

SEMICONDUCTORS

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**LET US START FROM THE
MICROSCOPIC
ORIGIN OF RESISTANCE**



**MANY BODY PHYSICS
COLLECTIVE PHENOMENA**

Consider Yourself an Electron



You can run without problem from a point A
to Point B without any problem

Consider Yourself an Electron



You CAN'T run without problem from a point A to Point B...You have to jostle, **may even collide with others** and find your way to reach

Consider Yourself an Electron

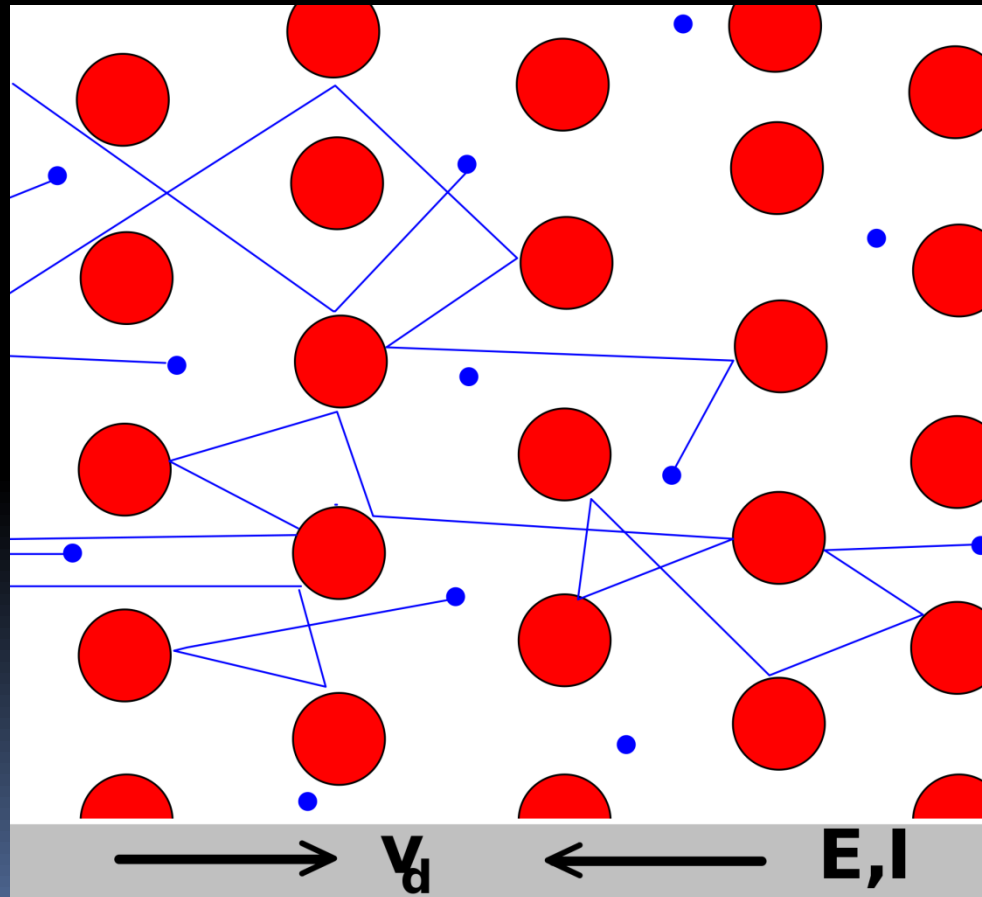


If the Members of the parade stand still then
you can run inhibited within the gaps
However, if they start to walk/rotate/move
sideways then your movement will be restricted

Origin of Electrical Resistance



Scattering of Charge Carriers with Lattice Ions and Impurities inside the Solid

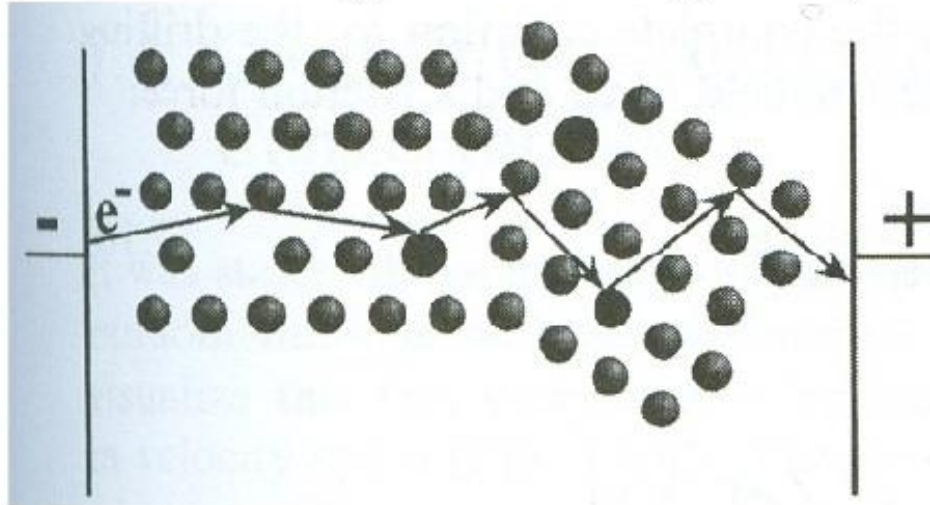


Origin of Electrical Resistance

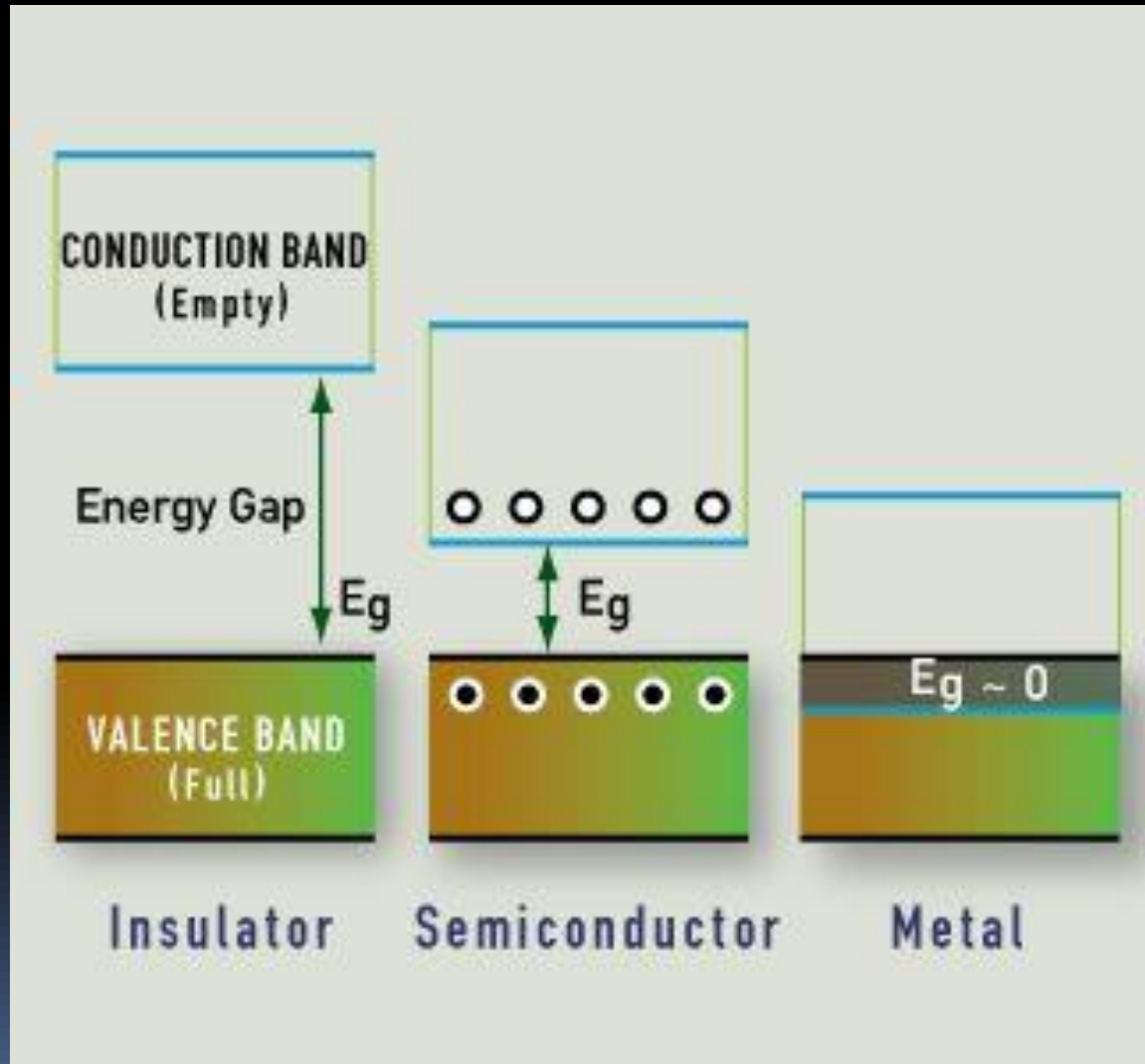


Scattering: origin of resistivity/conductivity

- Primary Scattering Events
 1. Thermal defects ($k_b T$ at room temperature is about 25 meV).
 2. Atomic Defects (impurities/dopants)
 3. 2d and/or 3d defects (grain boundaries, particles, dislocations)



Insulators, Semiconductors And Metals



Full and Partially Full Halls



⇐ All seats are full.
People have less
mobility to move



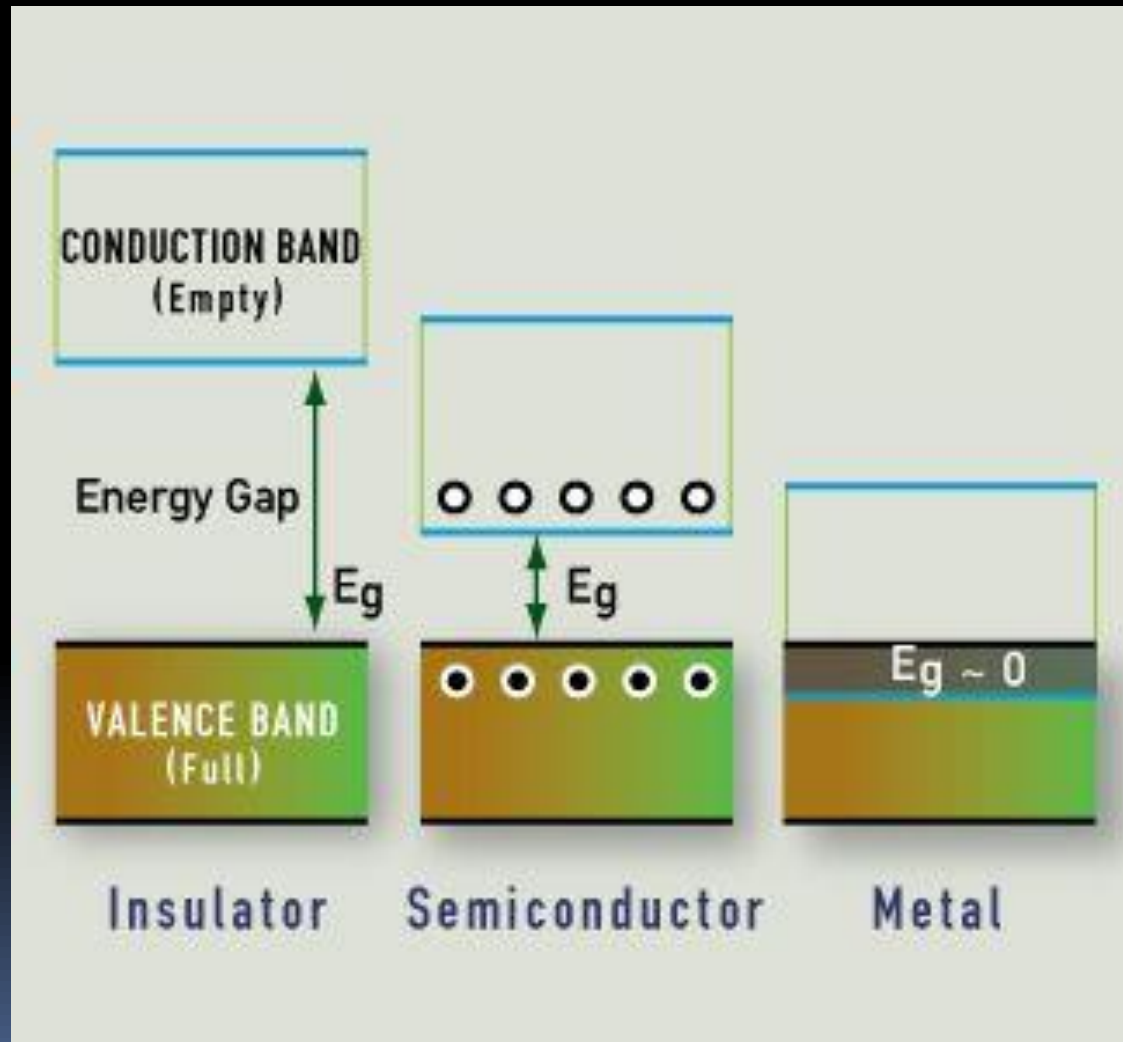
Hall filled hall. People
⇒
have much higher
mobility to move from
one place to other



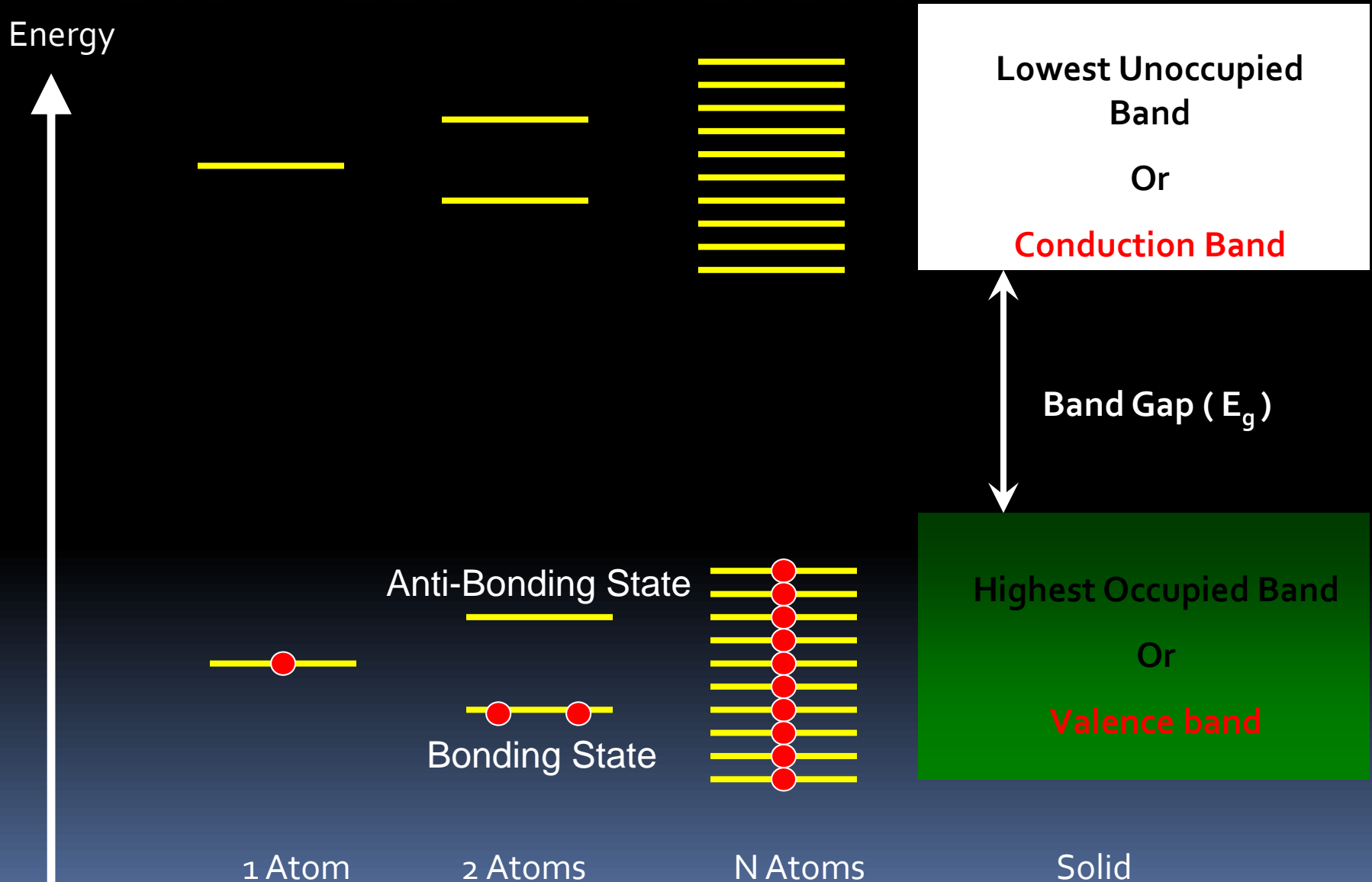
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Insulators, Semiconductors And Metals



FORMATION OF ENERGY BANDS



$$E_c, E_v \text{ \& } E_G = (E_c - E_v)$$

Energy



E_c (Lowest Edge of the Conduction Band) \Rightarrow

E_v (Highest Edge of the Valence Band) \Rightarrow

Lowest Unoccupied
Band

Or

Conduction Band

Band Gap (E_G)

Highest Occupied Band

Or

Valence band

Solid

FERMI DIRAC PROBABILITY DISTRIBUTION

$f(E)$ = The probability that an electron will occupy a level with energy E .

$1-f(E)$ = The probability that a level with energy E remains unoccupied.

$$f(E) = \frac{1}{1 + e^{\frac{(E-E_F)}{k_B T}}}$$

Consequence of Quantum Nature of Electrons : Two Electrons are indistinguishable

DEFINITION OF FERMI LEVEL

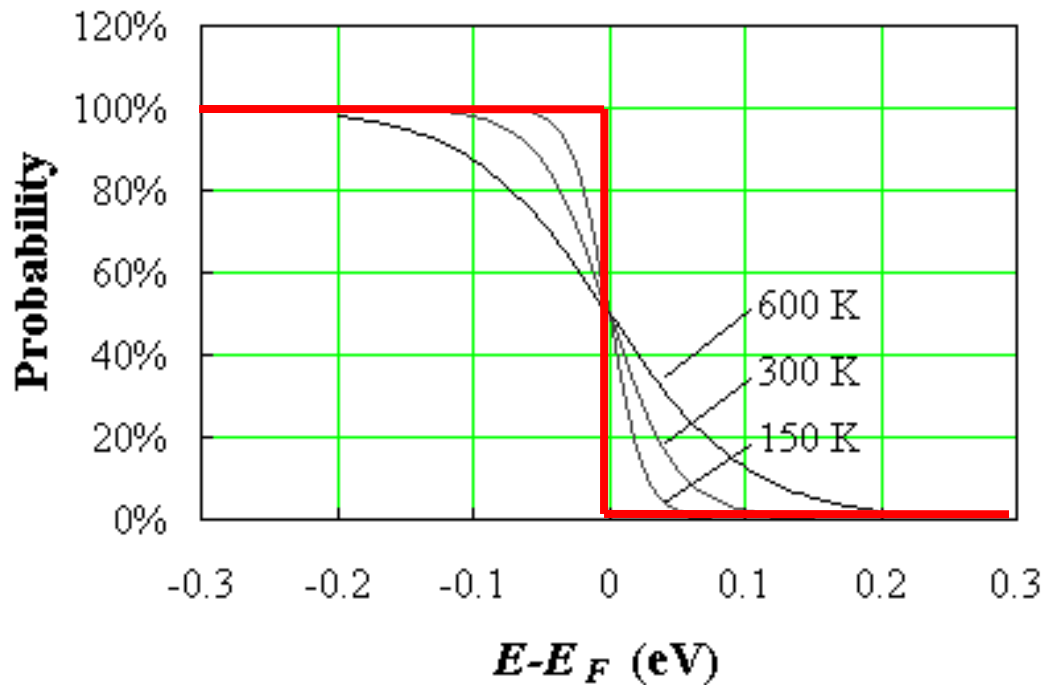
$$f(E) = \frac{1}{1 + e^{\frac{(E - E_F)}{k_B T}}}$$

Probability that electron will occupy an energy level at $E = E_F$ is $\frac{1}{2}$ at all temperature above $T=0$ K.



Very Definition of the
Fermi Level (or Chemical Potential)

FERMI DIRAC PROBABILITY DISTRIBUTION



@ $T = 0$ K

Fermi Dirac
Distribution

$$f(E) = \frac{1}{1 + e^{\frac{(E - E_F)}{k_B T}}}$$

FERMI DIRAC PROBABILITY DISTRIBUTION

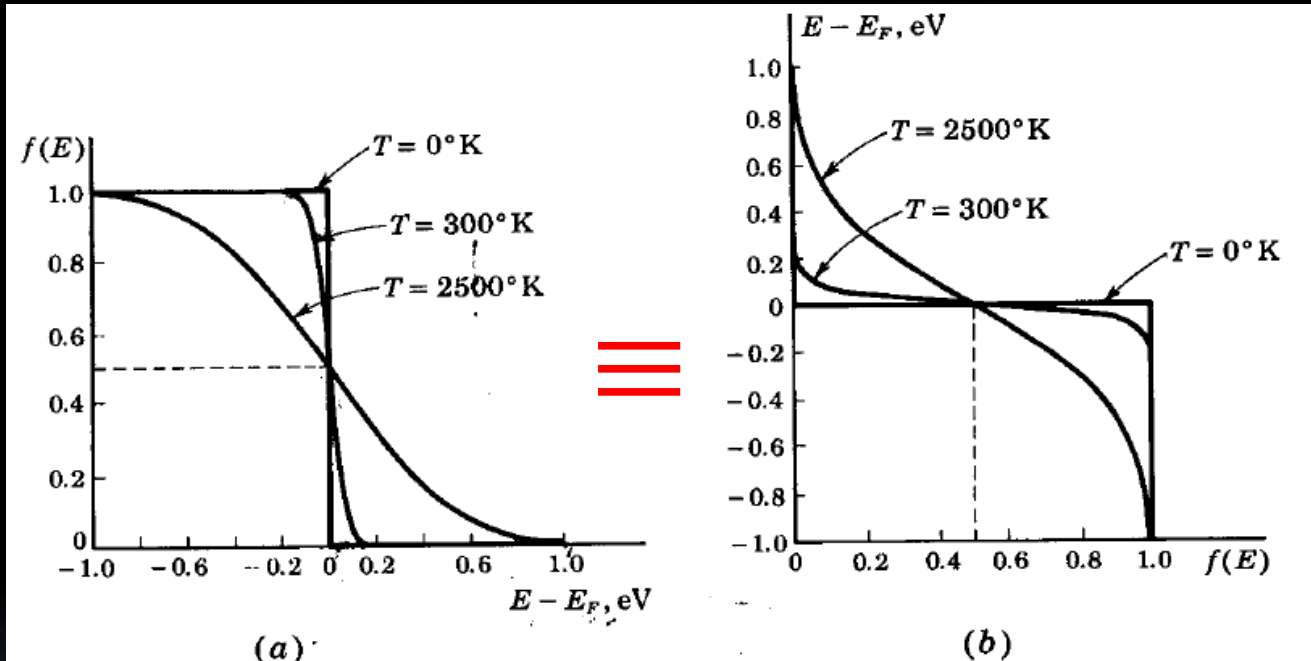
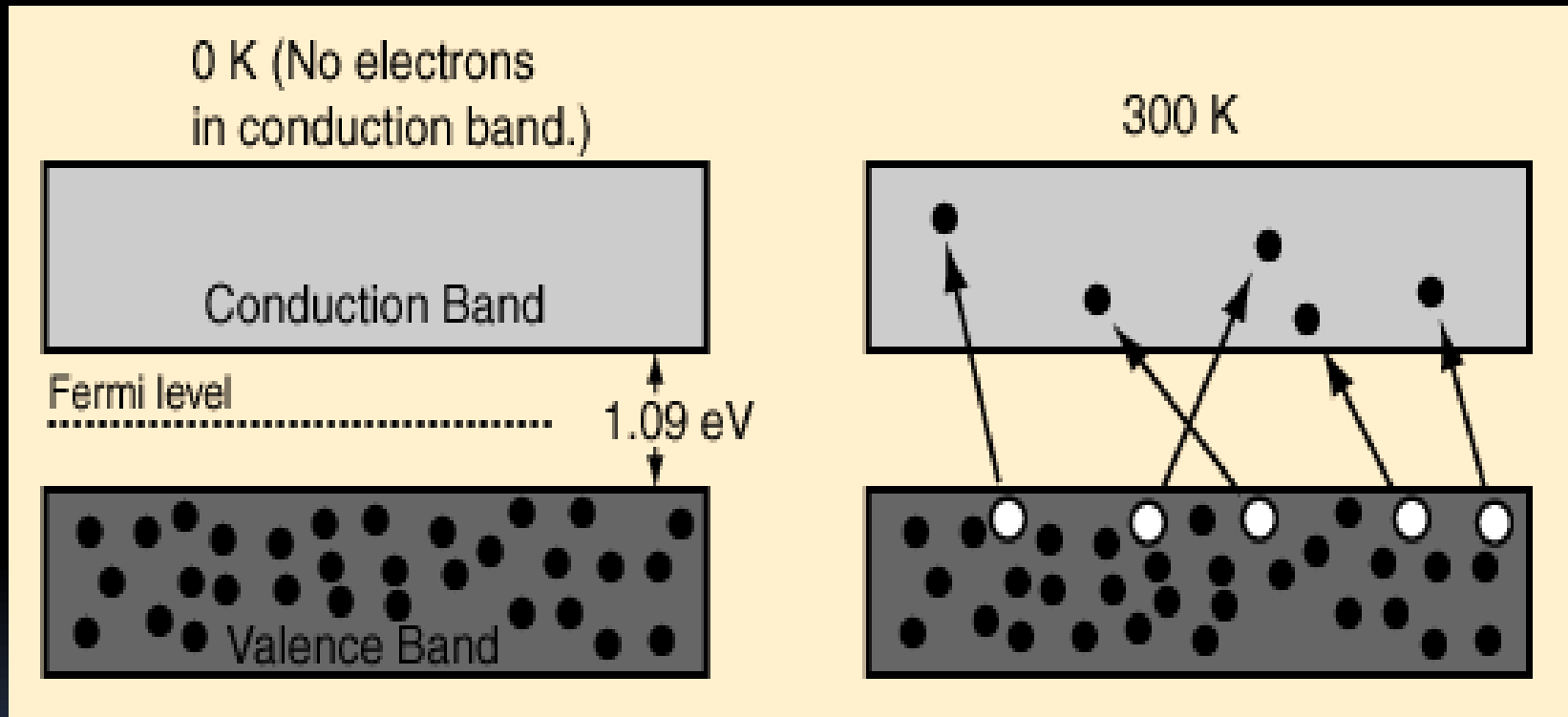


Fig. 19-3 The Fermi-Dirac function $f(E)$ gives the probability that a state of energy E is occupied.

Fermi Dirac
Distribution

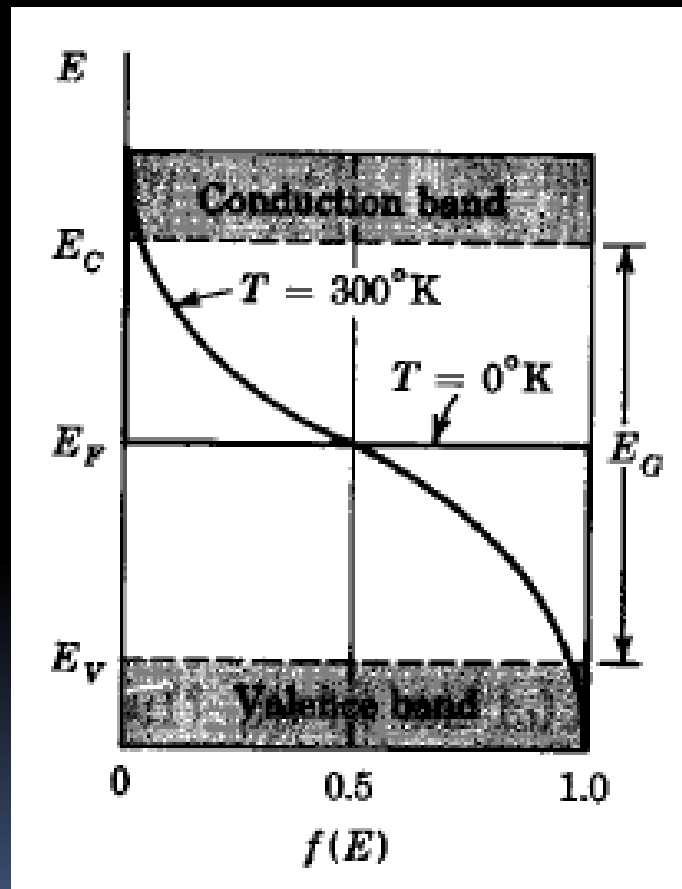
$$f(E) = \frac{1}{1 + e^{\frac{E - E_F}{k_B T}}}$$

THERMAL GENERATION OF FREE ELECTRONS



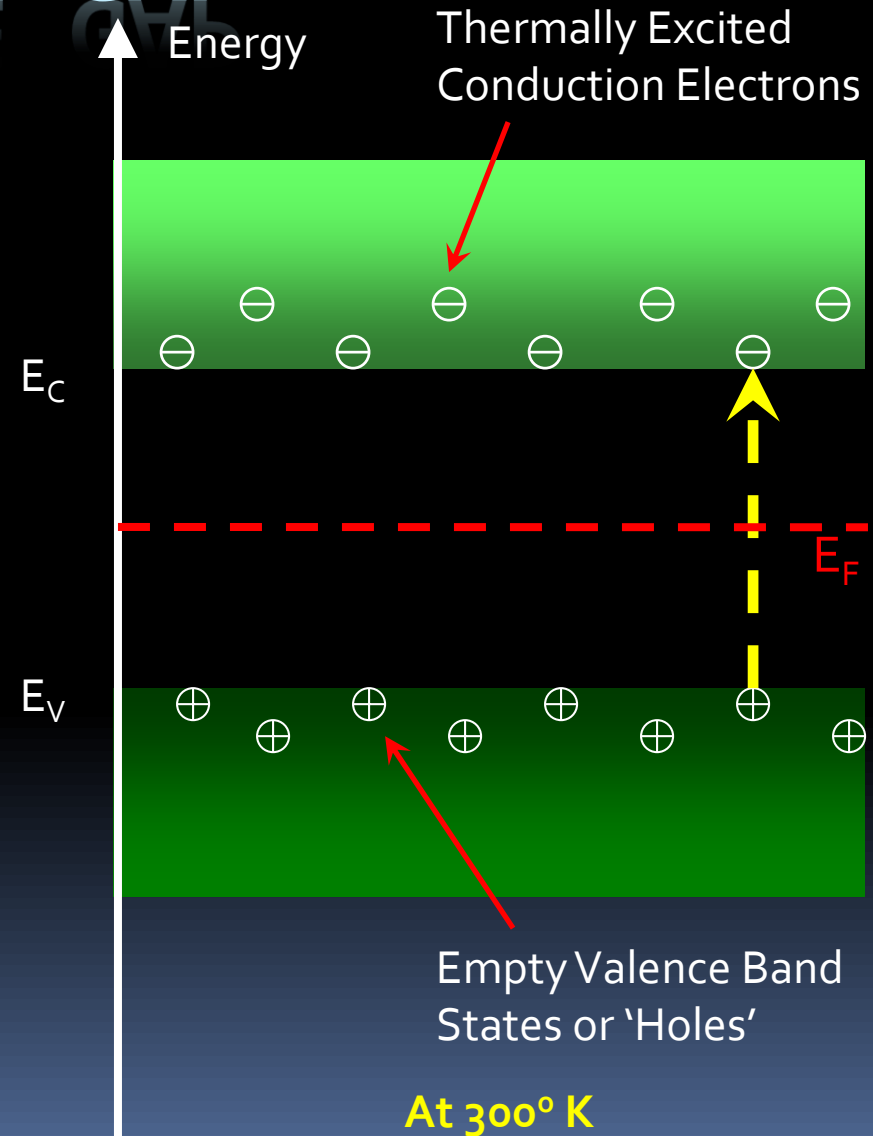
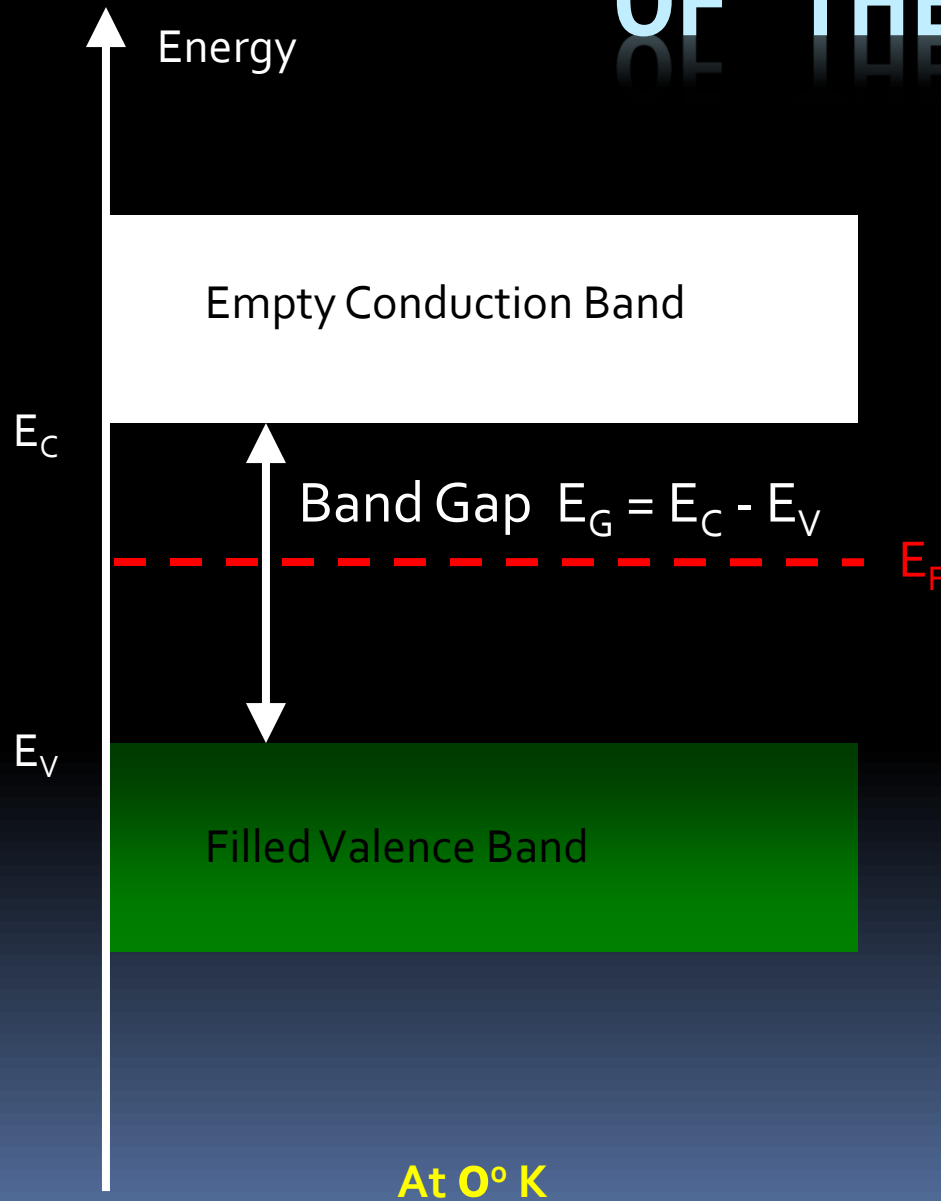
Semiconductor is a perfect insulator at zero Kelvin

SPREAD OF FERMI DIRAC DISTRIBUTION WITH TEMPERATURE



From : fig 19.8 Millman & Halkias, Integrated Electronics.

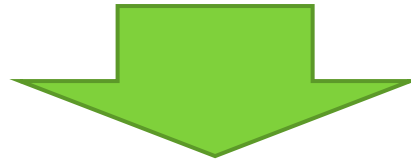
WHY FERMI LEVEL IS AT THE MIDDLE OF THE GAP



DENSITY OF ELECTRONS

$N(E)$ = Density of available Energy State

$f(E)$ = Probability that an electron will occupy an energy level E



Density of Electrons $n(E) \sim N(E)f(E)$

Probability of
Occupation of a State
with Energy E



Number of Available
Energy States

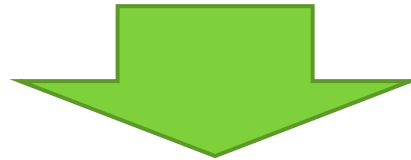


Actual Electron
Population of that
State with Energy E

NUMBER OF STUDENTS OCCUPYING THE SEATS AT RAMAN HALL

$N(E)$ = Available Bench Seats in the first year (Say 50)

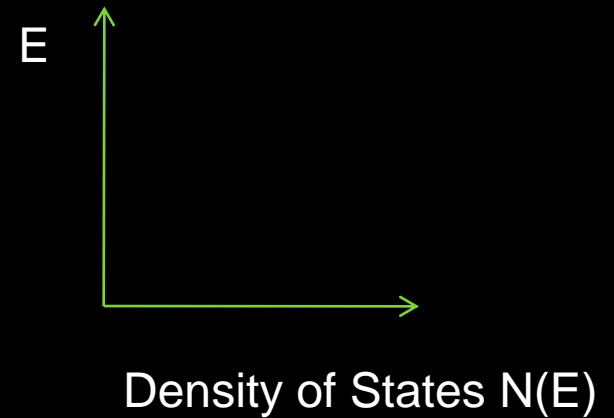
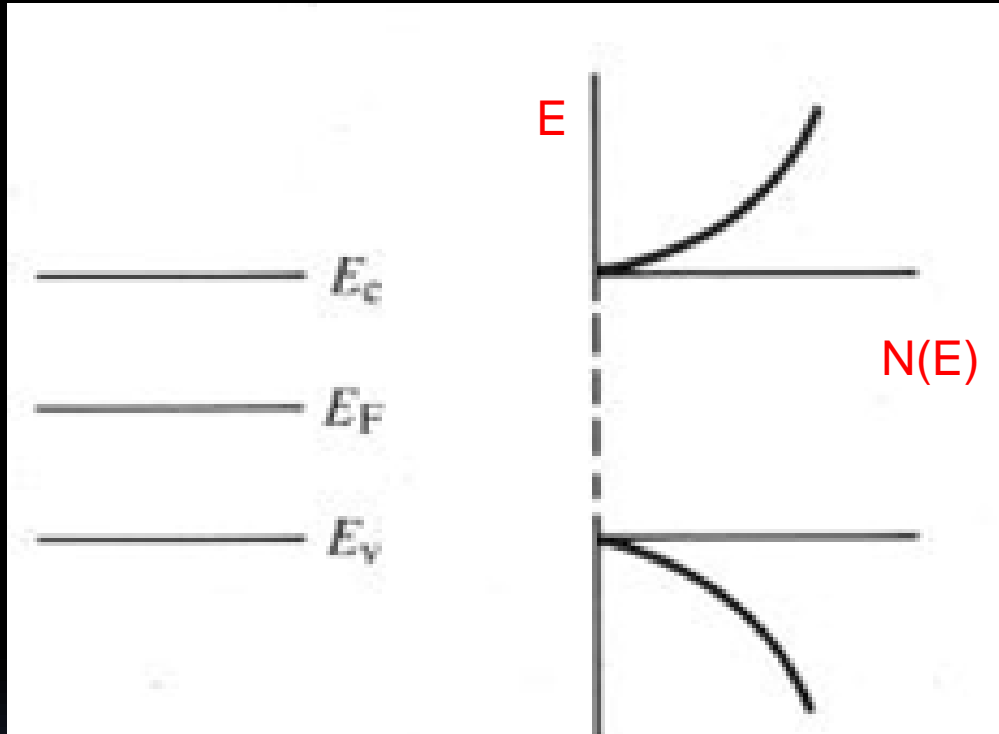
$f(E)$ = Probability that a student is practicing for KARVA ! (Say 60%)



Number of Students sitting on the bench $n(E) \sim N(E)f(E)$

$$n(E) = 50 \times (1 - 0.6) = 20$$

$N(E) \sim$ DENSITY OF AVAILABLE STATES

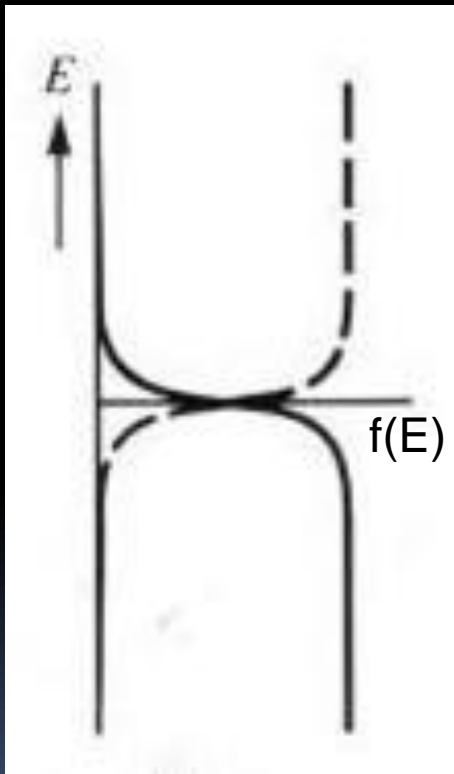


For 3D Bulk Semiconductors only

$$N_C(E) = \gamma(E - E_C)^{1/2}$$

$$N_V(E) = \gamma(E_V - E)^{1/2}$$

FERMI DIRAC PROBABILITY DISTRIBUTION OF ELECTRONS & HOLES

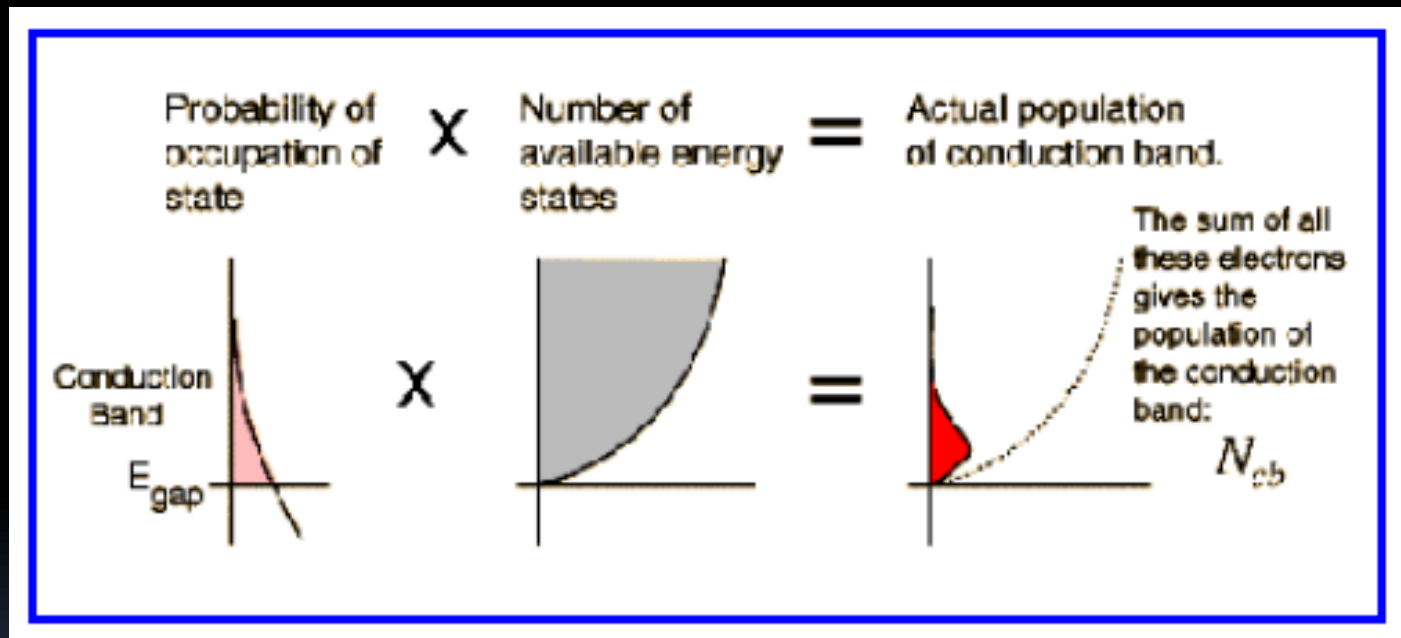


Absence of Electron \equiv Presence of Hole

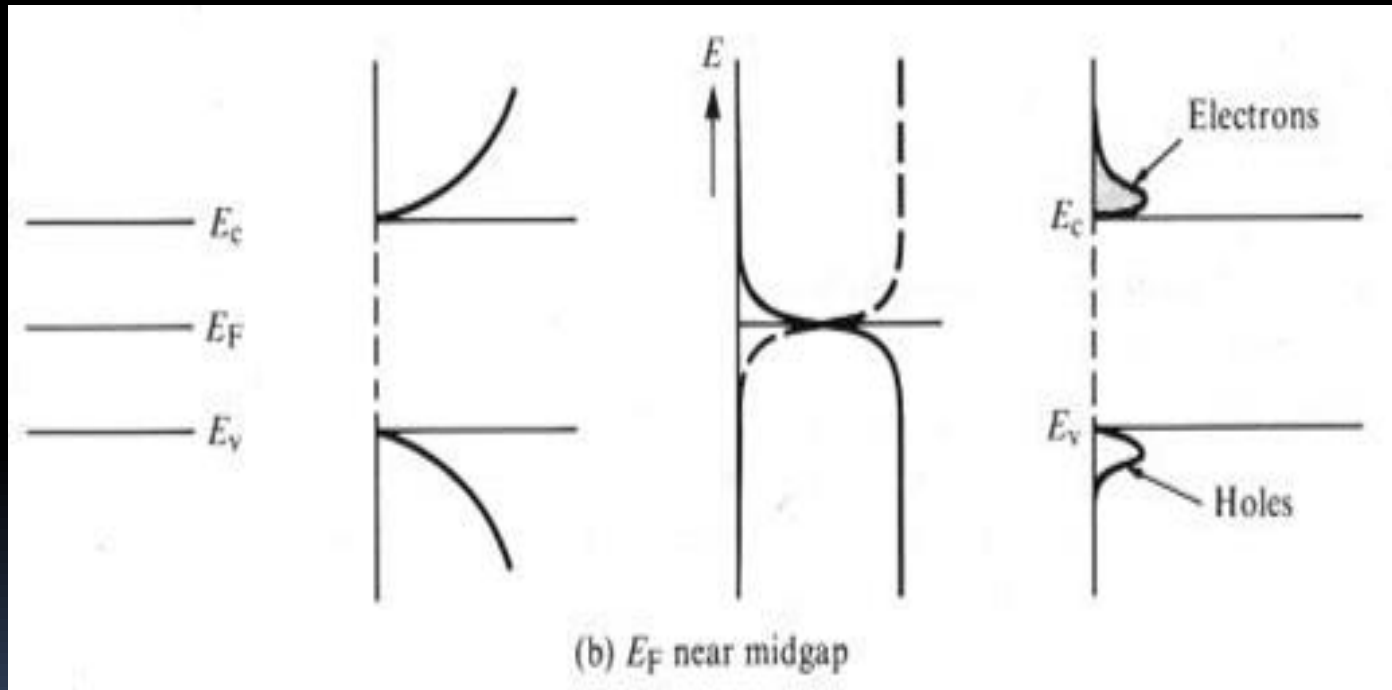
$1 - f(E)$ = The probability that a Hole will occupy an energy level with energy E .

Solid Line $\sim f(E)$; Dotted Line $\sim 1-f(E)$

DENSITY OF ELECTRONS



FERMI DIRAC PROBABILITY DISTRIBUTION AND DENSITY OF ELECTRONS



Band Gap

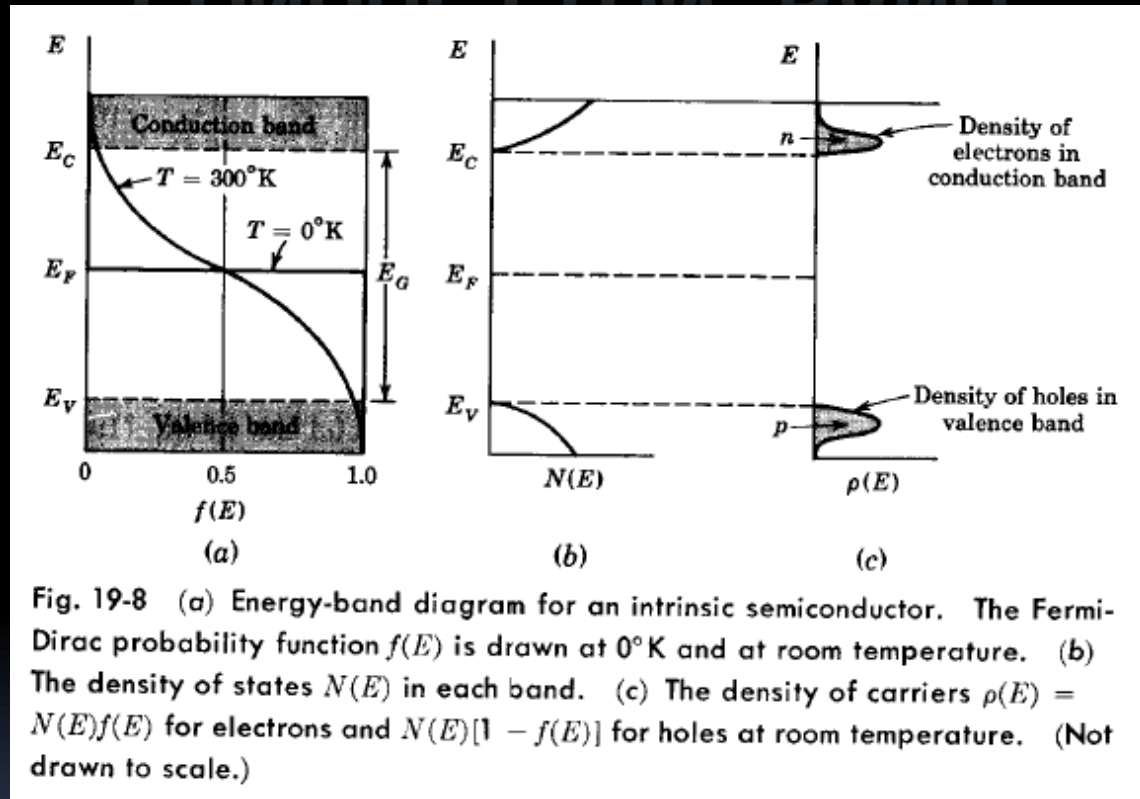
Density of State

F(E)

Concentration of e & h

Intrinsic Semiconductor Only

DENSITY OF ELECTRONS IN THE CONDUCTION BAND

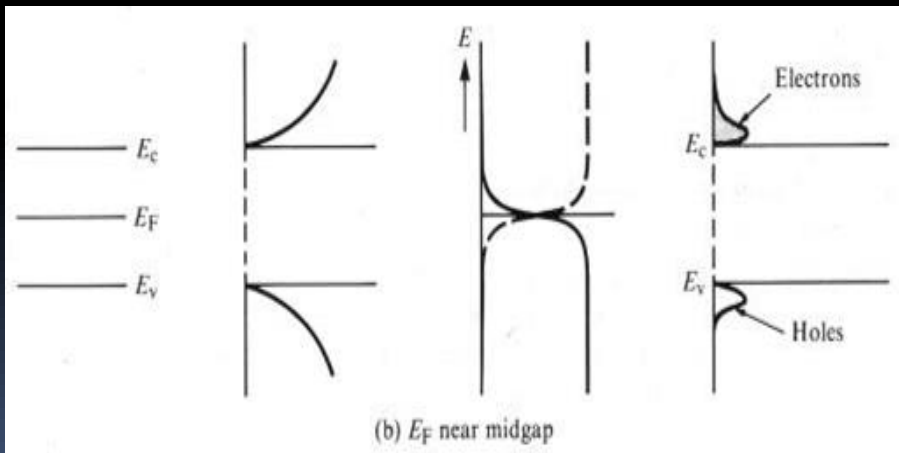


$$n_C = \int_{E_C}^{\infty} N(E) f(E) dE$$

DENSITY OF ELECTRONS IN THE CONDUCTION BAND

$$N_C(E) = \gamma(E - E_C)^{\frac{1}{2}}; \quad \gamma = \frac{4\pi}{h^3} (2m_e^*)^{3/2} (1.602 \times 10^{-19})^{3/2} \text{ in } (\text{m}^{-3})(\text{eV})^{-3/2}$$

$$f(E) = \frac{1}{1 + e^{\frac{(E - E_F)}{k_B T}}}$$



For $E \geq E_C$

$$n_C = \int_{E_C}^{\infty} N(E) f(E) dE$$

DENSITY OF ELECTRONS IN THE CONDUCTION BAND

In case of $E \geq E_C$ & $(E - E_F) \gg kT$

$$f(E) = \frac{1}{1 + e^{\frac{(E-E_F)}{k_B T}}} \approx e^{-\frac{(E-E_F)}{k_B T}}$$

$$n_C = \int_{E_C}^{\infty} N(E) f(E) dE = N_C e^{-\frac{(E_C-E_F)}{k_B T}}$$
$$N_C = 2 \left(\frac{2\pi m_e^* k_B T}{h^2} \right)^{3/2} (1.602 \times 10^{-19})^{3/2}$$

DENSITY OF HOLES IN THE VALENCE BAND

$$E \leq E_V \text{ \& } (E - E_F) \gg kT$$

$$f_p = 1 - \frac{1}{1 + \exp\left(\frac{E - E_F}{k_B T}\right)} = \frac{1}{1 + \exp\left(\frac{E_F - E}{k_B T}\right)}$$



$$n = N_c \exp\left(\frac{E_F - E_c}{k_B T}\right)$$

$$p = N_v \exp\left(\frac{E_v - E_F}{k_B T}\right)$$

PHYSICAL CONSTANTS

Boltzman Constant (k_B) = $8.617\,343(15) \times 10^{-5}$ eV/Kelvin

Plank Constant (h) = $6.626\,068\,96(33) \times 10^{-34}$ Joule - second

m_h^* = Effective Mass of Hole in Si = $0.16_L \times m_e$

$m_e = 9.109\,382\,15(45) \times 10^{-31}$ kg

Fundamental Physical Constants

<http://physics.nist.gov/cuu/Constants/index.html>

THINK!

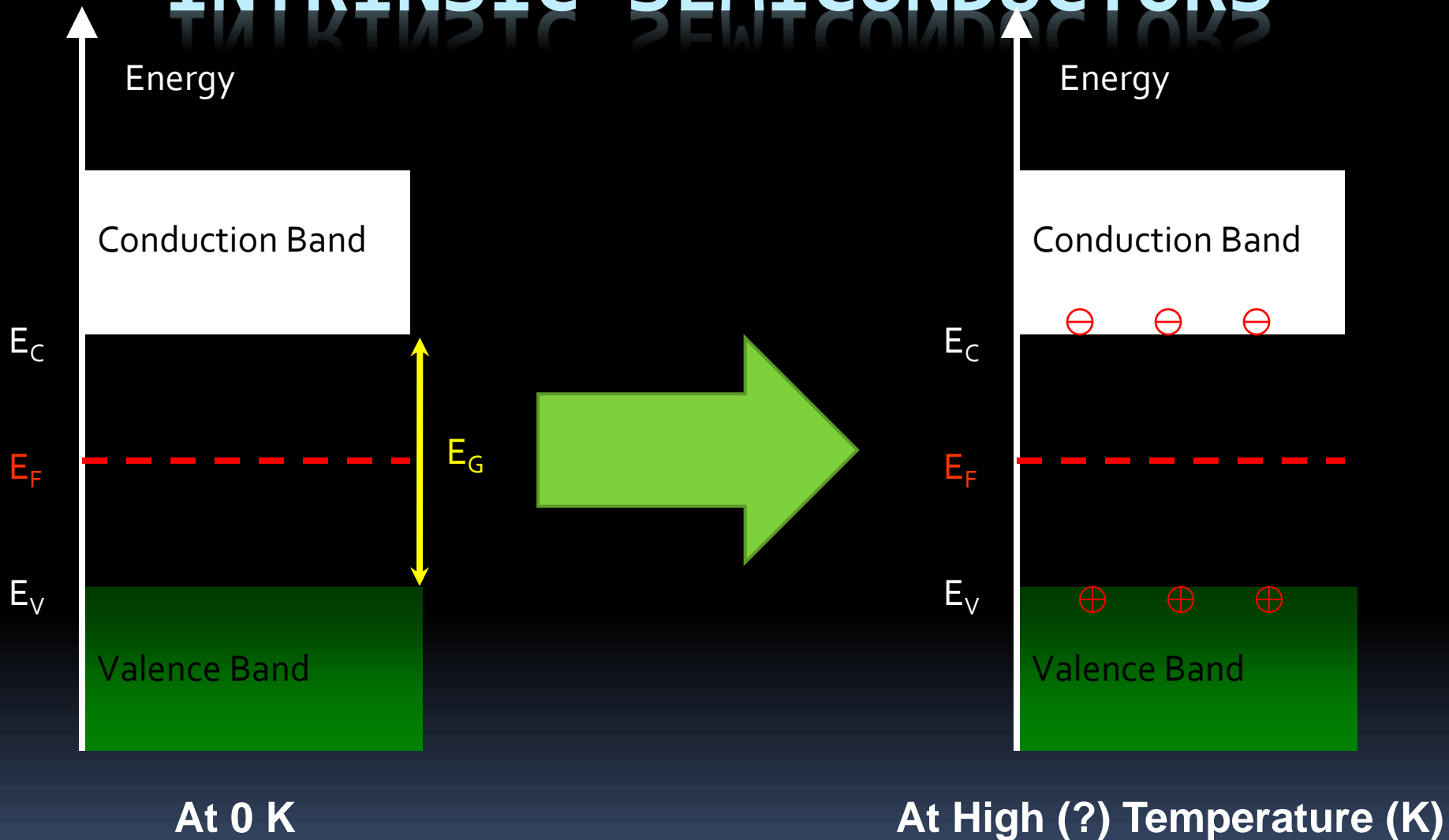
In case of intrinsic semiconductor $n_i = p_i$ then,

Does $n_C = p_V$ imply that $E_F = \frac{E_C + E_V}{2}$?

n_C is the number of electrons in the conduction band

p_V is the number of holes in the valence band

INTRINSIC SEMICONDUCTORS



Intrinsic Semiconductors \Rightarrow Current carrying carriers are generated only by thermal excitation across the band gap.

THINK !

Do we need $(k_B T) \geq E_G$ for Thermal Excitation ?



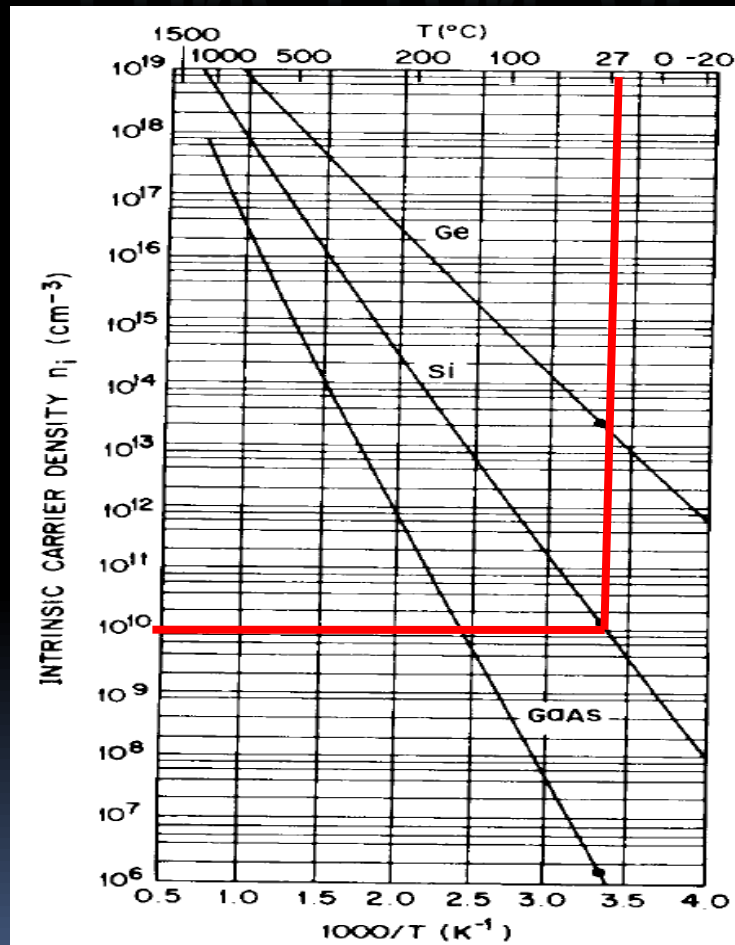
E_G of Silicon at 300 K is 1.12 eV

$k_B T = 0.025 \text{ eV}$ at 300 K



How some electrons are excited across the band gap at Room Temperature ?

INTRINSIC ELECTRON DENSITY AS A FUNCTION OF TEMPERATURE



Intrinsic (n_i) electron density around room temperature is $\sim 10^{10}$ cm⁻³ in Silicon



Small

Also $n_i = p_i$

GROUP IV IN THE PERIODIC TABLE

GROUP

IA

1

1

H

Hydrogen

1.00794

IIA

3

2

Li

Lithium

6.941

4

2

Be

Beryllium

9.01218

IIIB

5

3

B

Boron

10.811

6

3

C

Carbon

12.011

7

3

N

Nitrogen

14.00674

8

3

O

Oxygen

15.9994

9

3

F

Fluorine

18.99840

VIII

10

3

Ne

Neon

20.1797

IIIB

13

4

Al

Aluminum

26.98154

14

4

Si

Silicon

28.0855

15

4

P

Phosphorus

30.97376

16

4

S

Sulfur

32.066

VIIA

17

4

Cl

Chlorine

35.4527

18

4

Ar

Argon

39.948

IIIB

31

5

Ga

Gallium

69.723

32

5

Ge

Germanium

72.61

33

5

As

Arsenic

74.92160

VIIA

34

5

Se

Selenium

78.96

35

5

Br

Bromine

79.904

VIII

36

5

Kr

Krypton

83.80

IIIB

49

6

In

Indium

114.818

50

6

Sn

Tin

118.710

51

6

Sb

Antimony

121.760

VIIA

52

6

Te

Tellurium

127.60

53

6

I

Iodine

126.90447

VIII

54

6

Xe

Xenon

131.29

IIIB

81

7

Tl

Thallium

204.3833

82

7

Pb

Lead

207.2

83

7

Bi

Bismuth

208.98038

VIIA

84

7

Po

Polonium

(209)

85

7

At

Astatine

(210)

VIII

86

7

Rn

Radon

(222)

IIIB

104

8

Rf

Rutherfordium

(261)

105

8

Db

Dubnium

(262)

106

8

Sg

Seaborgium

(263)

VIIA

107

8

Bh

Bohrium

(264)

108

8

Hs

Hassium

(265)

VIII

109

8

Mt

Meitnerium

(268)

110

8

Uun

Ununium

(269)

111

8

Uuu

Unununium

(272)

112

8

Uub

Ununbium

(272)

IIIB

104

9

Rf

Rutherfordium

(261)

105

9

Db

Dubnium

(262)

106

9

Sg

Seaborgium

(263)

VIIA

107

9

Bh

Bohrium

(264)

108

9

Hs

Hassium

(265)

VIII

109

9

Mt

Meitnerium

(268)

110

9

Uun

Ununium

(269)

111

9

Uuu

Unununium

(272)

112

9

Uub

Ununbium

(272)

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110

11

Uun

Ununium

(269)

111

11

Uuu

Unununium

(272)

112

11

Uub

Ununbium

(272)

IIIB

104

12

Rf

Rutherfordium

(261)

105

12

Db

Dubnium

(262)

106

12

Sg

Seaborgium

(263)

VIIA

107

12

Bh

Bohrium

(264)

108

12

Hs

Hassium

(265)

VIII

109

12

Mt

Meitnerium

(268)

110

12

Uun

Ununium

(269)

111

12

Uuu

Unununium

(272)

112

12

Uub

Ununbium

(272)

IIIB

104

13

Rf

Rutherfordium

(261)

105

13

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Unununium

(272)

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Uub

Ununbium

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109

15

Mt

Meitnerium

(268)

110

15

Uun

Ununium

(269)

111

15

Uuu

Unununium

(272)

112

15

Uub

Ununbium

(272)

IIIB

104

16

Rf

Rutherfordium

(261)

105

16

Db

Dubnium

(262)

106

16

Sg

Seaborgium

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VIIA

107

16

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16

Hs

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(265)

VIII

109

16

Mt

Meitnerium

(268)

110

16

Uun

Ununium

(269)

111

16

Uuu

Unununium

(272)

112

16

Uub

Ununbium

(272)

IIIB

104

17

Rf

Rutherfordium

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105

17

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17

Mt

Meitnerium

(268)

110

17

Uun

Ununium

(269)

111

17

Uuu

Unununium

(272)

112

17

Uub

Ununbium

(272)

IIIB

104

18

Rf

Rutherfordium

(261)

105

18

Db

Dubnium

(262)

106

18

Sg

Seaborgium

(263)

VIIA

107

18

Bh

Bohrium

(264)

108

18

Hs

Hassium

(265)

VIII

109

18

Mt

Meitnerium

(268)

110

18

Uun

Ununium

(269)

111

18

Uuu

Unununium

(272)

112

18

Uub

Ununbium

(272)

IIIB

104

19

Rf

Rutherfordium

(261)

105

19

Db

Dubnium

(262)

106

19

Sg

Seaborgium

(263)

VIIA

107

19

Bh

Bohrium

(264)

108

19

Hs

Hassium

(265)

VIII

109

19

Mt

Meitnerium

(268)

110

19

Uun

Ununium

(269)

111

19

Uuu

Unununium

(272)

112

19

Uub

Ununbium

(272)

IIIB

104

20

Rf

Rutherfordium

(261)

105

20

Db

Dubnium

(262)

106

20

Sg

Seaborgium

(263)

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107

20

Bh

Bohrium

(264)

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Hs

Hassium

(265)

VIII

109

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Mt

Meitnerium

(268)

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20

Uun

Ununium

(269)

111

20

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Unununium

(272)

112

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IIIB

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Rf

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(272)

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Rutherfordium

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Db

Dubnium

(262)

106

22

Sg

Seaborgium

(263)

VIIA

107

22

Bh

Bohrium

(264)

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Hassium

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VIII

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Mt

Meitnerium

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22

Uun

Ununium

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Uuu

Unununium

(272)

112

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Uub

Ununbium

(272)

IIIB

104

23

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Rutherfordium

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23

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Dubnium

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(263)

VIIA

107

23

Bh

Bohrium

(264)

108

23

Hs

Hassium

(265)

VIII

109

23

Mt

Meitnerium

(268)

110

23

Uun

Ununium

(269)

111

23

Uuu

Unununium

(272)

112

23

Uub

Ununbium

(272)

IIIB

104

24

Rf

Rutherfordium

(261)

105

24

Db

Dubnium

(262)

106

24

Sg

Seaborgium

(263)

VIIA

107

24

Bh

Bohrium

(264)

108

24

Hs

Hassium

(265)

VIII

109

24

Mt

Meitnerium

(268)

110

24

Uun

Ununium

(269)

111

24

Uuu

Unununium

(272)

112

24

Uub

Ununbium

(272)

IIIB

104

25

Rf

Rutherfordium

(261)

105

25

Db

Dubnium

(262)

106

25

Sg

Seaborgium

(263)

VIIA

107

25

Bh

Bohrium

(264)

108

25

Hs

Hassium

(265)

VIII

109

25

Mt

Meitnerium

(268)

110

25

Uun

Ununium

(269)

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COVALENT & IONIC COMPOUNDS

Covalent and Ionic Crystals			
Covalent (Group IV)	III-V Compounds	II-VI Compounds	I-VII Compounds (Ionic)
C (diamond)	BN	BeO	LiF
	BP	BeS	LiCl
	AlN	MgO	NaF
Si	AlP	MgS	NaCl
	GaP	CaO	KF
	AlAs	CaS	KCl
		ZnS	LiBr
Ge	GaAs	CaSe	KBr
	InP	ZnSe	NaI
	AlSb	CdS	RbCl
	GaSb	CdSe	RbBr
	InAs	ZnTe	KF
α -Sn	InSb	CdTe	RbI
			CsBr

Table 1.2 J. P. McKelvey, Solid State and Semiconductor Physics

INTRINSIC GERMANIUM

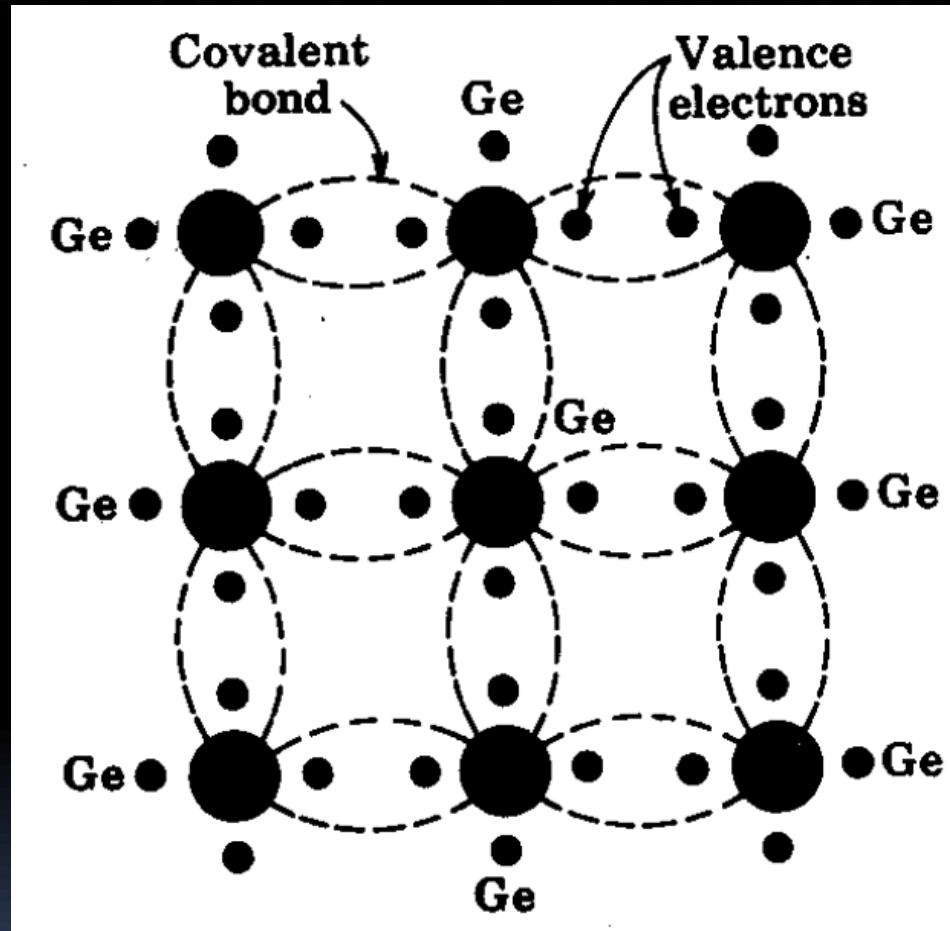


Fig 2.3 Millman &
Halkias, Integrated
Electronics

4 fold coordinated Ge atoms.
Each atom share 8 valence electrons.

DONOR IMPURITIES IN GERMANIUM

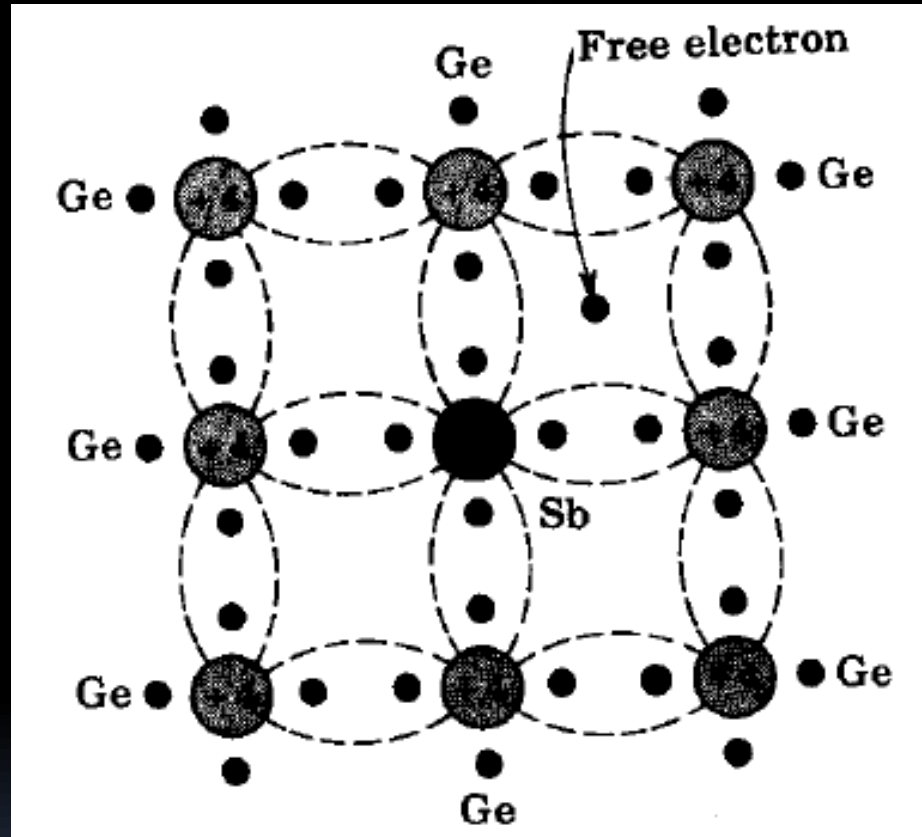


Fig 2.4 Millman &
Halkias, Intergrated
Electronics

Donor : An impurity atom which becomes positively charged after donating an electron to the material (e.g pentavalent Sb, As in Si); Donors contribute free electrons.

ACCEPTOR IMPURITIES IN GERMANIUM

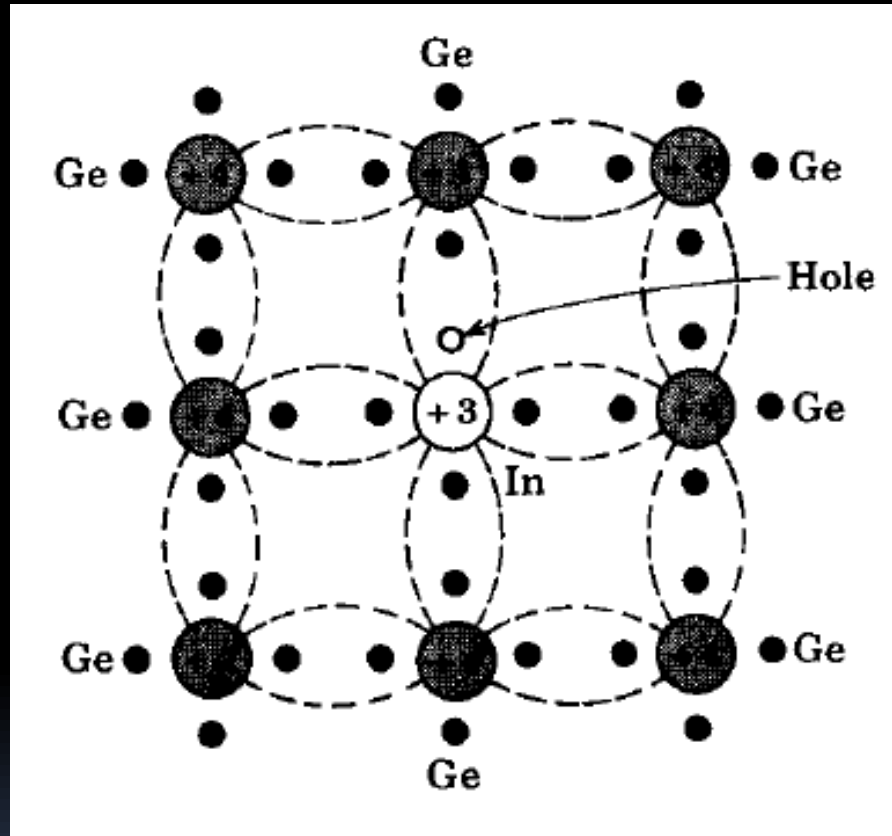
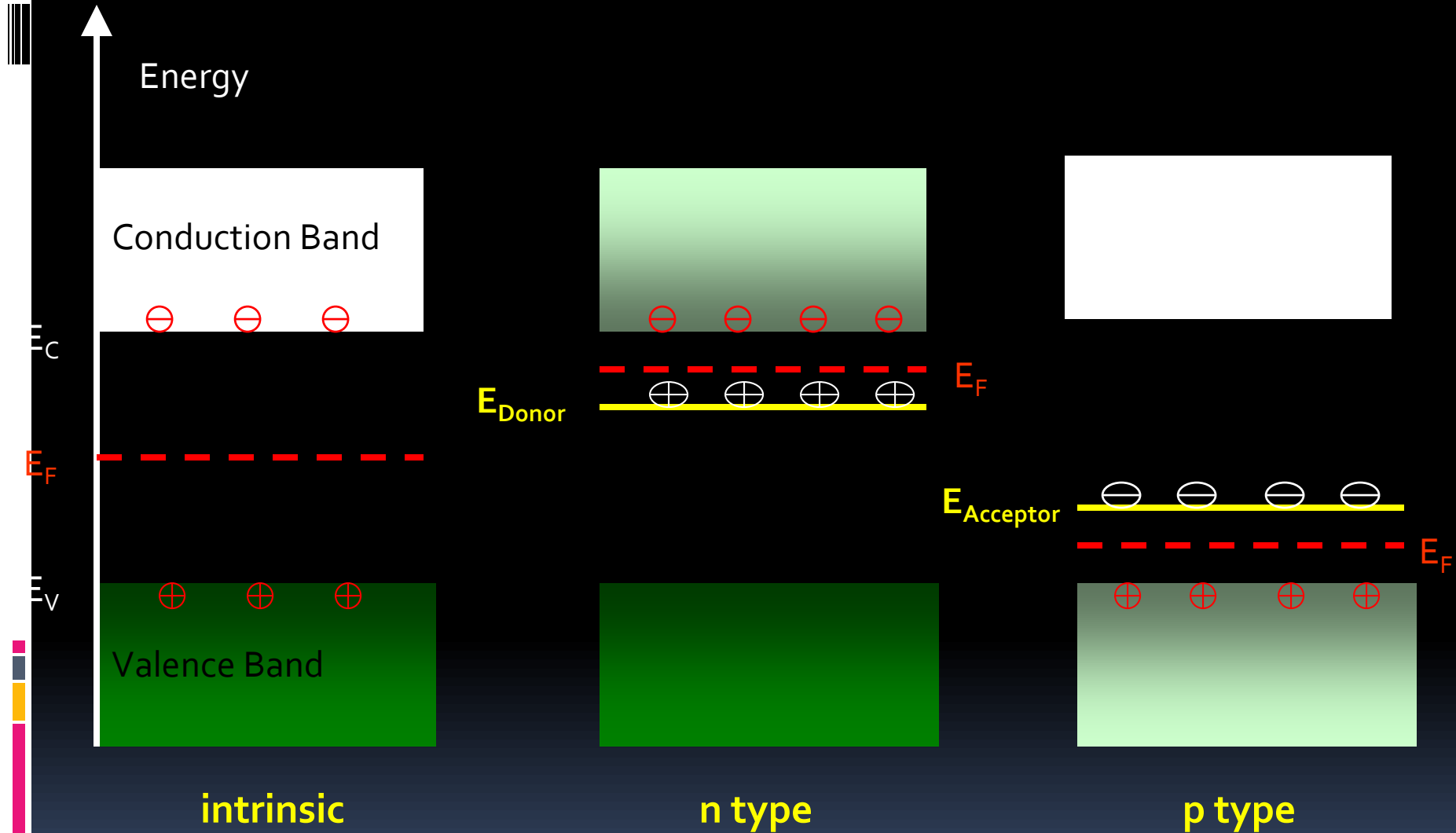


Fig 2.6 Millman &
Halkias, Integrated
Electronics

Acceptor : An impurity atom which becomes negatively charged after accepting an electron from the material (e.g trivalent Al in Si); Acceptor contribute free holes.

n and p type semiconductors



Fermi level E_F Determines the Occupation of any State

n and p type semiconductors

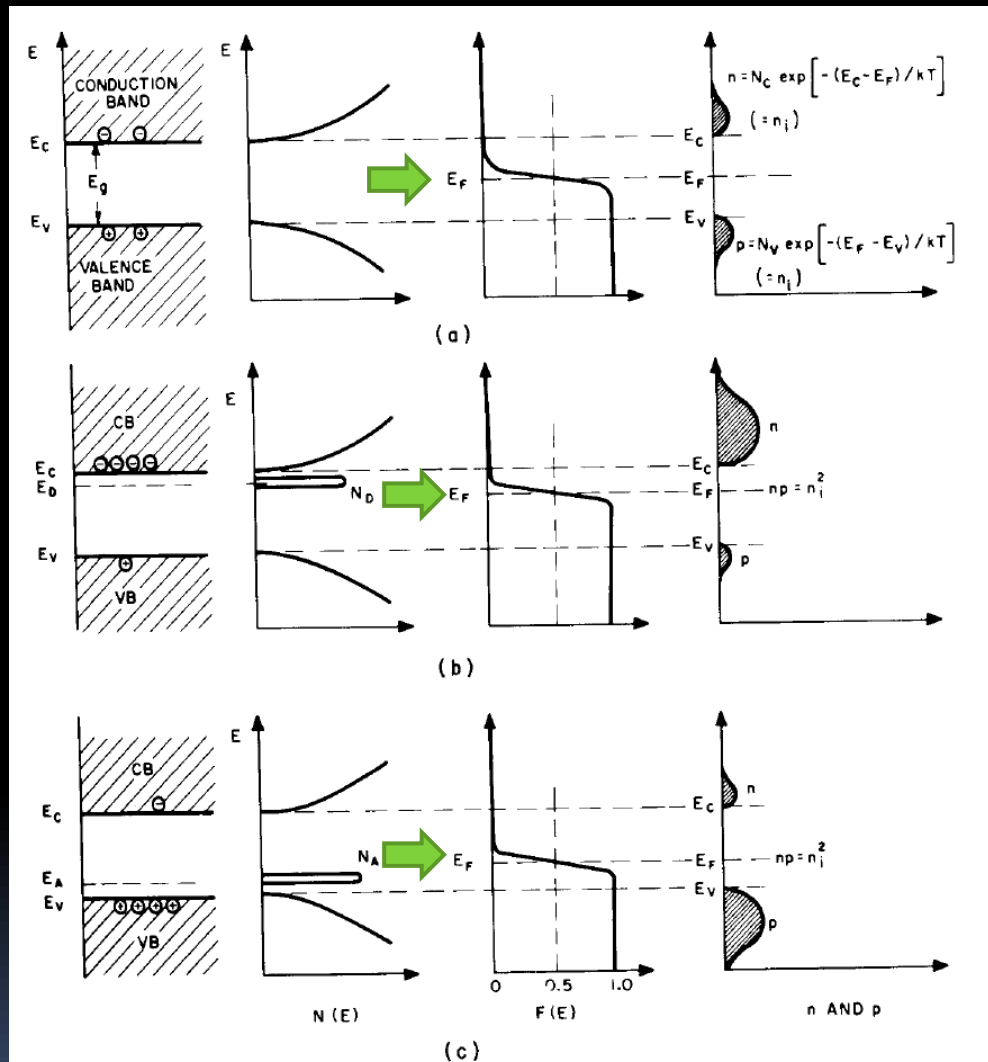
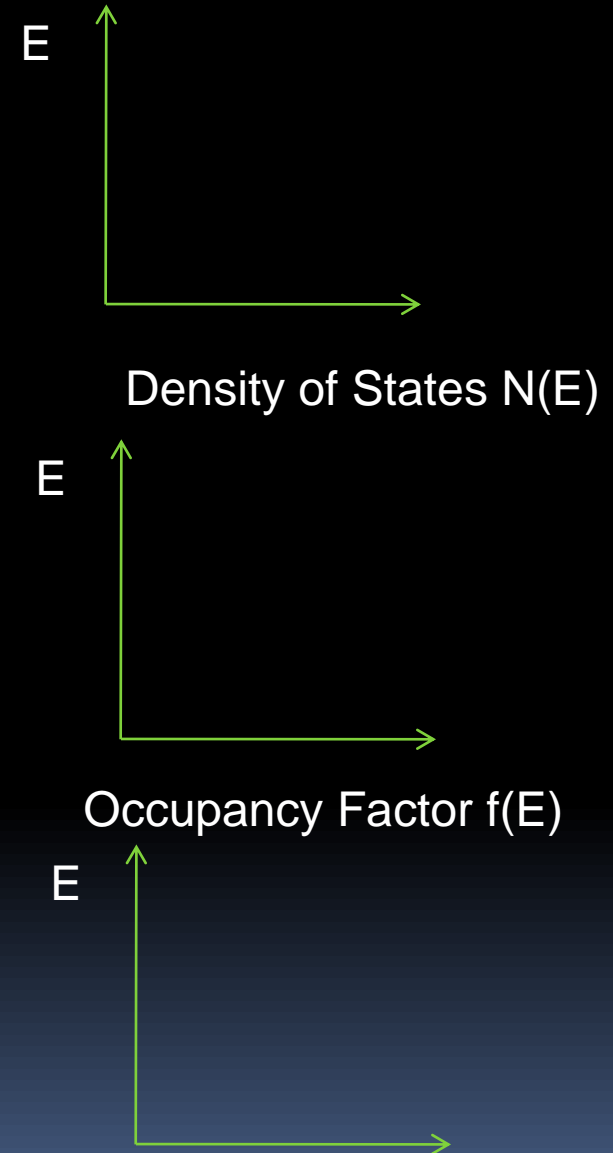


Fig. 14 Schematic band diagram, density of states, Fermi-Dirac distribution, and the carrier concentrations for (a) intrinsic, (b) n-type, and (c) p-type semiconductors at thermal equilibrium. Note that $pn = n_i^2$ for all three cases.



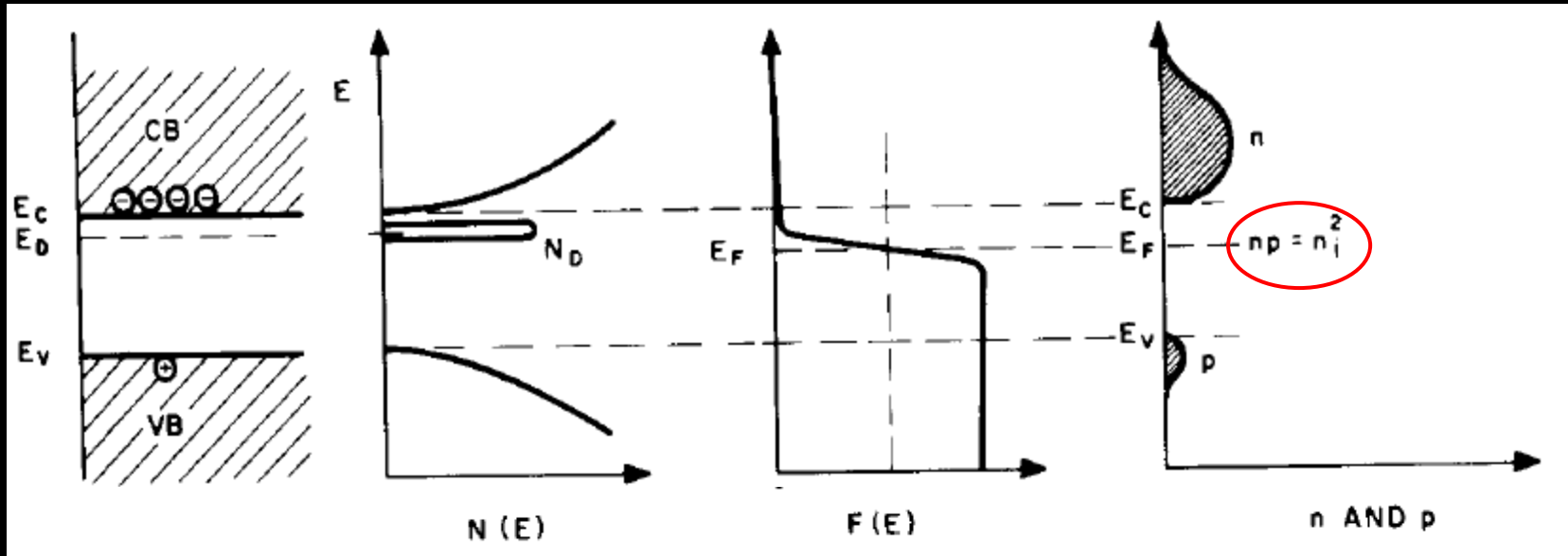
**TOTAL ELECTRON DENSITY =
INTRINSIC ELECTRON DENSITY
+ DOPED ELECTRON DENSITY**

Intrinsic electron density \Rightarrow Thermally excited electrons across the band gap

Doped electron density \Rightarrow contributed by donor impurities.

MAJORITY & MINORITY CARRIERS

n type Semiconductors

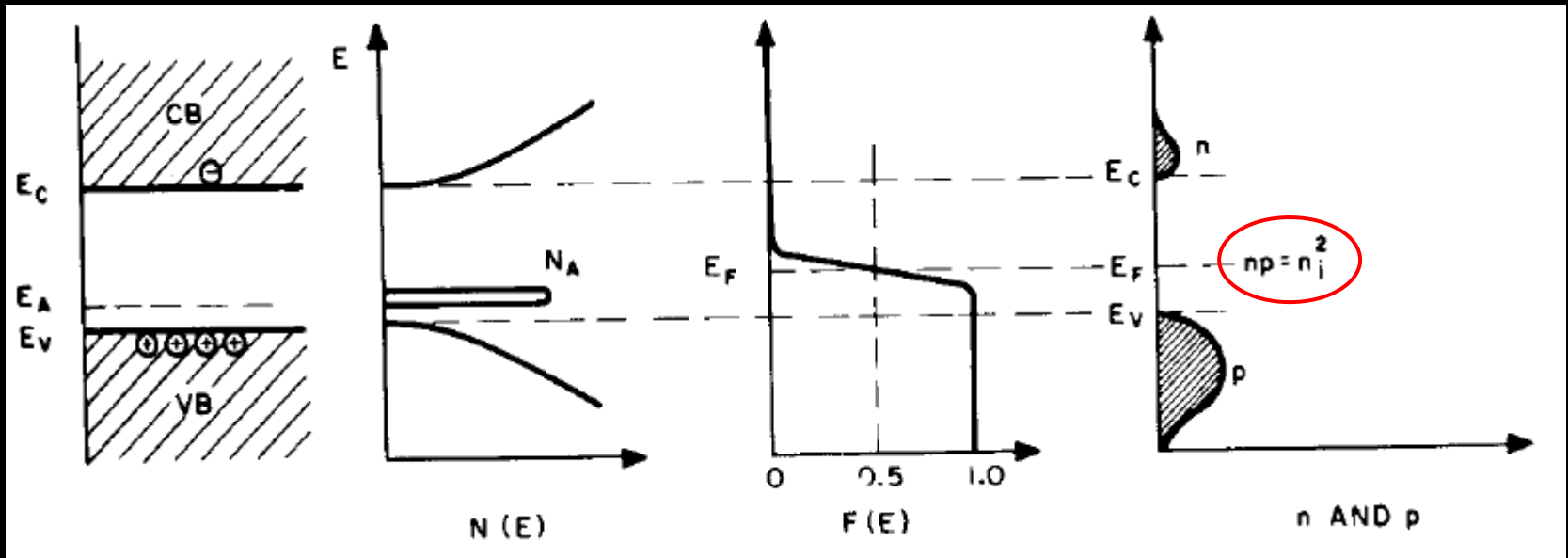


More number of free Electrons (n) \Rightarrow **Electrons are Majority carriers.**

Less number of free Holes (p) \Rightarrow **Holes are Minority carriers**

MAJORITY & MINORITY CARRIERS

p type Semiconductors



More number of free Holes (p) \Rightarrow Holes are Majority carriers.

Less number of free Electrons (n) \Rightarrow Electrons are Minority carriers

Law of Mass Action

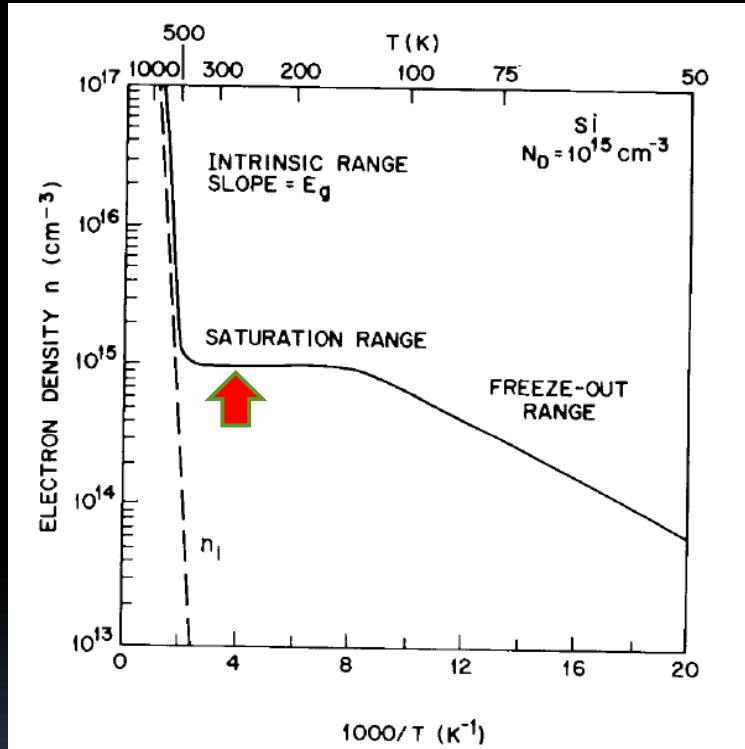


In CHEMISTRY, Law of Mass Action determines the Relative Concentration of ions in weakly ionized electrolytes

$$n.p = 4 \left(\frac{2\pi (m_e^* m_h^*)^{1/2} k_B T}{h^2} \right) e^{-\frac{E_G}{k_B T}} = n_i^2$$

Semiconductors can be thought of as weakly ionized electrolytes where a chemical bond having dissociation energy $E_G = (E_C - E_V)$ dissociates to produce an electron in the conduction band and a hole in the valence band.

ELECTRON CONCENTRATION IN DONOR IMPURITY DOPED SILICON

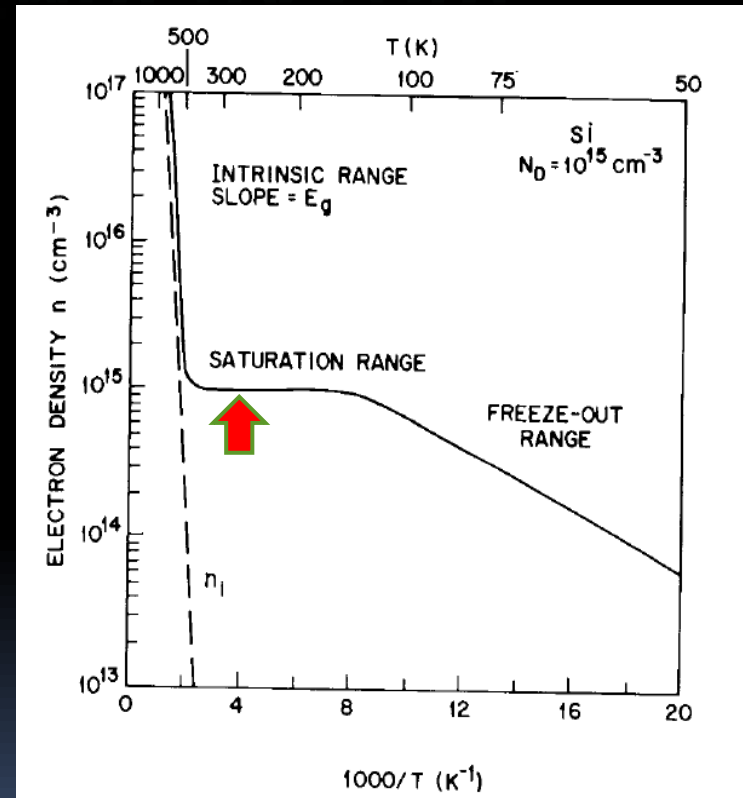
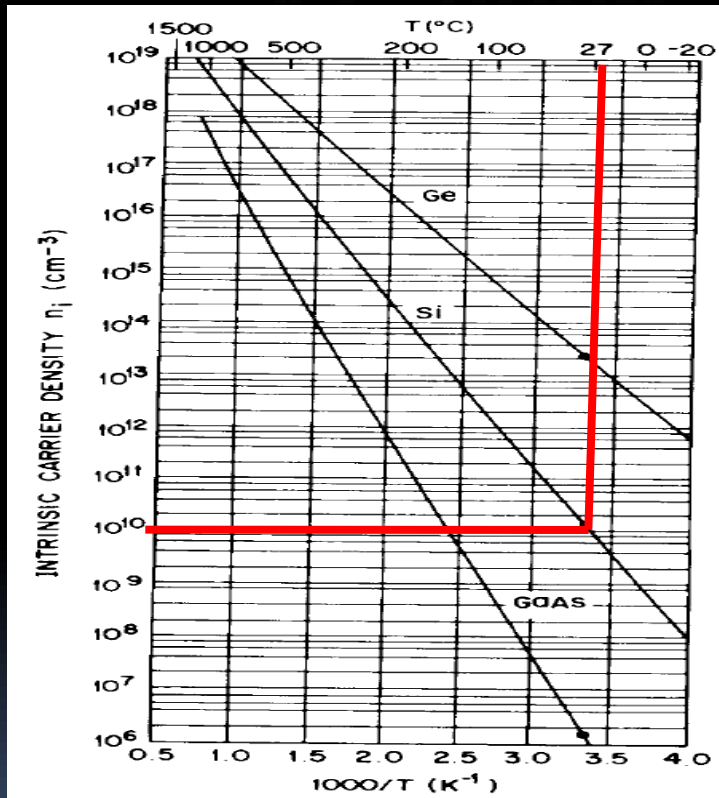


Impurity
concentration
 $\sim 10^{15} / \text{cm}^3$

Room
Temperature
 $\sim 300 \text{ K}$

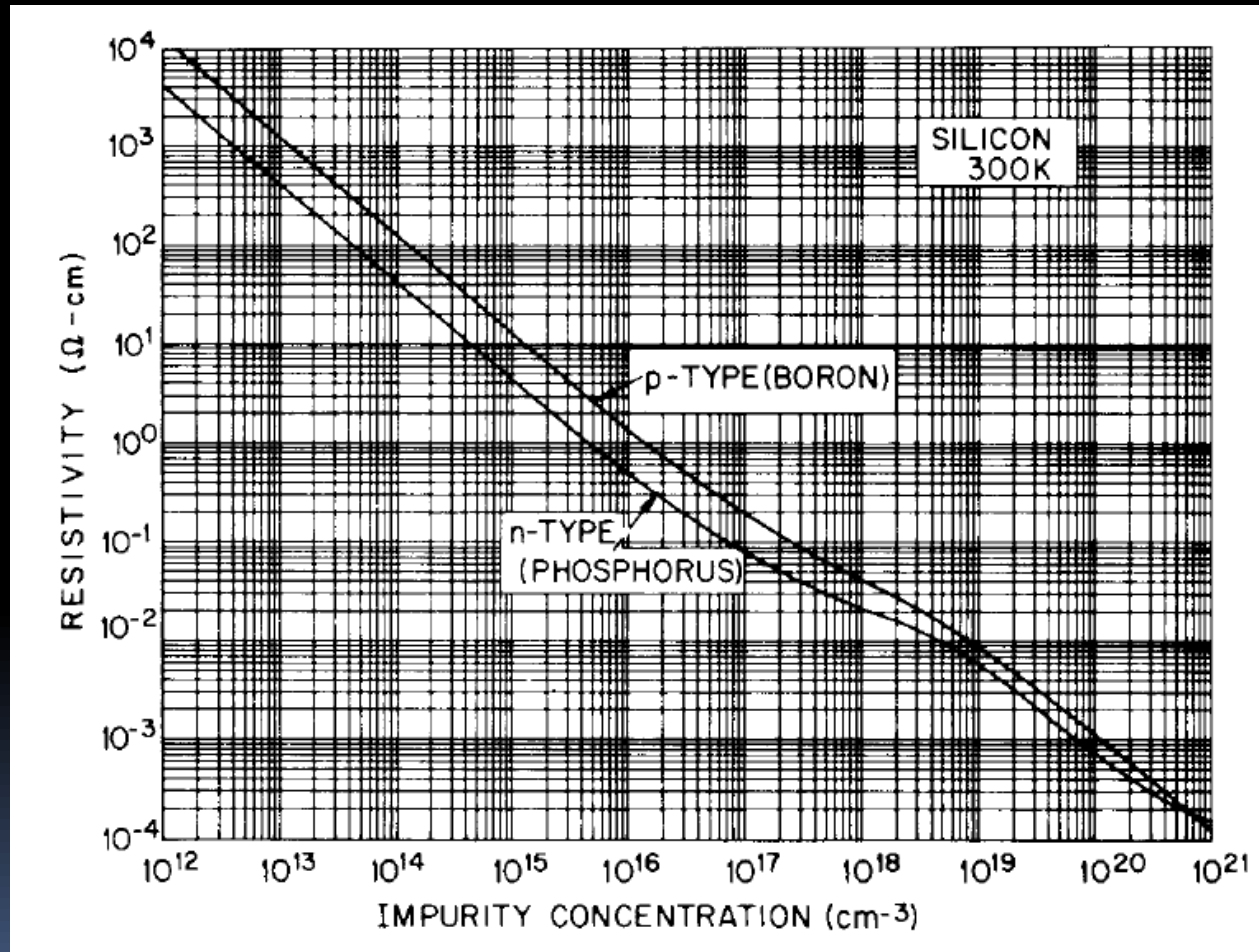
Total Electron Density = INTRINSIC CARRIER DENSITY
+ DOPANT IMPURITY generated Electron Density.

INCREASE IN ELECTRON CONCENTRATION IN DOPED SILICON AT ROOM TEMPERATURE

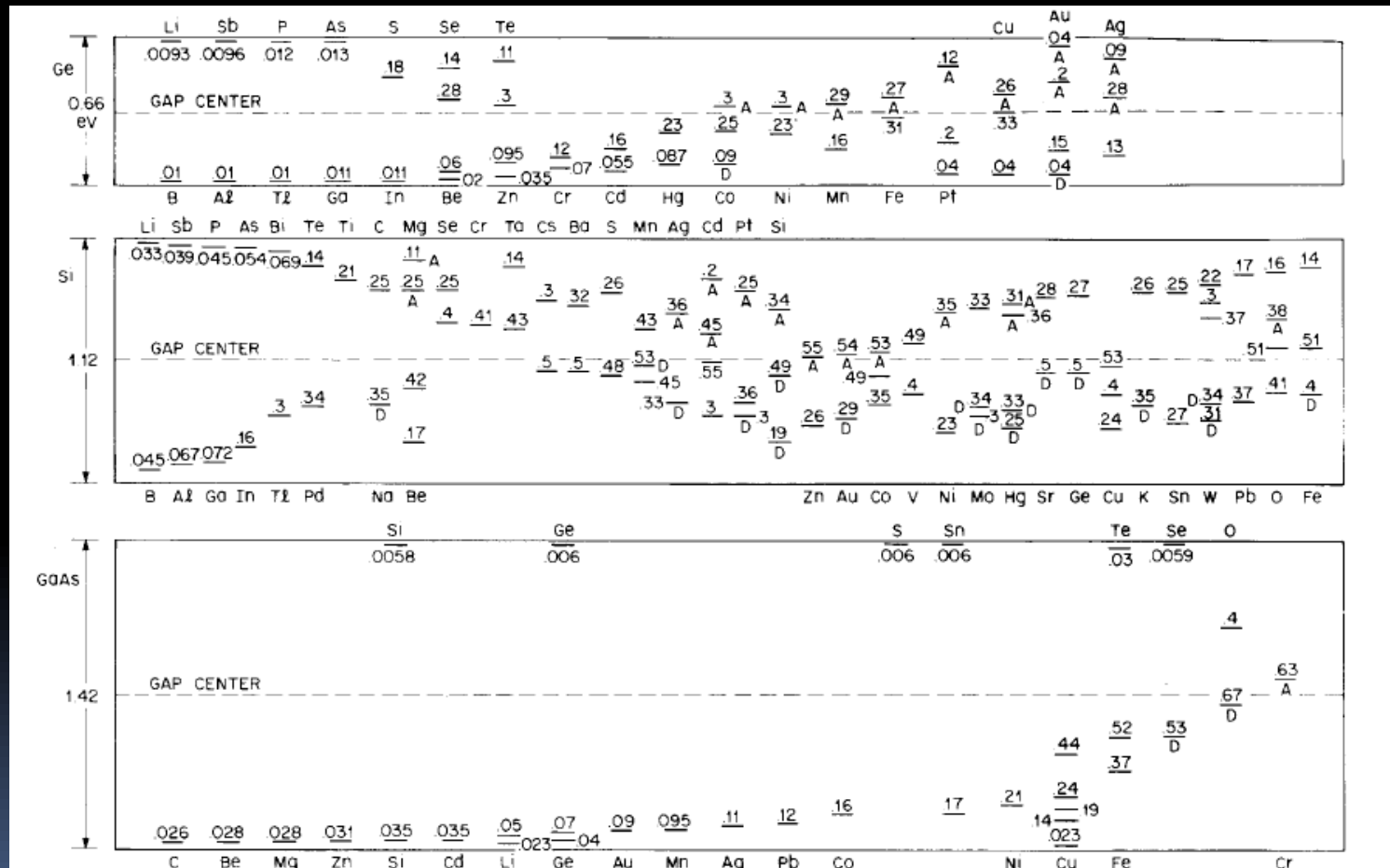


From $\sim 10^{10}$ to $\sim 10^{15}$ per cm^3 !

RESISTIVITY DECREASES WITH INCREASED DOPING



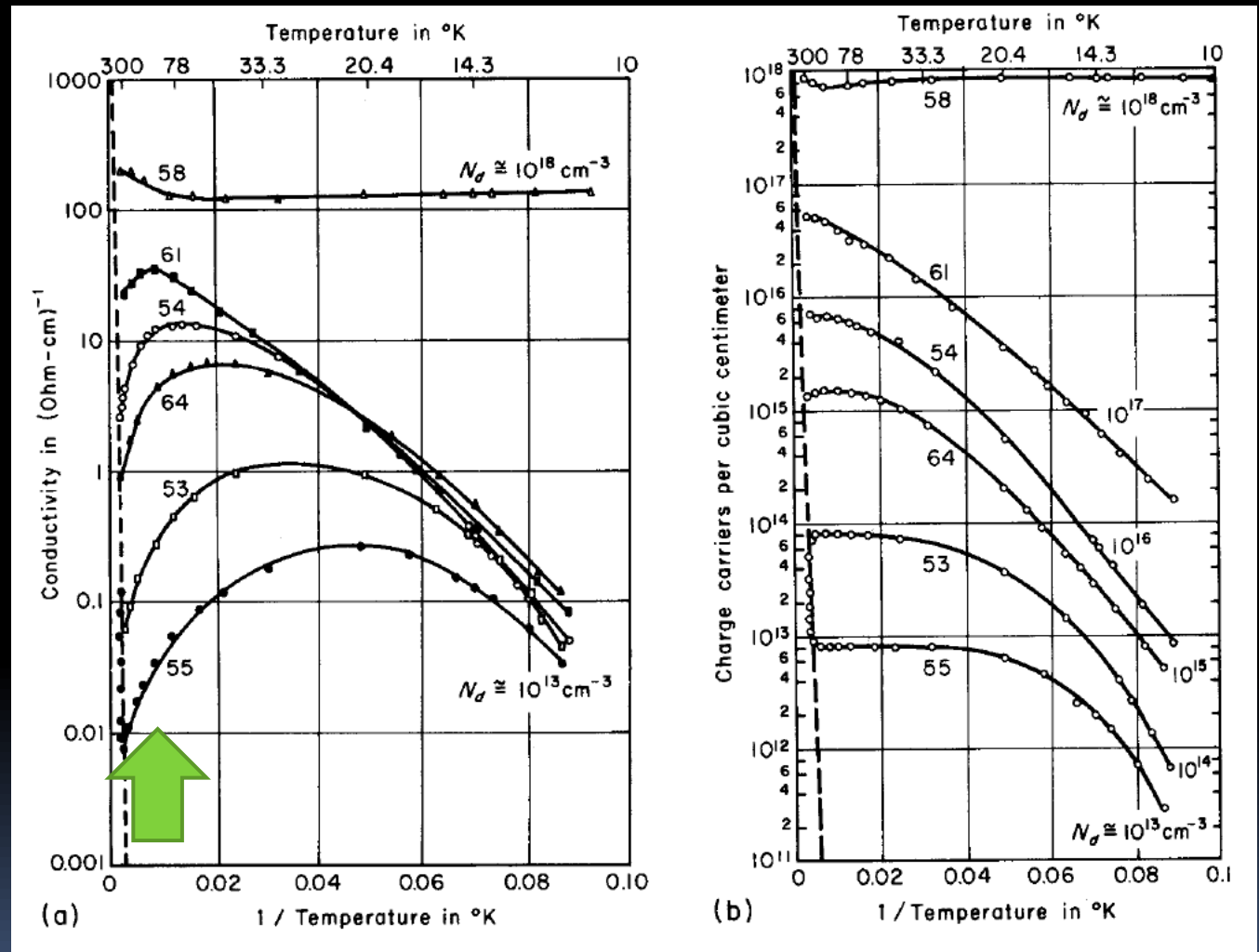
DIFFERENT KIND OF IMPURITY DOPANTS



WHAT IS WRONG IN ASKING.....

**“WHAT IS THE DIFFERENCE BETWEEN
A METAL AND A SEMICONDUCTOR IN
TERMS OF THEIR TEMPERATURE
DEPENDENCE OF THE RESISTANCE ?”**

CONDUCTIVITY OF GERMANIUM ALSO DECREASES WITH INCREASING TEMPERATURE



Ch.1 , S. M Sze, Physics
of Semiconductor
Devices

Conductivity again increases as soon as the intrinsic generation of carriers dominates the impurity doped carriers.

TEMPERATURE DEPENDENCE OF RESISTIVITY

Around room temperature and below room temperature the resistivity of Si can increase with increasing temperature like metals !!!!

Only when at very high temperature (say) $> 500^{\circ}\text{C}$ the **Silicon becomes Intrinsic** and resistivity then decreases with increasing temperature.