



Neutron scattering: an experimental perspective (with side-notes on history and philosophy)

Allen Scheie

6/25/2025

LA-UR-25-25909

Outline:

1. History of neutron scattering
2. Modern neutron scattering
3. Current uses of neutron scattering: exploring the unknown

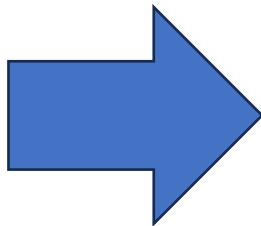
Part 1:

History of neutron scattering

1939: USA begins secret project to research nuclear bomb



ahf.nucrarmuseum.org

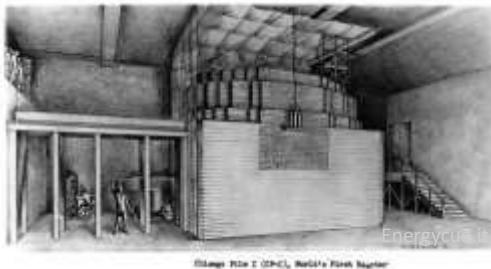


(Unofficial) MED emblem, 1946

Osti.gov

Nuclear reactors built to produce plutonium

Chicago pile 1 (1942)



0.5 W



Oak Ridge X-10 (1943)



1 MW



Hanford B (1944)



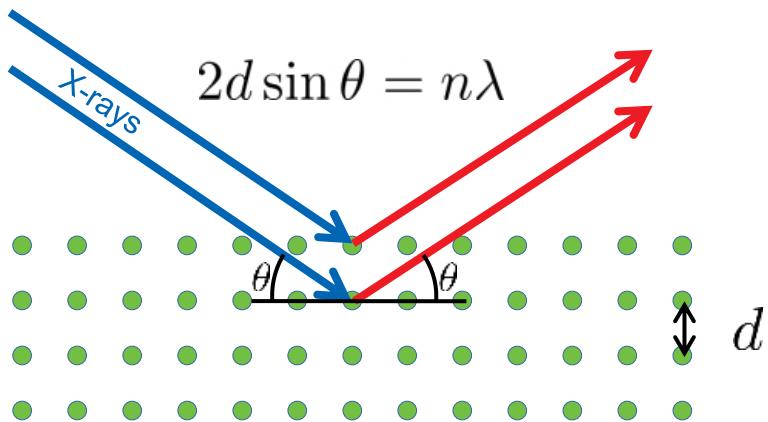
250 MW

After the war, X-10 kept running: what else can we do with a nuclear reactor?



One idea: crystallography

- Von Laue & Paul Ewald: X-ray diffraction on copper sulfate (1912)
 - Technique refined by William Braggs (Nobel Prize 1915), Scherrer, Debye, Hull, etc.



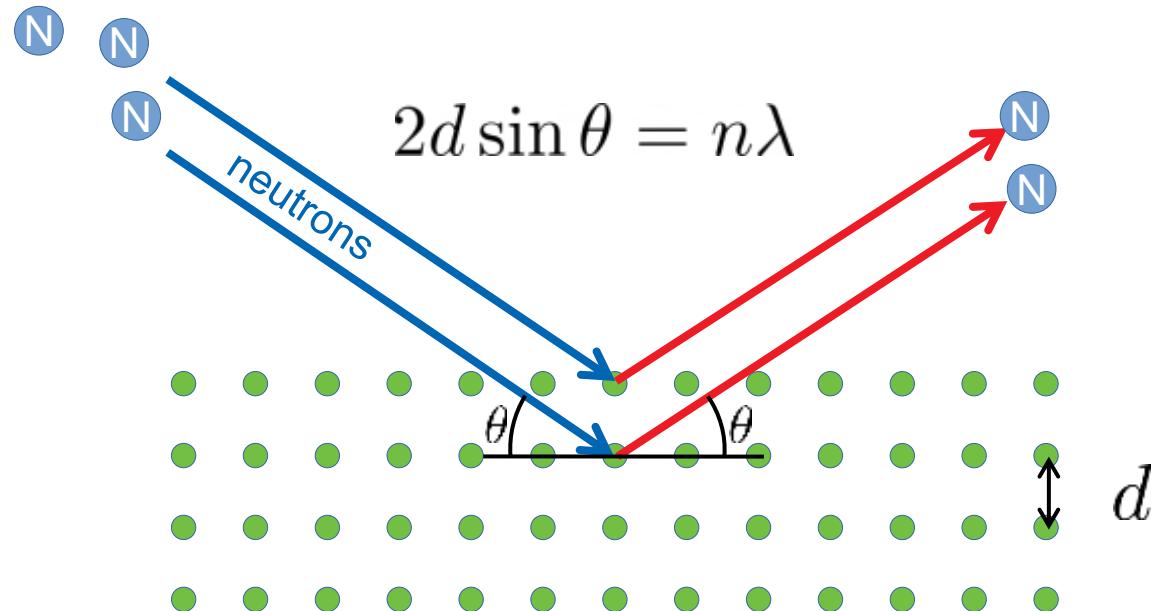
- Moderated neutrons have a similar wavelength to X-rays
(by DeBroglie's relation)

Thermal neutron wavelength:

$$\lambda = 1.8 \text{ \AA}$$

Close to typical atom spacing in a crystal

1944: Ernest Wollan measures neutron diffraction on gypsum and NaCl (table salt)

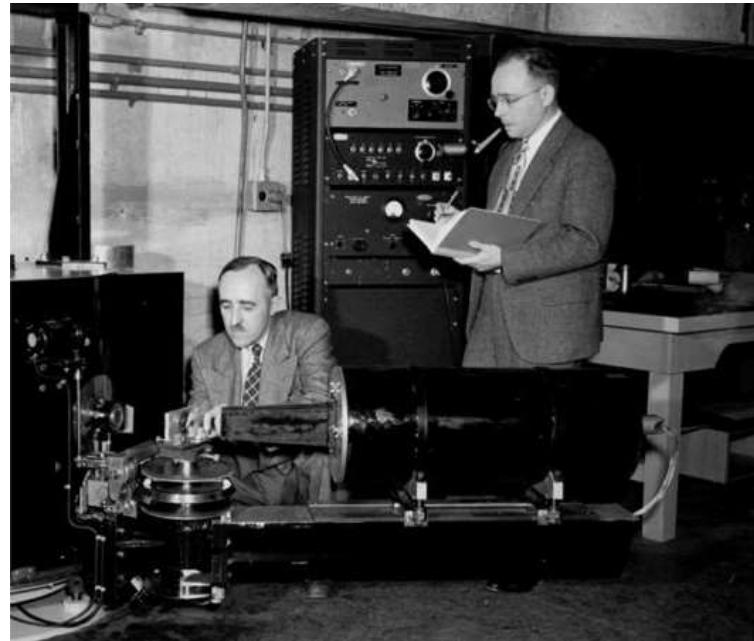


1946: Shull joins Wollan to pursue technique

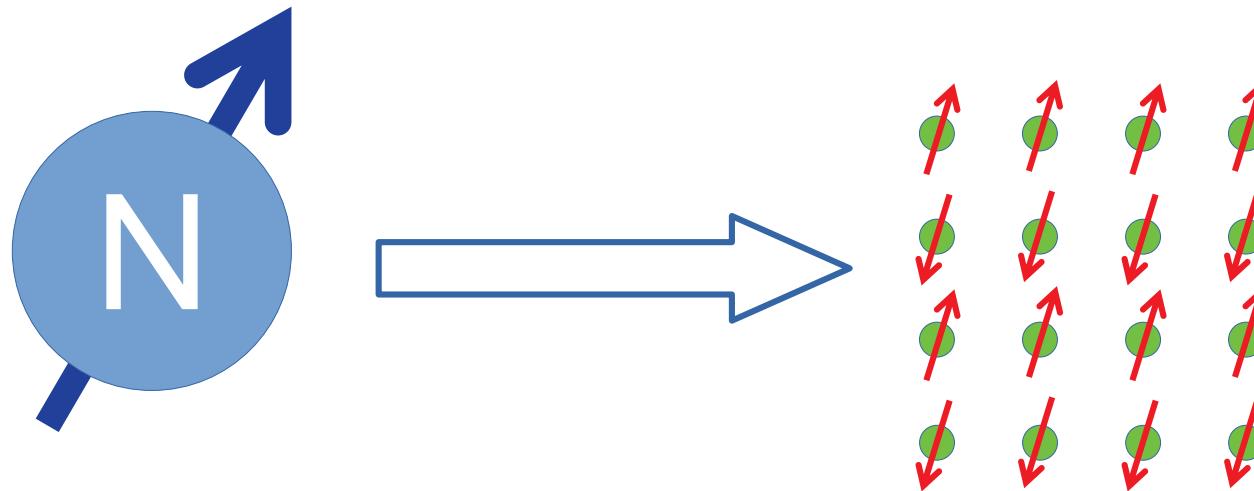
X-10 site (now ORNL)



Clifford Shull and Ernest Wollan
at the X-10 graphite reactor



Neutrons are chargless (pass through materials easily) but magnetic: so they scatter off magnetic patterns



1949: neutrons reveal antiferromagnetism in MnO

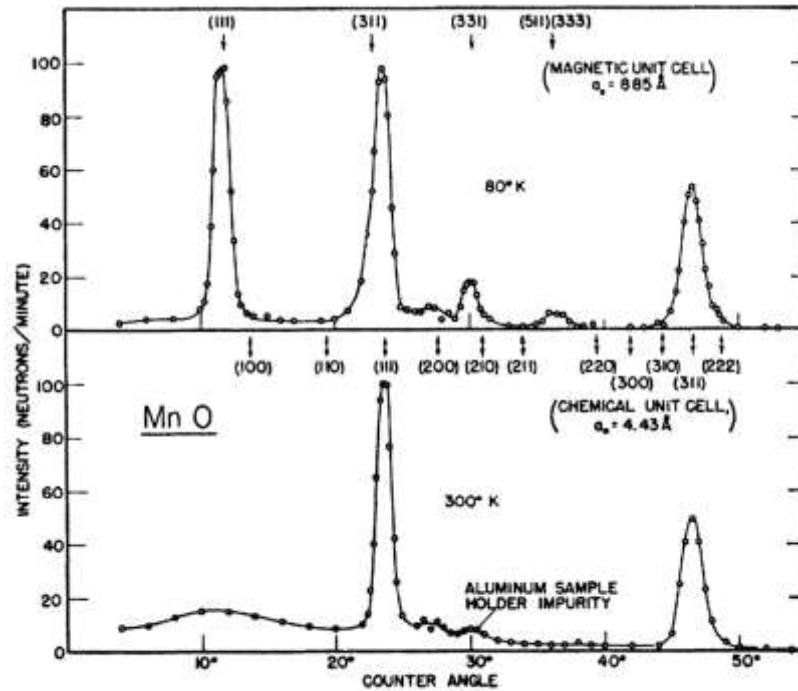
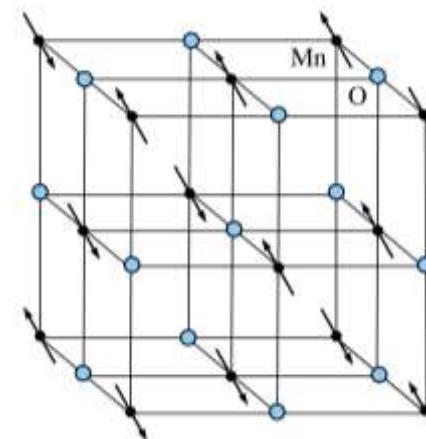
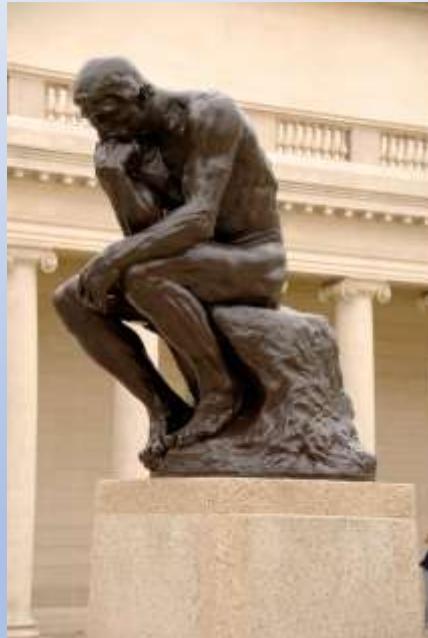


FIG. 1. Neutron diffraction patterns for MnO at room temperature and at 80°K.



Philosophical interlude:





Antiferromagnetism was first clear example of *spontaneous symmetry breaking*

- Global symmetry preserved, local symmetry broken
- Only occurs in thermodynamic limit (large system size)
 - Quantum theory predicts a superposition of two AFM states:
$$\frac{1}{\sqrt{2}}(|\uparrow\downarrow\uparrow\downarrow\dots\rangle \pm |\downarrow\uparrow\downarrow\uparrow\dots\rangle)$$
 - Seems to be related to decoherence?
- Now invoked all over physics (e.g. QCD)

Back to the physics.

Neutrons vs. X-rays: neutrons have much smaller energy at the same wavelength

Thermal neutron energy
($\lambda=1.8 \text{ \AA}$):

25 meV

10^5
difference

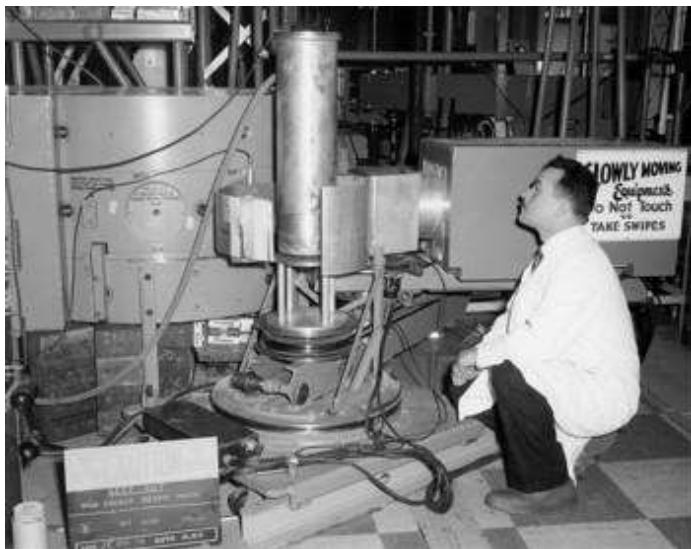
X-ray energy
($\lambda=1.8 \text{ \AA}$):

6.9 keV



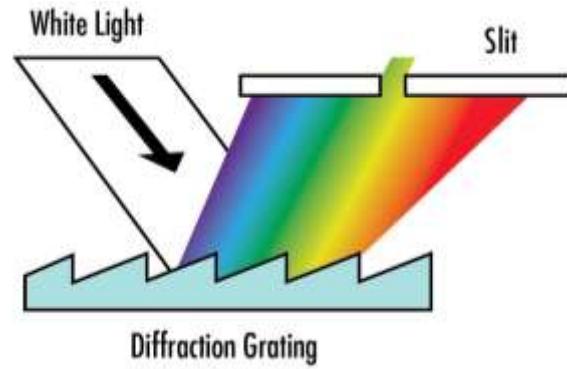
Can energy resolve to measure low-energy features?

Bertram Brockhouse: builds triple-axis spectrometer in 1955 at Chalk River reactor in Canada



Cins.ca

crystal diffraction acts like a diffraction grating



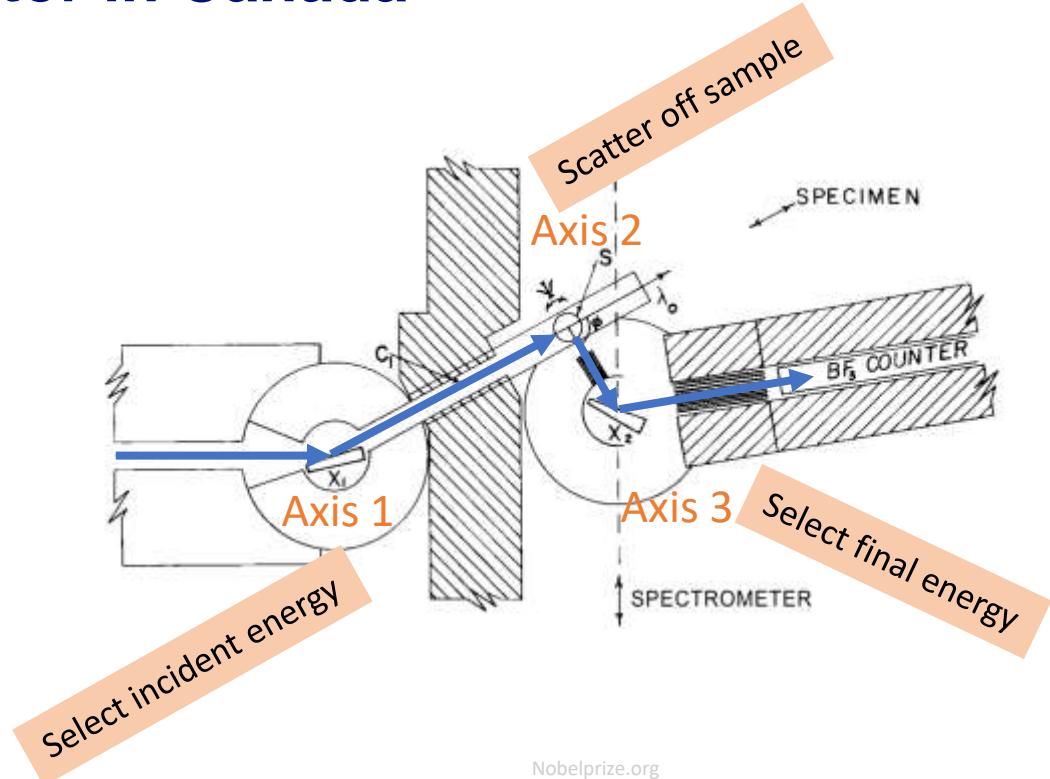
Selecting a wavelength = selecting an energy

$$\lambda = \frac{h}{mv} \quad E = \frac{1}{2}mv^2$$

Bertram Brockhouse: builds triple-axis spectrometer in 1955 at Chalk River reactor in Canada

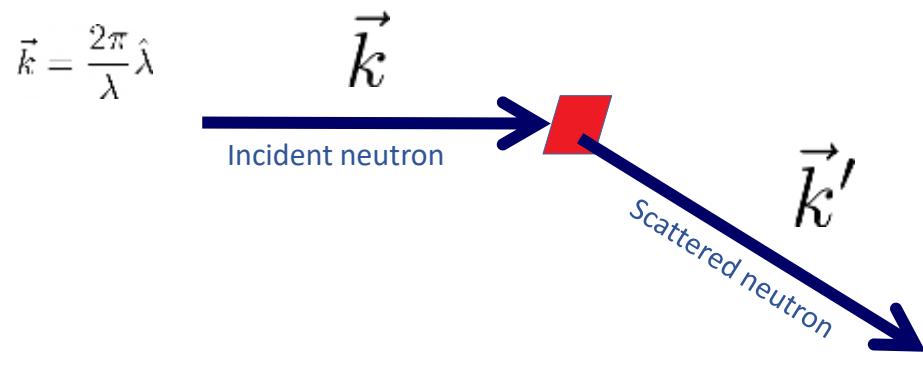


Cins.ca



Nobelprize.org

Conservation of momentum and energy:



$$\vec{k} = \frac{2\pi}{\lambda} \hat{\lambda}$$

Momentum:

$$\vec{Q} = \vec{k} - \vec{k}'$$

$$\vec{k} = \frac{m\vec{v}}{\hbar} \quad \vec{Q} = \frac{1}{\hbar}(m\vec{v} - m\vec{v}')$$

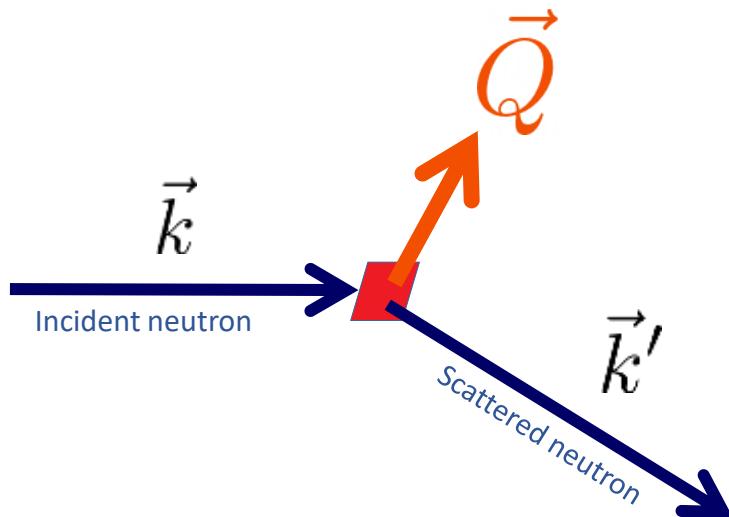
Energy:

$$\Delta E = E_1 - E_2$$

also written

$$\hbar\omega = \frac{\hbar^2}{2m}(k^2 - k'^2)$$

Energy and momentum can be deposited in the sample, and it can be measured!



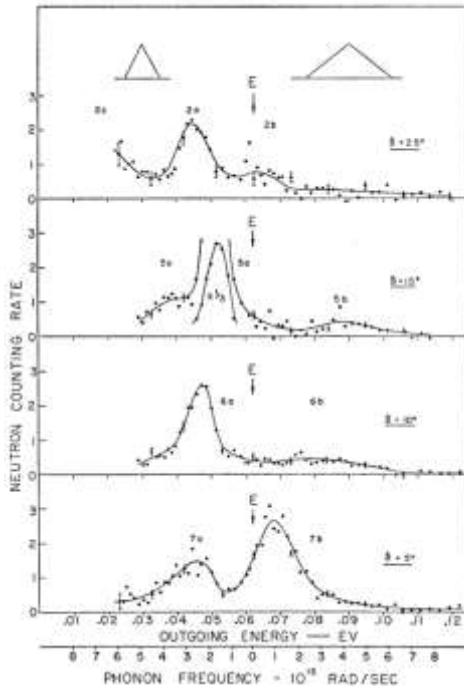
Momentum:

$$\vec{Q} = \vec{k} - \vec{k}'$$

Energy:

$$\hbar\omega = \frac{\hbar^2}{2m}(k^2 - k'^2)$$

Brockhouse measured a dispersion curve for the lattice vibrations in aluminum



Scattering of Neutrons by Phonons in an Aluminum Single Crystal

B. N. BROCKHOUSE AND A. T. STEWART

Physics Division, Atomic Energy of Canada, Limited,
Chalk River, Ontario, Canada

(Received August 29, 1955)

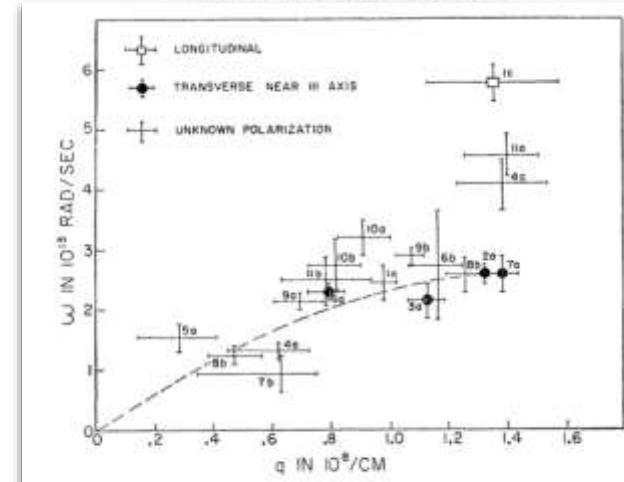
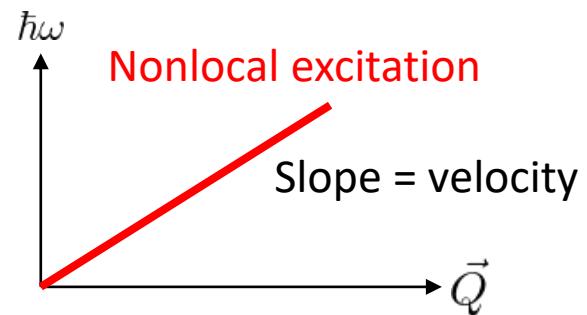
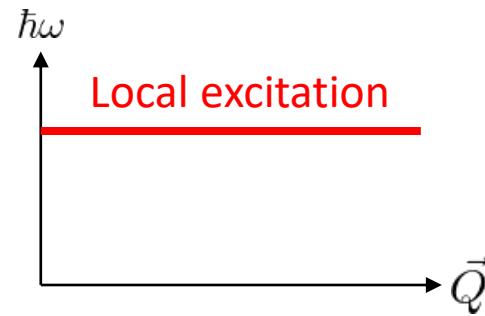
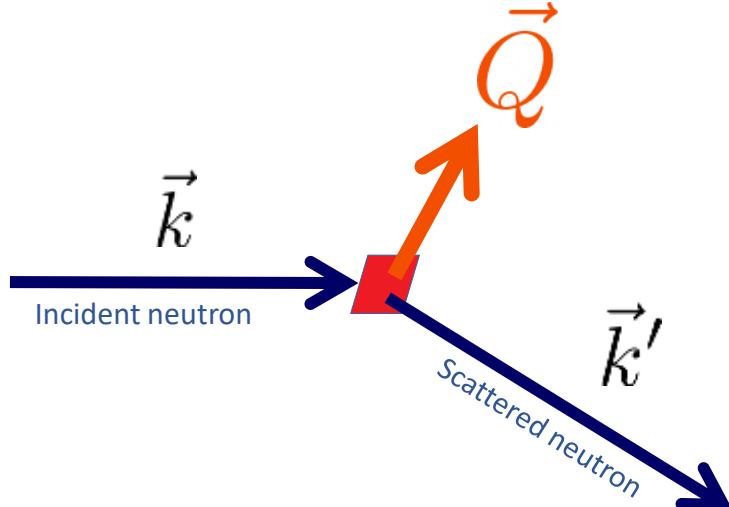
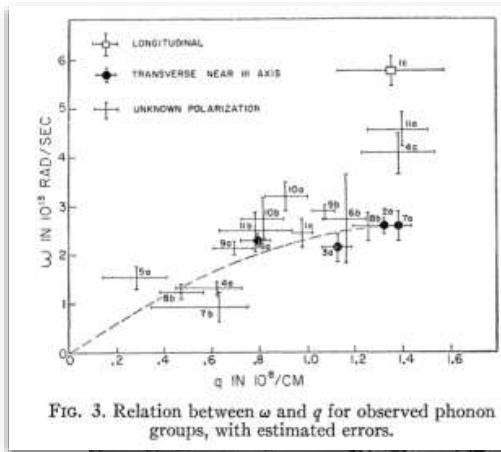


FIG. 3. Relation between ω and q for observed phonon groups, with estimated errors.

Dispersion curve: relation between energy and momentum



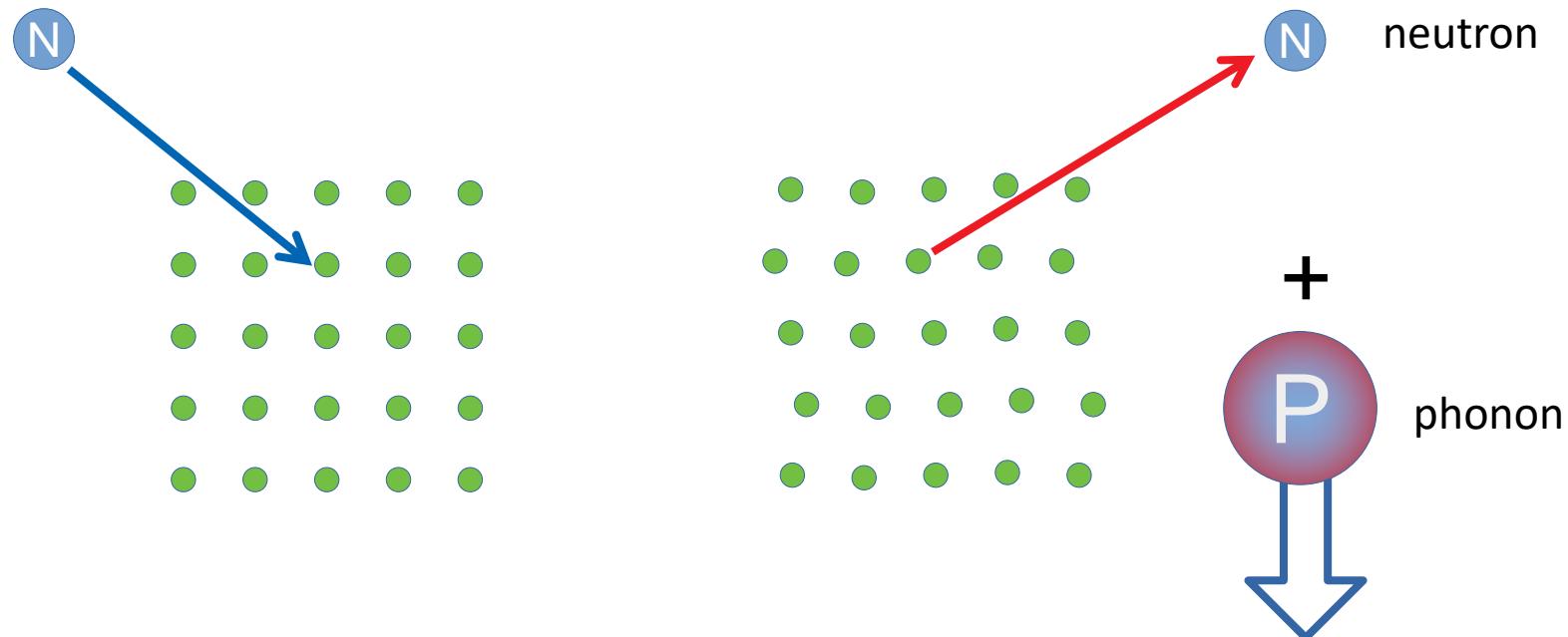
First direct evidence of “quasiparticles”



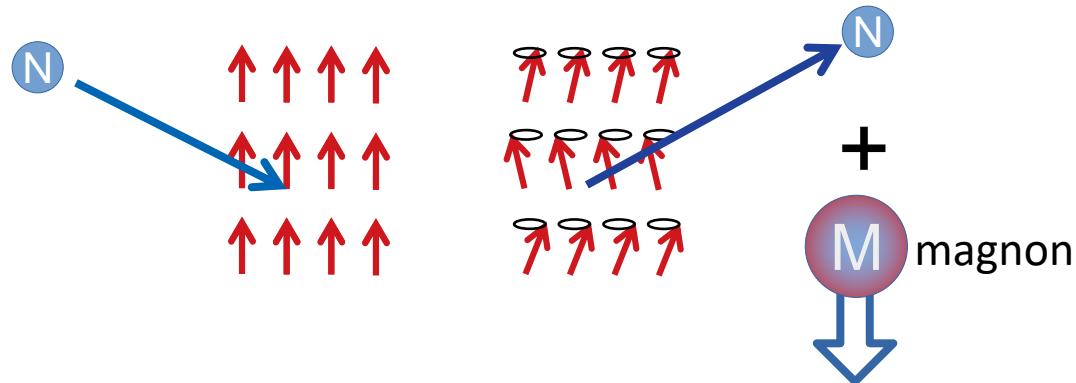
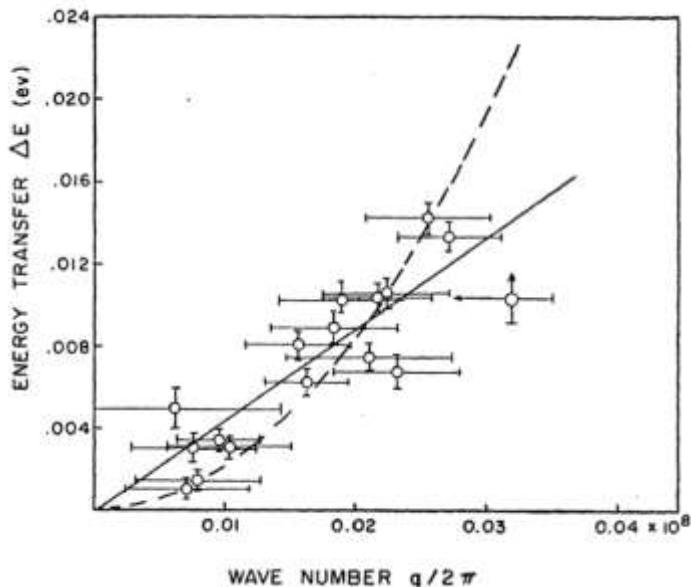
“The results permitted refutation of the simple models for interatomic forces then extant... these early results show surprisingly long-range behavior for the interatomic force system.”

Brockhouse, Nobel Prize lecture (1995)

Phonons: collective wave of atomic vibrations, we use the mathematics of particle physics to describe it

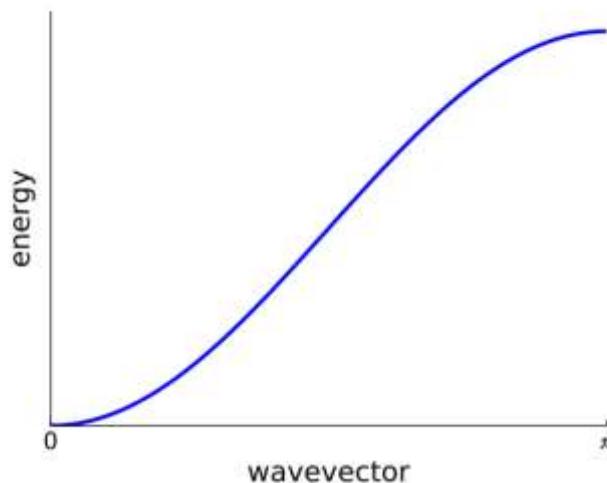


Brockhouse also measured magnetic quasiparticles (“Magnons”) in Fe_3O_4 (1957)

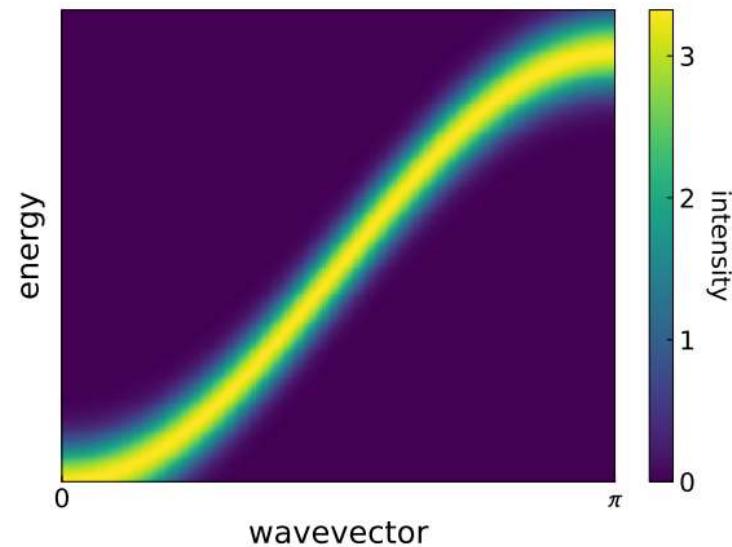


Quasiparticles were originally considered a calculation tool; neutron scattering showed they are *real*

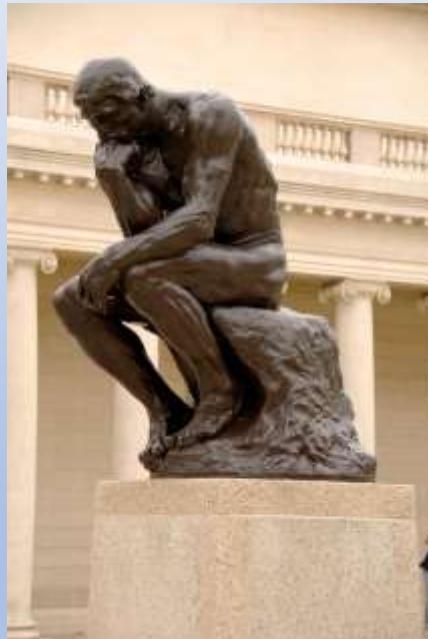
Theoretical dispersion



Measured neutron signal



Philosophical interlude:





Are quasiparticles real?

- We treat these as “emergent reality”
 - Can take on properties that are forbidden for individual parts

Examples: quasiparticles, 2nd law of thermodynamics, organic life

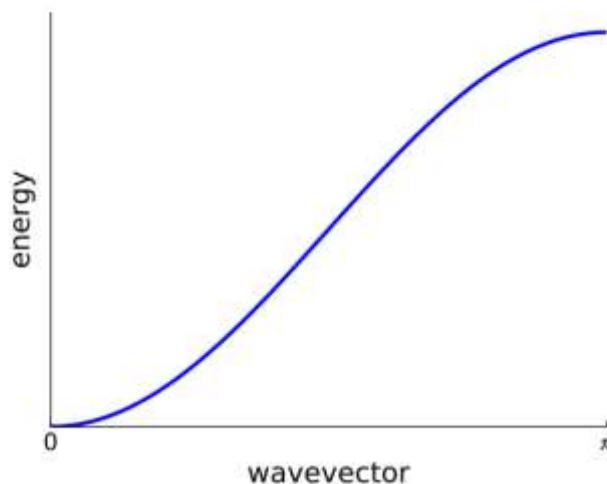
- Most of my colleagues (myself included) consider quasiparticles to be “real” objects.
- ALL our theories are incomplete approximations of underlying reality.
 - “quasiparticles” differ from “particles” in degree, not in kind.

“In quantum physics there is no logical way to distinguish a real particle from an excited state of the system that behaves like one. We therefore use the same word for both.”

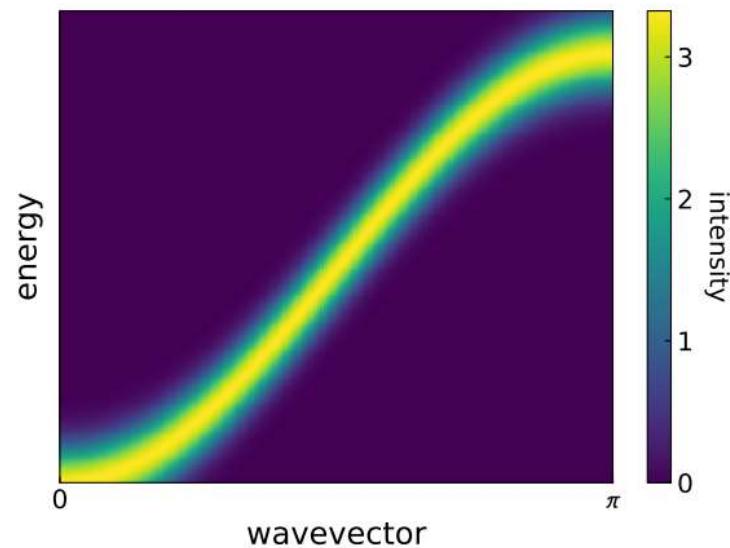
- R. Laughlin (Nobel Lecture, 1998)

Quasiparticles were originally considered a calculation tool; neutron scattering showed they are real

Theoretical dispersion



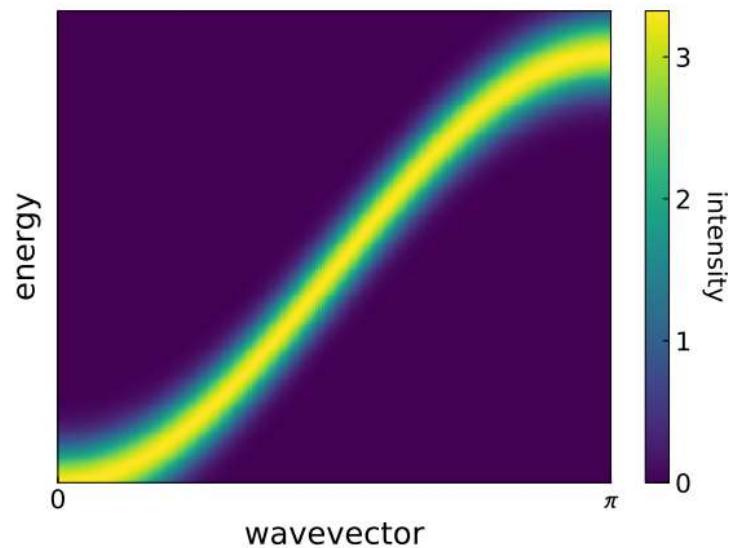
Measured neutron signal



Quasiparticles were originally considered a calculation tool; neutron scattering showed they are real

Scattered intensity is proportional to the
Fourier transform of spin correlation

$$\frac{d^2\sigma}{d\Omega d\omega} \propto \int \sum_l d^{-i\vec{Q}\cdot\vec{l}} \langle S_0(0)S_l(t) \rangle e^{-i\omega t} dt$$



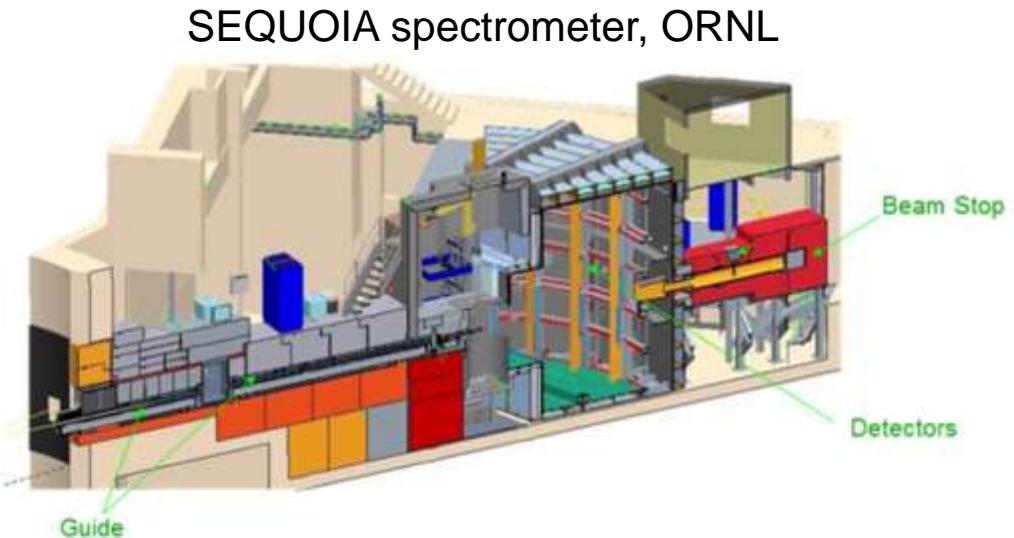
Part 2:

What neutrons are used for now

Old... vs New



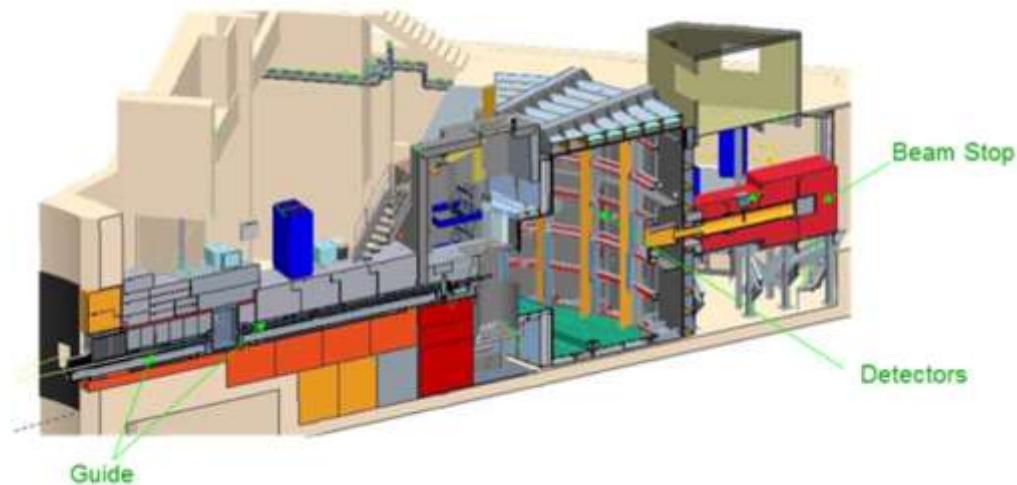
1 detector



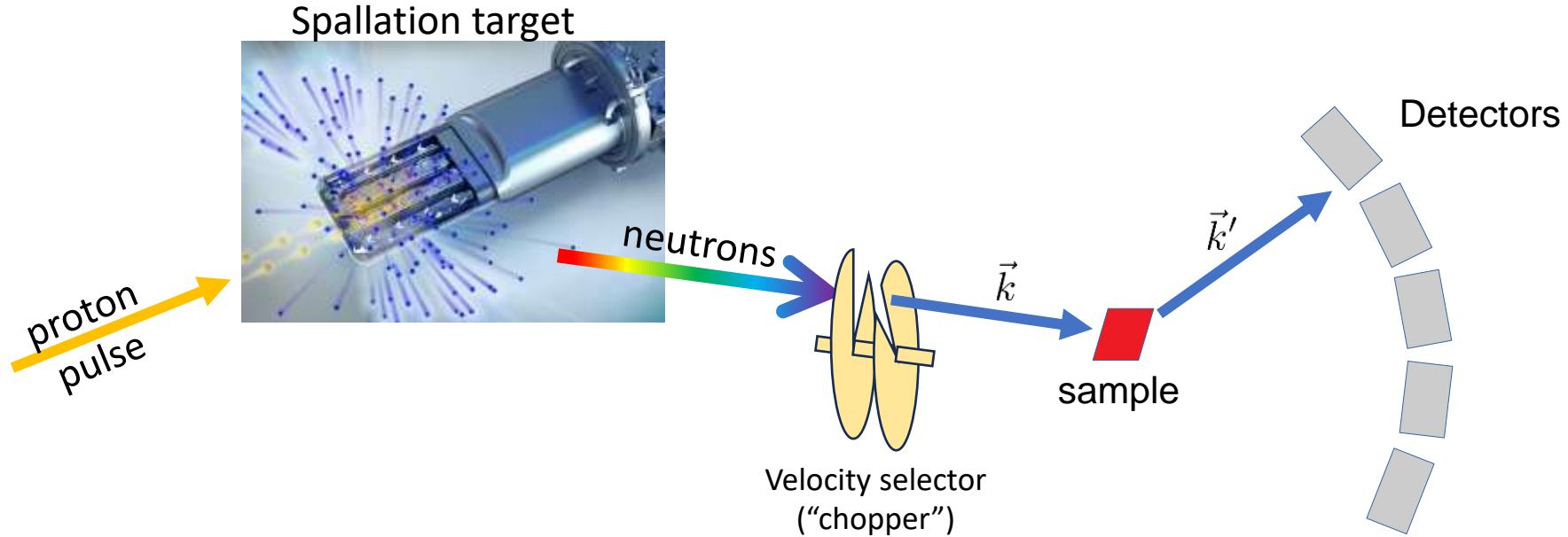
144,000 detector pixels

Time-of-flight technique: hit sample with a pulse of neutrons, and measure neutron velocity (more efficient)

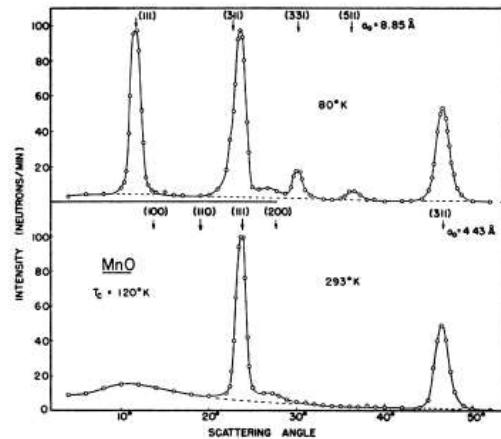
SEQUOIA spectrometer, ORNL



Time-of-flight technique: hit sample with a pulse of neutrons, and measure neutron velocity (more efficient)

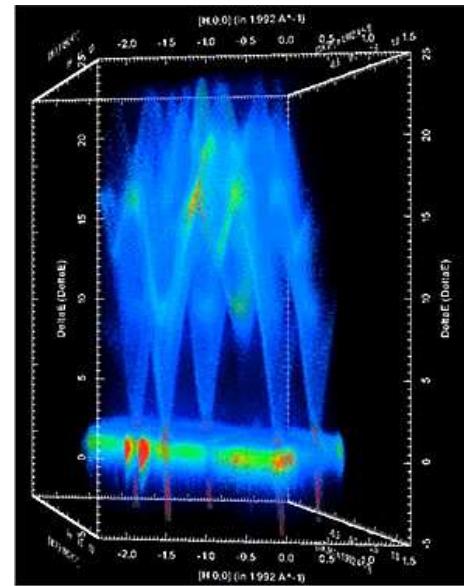


Old: 1D plots



New: Map out four-dimensional space:

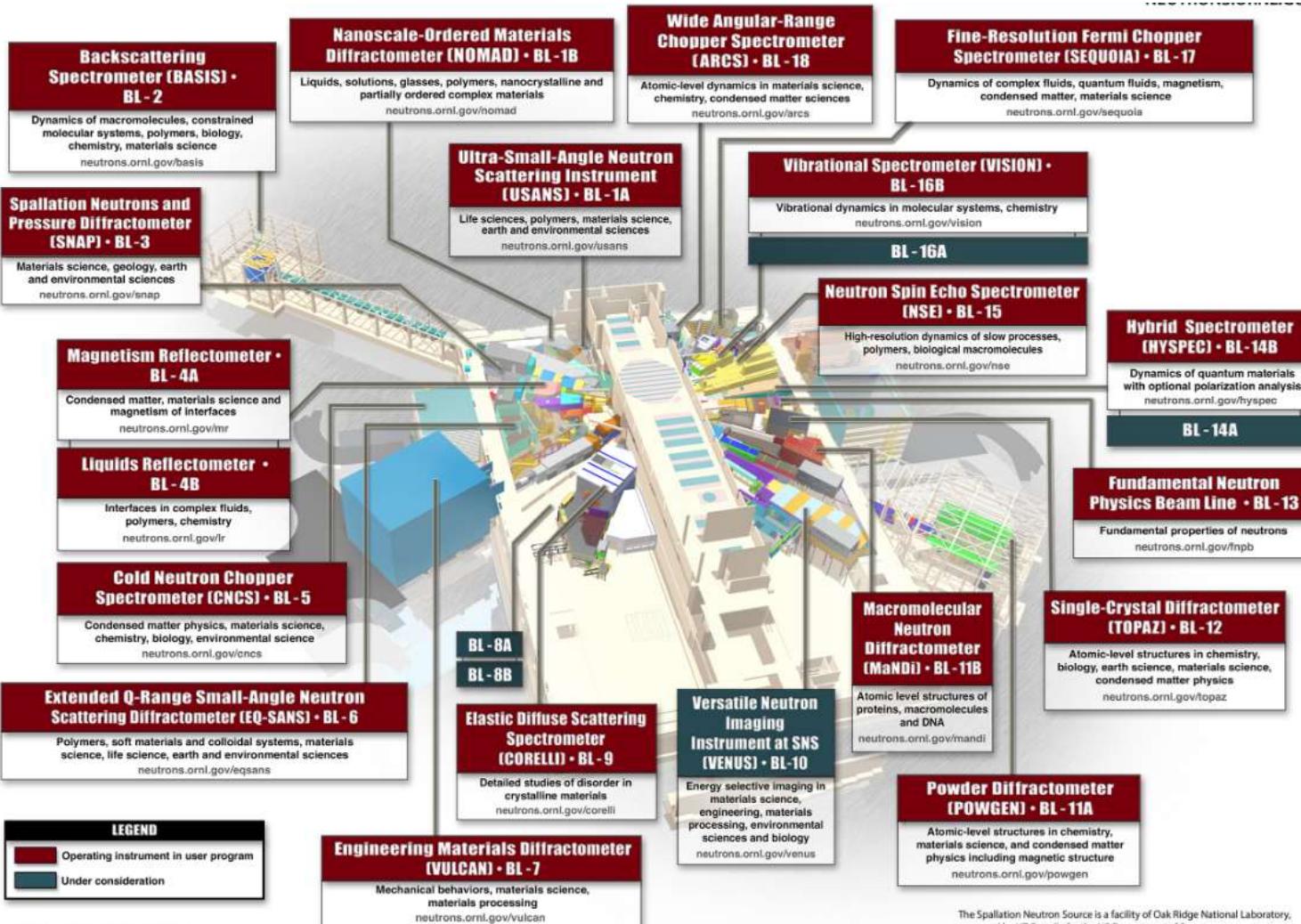
Three spatial dimensions,
Energy transfer, Intensity



European-spallationsource.se

Modern neutron facilities:





Neutron instruments worldwide:



"Neutrons for the Nation" report, Aps.org

Neutron facilities in the US:



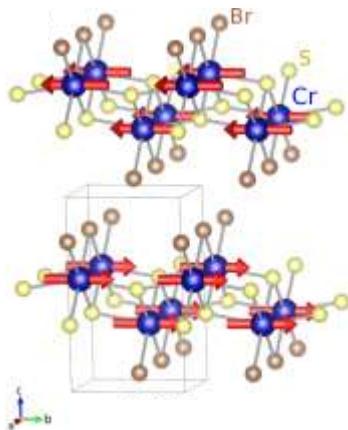
Oak Ridge, TN



Gaithersburg, MD

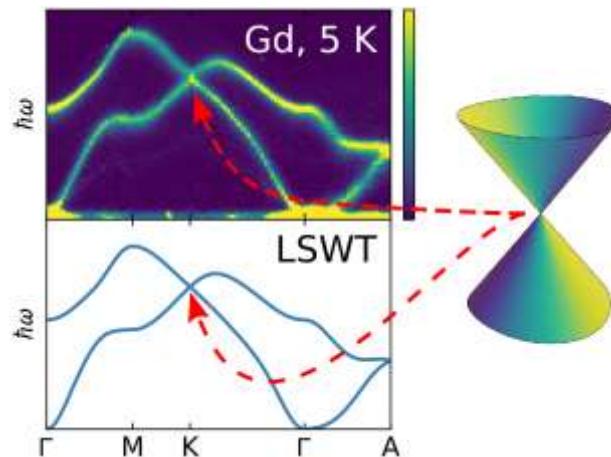
“Bread and butter” neutron experiments:

Static magnetic order



Scheie et al, *Adv. Science* (2022)

Magnetic dynamics

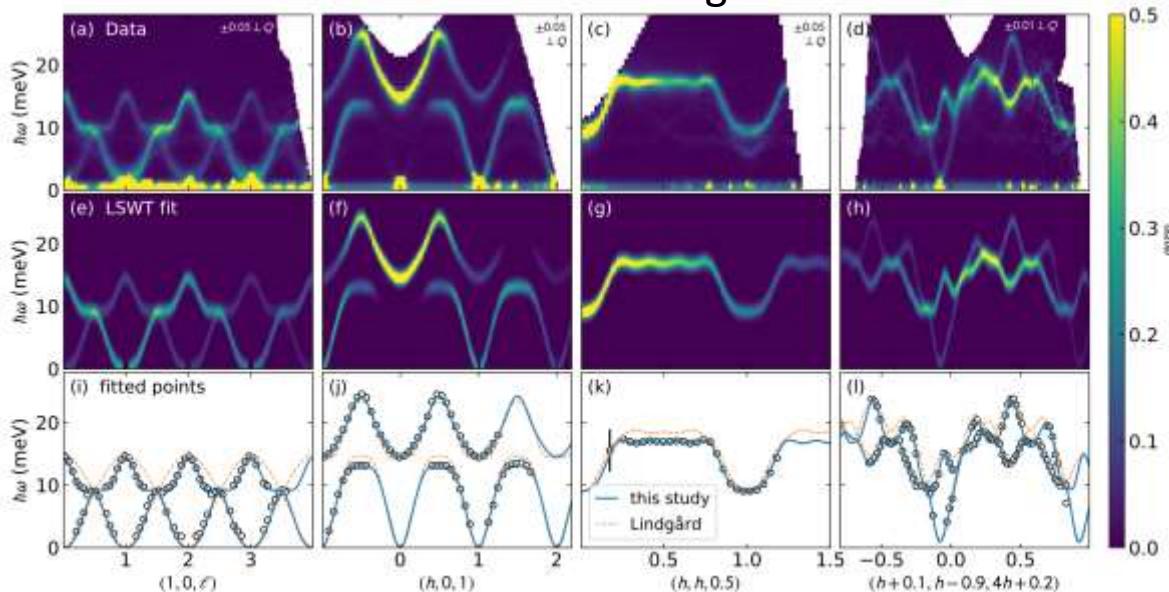


Scheie et al, *PRL* (2022)

Example: modeling and fitting the magnetic interactions of elemental gadolinium

Neutron data and fitted magnon model

data
model



Fitted exchange parameters

$J_1 = -138 \pm 8$	$J_{10} = -2 \pm 2$	$J_{19} = -5 \pm 3$
$J_2 = -174 \pm 4$	$J_{11} = 0 \pm 20$	$J_{20} = -5.6 \pm 1.3$
$J_3 = 50 \pm 20$	$J_{12} = -25 \pm 6$	$J_{21} = 5 \pm 9$
$J_4 = 41 \pm 6$	$J_{13} = 0 \pm 2$	$J_{22} = -5 \pm 7$
$J_5 = -10 \pm 20$	$J_{14} = -8 \pm 10$	$J_{23} = 29 \pm 7$
$J_6 = -14 \pm 4$	$J_{15} = 4 \pm 2$	$J_{24} = -5 \pm 9$
$J_7 = -4 \pm 2$	$J_{16} = 1 \pm 7$	$J_{25} = 1 \pm 8$
$J_8 = 0 \pm 3$	$J_{17} = 4 \pm 8$	$J_{26} = 7 \pm 2$
$J_9 = -10 \pm 20$	$J_{18} = 10 \pm 20$	

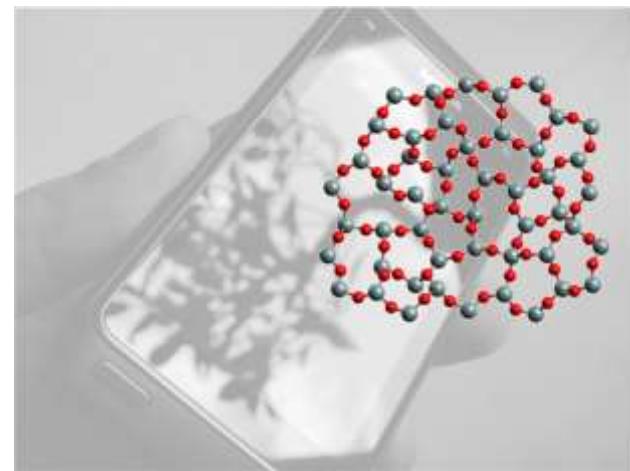
Other “Bread and butter” neutron experiments:

Locations of hydrogen



Kneller et al, *JCB* **295**, 50 (2020)

Contrast between elements



Shi et al, *N. Comm.* **14**, 13 (2023)

Part 3

Current uses of neutron scattering: exploring the unknown



Condensed matter theorists propose exotic states of matter—we try to find them.

Quantum spin liquid (QSL): a massively entangled magnetic state potentially useful for quantum computers

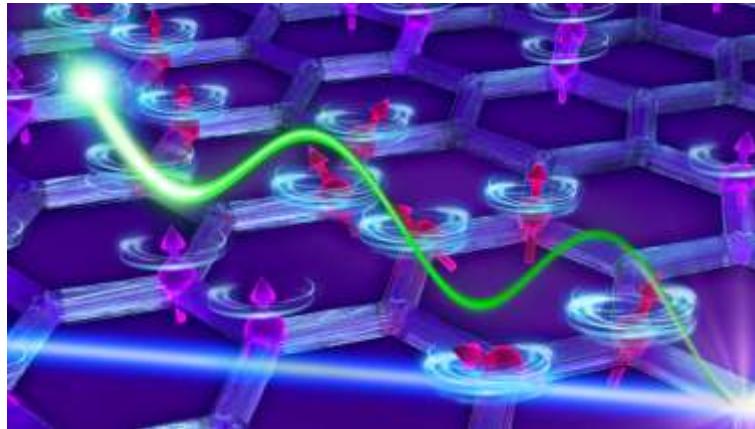
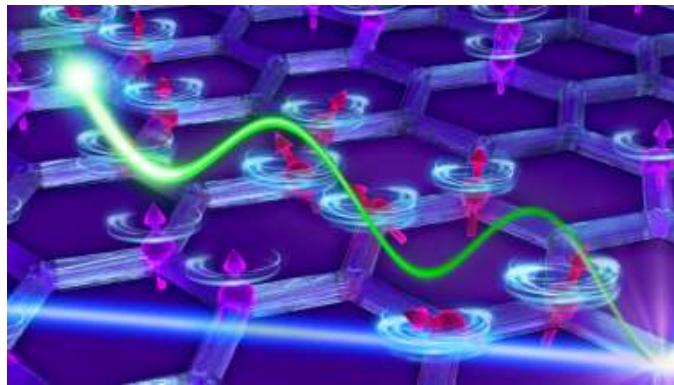


Image: neutrons.ornl.gov

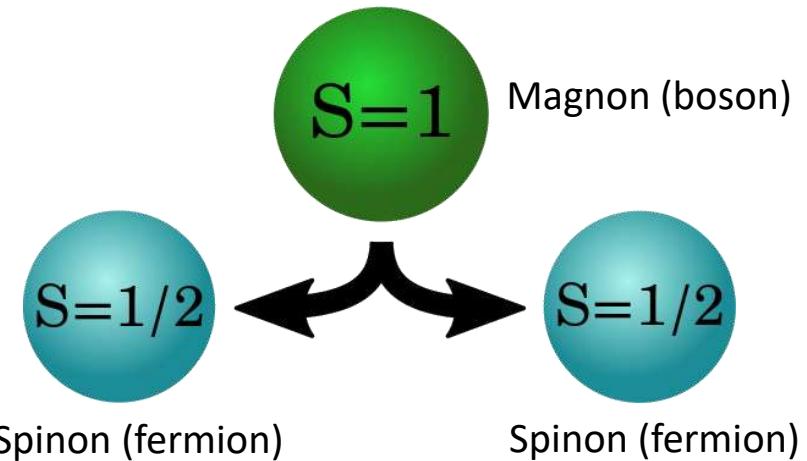
Theoretically, they should exist.
...but they are frustratingly difficult to identify experimentally.

Key features of QSL ground state:

(1) Long range entanglement:



(2) Fractional quasiparticles:

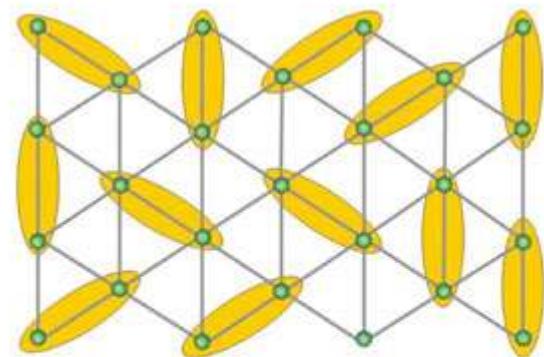


Original idea for a quantum spin liquid (QSL): Triangular lattice Heisenberg antiferromagnet

RESONATING VALENCE BONDS: A NEW KIND OF INSULATOR?*

P. W. Anderson
Bell Laboratories, Murray Hill, New Jersey 07974
and
Cavendish Laboratory, Cambridge, England

(Received December 5, 1972; Invited**)



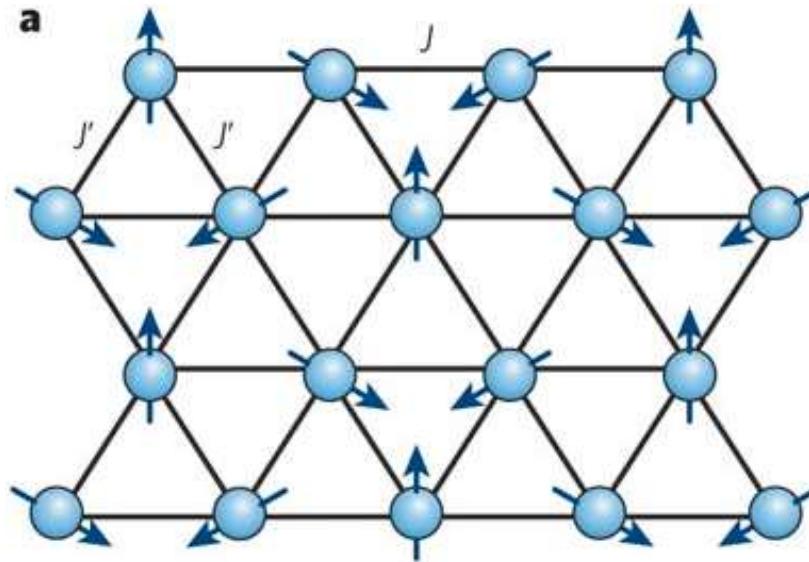
$$\text{Oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Image: Universität Augsburg, EP VI/EKM

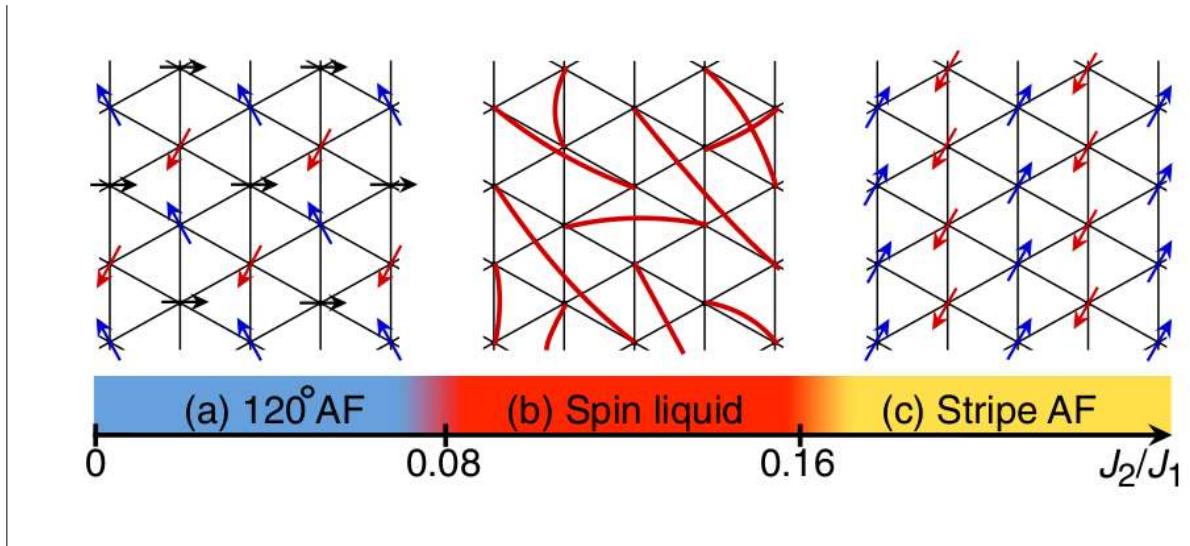
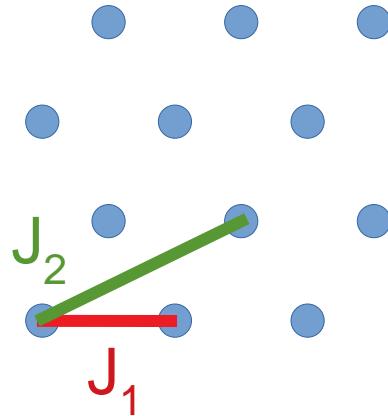
However, subsequent calculations showed that the Heisenberg triangular lattice antiferromagnet magnetically orders at low temperatures.

120° order

Capriotti et al, PRL (1999)
White and Chernyshev, PRL (2007)



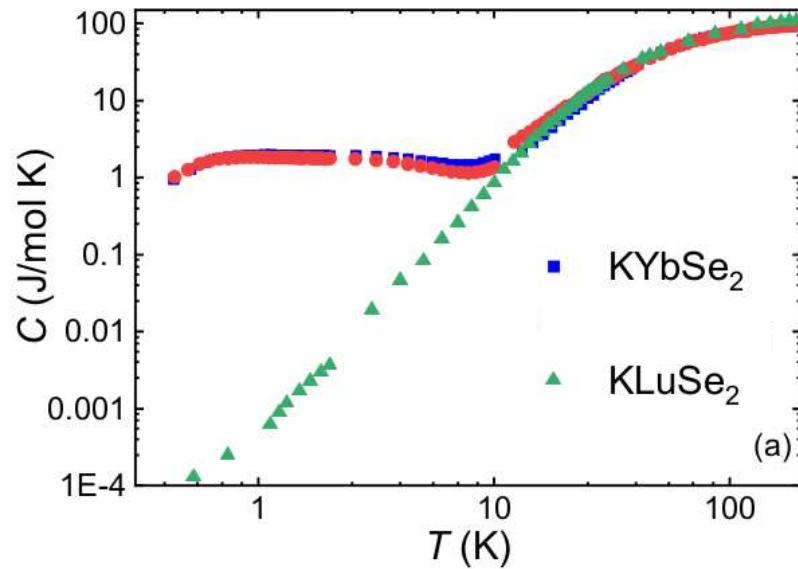
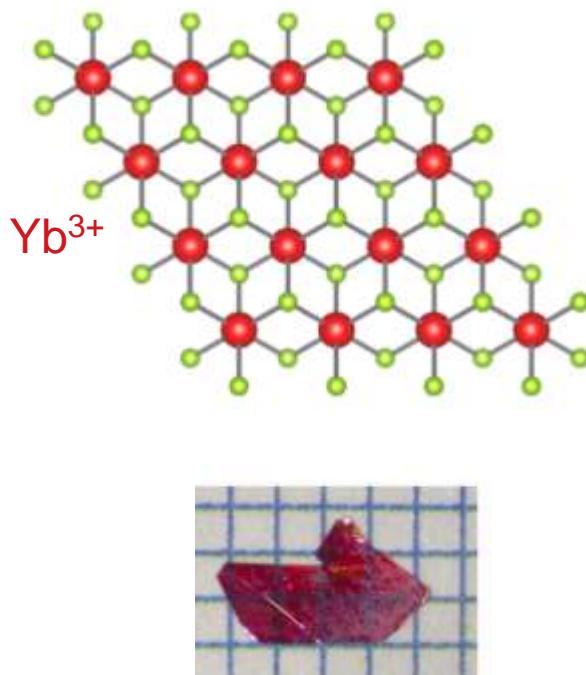
Instead, a small second neighbor exchange J_2
 $(0.06 \lesssim J_2/J_1 \lesssim 0.16)$ induces a QSL phase



Iqbal et al, PRB (2016)
Zhu et al, PRB (2015)

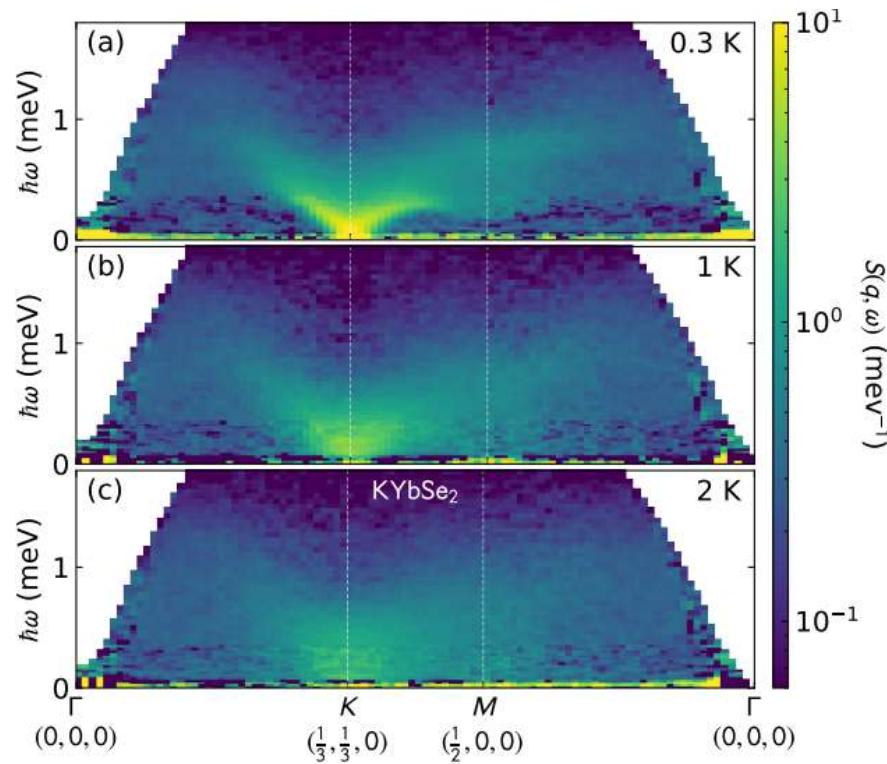
The precise nature of the QSL is debated

KYbSe_2 showed no magnetic order down to 400 mK. Maybe a QSL state?

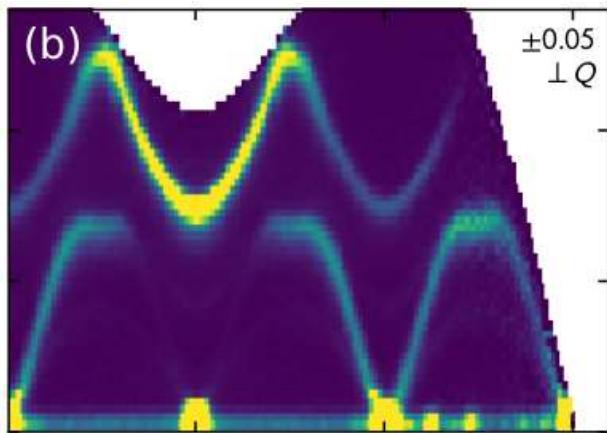


Xing, *APL Materials* (2021)

Neutron scattering experiment on KYbSe₂ @ CNCS



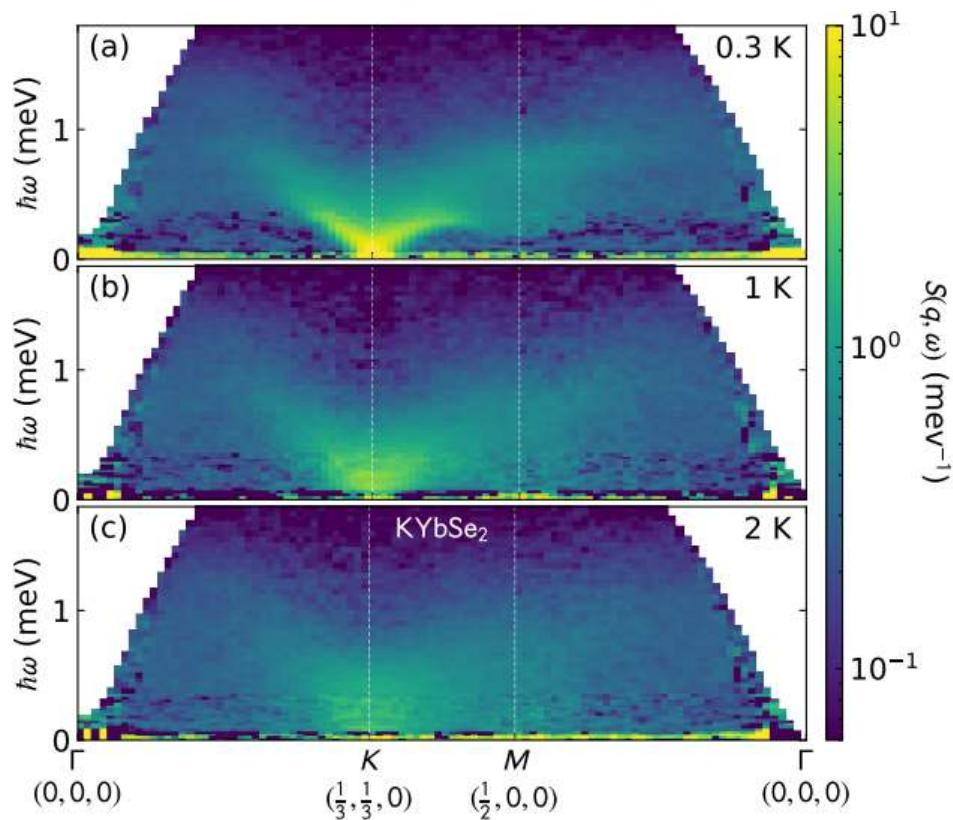
Ordinary magnetic material: well-defined modes



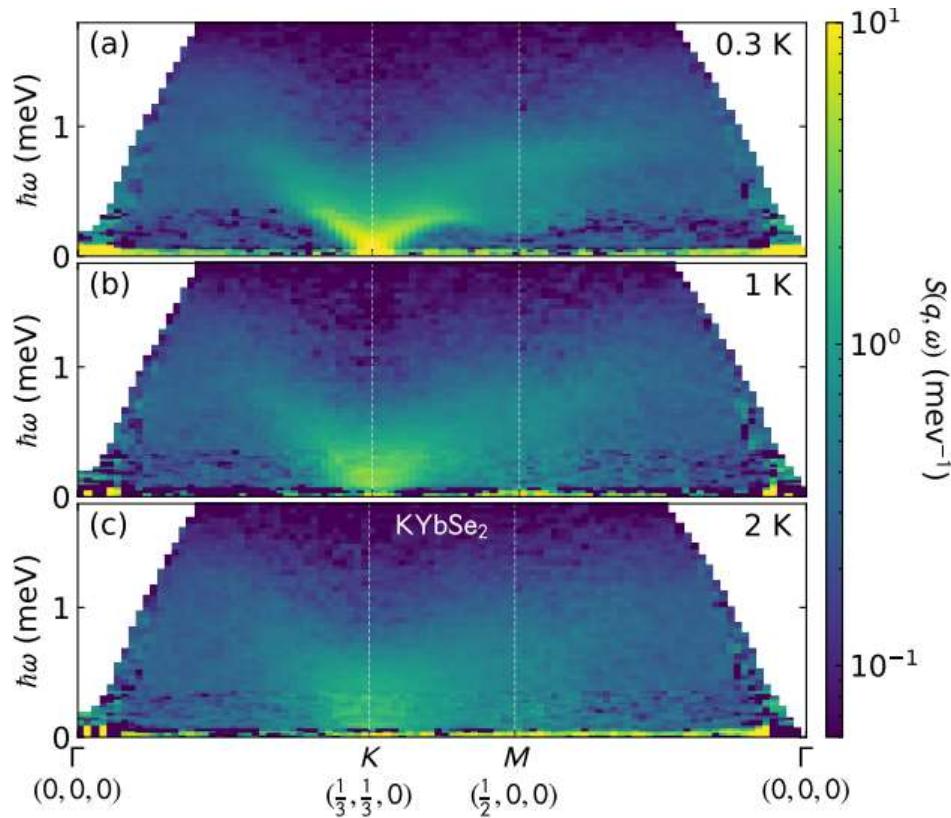
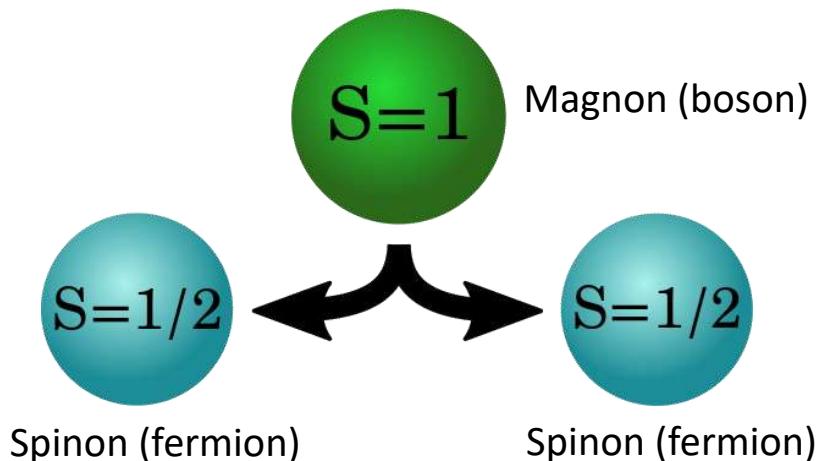
Elemental Gd, 5K

Scheie et al, PRB (2022)

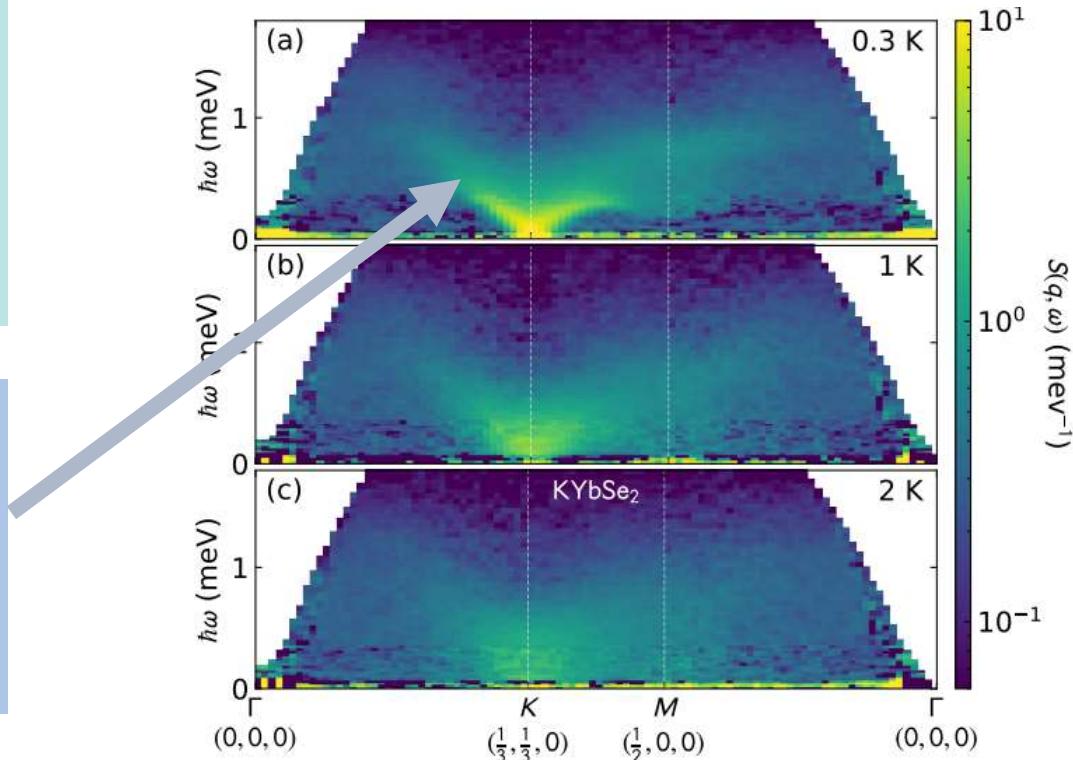
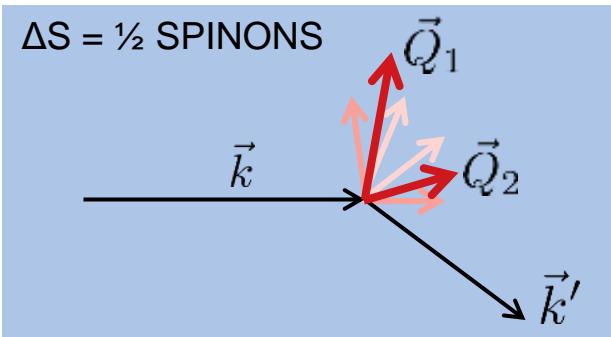
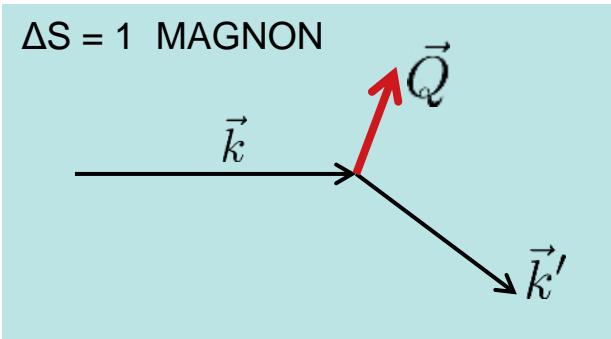
KYbSe₂: Diffuse continuum



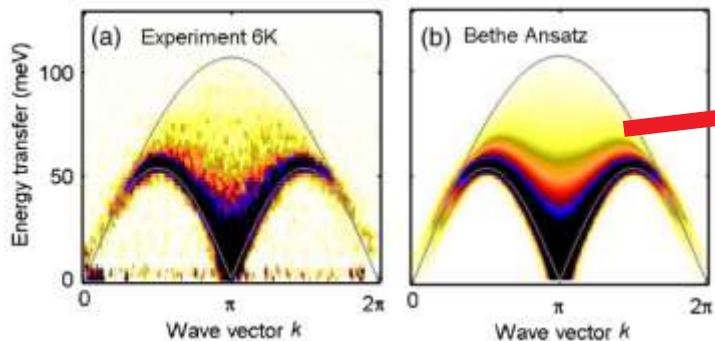
Continuum scattering can indicate fractional quasiparticles



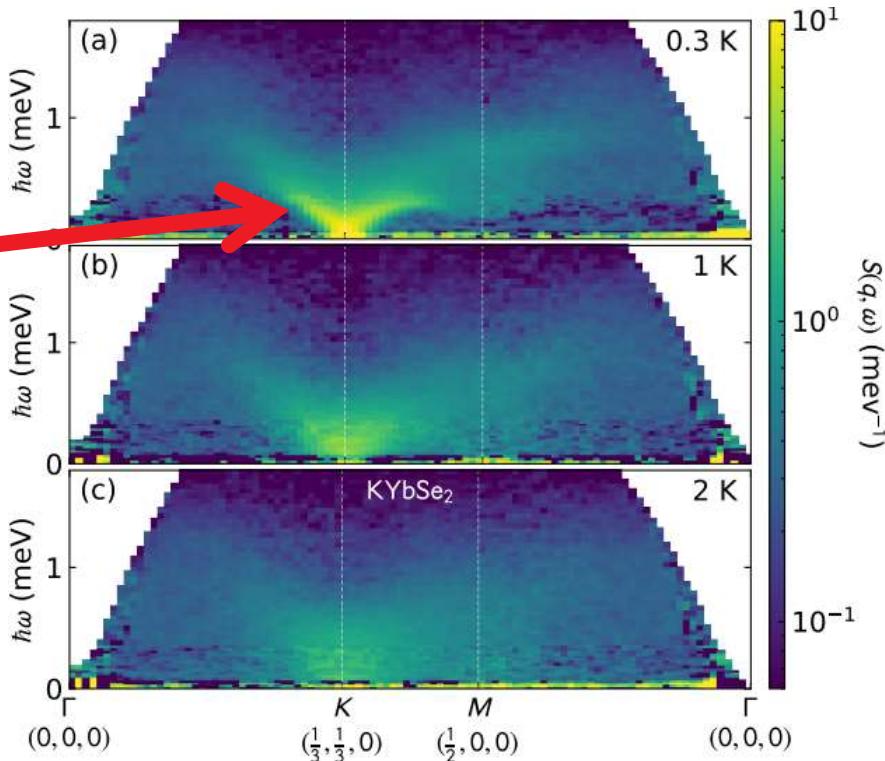
Continuum scattering can indicate fractional quasiparticles



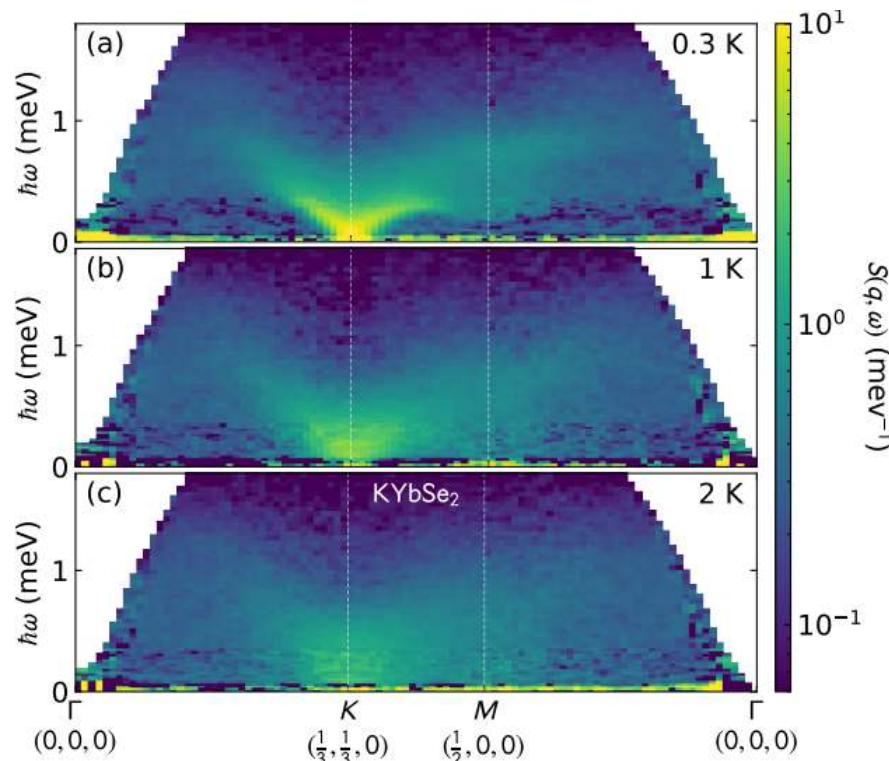
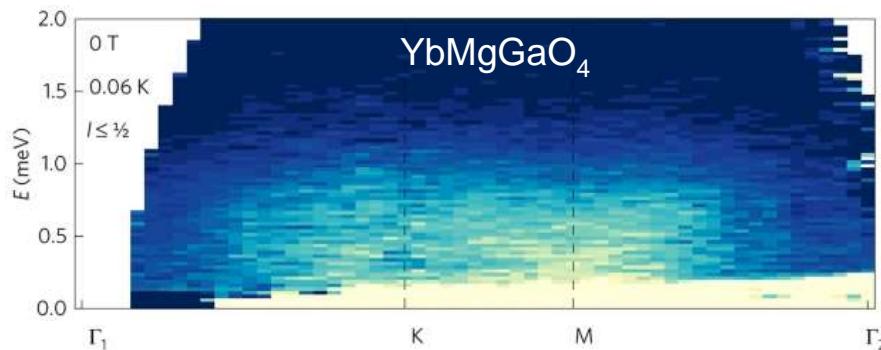
Continuum resembles the 1D spin chain spinon signal (known instance of fractionalization)...



Lake et al, Phys. Rev. Lett. 111, 137205 (2013)



But crystalline disorder also causes a diffuse continuum...



So... is KYbSe₂ a QSL?



We use two tools to answer this question:

- A) Magnetic Hamiltonian analysis
- B) Comparison to theoretical models

So... is KYbSe₂ a QSL?

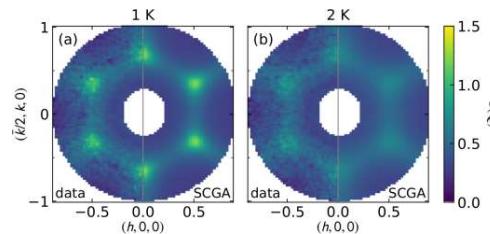
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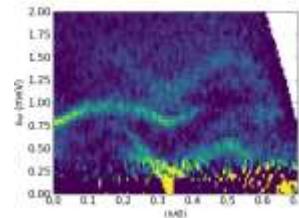
- A) Magnetic Hamiltonian analysis
- B) Comparison to theoretical models

Three methods to fit Hamiltonian:

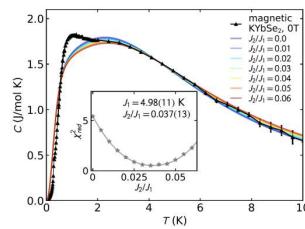
1. Paramagnetic diffuse scattering



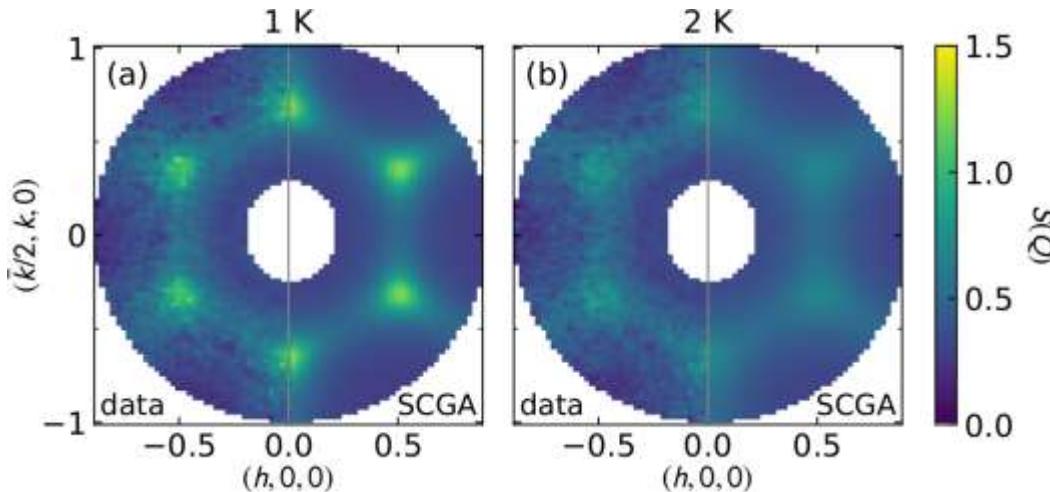
2. Nonlinear spin waves



3. Heat capacity expansion



Method 1: Onsager Reaction Field theory fitting the high temperature diffuse scattering



Fitted values (in K):

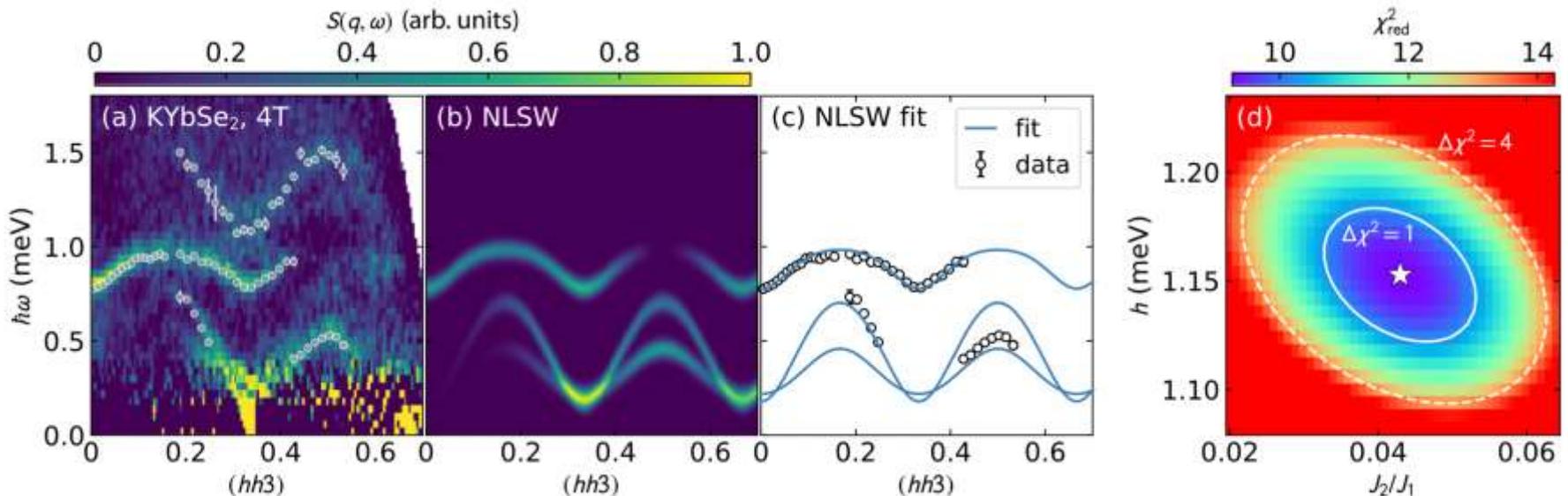
$$J_1 = \begin{pmatrix} 2.34(10) & 0.016(7) & 0.00(4) \\ 0.016(7) & 2.34(10) & 0.00(6) \\ 0.00(4) & 0.00(6) & 2.28(10) \end{pmatrix}$$

$$J_2 = 0.11(2)$$

(Classical method, so overall energy scale is wrong.)

Result: nearly isotropic exchange, with $J_2/J_1 = 0.047(7)$

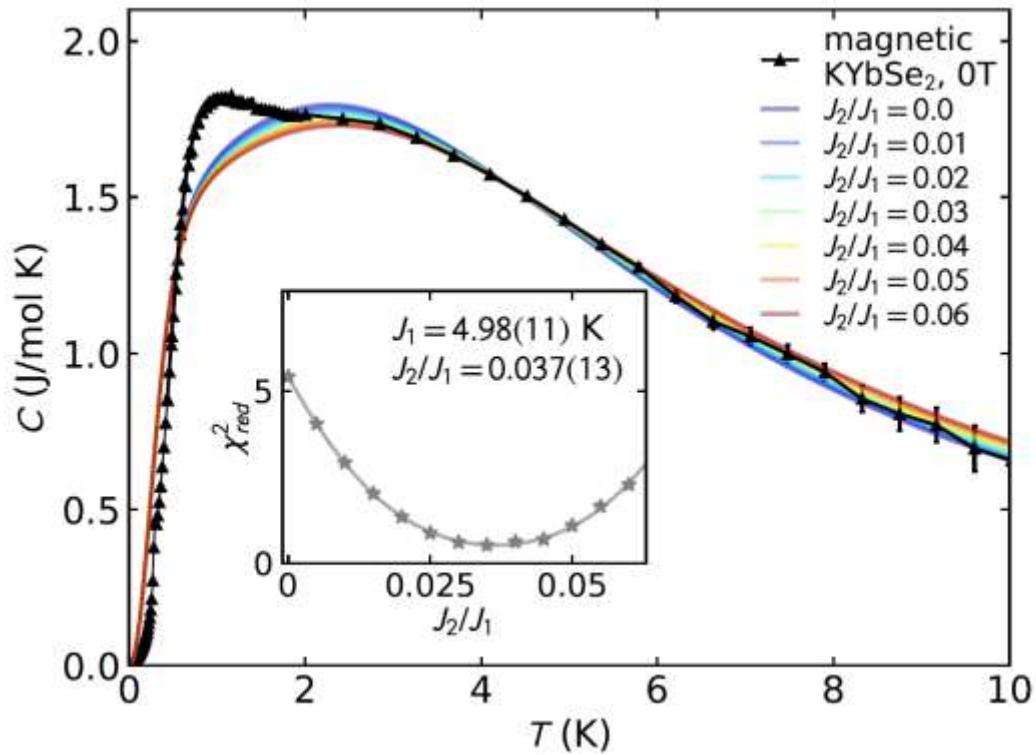
Method 2: Nonlinear spin wave fits in the 1/3 magnetization plateau



Result: $J_2/J_1 = 0.043 \pm 0.010$

Agrees with $J_2/J_1 = 0.047(7)$ from diffuse scattering fits

Method 3: Heat capacity expansion fits



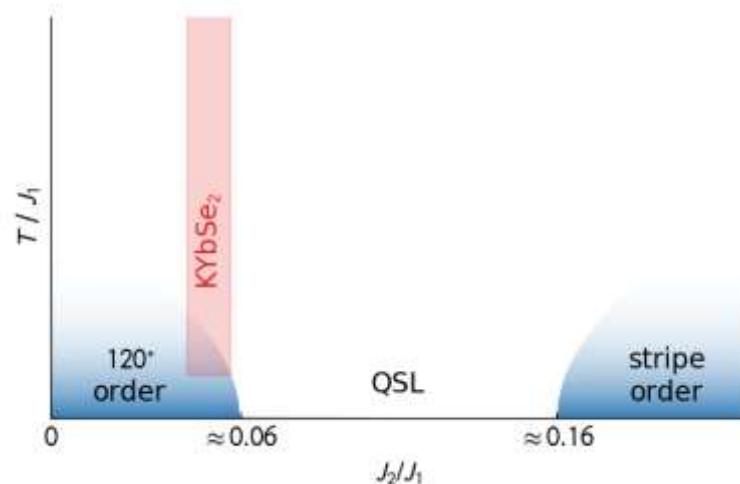
Fit between 2 K and 8 K

Courtesy Laura Messio @ Sorbonne University

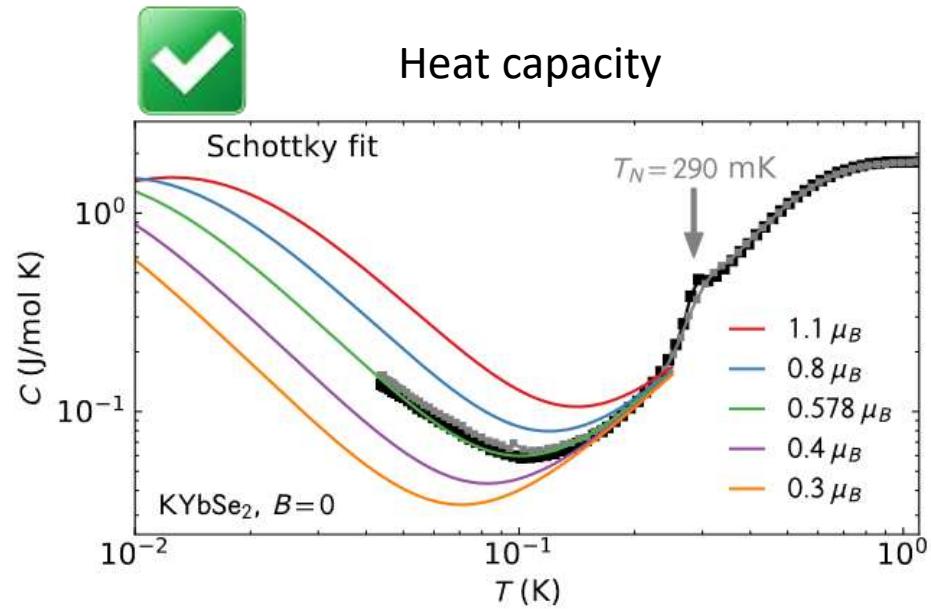
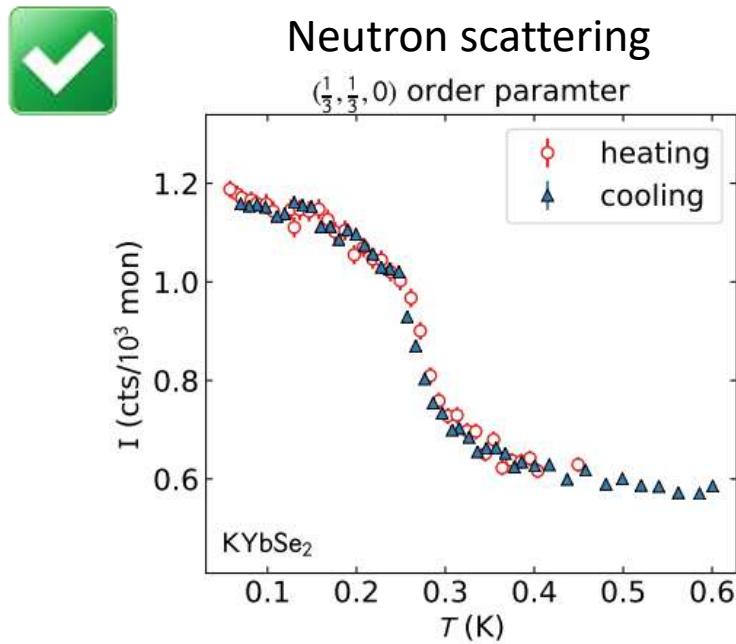
Result:
 $J_2/J_1 = 0.037 \pm 0.013$

All methods agree to within uncertainty: KYbSe₂ is very close to the QSL phase

Theoretical technique	J_1 (meV)	J_2/J_1
Onsager reaction field	NA	0.047 ± 0.007
Nonlinear spin waves	0.456 ± 0.013	0.043 ± 0.010
Heat capacity	0.429 ± 0.010	0.037 ± 0.013
Weighted mean:	0.438 ± 0.008	0.044 ± 0.005



However, if our model is correct, we should see 120° magnetic order at low temperatures. This was subsequently observed:



So... is KYbSe₂ a QSL?

We use two tools to answer this question:

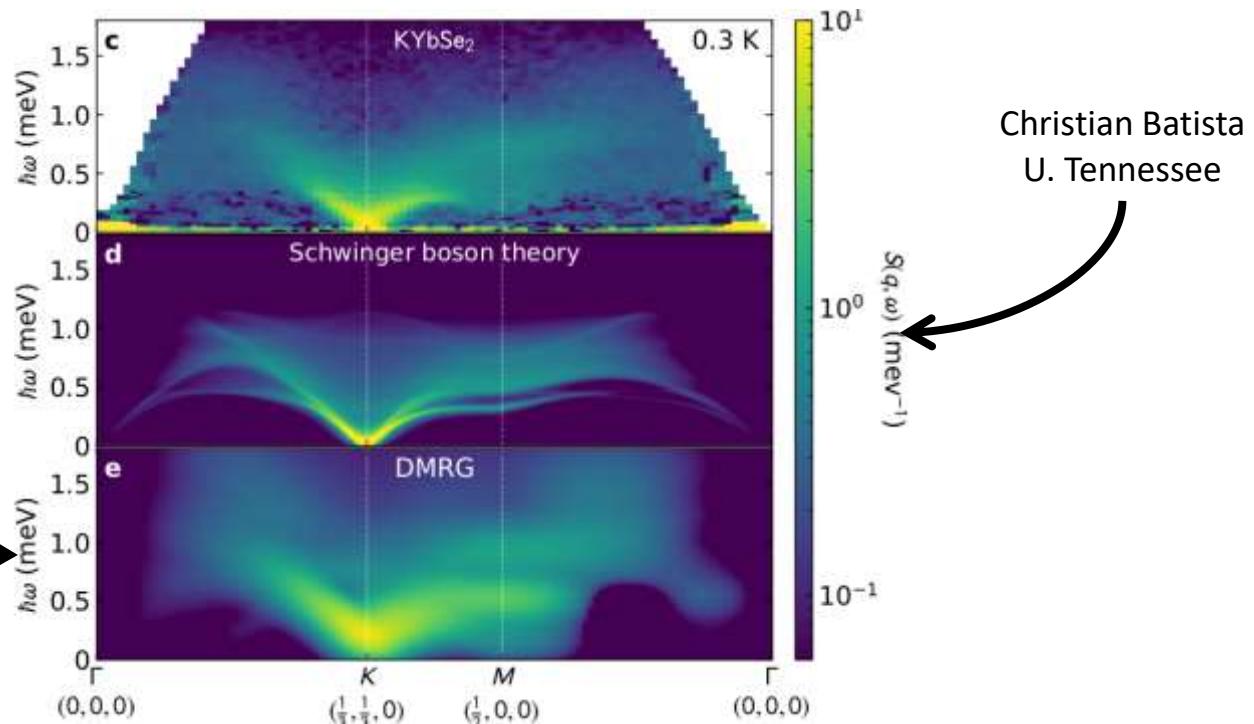
A) Magnetic Hamiltonian analysis



B) Comparison to theoretical models

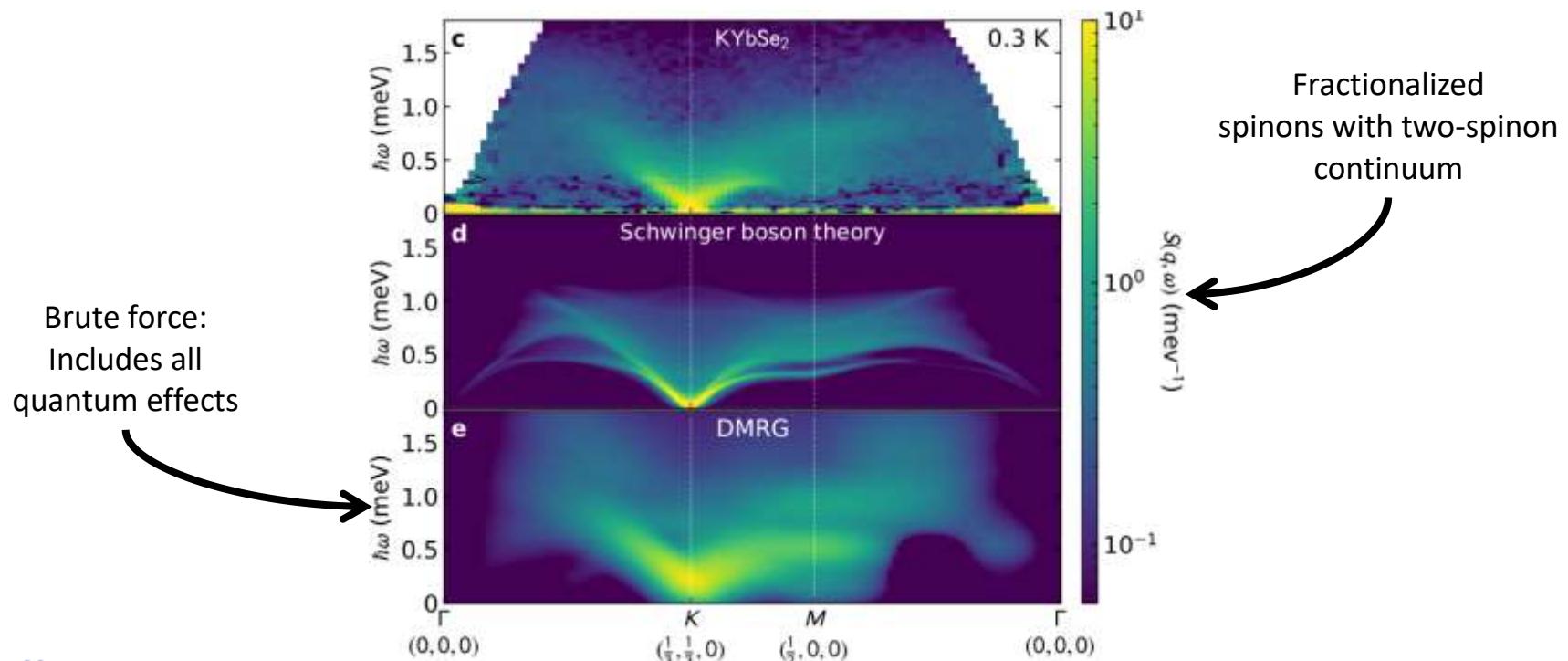
Theoretical calculations for a 2D triangular lattice with the fitted Hamiltonian closely resemble our data

Joel Moore
UC Berkeley

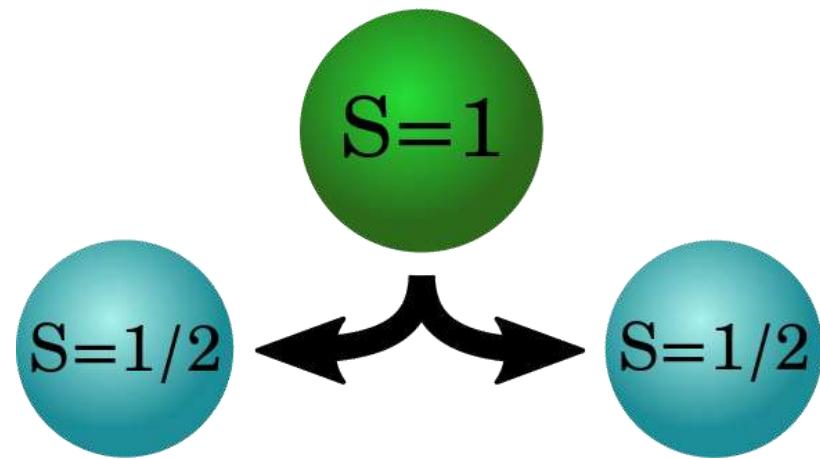
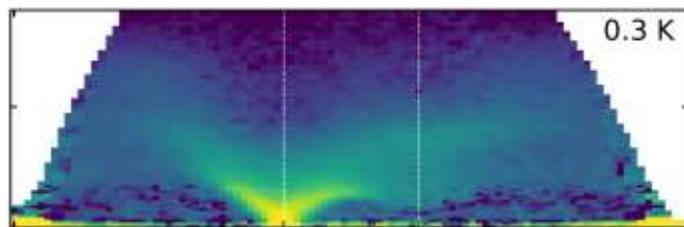


Christian Batista
U. Tennessee

Theoretical calculations for a 2D triangular lattice with the fitted Hamiltonian closely resemble our data

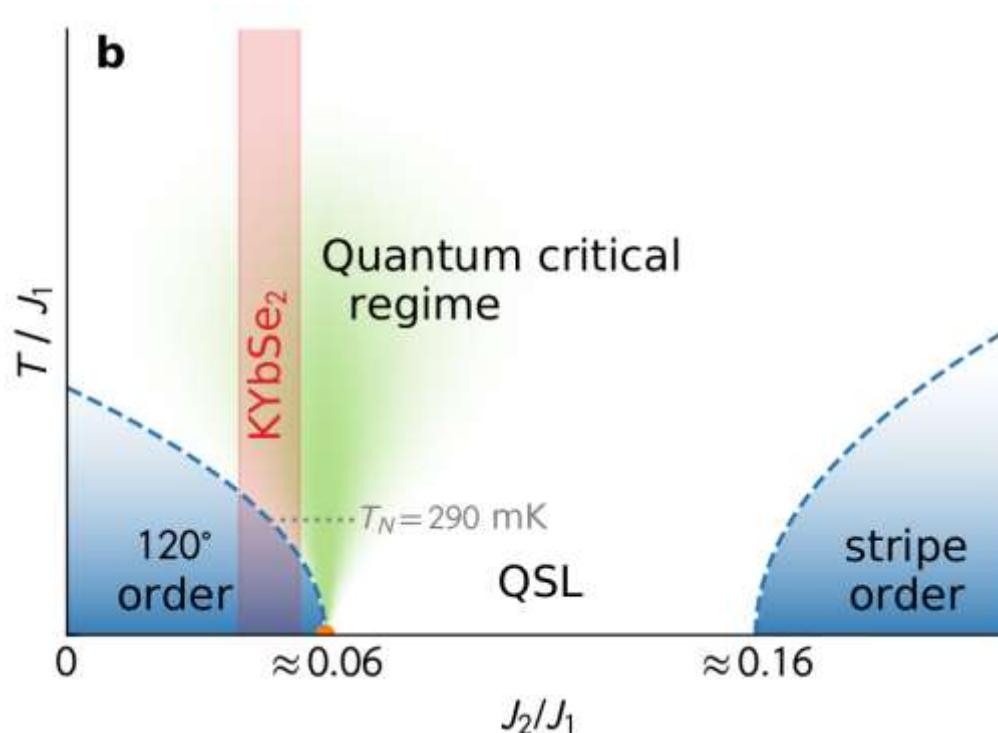


Although KYbSe₂ orders magnetically, its inelastic spectrum is dominated by fractionalized spinons



So is KYbSe₂ a QSL?

So is KYbSe₂ a QSL? It's a proximate QSL



Another project: measuring quantum entanglement using neutron scattering

To really understand quantum materials, we need a probe which can “see the quantumness” in a material.

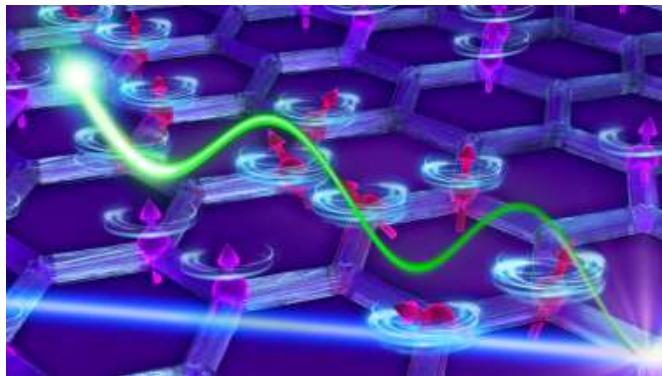
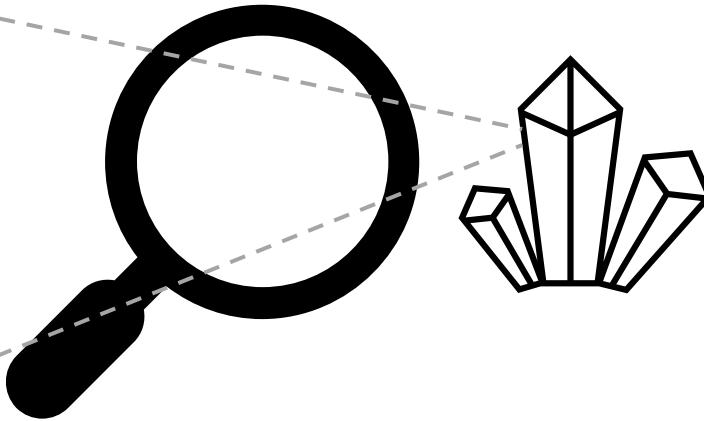
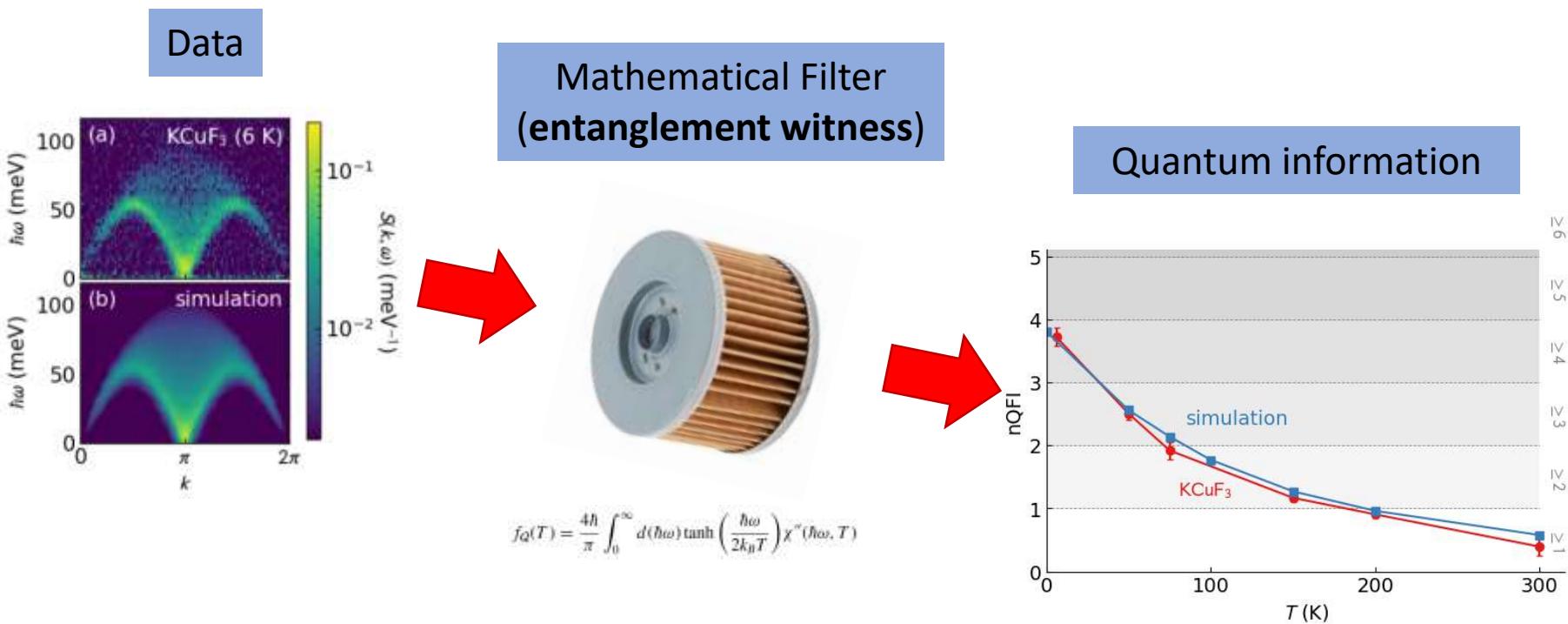


Image: neutrons.ornl.gov



entanglement witnesses



First “entanglement witness” proposed by John Stewart Bell in 1964, based on two-point (i.e., two-spin) correlation

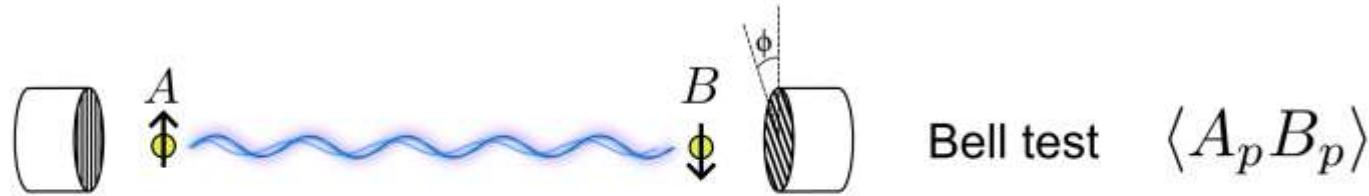
$$|\langle S_1^a S_2^b \rangle - \langle S_1^a S_2^c \rangle| \leq 1 + \langle S_1^b S_2^c \rangle$$

Classically, the inequality is always satisfied.
In the presence of quantum entanglement,
the inequality is violated.



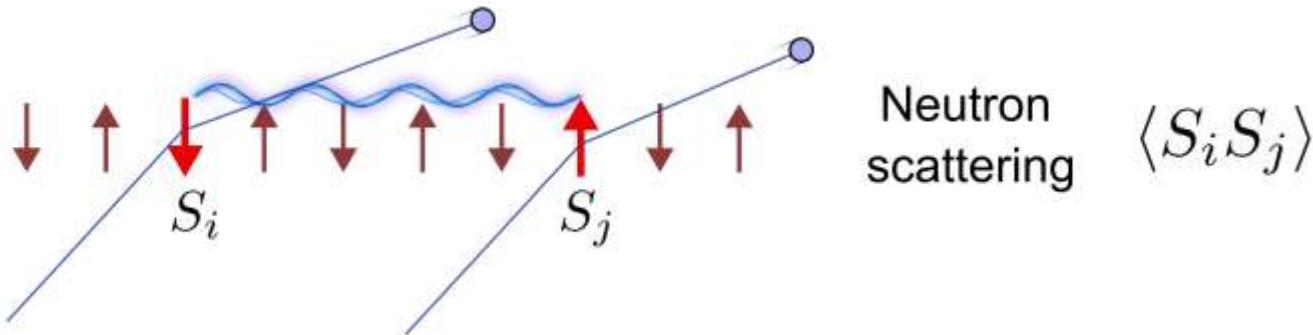
Wikipedia
Bell, *Rev. Mod. Phys.* (1966)

Entanglement information is encoded in two-point correlations:



Bell test

$$\langle A_p B_p \rangle$$



Neutron
scattering

$$\langle S_i S_j \rangle$$

Entanglement information is encoded in the neutron structure factor; the trick is knowing how to extract it.

$$S(\vec{q}, \omega) \propto \int e^{i\vec{q} \cdot \vec{r}} \langle \vec{S}_{i,\alpha} \vec{S}_{j,\beta} \rangle$$

We recently showed it is possible to measure quantum spin entanglement called Quantum Fisher Information (QFI) using neutron scattering

REVIEW *Adv. Quantum Technol.*, 2024, 2400196
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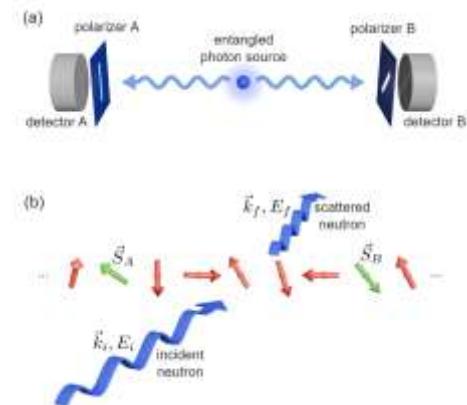
Witnessing Entanglement and Quantum Correlations in Condensed Matter: A Review

Pontus Laurell,^a Allen Scheie,^a Elbio Dagotto,^{c,d} and D. Alan Tennant^{a*}

Materials Today Quantum
Volume 5, March 2025, 100020

Tutorial: Extracting entanglement signatures from neutron spectroscopy

Allen Scheie^a, Pontus Laurell^b, Wolfgang Simeth^a, Elbio Dagotto^{c,d}, D. Alan Tennant^{c,e}



Quantum Fisher Information:

defines how precisely an operator \mathcal{O} can be measured

$$f_Q = 4 \left(\langle \psi | \mathcal{O}^\dagger \mathcal{O} | \psi \rangle - \langle \psi | \mathcal{O} | \psi \rangle^2 \right) \quad (\text{pure state QFI})$$

By the Cramér-Rao bound, this gives a lower bound on the number of entangled objects:

$$f_Q [\rho] > m (\lambda_{\max} - \lambda_{\min})^2$$

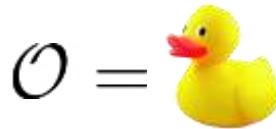
**m = entanglement depth,
i.e. ($m+1$) partite entanglement**

Quantum Fisher Information:

defines how precisely an operator \mathcal{O} can be measured

$$f_{\mathcal{Q}} = 4 \left(\langle \psi | \mathcal{O}^\dagger \mathcal{O} | \psi \rangle - \langle \psi | \mathcal{O} | \psi \rangle^2 \right) \quad (\text{pure state QFI})$$

\mathcal{O} can be *anything**:



$$f_{\mathcal{Q}} = 4 \left(\langle \begin{array}{c} \text{duck} \\ \dagger \end{array} | \begin{array}{c} \text{duck} \end{array} \rangle - \langle \begin{array}{c} \text{duck} \end{array} \rangle^2 \right)$$

*must be Hermitian

Quantum Fisher Information:

defines how precisely an operator \mathcal{O} can be measured

$$f_{\mathcal{Q}} = 4 \left(\langle \psi | \mathcal{O}^\dagger \mathcal{O} | \psi \rangle - \langle \psi | \mathcal{O} | \psi \rangle^2 \right) \quad (\text{pure state QFI})$$

For magnetic neutron scattering: $\mathcal{O} = S_\alpha(\mathbf{Q}) = \sum_j S^\alpha{}_j e^{i\mathbf{Q}\cdot\mathbf{r}_j}$

$$f_{\mathcal{Q}} = 4 \left(\underbrace{\langle \psi | S_\alpha^\dagger(\mathbf{Q}) S_\alpha(\mathbf{Q}) | \psi \rangle}_{\text{Total scattering}} - \underbrace{\langle \psi | S_\alpha(\mathbf{Q}) | \psi \rangle^2}_{\text{elastic scattering}} \right)$$

This is experimentally measurable!

Hauke et al provided a generalization to T>0:

$$f_Q = 2 \sum_{k,l} \frac{1}{Z} \frac{(e^{-\beta E_k} - e^{-\beta E_l})^2}{e^{-\beta E_k} + e^{-\beta E_l}} |\langle k | \hat{O} | l \rangle|^2 \quad (\text{thermal state QFI})$$

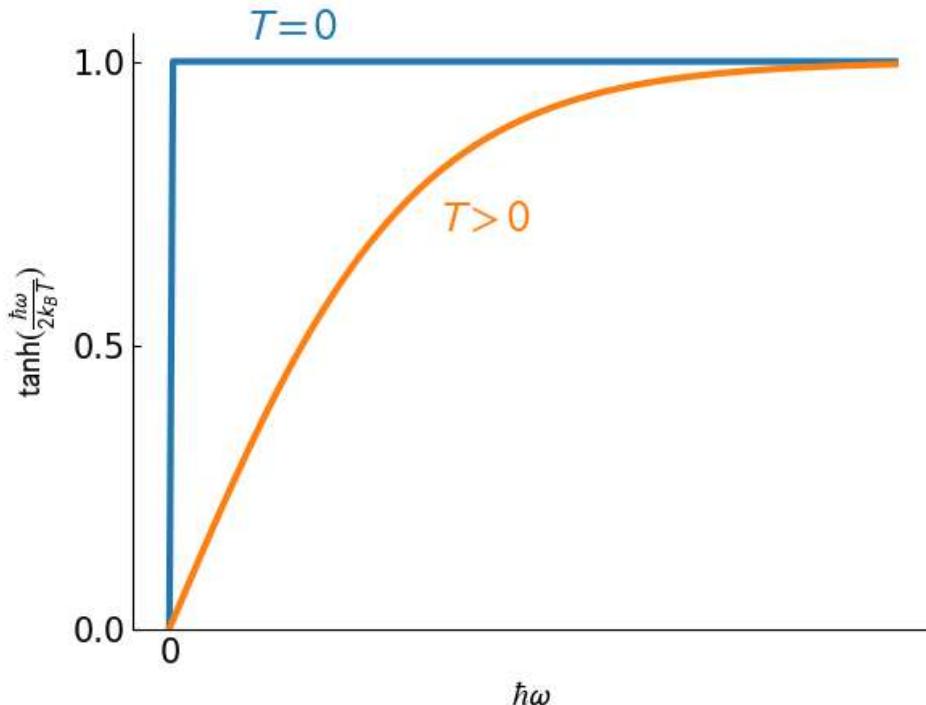
Can be related to dynamic susceptibility

$$\chi(\omega, T) = \frac{i}{\hbar} \int_0^\infty dt e^{i\omega t} \langle [\hat{O}(t), \hat{O}] \rangle$$

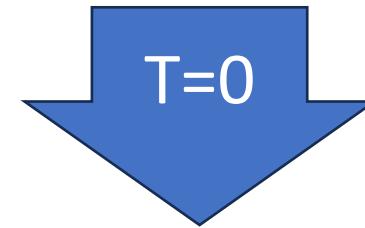


$$f_Q(T) = \frac{4\hbar}{\pi} \int_0^\infty d(\hbar\omega) \tanh\left(\frac{\hbar\omega}{2k_B T}\right) \chi''(\hbar\omega, T)$$

At T=0, thermal state QFI reduces to pure state QFI

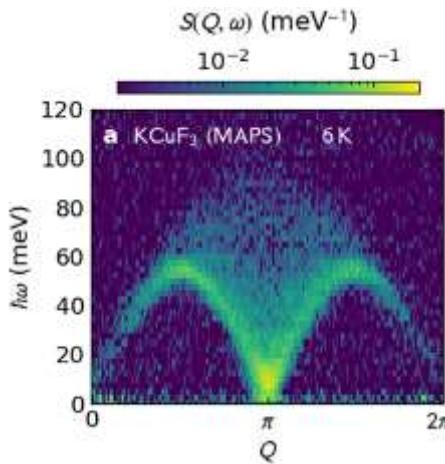


$$f_Q(T) = \frac{4\hbar}{\pi} \int_0^{\infty} d(\hbar\omega) \tanh\left(\frac{\hbar\omega}{2k_B T}\right) \chi''(\hbar\omega, T)$$



$$f_Q = 4 \left(\langle S^2 \rangle - \langle S \rangle^2 \right)$$

Quantum Fisher Information gives a lower bound on the number of entangled spins, and is measurable with neutron scattering



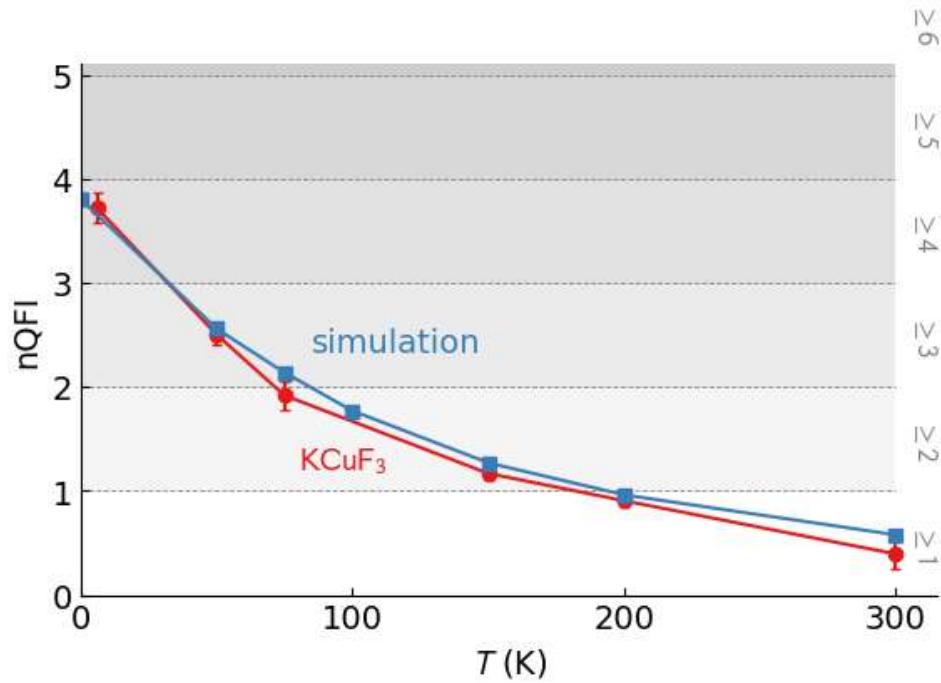
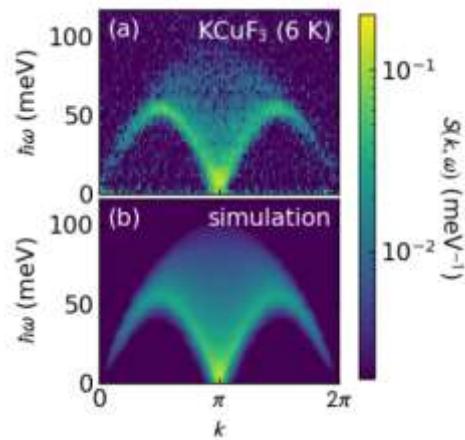
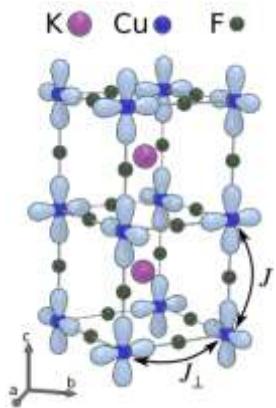
$$f_{\mathcal{Q}} [\mathbf{Q}, T] = 4 \int_0^{\infty} d(\hbar\omega) \left[\tanh \left(\frac{\hbar\omega}{2k_B T} \right) \left(1 - e^{-\hbar\omega/k_B T} \right) S_{\alpha\alpha}(\mathbf{Q}, \omega) \right]$$

if a system has nQFI > m ,
it is at least $(m+1)$ -partite entangled.

By the Cramer Rao bound:

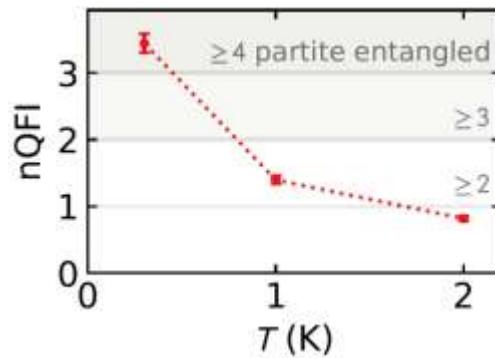
$$\text{nQFI} = \frac{f_{\mathcal{Q}}}{4S^2} > m \quad m = \begin{matrix} \text{entanglement} \\ \text{depth} \end{matrix}$$

We tested QFI on 1D spin chain KCuF_3 , witnessing ≥ 4 partite entanglement at low temperatures

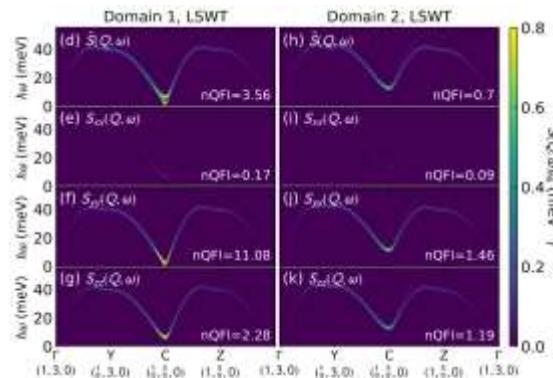


We've also measured quantum entanglement in

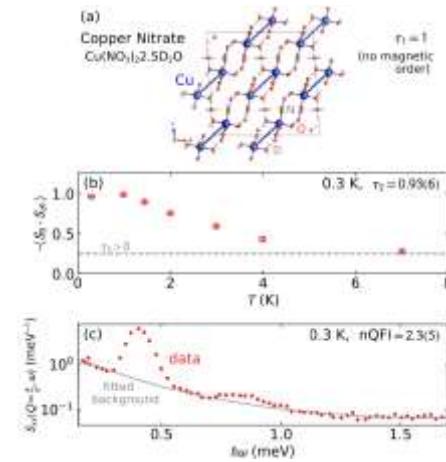
KYbSe₂



NiPS₃



Copper nitrate
(dimer)



Is it crazy that integrating conventional neutron scattering tells us how many particles are entangled?

Some trivial examples demonstrate that nQFI does scale with entanglement:

$$\text{nQFI} = 4 \left(\langle \psi | \mathcal{O}^\dagger \mathcal{O} | \psi \rangle - \langle \psi | \mathcal{O} | \psi \rangle^2 \right)$$

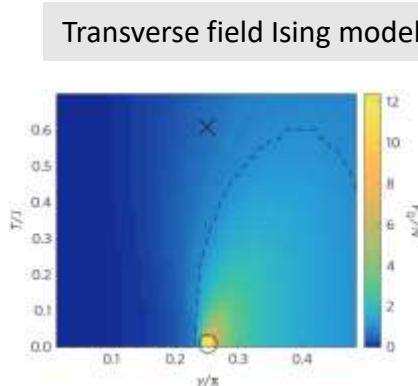
wavefunction	operator	nQFI
Free spin $\frac{1}{\sqrt{2}}(\uparrow\rangle + \downarrow\rangle)$	$\mathcal{O} = S^z$	1
Singlet $\frac{1}{\sqrt{2}}(\uparrow\downarrow\rangle - \downarrow\uparrow\rangle)$	$\mathcal{O} = \sum_j S_j^z e^{i\pi r_j}$	2
GHZ $\frac{1}{\sqrt{2}}(\uparrow\uparrow\uparrow\rangle - \downarrow\downarrow\downarrow\rangle)$	$\mathcal{O} = \sum_j S_j^z$	3
Cat state $\frac{1}{\sqrt{2}}(\uparrow\downarrow\dots N\rangle - \downarrow\uparrow\dots N\rangle)$	$\mathcal{O} = \sum_j S_j^z e^{i\pi r_j}$	N

There's a way in which the QFI measures the “cat-ness” of a ground state

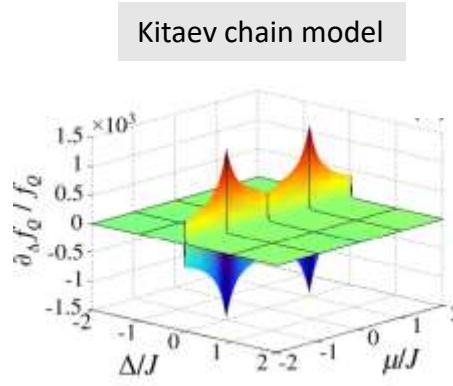
$$\frac{1}{\sqrt{2}} \left(| \text{cat} \rangle - | \text{bird} \rangle \right)$$

Q: But what does QFI really tell us?

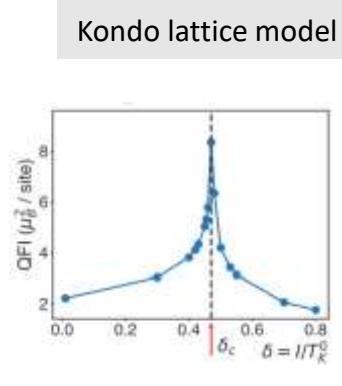
- Allows us to rule out trivial un-entangled states.
 - Doesn't necessarily rule out a magnon spectrum
- Diverging QFI seems to be related to quantum criticality



Hauke, N. Phys. (2016)



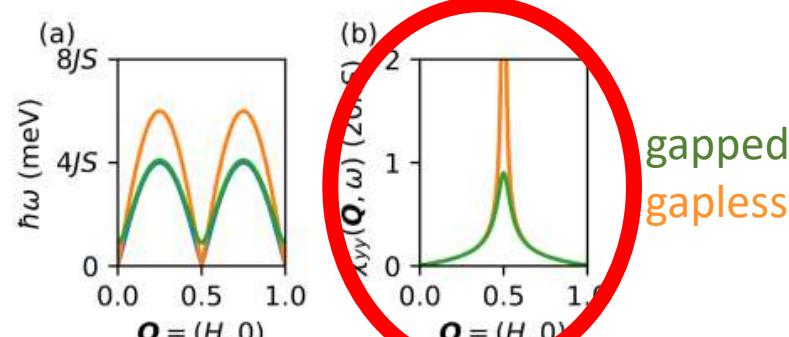
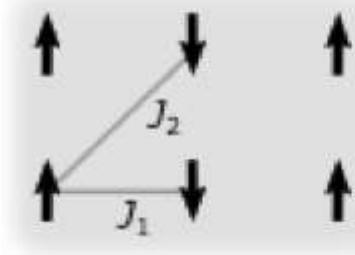
Pezze, PRL (2017)



Fang, ArXiv:2402.18552 (2024)

Implication: many conventional antiferromagnets are highly entangled states!

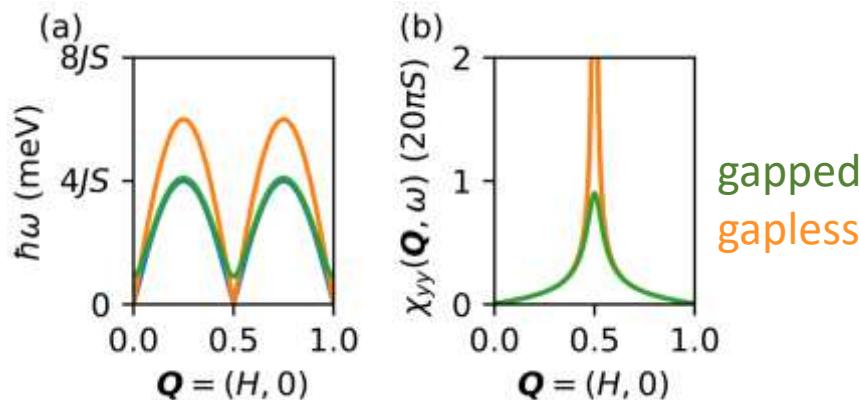
2D Square lattice antiferromagnet: divergent inelastic intensity at ordering vector



Implies divergent QFI $f_Q = 4 (\langle \psi | S_\alpha^\dagger(\mathbf{Q}) S_\alpha(\mathbf{Q}) | \psi \rangle - \langle \psi | S_\alpha(\mathbf{Q}) | \psi \rangle^2)$

What's going on?

The Bogoliubov transform also diverges at the ordering wavevector, so the diagonal basis is dramatically rotated from the Neel basis (implying entanglement!)



$$u_{\mathbf{q}} = \cosh(\theta_{\mathbf{q}}) = \sqrt{\frac{A_{\mathbf{q}} + \hbar\omega_{\mathbf{q}}}{2\hbar\omega_{\mathbf{q}}}}$$

$$v_{\mathbf{q}} = \sinh(\theta_{\mathbf{q}}) = s \cdot \sqrt{\frac{A_{\mathbf{q}} - \hbar\omega_{\mathbf{q}}}{2\hbar\omega_{\mathbf{q}}}}$$



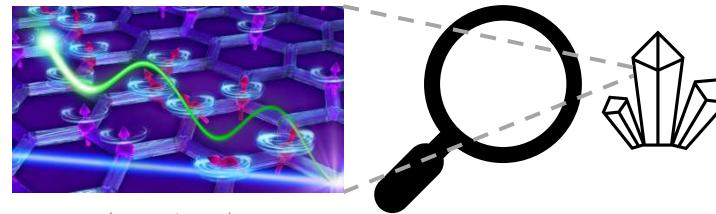
Forthcoming paper with Wolfgang Simeth
(LANL postdoc)

Summary of entanglement witnesses:

- We've successfully witnessed solid state entanglement using spectroscopy.
- These are *model-independent*: don't have to know or solve the underlying Hamiltonian.
- This doesn't tell us everything about the underlying quantum state, but it gives some valuable information.

But really, these are just baby steps.

Scheie, Laurell et al, PRB (2021)
Laurell, Scheie et al, PRL (2021)
Laruell et al, *Adv. Quant. Tech.* (2024)
Scheie et al, *Mat. Tod. Quant.* (2025)



Closing thoughts:

- Neutron scattering has historically revolutionized the way we think about condensed matter.
- This will continue to be the case: neutron scattering is one of the most direct ways to probe atomic scale quantum properties.
 - Quantum spin liquids
 - Many-body quantum entanglement and decoherence
 - Other states: unconventional superconductivity