1}\vartheta_\ell} d\vartheta_\ell \bigwedge_{k \notin J(x_0)} dr_k \\
&= \frac{1}{|s|^{2\left(1-\frac{k_{j_0}+1}{a_{j_0}}\right)}} \int_{V_s} \prod_{j \in J(x_0)} \frac{1}{a_j^2} e^{2\left(\frac{k_j+1}{a_j} - \frac{k_{j_0}+1}{a_{j_0}}\right)\rho_j} d\rho_j \bigwedge_{\ell \in J(x_0)^*} e^{2\sqrt{-1}\vartheta_\ell} d\vartheta_\ell \bigwedge_{k \notin J(x_0)} dr_k \\
&\sim \frac{1}{|s|^{2\left(1-\frac{k_{j_0}+1}{a_{j_0}}\right)}} \int_{V_s} \prod_{j \in J(x_0)} e^{2\left(\frac{k_j+1}{a_j} - \frac{k_{j_0}+1}{a_{j_0}}\right)\rho_j} d\rho_j.
\end{aligned}$$

We claim that

$$\int_{V_s} \prod_{j \in J(x_0)} e^{2\left(\frac{k_j+1}{a_j} - \frac{k_{j_0}+1}{a_{j_0}}\right)\rho_j} d\rho_j \sim (-\log |s|)^{N_0-1},$$

where N_0 is the cardinality of the set $\{j \in J(x_0) : \frac{k_j+1}{a_j} = \frac{k_{j_0}+1}{a_{j_0}}\}$. For simplicity, take $J(x_0) = \{1, 2, 3\}$, and $j_0 = 3$. We will consider three cases:

Case 1: $N_0 = 3$ and

$$\frac{k_1 + 1}{a_1} = \frac{k_2 + 1}{a_2} = \frac{k_3 + 1}{a_3}.$$

The region of integration R is described by $\rho_1 + \rho_2 + \rho_3 = \log |s|$ and $\rho_1, \rho_2, \rho_3 \leq 0$. Since $\log |s| \leq \rho_3 \leq 0$, we can write

$$R = \{(\rho_1, \rho_2) \in \mathbb{R}^2 : \rho_1 \leq 0, \rho_2 \leq 0, \log |s| \leq \rho_1 + \rho_2 \leq 0\},$$

which describes a triangle in the plane whose area is $\frac{1}{2}(\log |s|)^2 = \frac{1}{2}(-\log |s|)^2 \sim (-\log |s|)^2$.

Case 2: $N_0 = 2$, where we may assume that $\frac{k_1+1}{a_1} > \frac{k_3+1}{a_3}$ and $\frac{k_2+1}{a_2} = \frac{k_3+1}{a_3}$. Set $\alpha = 2\left(\frac{k_1+1}{a_1} - \frac{k_3+1}{a_3}\right) > 0$. Then

$$\begin{aligned} \int_R e^{\alpha \rho_1} d\rho_1 d\rho_2 &= \int_{\log |s|}^0 \int_{\log |s| - \rho_1}^0 e^{\alpha \rho_1} d\rho_2 d\rho_1 \\ &= \int_{\log |s|}^0 (\rho_1 - \log |s|) e^{\alpha \rho_1} d\rho_1 \\ &= -\frac{1}{\alpha} \log |s| + \frac{1}{\alpha^2} (e^{\alpha \log |s|} - 1) \\ &= -\frac{1}{\alpha} \log |s| + \frac{1}{\alpha^2} (|s|^\alpha - 1) \\ &\sim -\log |s| \end{aligned}$$

Case 3: $N_0 = 1$, where we may assume that $\frac{k_1+1}{a_1} > \frac{k_2+1}{a_2} > \frac{k_3+1}{a_3}$. Set $\alpha_i := \frac{k_i+1}{a_i} - \frac{k_3+1}{a_3} > 0$ for $i \in \{1, 2\}$. We compute

$$\begin{aligned}
\int_R e^{\alpha_1 \rho_1 + \alpha_2 \rho_2} d\rho_1 d\rho_2 &= \int_{\log |s|}^0 \int_{\log |s| - \rho_1}^0 e^{\alpha_1 \rho_1 + \alpha_2 \rho_2} d\rho_2 d\rho_1 \\
&= \int_{\log |s|}^0 \left[\frac{1}{\alpha_2} e^{\alpha_1 \rho_1 + \alpha_2 \rho_2} \right]_{\log |s| - \rho_1}^0 d\rho_1 \\
&= \int_{\log |s|}^0 \left[\frac{1}{\alpha_2} e^{\alpha_1 \rho_1} - \frac{1}{\alpha_2} e^{\alpha_1 \rho_1} e^{\alpha_2 \log |s| - \alpha_2 \rho_1} \right] d\rho_1 \\
&= \int_{\log |s|}^0 \frac{1}{\alpha_2} e^{\alpha_1 \rho_1} d\rho_1 - \int_{\log |s|}^0 \frac{1}{\alpha_2} |s|^{\alpha_2} e^{(\alpha_1 - \alpha_2) \rho_1} d\rho_1 \\
&= \left[\frac{1}{\alpha_1 \alpha_2} e^{\alpha_1 \rho_1} \right]_{\log |s|}^0 - \frac{1}{\alpha_2} |s|^{\alpha_2} \int_{\log |s|}^0 e^{(\alpha_1 - \alpha_2) \rho_1} d\rho_1 \\
&= \frac{1}{\alpha_1 \alpha_2} - \frac{1}{\alpha_1 \alpha_2} |s|^{\alpha_1} - \frac{1}{\alpha_2} |s|^{\alpha_2} \left[\frac{1}{\alpha_1 - \alpha_2} - |s|^{\alpha_1 - \alpha_2} \right] \\
&= \frac{1}{\alpha_1 \alpha_2} - \frac{1}{\alpha_1 \alpha_2} |s|^{\alpha_1} - \frac{1}{\alpha_2 (\alpha_1 - \alpha_2)} |s|^{\alpha_2} + \frac{1}{\alpha_2} |s|^{\alpha_1},
\end{aligned}$$

where all of these terms are bounded.

Therefore,

$$\int_{V_s} \prod_{j \in J(x_0)} |z_j|^{2k_j} \xi_0 \wedge \overline{\xi_0} \sim |s|^{-2 \left(1 - \frac{k_{j_0}+1}{a_{j_0}} \right)} (-\log |s|)^{N_0-1},$$

and we set $\beta := \frac{k_{j_0}+1}{a_{j_0}}$, and let N be the maximal number such that there are N divisors E_j with $\frac{k_j+1}{a_j} = \beta$. Since X is compact, cover X by a finite number of charts U_α . Let ϕ_α be a partition of unity subordinate to (U_α) . Then

$$\frac{\int_{\tilde{X}_s} \pi^* \mu}{\sqrt{-1} ds \wedge d\bar{s}} = \sum_\alpha \frac{\int_{\tilde{X}_s} \phi_\alpha \pi^* \mu}{\sqrt{-1} ds \wedge d\bar{s}} = \sum_\alpha \frac{\int_{\tilde{X}_s \cap U_\alpha} \phi_\alpha \pi^* \mu}{\sqrt{-1} ds \wedge d\bar{s}},$$

completing the proof. \square

REFERENCES

- [1] Zhang, Y., *Collapsing Limits of the Kähler–Ricci flow and the Continuity Method*, arXiv:1705.01434 v3, (2018).

We now compute the full expansion: That is, we compute the integrals

$$\begin{aligned}
 & \int_{V_s} \prod_{j \in J(x_0)} |z_j|^{2k_j} \xi_0 \wedge \overline{\xi_0} \\
 &= \int_{V_s} |s|^{-2} \prod_{j \in J(x_0)} e^{2\left(\frac{k_j+1}{a_j}\right)\rho_j} d\rho_j \\
 &= \int_{V_s} e^{-2\sum_{j \in J(x_0)} \rho_j} \prod_{j \in J(x_0)} e^{2\left(\frac{k_j+1}{a_j}\right)\rho_j} d\rho_j.
 \end{aligned}$$

where we omit the $1/a_j^2$ term. Let $\alpha_j := 2\left(\frac{k_j+1}{a_j} - 1\right)$. Then the integral simplifies to

$$\int_{V_s} e^{\sum_{j \in J(x_0)} \alpha_j \rho_j} d\rho_j.$$

Let $|J(x_0)| = \ell$, and recall that the region of integration is described by

$$\begin{aligned}
 V_s &= \left\{ (\rho_1, \dots, \rho_\ell) \in (\mathbb{R}_{<0})^\ell : \sum_{j=1}^{\ell} \rho_j = \log |s| \right\} \\
 &= \left\{ (\rho_1, \dots, \rho_{\ell-1}) \in (\mathbb{R}_{<0})^{\ell-1} : \log |s| \leq \sum_{j=1}^{\ell-1} \rho_j \leq 0 \right\}.
 \end{aligned}$$

We claim that

$$\int_{V_s} e^{\alpha_1 \rho_1 + \dots + \alpha_\ell \rho_\ell} d\rho_1 \dots d\rho_\ell = \frac{1}{\prod_{k=1}^{\ell} \alpha_k} + (-1)^\ell \sum_{k=1}^{\ell} \frac{1}{\alpha_k \prod_{1 \leq j \neq k \leq \ell} A_{kj}} |s|^{\alpha_k}, \quad A_{kj} := \alpha_k - \alpha_j.$$

Proof. Proceed inductively on ℓ . For $\ell = 2$, we compute

$$\begin{aligned}
\int_{V_s} e^{\alpha_1 \rho_1 + \alpha_2 \rho_2} d\rho_1 d\rho_2 &= \int_{\log |s|}^0 \int_{\log |s| - \rho_2}^0 e^{\alpha_1 \rho_1 + \alpha_2 \rho_2} d\rho_1 d\rho_2 \\
&= \int_{\log |s|}^0 \left[\frac{1}{\alpha_1} e^{\alpha_1 \rho_1 + \alpha_2 \rho_2} \right]_{\log |s| - \rho_2}^0 d\rho_2 \\
&= \int_{\log |s|}^0 \left[\frac{1}{\alpha_1} e^{\alpha_2 \rho_2} - \frac{1}{\alpha_1} e^{\alpha_1 (\log |s| - \rho_2) + \alpha_2 \rho_2} \right] d\rho_2 \\
&= \int_{\log |s|}^0 \left[\frac{1}{\alpha_1} e^{\alpha_2 \rho_2} - \frac{1}{\alpha_1} e^{\log |s|^{\alpha_1} + (\alpha_2 - \alpha_1) \rho_2} \right] d\rho_2 \\
&= \int_{\log |s|}^0 \left[\frac{1}{\alpha_1} e^{\alpha_2 \rho_2} - \frac{1}{\alpha_1} |s|^{\alpha_1} e^{(\alpha_2 - \alpha_1) \rho_2} \right] d\rho_2 \\
&= \left[\frac{1}{\alpha_1 \alpha_2} e^{\alpha_2 \rho_2} - \frac{1}{\alpha_1 (\alpha_2 - \alpha_1)} |s|^{\alpha_1} e^{(\alpha_2 - \alpha_1) \rho_2} \right]_{\log |s|}^0 \\
&= \left[\frac{1}{\alpha_1 \alpha_2} - \frac{1}{\alpha_1 (\alpha_2 - \alpha_1)} |s|^{\alpha_1} \right] - \left[\frac{1}{\alpha_1 \alpha_2} |s|^{\alpha_2} - \frac{1}{\alpha_1 (\alpha_2 - \alpha_1)} |s|^{\alpha_1} |s|^{\alpha_2 - \alpha_1} \right] \\
&= \frac{1}{\alpha_1 \alpha_2} + \frac{1}{\alpha_1 (\alpha_1 - \alpha_2)} |s|^{\alpha_1} - \frac{1}{\alpha_1 \alpha_2} |s|^{\alpha_2} + \frac{1}{\alpha_1 (\alpha_2 - \alpha_1)} |s|^{\alpha_2} \\
&= \frac{1}{\alpha_1 \alpha_2} + \frac{1}{\alpha_1 (\alpha_1 - \alpha_2)} |s|^{\alpha_1} + \frac{1}{\alpha_2 (\alpha_2 - \alpha_1)} |s|^{\alpha_2} \\
&= \frac{1}{\alpha_1 \alpha_2} + \frac{1}{\alpha_1 A_{12}} |s|^{\alpha_1} + \frac{1}{\alpha_2 A_{21}} |s|^{\alpha_2}.
\end{aligned}$$

Suppose the claim is true for ℓ , i.e.,

$$\int_{V_s} e^{\alpha_1 \rho_1 + \dots + \alpha_\ell \rho_\ell} d\rho_1 \dots d\rho_\ell = \frac{1}{\prod_{k=1}^\ell \alpha_k} + (-1)^\ell \sum_{k=1}^\ell \frac{1}{\alpha_k \prod_{1 \leq j \neq k \leq \ell} A_{kj}} |s|^{\alpha_k}, \quad A_{kj} := \alpha_k - \alpha_j.$$

We will show the result holds for $\ell + 1$. To this end,

□

Proof. Evaluate the integral directly:

$$\begin{aligned}
 & \int_{V_s} e^{\alpha_1 \rho_1 + \dots + \alpha_\ell \rho_\ell} d\rho_1 \dots d\rho_\ell \\
 = & \int_{\log |s|}^0 \int_{\log |s| - \rho_{\ell-1}}^0 \dots \int_{\log |s| - (\rho_2 + \dots + \rho_\ell)}^0 e^{\alpha_1 \rho_1 + \dots + \alpha_\ell \rho_\ell} d\rho_1 \dots d\rho_\ell \\
 = & \int_{\log |s|}^0 \int_{\log |s| - \rho_{\ell-1}}^0 \dots \int_{\log |s| - (\rho_3 + \dots + \rho_{\ell-1})}^0 \left[\frac{1}{\alpha_1} e^{\alpha_1 \rho_1 + \dots + \alpha_\ell \rho_\ell} \right]_{\log |s| - (\rho_2 + \dots + \rho_\ell)}^0 d\rho_2 \dots d\rho_\ell \\
 = & \int_{\log |s|}^0 \int_{\log |s| - \rho_{\ell-1}}^0 \dots \int_{\log |s| - (\rho_3 + \dots + \rho_{\ell-1})}^0 \frac{1}{\alpha_1} e^{\alpha_2 \rho_2 + \dots + \alpha_\ell \rho_\ell} d\rho_2 \dots d\rho_\ell \\
 & - \int_{\log |s|}^0 \int_{\log |s| - \rho_{\ell-1}}^0 \dots \int_{\log |s| - (\rho_3 + \dots + \rho_{\ell-1})}^0 \frac{1}{\alpha_1} |s|^{\alpha_1} e^{A_{21} \rho_2 + A_{31} \rho_3 + \dots + A_{(\ell-1)1} \rho_{\ell-1} + \alpha_\ell \rho_\ell} d\rho_2 \dots d\rho_\ell.
 \end{aligned}$$

Calculate:

$$\begin{aligned}
 & \int_{\log |s| - (\rho_3 + \dots + \rho_{\ell-1})}^0 \frac{1}{\alpha_1} e^{\alpha_2 \rho_2 + \dots + \alpha_\ell \rho_\ell} d\rho_2 \\
 = & \left[\frac{1}{\alpha_1 \alpha_2} e^{\alpha_3 \rho_3 + \dots + \alpha_\ell \rho_\ell} e^{\alpha_2 \rho_2} \right]_{\log |s| - (\rho_3 + \dots + \rho_{\ell-1})}^0 \\
 = & \frac{1}{\alpha_1 \alpha_2} e^{\alpha_3 \rho_3 + \dots + \alpha_\ell \rho_\ell} - \frac{1}{\alpha_1 \alpha_2} e^{A_{32} \rho_3 + \dots + A_{(\ell-1)2} \rho_{\ell-1} + \alpha_\ell \rho_\ell} |s|^{\alpha_2}
 \end{aligned}$$

And:

$$\begin{aligned}
 & \int_{\log |s| - (\rho_3 + \dots + \rho_{\ell-1})}^0 \frac{1}{\alpha_1} |s|^{\alpha_1} e^{A_{21} \rho_2 + A_{31} \rho_3 + \dots + A_{(\ell-1)1} \rho_{\ell-1} + \alpha_\ell \rho_\ell} d\rho_2 \\
 = & \frac{1}{\alpha_1 A_{21}} |s|^{\alpha_1} e^{A_{31} \rho_3 + \dots + A_{(\ell-1)1} \rho_{\ell-1} + \alpha_\ell \rho_\ell} \left(1 - |s|^{A_{21}} e^{-A_{21}(\rho_3 + \dots + \rho_{\ell-1})} \right) \\
 = & \frac{1}{\alpha_1 A_{21}} |s|^{\alpha_1} e^{A_{31} \rho_3 + \dots + A_{(\ell-1)1} \rho_{\ell-1} + \alpha_\ell \rho_\ell} - \frac{1}{\alpha_1 A_{21}} |s|^{\alpha_2} e^{A_{32} \rho_3 + \dots + A_{(\ell-1)2} \rho_{\ell-1} + \alpha_\ell \rho_\ell}
 \end{aligned}$$

Combining the expressions gives:

$$\begin{aligned}
 & \frac{1}{\alpha_1 \alpha_2} e^{\alpha_3 \rho_3 + \dots + \alpha_\ell \rho_\ell} + \frac{1}{\alpha_1 A_{21}} |s|^{\alpha_1} e^{A_{31} \rho_3 + \dots + A_{(\ell-1)1} \rho_{\ell-1} + \alpha_\ell \rho_\ell} \\
 & - \left(\frac{1}{\alpha_1 \alpha_2} + \frac{1}{\alpha_1 A_{21}} \right) |s|^{\alpha_2} e^{A_{32} \rho_3 + \dots + A_{(\ell-1)2} \rho_{\ell-1} + \alpha_\ell \rho_\ell} \\
 = & \frac{1}{\alpha_1 \alpha_2} e^{\alpha_3 \rho_3 + \dots + \alpha_\ell \rho_\ell} + \frac{1}{\alpha_1 A_{21}} |s|^{\alpha_1} e^{A_{31} \rho_3 + \dots + A_{(\ell-1)1} \rho_{\ell-1} + \alpha_\ell \rho_\ell} \\
 & - \frac{2\alpha_2 - \alpha_1}{\alpha_1 \alpha_2 (\alpha_2 - \alpha_1)} |s|^{\alpha_2} e^{A_{32} \rho_3 + \dots + A_{(\ell-1)2} \rho_{\ell-1} + \alpha_\ell \rho_\ell}
 \end{aligned}$$

□

For $\ell = 3$, we get

$$= \frac{1}{\alpha_1 \alpha_2 \alpha_3} - \frac{1}{\alpha_1(\alpha_1 - \alpha_2)(\alpha_1 - \alpha_3)} |s|^{\alpha_1} + \frac{1}{\alpha_2(\alpha_1 - \alpha_2)(\alpha_2 - \alpha_3)} |s|^{\alpha_2} - \frac{1}{\alpha_3(\alpha_1 - \alpha_3)(\alpha_2 - \alpha_3)} |s|^{\alpha_3},$$

and for $\ell = 4$, we get

$$= \frac{1}{\alpha_1 \alpha_2 \alpha_3 \alpha_4} + \frac{1}{\alpha_1 A_{12} A_{13} A_{14}} |s|^{\alpha_1} - \frac{1}{\alpha_2 A_{12} A_{23} A_{24}} |s|^{\alpha_2} + \frac{1}{\alpha_3 A_{13} A_{23} A_{34}} |s|^{\alpha_3} - \frac{1}{\alpha_4 A_{14} A_{24} A_{34}} |s|^{\alpha_4}.$$