FAU

Friedrich-Alexander-Universität
Faculty of Engineering

Department Materials Science

WW8: Materials Simulation

Practical: Phase-Field Method

Basics and Application in Materials Science

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Contents

1 Introduction

Phase-field simulation is a versatile application in the toolbox of materials simulation. It is often used for simulations of phase transitions, dislocation evolution, fracture simulations etc. The following practicals aim is to get a practical introduction into the subject. In two separate tasks a 1D single-component solidification simulation, calculation of a bulk energy density coefficient and gradient energy density coefficient are going to be conducted.

2 Task 1: 1D Single-Component Solidification

To simulate a 1D single-component solidification we utilize the following energy density:

$$F = \int \left(f_0 \phi^2 (1 - \phi)^2 + \frac{K_\phi}{2} |\nabla \phi|^2 \right) d\vec{r} \tag{1}$$

As solidification is a a non-conservative process, the kinetics are governed by the Allen-Cahn equation (for the homogeneous, isotropic case).

$$\frac{\partial \phi}{\partial t} = -L \frac{\delta F}{\delta \phi} \tag{2}$$

Here the functional derivative $\frac{\delta F}{\delta \phi}$ of an energy functional eq. 3a can be evaluated as eq. 3b.

$$F = \int f(\vec{r}, \phi, \nabla \phi) d\vec{r} \tag{3a}$$

$$\frac{\delta F}{\delta \phi} = \frac{\partial f}{\partial \phi} - \nabla \cdot \frac{\partial f}{\partial (\nabla \phi)} \tag{3b}$$

Solving the functional derivative from eq. 2 yields:

$$\frac{\delta F}{\delta \phi} = 2 f_0 \phi (1 - \phi)^2 + 2 f_0 \phi^2 (\phi - 1) - K_\phi \nabla^2 \phi \Leftrightarrow \tag{4a}$$

$$\Leftrightarrow 2\phi^3 - 4\phi^2 + 2\phi - K_\phi \nabla^2 \phi \tag{4b}$$

Inserting the result into eq. 2 results in:

$$\frac{\partial \phi}{\partial t} = -L\big[f_0(2\phi^3 - 4\phi^2 + 2\phi) - K_\phi \nabla^2 \phi\big] \tag{5}$$

As the governing equation for solving the PDE we apply Neumann boundary conditions for the left $\phi_0 = \phi_1$ and the right side $\phi_n = \phi_{n-1}$ and discretize the partial derivative $\frac{\partial \phi}{\partial t}$ as:

$$\frac{\partial \phi}{\partial t} \approx \frac{\phi_{next} - \phi}{\Delta t} = -L\big[f_0(2\phi^3 - 4\phi^2 + 2\phi) - K_\phi \nabla^2 \phi\big] \Leftrightarrow \tag{6a}$$

$$\Leftrightarrow \phi_{next} = -\Delta t L\big[f_0(2\phi^3 - 4\phi^2 + 2\phi) - K_\phi \nabla^2 \phi\big] \tag{6b}$$

and $\nabla^2 \phi(x)$ is discretized according to the (forward) finite differences scheme:

$$\nabla^2 \phi(x_i) \approx \frac{\phi(x_{i+2}) - 2\phi(x_{i+1}) + \phi(x_i)}{\Delta x^2} \tag{7}$$

where x_i is the node at which the value is computed and Δx is the node distance. The initial result is plotted in fig. 1. As can be seen in fig. 2 f_0 seems to be reciprocally proportional energy functional over time. With decreasing value the transition region (where ϕ changes from 0 to 1) becomes broader. The opposite effect can be observed in fig. 3 with variation of K_ϕ. Additional the energy density functional seems to become broader but only changes its peak value for increased K-values. L does not seem to have any influence on the ϕ-field, the energy or energy density - see fig. 4 .
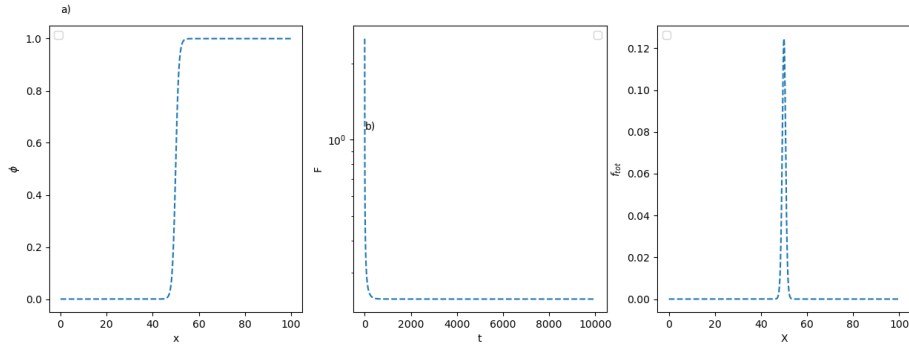

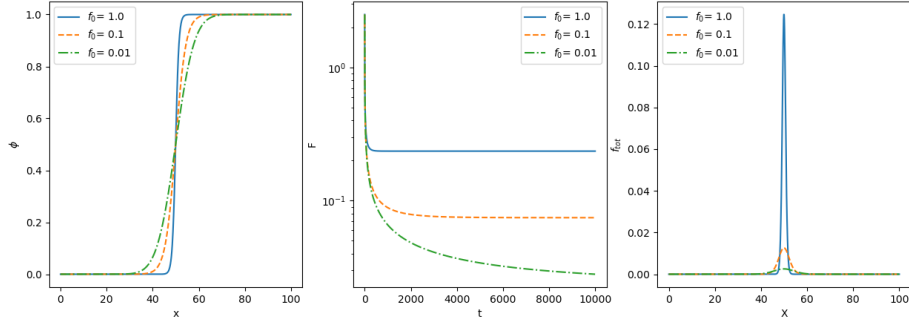Figure 1: Initial results of the 1D single-component solidification simulation.

Figure 2: Results of the 1D single-component solidification simulation with various f_0 values.
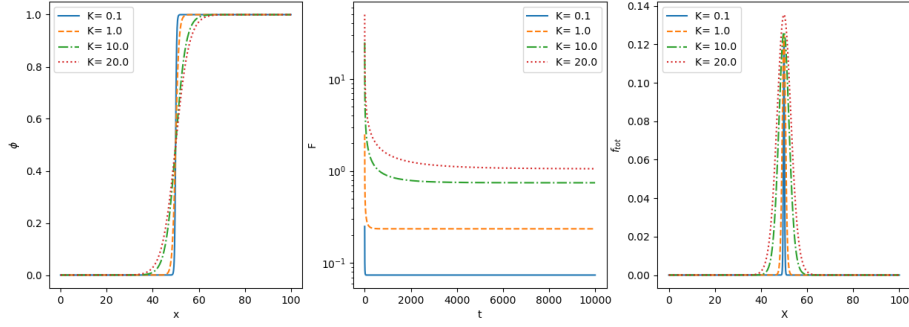

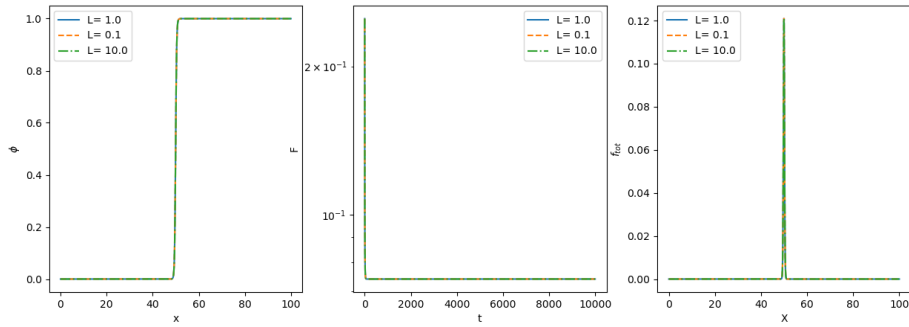Figure 3: Results of the 1D single-component solidification simulation with various K values.

Figure 4: Results of the 1D single-component solidification simulation with various L values.

3 Task 2: Determination of f_0

The equilibrium Al concentration of the disordered γ-Phase c_γ^e has been determined to be 0.165 and the Al concentration of the ordered γ'-Phase $c_{\gamma'}^e$ has been determined to be 0.230 - which is shown in fig. 5 [1]. The bulk energy density function is constructed as:

$$f_{bulk} = f_0 (c_{\gamma'}^e - c)^2 (c_\gamma^e - c)^2 \tag{8}$$

The f_0 needs to be fitted so that it results in an energy gradient of $\Delta f = 149 \frac{\text{J}}{\text{mol}}$ [2]. This is accomplished by choosing the right f_0 to reach the right value at the local maximum of eq. 8. The local maximum of the function can be found automatically by finding the second root of the derivative of eq. 8:

$$\frac{df_{bulk}}{dc} = -2 f_0 (c_{\gamma'}^e - c)(c - c_\gamma^e)^2 + 2 f_0 (c_{\gamma'}^e - c)^2 (c - c_\gamma^e) \tag{9}$$

Using the `scipy` package of Python this can be achieved using `scipy.optimize.root_scalar(func, x0=0.2)` where 0.2 is the initial starting point, looking for c_{max}. Here c_{max} is found at 0.1975. Solving 8 for f_0 results in a value of $\approx 133.6 \cdot 10^6 \frac{\text{J}}{\text{mol}}$. The resulting function of the corresponding parameters is plotted in 6.
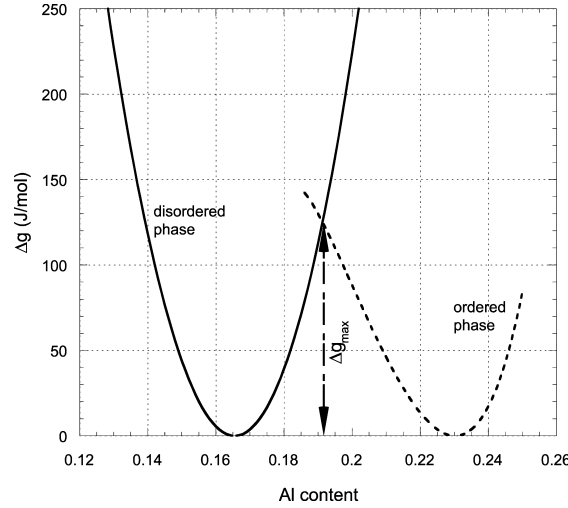

Figure 5: Calculated chemical free energy as a function of composition for both the disordered and ordered phase at 1300 K reproduced from [1].
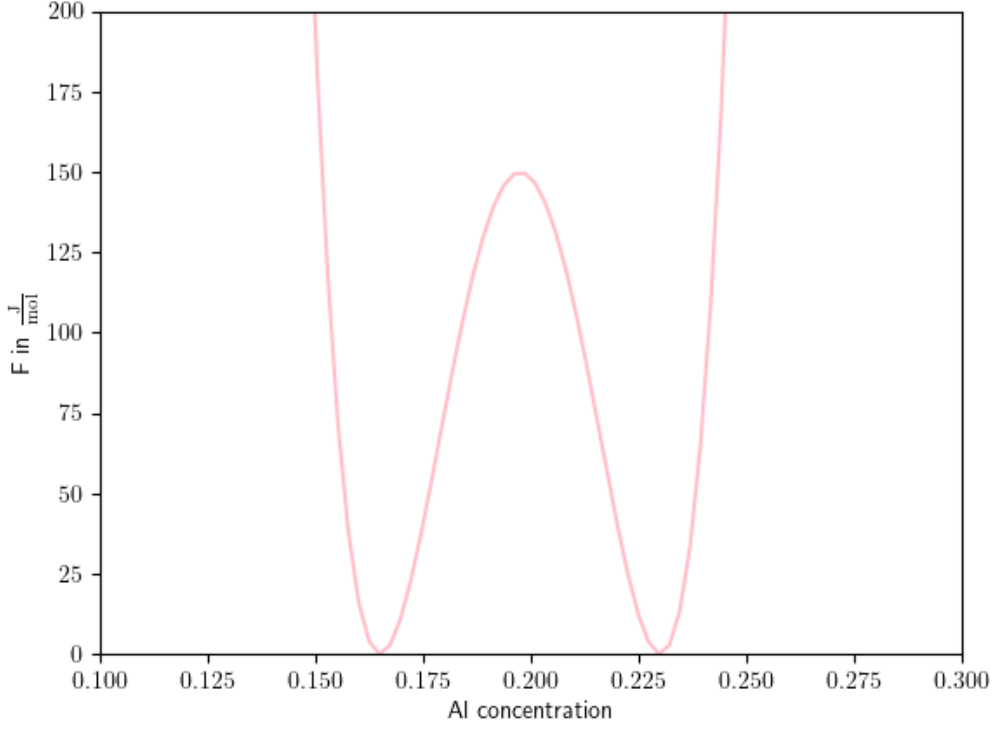
Figure 6: Plot of bulk energy of binary NiAl mixture with $c_\gamma^e = 0.165$, $c_{\gamma'}^e = 0.23$ and $f_0 = 133.6 \cdot 10^6 \frac{\text{J}}{\text{mol}}$.

4 Task 3: Determination of K_c

In this task the gradient energy density coefficient of a binary NiAl-system is calculated. Apart from the governing equation being the Cahn-Hillard equation (eq. 10) as we are dealing with a conserved field, the general algorithm is similar to sec. 2.

$$\frac{\partial c}{\partial t} = M \nabla^2 \frac{\delta F}{\delta c} \tag{10}$$

The functional for the binary system is evaluated as follows:

$$\frac{\delta F}{\delta c} = 2f_0(c_{\gamma'} - c)(c - c_\gamma)^2 + 2f_0(c - c_\gamma)(c_{\gamma'} - c)^2 - K_c \nabla^2 c \tag{11}$$

$\nabla^2(\cdot)$ is (both in eq. 11 and eq. 10) is discretized as in sec. 2 eq. 7. The following

coefficients are used molar Volumen $V_m = 10^{-5}\frac{m^3}{mol}$, interface mobility $M = 10^{-17}\frac{mol^2}{Jms}$, length segment $dx = 10^{-8}m$, $f_0 = 134 \cdot 10^6\frac{J}{mol}$ and a experimental interface energy of $F_{inte} = 5 \sim 50 \cdot 10^{-3}\frac{J}{m}$.

As M is function of molar quantity of substance for calculating the energy we have to add define an adjusted $M_{new} = MV_m^2$. Running the simulation for various values of K_c results in fig. 7. Printing the energies alongside the K_c values shows that to achieve $F_{inte} < 0.5\frac{J}{mol}$ we need to use $K_c < 3.5 \cdot 10^{-6}\frac{J}{m}$.
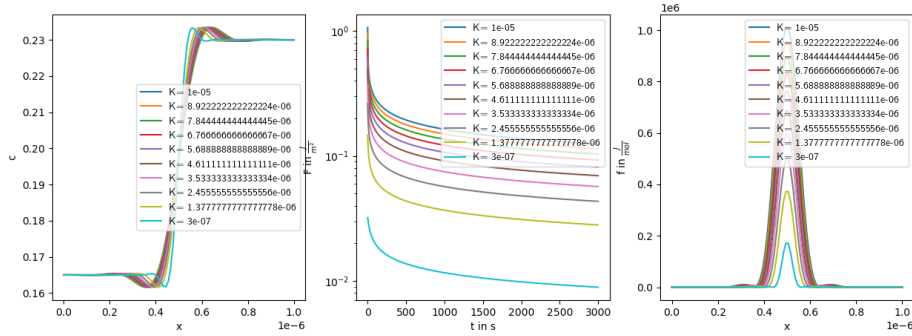

Figure 7: Results of the conserved binary NiAl-system simulation with various K_c values.

For expanding the calculation of the concentrations we need to implement eq. 7 in 2D:

$$\nabla^2 c_{i,j} = \frac{c_{i,j-1} + c_{i,j+1} - 4c_{i,j} + c_{i-1,j} + c_{i+1,j}}{\Delta x^2} \tag{12}$$

where the concentrations are mapped on a $n \times n$-Matrix. Running the simulation for $1000dt$ results in fig. 8. Here we applied an initially random concentration with $c_{init} = 0.195$ as a mean value for a gauss distributed matrix using the `numpy.random.normal(loc=0.195, scale=0.1, size=(N,N))` function. Over time we see, the distribution becomes smooth, while maintaining the composition of several domains of γ
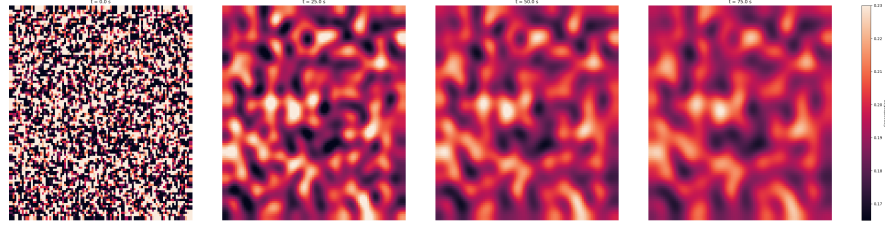γ'-Phases.

Figure 8: Heatmaps of concentrations over time calculated by conserved phase-field method from an initially randomly distributed concentration map.

References

1. Zhu, J. Z., Liu, Z. K., Vaithyanathan, V. & Chen, L. Q. Linking Phase-Field Model to CALPHAD: Application to Precipitate Shape Evolution in Ni-base Alloys. *Scripta Materialia* **46,** 401–406. ISSN: 1359-6462. `https://www.sciencedirect.com/science/article/pii/S1359646202000131` (2024) (Mar. 11, 2002).

2. Zaiser, M. *Phase-Field Method: Basics and Application in Materials Science*