

# Numerical Solutions & Analysis for PDEs&ODEs

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## Abstract

□ **Infinite domain problems** can be applied broadly and effectively in our daily life such as temperature distribution in physical phenomena and option pricing in financial mathematics.

□ **Sturm–Liouville theory** is the theory of real second-order linear ordinary differential equations of a certain form. Sturm–Liouville equations have been widely researched and its numerical solutions can have good properties.

□ **Partial Differential Equations (PDEs)** can be applied broadly and effectively in our daily life such as temperature distribution in physical phenomena and option pricing in financial mathematics.

□ **The heat equation is a parabolic PDE** that describes the distribution of heat (or variation in temperature) in a given region over time.

## Objectives & Applications

□ Validate some finite difference methods to solve one dimensional Stum-Liouville problems and use them to solve one dimensiional finite domain problems.

□ Transform a classical infinite domain ordinary differencial equations into a Sturm-Liouville problem, generate a method for numerical solutions, analyze and minimize the error.

□ Develop and validate some finite difference methods to solve parabolic PDEs such as the heat equation numerically. And we build a model to solve the special clothing design for high temperature work.

## Models

□ **Infinite domain problem in 1D**

$$u''(x)=f(x),a<x<+\infty$$

□ **Boundary conditons**

$$u(a)=ua,\lim_{x\rightarrow\infty}u(x)=u\inf$$

□ **Sturm-Liouville problem**

$$\frac{d}{dx}(p(x)\frac{du(x)}{dx})+q(x)u(x)=f(x)$$

□ **Boundary conditons**

$$u(a)=\alpha,u(b)=\beta$$

□ **Heat Equation in 1D**

$$\frac{\partial u}{\partial t}=a^2\frac{\partial^2 u}{\partial t^2},a<x<b,0<t$$

□ **Boundary conditions**

$$u(a,t)=g_1(t),u(b,t)=g_2(t)$$

□ **Initial Condition:**

$$u(x,0)=u_0(x)$$

## Methods

□ **Finite difference method**

$$\frac{u_{i+1}-2u_i+u_{i-1}}{h^2}=f(x_i)$$

,where  $u_{i+1}=u_i+h$

Applying Taylor expansion, we can see the pattern is second order convergent.

□ **Splitting method**

Split the interval  $[a,+\infty)$  into two intervals:

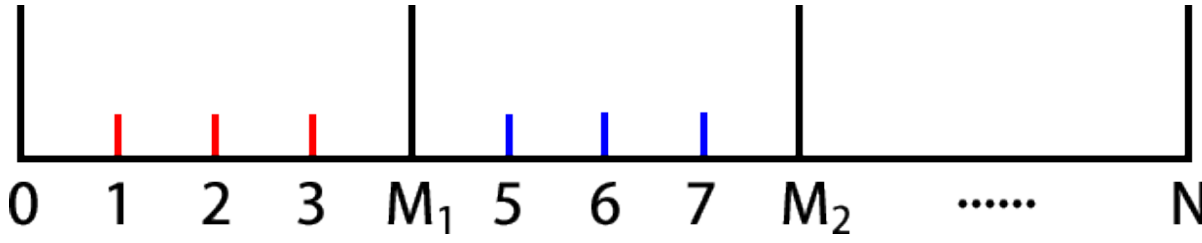
$[a,b]$  and  $[b,+\infty)$ , and in the second interval,

Let  $t=\frac{1}{x}$ , then the original equation  $u''(x)=f(x),a<x<+\infty$

transforms to 
$$\frac{d^2u(t)}{dt^2}+\frac{2}{t}\frac{du(t)}{dt}=\frac{f(\frac{1}{t})}{t^4}$$

,which is a transformation of a Sturm-Liouville equation having boundary  $u(0)=\lim_{x\rightarrow\infty}u(x)=u\inf$

□ **Finite-Difference Approximations to the Heat Conduction Equation**



$$A_1u_{0,j+1}+B_1u_{1,j+1}+C_1u_{2,j+1}=v'_{1,j}$$
$$\dots=\dots$$

$$A_{M_1-1}u_{M_1-2,j+1}+B_{M_1-1}u_{M_1-1,j+1}+C_{M_1-1}u_{M_1,j+1}=v_{M_1-1,j}$$

$$A_{M_1}u_{M_1-1,j+1}+B_{M_1}u_{M_1,j+1}+C_{M_1}u_{M_1+1,j+1}=v_{M_1,j}$$

$$A_{M_1+1}u_{M_1,j+1}+B_{M_1+1}u_{M_1+1,j+1}+C_{M_1+1}u_{M_1+2,j+1}=v_{M_1+1,j}$$

$\dots=\dots$

$$A_{M_2}u_{M_2-1,j+1}+B_{M_2}u_{M_2,j+1}+C_{M_2}u_{M_2+1,j+1}=v_{M_2,j}$$

$\dots=\dots$

$$A_Nu_{N-1,j+1}+B_Nu_{N,j+1}+C_Nu_{N+1,j+1}=v'_{N,j}$$

$$\begin{bmatrix} B_1 & C_1 \\ A_2 & B_2 & C_2 \\ & A_3 & B_3 & C_3 \\ & & A_{M_1} & B_{M_1} & C_{M_1} \\ & & & \ddots & \ddots & \ddots \\ & & & & A_{M_2} & B_{M_2} & C_{M_2} \\ & & & & & \ddots & \ddots & \ddots \\ & & & & & & A_N & B_N \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \\ u_3 \\ \vdots \\ u_{M_1} \\ \vdots \\ u_{M_2} \\ \vdots \\ u_N \end{bmatrix} = \begin{bmatrix} v_1 \\ v_2 \\ v_3 \\ \vdots \\ v_M \\ \vdots \\ v_{M_2} \\ \vdots \\ v_N \end{bmatrix}$$

## Results

□ **Infinite domain problem**

**Splitting method**

Solve the equation for

$$u(x)=1+e^{-x}$$

Boundaries :

$$u(1)=1+1/e,\lim_{x\rightarrow\infty}u=1$$

Choose b=10

Switch n1,n2

Plot infinite norm of error  
Against the maximum of h.

The plot proves the error  
converges by second order.

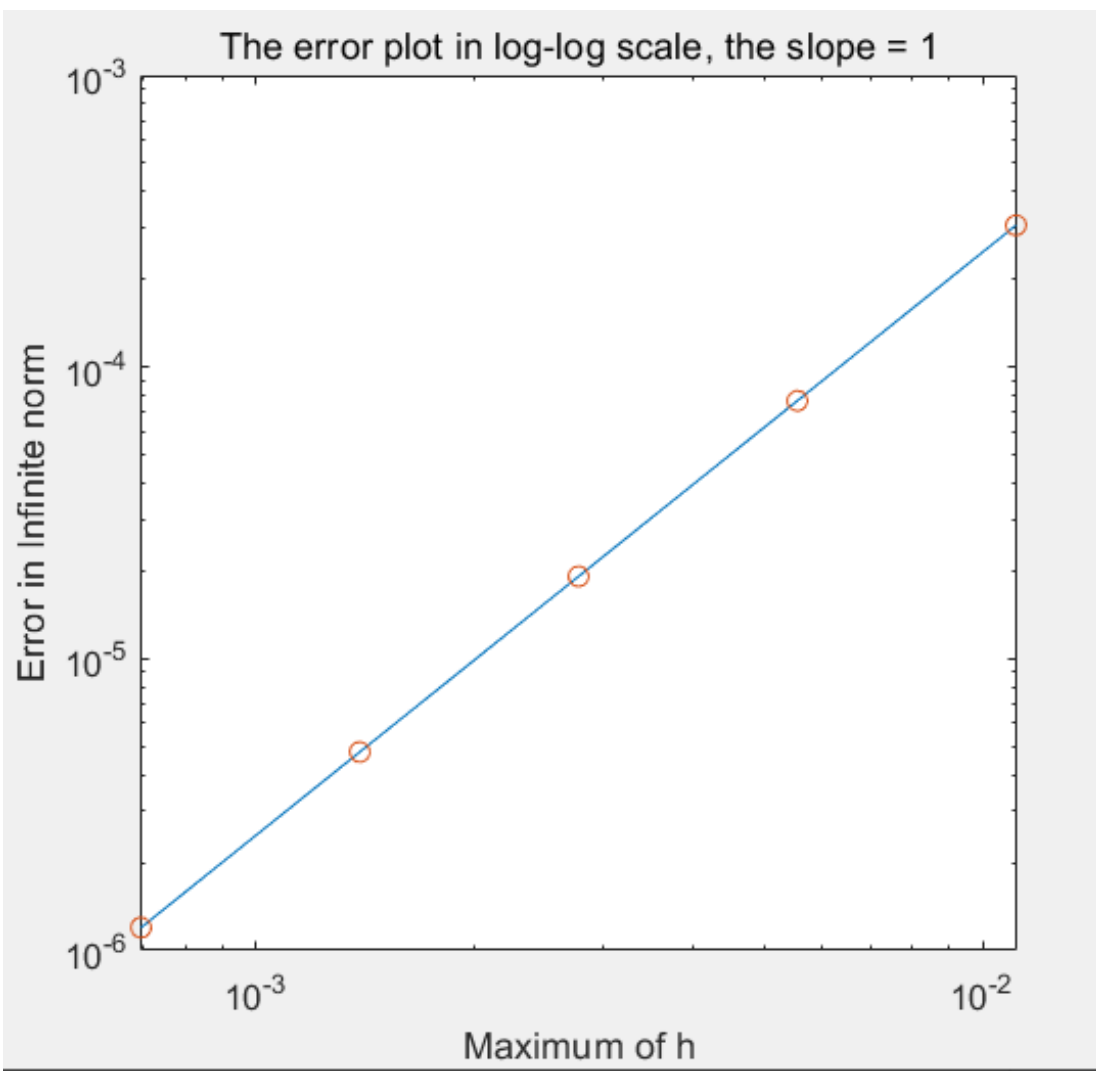


Figure 1. Error plot in log-log scale

## Results

□ **Error analysis**

Three parts of different truncation error

Interval  $[a,b]$ :  $e_1=\frac{1}{12}u''(x)h_1^2$

Interval  $[b,+\infty)$ :  $e_2=\frac{1}{12}u'''(t)h_2^2+\frac{1}{3t}u'''(t)h_2^2$

Point b:  $e_3=-\frac{1}{6}\alpha_1u'''(b)h_1^2+\frac{1}{6}\alpha_3u'''(b)\left(\frac{b^2h_2}{1-bh_2}\right)^2$

Compare all the three errors,  
we prove in theory the error  
converges by second order.

Switch b to see when  
the error is the smallest.

Fix n1=100, n2=50, a=1,

Switch b between 0 and 100,

We see at point b=4.49 the  
error reaches its bottom.

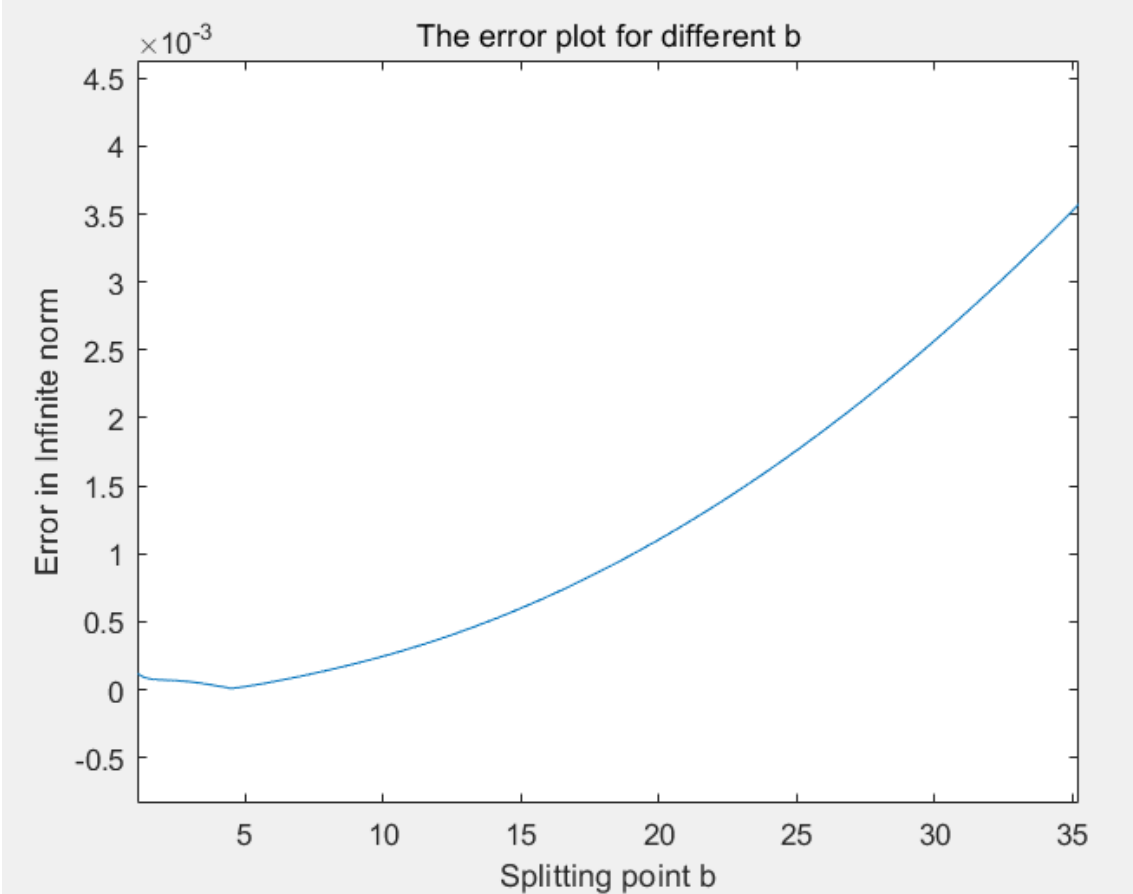


Figure 2. Error plot for different b

□ **The numerical solution of heat equation**

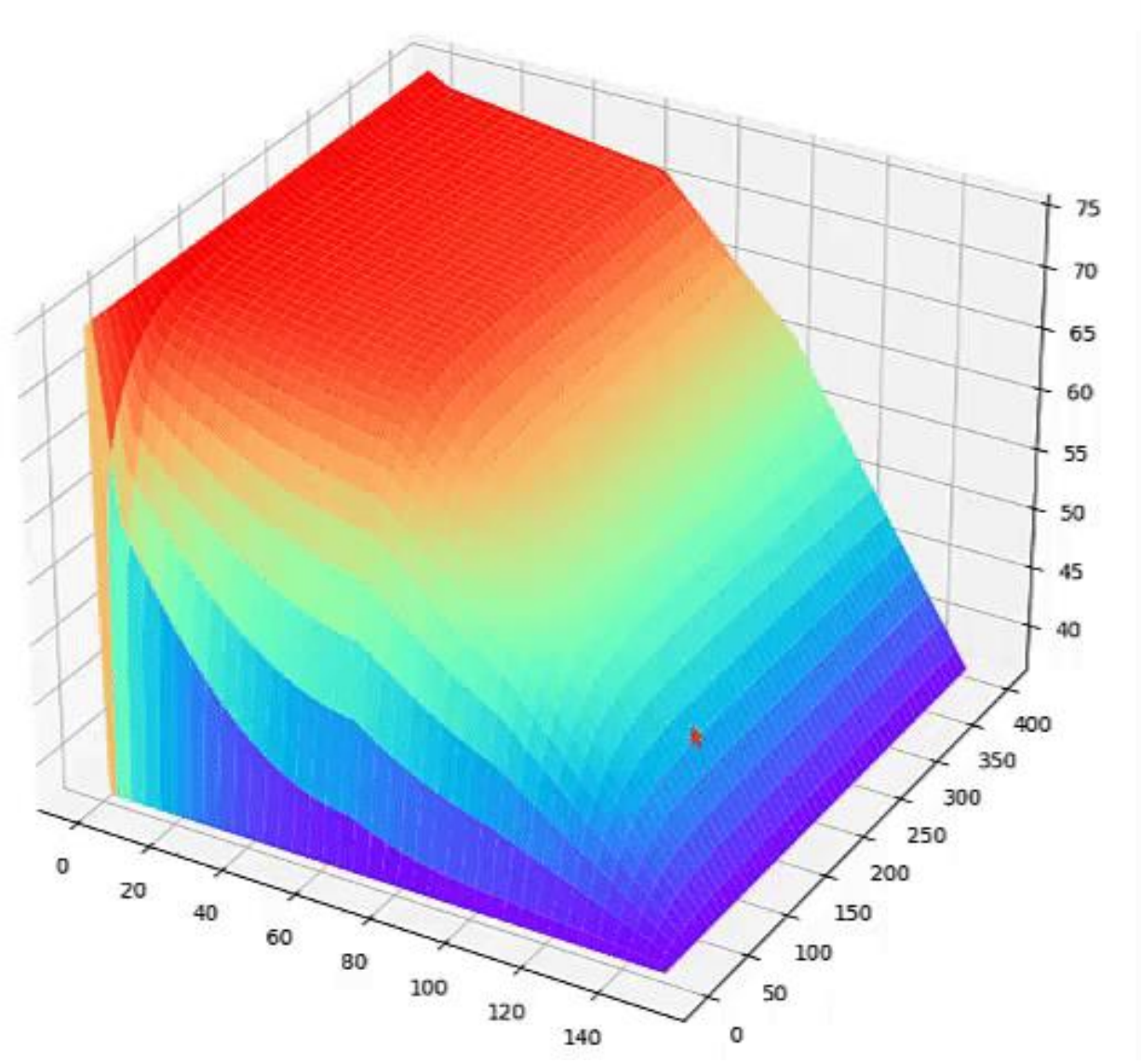


Figure 3. Temperature changes with time goes

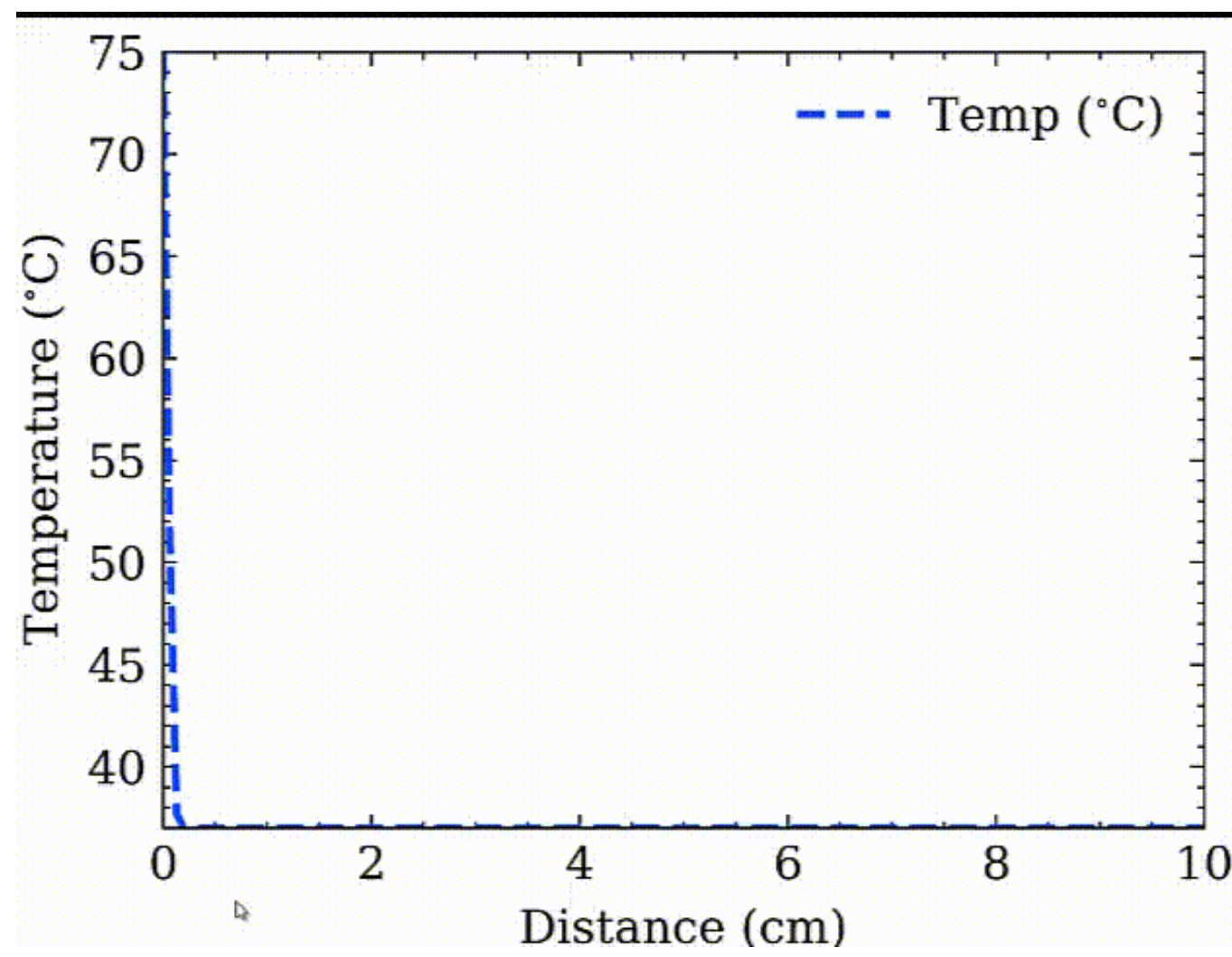


Figure 4. Temperature distribution with respect to time and space

## Future Work

□ Carry out further and deeper error analysis and work out a general code to generate the numerical solution for the infinite domain problem which minimize the overall error.

□ Extend the splitting method further into infinite domain partial differential equations and utilize polar coordinates to generate a solution.

□ Combine our mathematical learning with financial problems like put option with infinite domain.

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