

PH3244
Experiment - 1
Measurement of Suseptibility of a liquid or a solution by
Quincke's method

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Synopsis

In this experiment we try to measure the suseptibility of given substance by Quincke's method.

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I. THEORY AND PROCEDURE

A. Theory

The Quincke's method is used to determine magnetic susceptibility of diamagnetic or paramagnetic substances in the form of a liquid or an aqueous solution. When an object is placed in a magnetic field, a magnetic moment is induced in it. Magnetic susceptibility χ is the ratio of the magnetization I (magnetic moment per unit volume) to the applied magnetizing field intensity H . The magnetic moment can be measured either by force methods, which involve the measurement of the force exerted on the sample by an inhomogeneous magnetic field or induction methods where the voltage induced in an electrical circuit is measured by varying magnetic moment. The Quincke's method like the Gouy's method belongs to the former class. The force f on the sample is negative of the gradient of the change in energy density when the sample is placed,

$$f = \frac{d}{dx} \left[\frac{1}{2} \mu_0 (\mu_r - \mu_{ra}) H^2 \right] = \frac{1}{2} \mu_0 (\chi - \chi_a) \frac{d}{dx} H^2 \quad (1)$$

Here μ_0 is permeability of the free space and μ_r, χ and μ_{ra}, χ_a are respectively relative permeability and susceptibility of the sample and the air which the sample displaces. The force acting on an element of area A and length dx of the liquid column is $fAdx$, so the total force F on the liquid is

$$F = A \int f dx = \frac{A\mu_0}{2} (\chi - \chi_a) (H^2 - H_0^2) \quad (2)$$

where the integral is taken over the whole liquid. This means that H is equal to the field at the liquid surface between the poles of the magnet and H_0 is the field at the other surface away from the magnet. The liquid (density ρ) moves under the action of this force until it is balanced by the pressure exerted over the area A due to a height difference h between the liquid surfaces in the two arms of the U-tube. It follows that

$$F = Ah(\rho - \rho_a)g$$

Or

$$\chi = \chi_a + \frac{2}{\mu_0} g(\rho - \rho_a) \frac{h}{(H^2 - H_0^2)} \quad (3)$$

In actual practice χ_a , density of air ρ_a and H_0 are negligible and can be ignored and the above expression simplifies to

$$\chi = \frac{2\rho gh}{\mu_0 H^2} \quad (4)$$

This equation shows that by plotting h as a function of H^2 , the susceptibility χ can be determined from the slope. The expressions in C.G.S units are given by

$$\chi = \chi_a + 2g(\rho - \rho_a) \frac{h}{(H^2 - H_0^2)} \quad (5)$$

$$\boxed{\chi = \frac{2\rho gh}{H^2}} \quad (6)$$

B. Procedure

A schematic diagram of Quincke's set up is shown in Fig.1. One limb of the glass U-tube is very narrow (about 2 – 3 mm in diameter) and the other one quite wide. The result is that a change in the level of the liquid in the narrow limb does not affect the level in the wider limb. The narrow limb is placed between the pole pieces of an electromagnet shown as N-S such that the meniscus of the liquid lies symmetrically between N-S. The length of the limb should be sufficient enough to keep the lower extreme end of this limb well outside the field of the magnet. The rise or fall h on applying the field is measured by means of a traveling microscope fitted with a micrometer scale of least count of order 10^{-3} cm.

1. Measurement of H

1. Fix the air gap between the pole pieces of the electromagnet to the minimum distance required to insert Quinck's tube without touching the pole pieces.
2. Measure the air gap. Each time the air gap changes, the graph will change.
3. Mount the Hall probe of the Digital Gaussmeter, DGM-102 in the wooden stand provided and place it at the centre of the air gap such that the surface of the probe is parallel to the pole pieces. The small black crystal in the probe is its transducer, so this part should be at the centre of the air gap.
4. Connect the leads of the Electromagnet to the Power Supply, bring the current potentiometer of the Power Supply to the minimum. Switch on the Power Supply and the Gaussmeter.
5. Slowly raise the current in the Power Supply and note the magnetic field reading in the Gaussmeter.
6. Plot the graph between the current and the magnetic field. This graph will be linear for small values of the current and then the slope will decrease as magnetic saturation occurs in the material of the pole pieces.

2. Measurement of h

1. Test and ensure that each unit (Electromagnet and Power Supply) is functioning properly.
2. Measure the density ρ of the specimen (liquid or solution) by specific gravity bottle. If the mass of empty bottle is w_1 , filled with specimen w_2 and filled with water w_3 , then

$$\rho = \rho_{\text{water}} \frac{w_2 - w_1}{w_3 - w_1} \quad (7)$$

3. *Scrupulous* cleaning of the tube is essential. *Thoroughly* clean the Quincke's tube, rinse it well with distilled water and dry it (preferably with dry compressed air). Do not use the tube for longer than one laboratory period without recleaning it.
4. Keep the Quincke's tube between the pole pieces of the magnet as shown in Fig.1. The length of the horizontal connecting limb should be sufficient to keep the wide limb out of the magnetic field.
5. Fill the liquid in the tube and set the meniscus centrally within the pole pieces as shown. Focus the microscope on the meniscus and take reading.

6. Apply the magnetic field H and note its value from the calibration, which is done earlier as an auxiliary experiment. Note whether the meniscus rises up or descends down. It rises up for paramagnetic liquids and solutions while descends down for diamagnetics. Readjust the microscope on the meniscus and take reading. The difference of these two readings gives h for the field H . The magnetic field between the poles of the magnet does not drop to zero even when the current is switched off. There is a residual magnetic field R which requires a correction.
7. Measure the displacement h as a function of applied field H by changing the magnet current in small steps. Plot a graph of h as a function of H^2 .

II. OBSERVATION

A. Measurement of h and H

| Current (A) | Main Scale (mm) | Circular Scale (0.01 mm) |
|----------------|--------------------|-----------------------------|
| 0.00 | 9.5 | 48 |
| 0.80 | 11.5 | 10 |
| 1.00 | 12.5 | 5 |
| 1.25 | 13.5 | 3 |
| 1.50 | 16.5 | 2 |
| 1.75 | 18.5 | 12 |
| 2.00 | 20.0 | 19 |
| 2.50 | 24.0 | 12 |
| 2.75 | 24.0 | 38 |
| 3.00 | 24.5 | 15 |
| 3.25 | 25.0 | 40 |
| 3.30 | 25.0 | 48 |

TABLE I: Δh measurement; Data taken on 21/01/26

| Current (A) | Field Strength (Gauss) |
|----------------|---------------------------|
| 0.00 | 423 |
| 0.10 | 729 |
| 0.20 | 1052 |
| 0.30 | 1400 |
| 0.40 | 1769 |
| 0.50 | 2150 |
| 0.60 | 2550 |
| 0.70 | 2970 |
| 0.80 | 3380 |
| 0.90 | 3790 |
| 1.00 | 4210 |
| 1.50 | 6340 |
| 2.10 | 8490 |
| 2.50 | 10230 |
| 3.00 | 11840 |
| 3.50 | 13260 |
| 3.55 | 13390 |

TABLE II: H measurement; Data taken on 21/01/26

add least count
 add 3M, conc and other data, also add the previous data too of that table to not include Cus04 ka data, why repeat add and toehr stuff toto

B. Comments and Sources of Error

1. sadf
2. adf

III. ANALYSIS

A. Error propagation

Below we summarise the error propagation used to obtain uncertainties on the predicted magnetic field at the height-measurement currents and on the derived quantity B^2 . The

notation is:

- I : current (measured where height h was taken).
- $B(I) = m_B I + c_B$: linear fit of magnetic field vs current (fit parameters m_B, c_B).
- σ_m, σ_c : standard uncertainties (one-sigma) of the fitted parameters m_B, c_B .
- $\text{Cov}(m, c)$: covariance between m_B and c_B from the fit (if available).
- $\hat{B} = B_{\text{pred}} = m_B I + c_B$: predicted B at a given I .
- $\sigma_{\hat{B}}$: uncertainty on the predicted B .
- $x = B^2$ and σ_x : the squared field and its uncertainty.
- Δh : measured meniscus rise, with uncertainty σ_h .

a. *Uncertainty on the predicted field \hat{B} .* Using standard linear propagation for a function $f(m, c) = mI + c$,

$$\text{Var}(\hat{B}) = \left(\frac{\partial f}{\partial m} \right)^2 \text{Var}(m) + \left(\frac{\partial f}{\partial c} \right)^2 \text{Var}(c) + 2 \left(\frac{\partial f}{\partial m} \right) \left(\frac{\partial f}{\partial c} \right) \text{Cov}(m, c).$$

Since $\partial f / \partial m = I$ and $\partial f / \partial c = 1$, this becomes

$$\sigma_{\hat{B}}^2 = I^2 \sigma_m^2 + \sigma_c^2 + 2I \text{Cov}(m, c).$$

If the covariance $\text{Cov}(m, c)$ is not available or is neglected, use the conservative approximation

$$\sigma_{\hat{B}} \approx \sqrt{I^2 \sigma_m^2 + \sigma_c^2}.$$

b. *Uncertainty on B^2 .* For $x = B^2$ and small relative uncertainty on B , linear propagation gives

$$\sigma_x = \left| \frac{d(B^2)}{dB} \right| \sigma_B = 2|B| \sigma_B.$$

Applied to the predicted field,

$$\sigma_{B^2} \approx 2 \hat{B} \sigma_{\hat{B}}.$$

(Use the sign of \hat{B} consistently; in practice $\hat{B} > 0$ here.)

c. *Uncertainty on the height measurement Δh .* Each height reading was formed from a main scale reading plus a circular-count reading. Treat the least-counts as uniform distribution uncertainties and combine in quadrature. If LC_{main} is the main-scale least count (mm) and each circular step equals δ_{circ} (mm) with a least-count LC_{circ} in counts, then

$$\sigma_h = \sqrt{\left(\frac{LC_{\text{main}}}{\sqrt{12}} \right)^2 + \left(\frac{\delta_{\text{circ}} \cdot LC_{\text{circ}}}{\sqrt{12}} \right)^2}.$$

In the code this was implemented as

$$\sigma_h = \sqrt{(h_{\text{lc}}/\sqrt{12})^2 + (0.01/\sqrt{12})^2},$$

where h_{lc} denotes the main-scale least count (in mm) and 0.01 mm is the circular-step size converted to mm.

d. Fitting Δh vs B^2 . With $x = B^2$ (predicted at the height currents) and $y = \Delta h$, each data point has uncertainties σ_x and $\sigma_y = \sigma_h$. Because both x and y have errors we used orthogonal distance regression (ODR) to fit the linear model

$$y = \alpha x + \beta,$$

and ODR returns best-fit parameters $\hat{\alpha}, \hat{\beta}$ together with their standard uncertainties $\sigma_\alpha, \sigma_\beta$. These reported $\sigma_\alpha, \sigma_\beta$ are used as the final uncertainties for the slope and intercept of Δh vs B^2 .

e. Remarks and alternatives

1. The propagation formula that includes $\text{Cov}(m, c)$ is the formally correct one; neglecting the covariance may under- or over-estimate $\sigma_{\hat{B}}$ depending on the sign and magnitude of $\text{Cov}(m, c)$. If available, use the full covariance matrix from the fit (ODR provides `cov_beta`).
2. An alternative to using the linear fit $B(I)$ for predicting \hat{B} is to interpolate measured $B(I)$ (e.g. using `scipy.interpolate`) and propagate interpolation uncertainty via bootstrap or measurement uncertainties. This avoids model bias at the cost of potentially larger pointwise noise.
3. A more rigorous (and often preferable) approach is to perform a *joint* fit of the two datasets by fitting a model of the form

$$\Delta h(I) = \alpha(m_B I + c_B)^2 + \beta,$$

and fitting simultaneously for α, β, m_B, c_B . This treats the relation between B and h consistently and yields the full covariance among all parameters; implementation is straightforward with ODR or a non-linear least squares package.

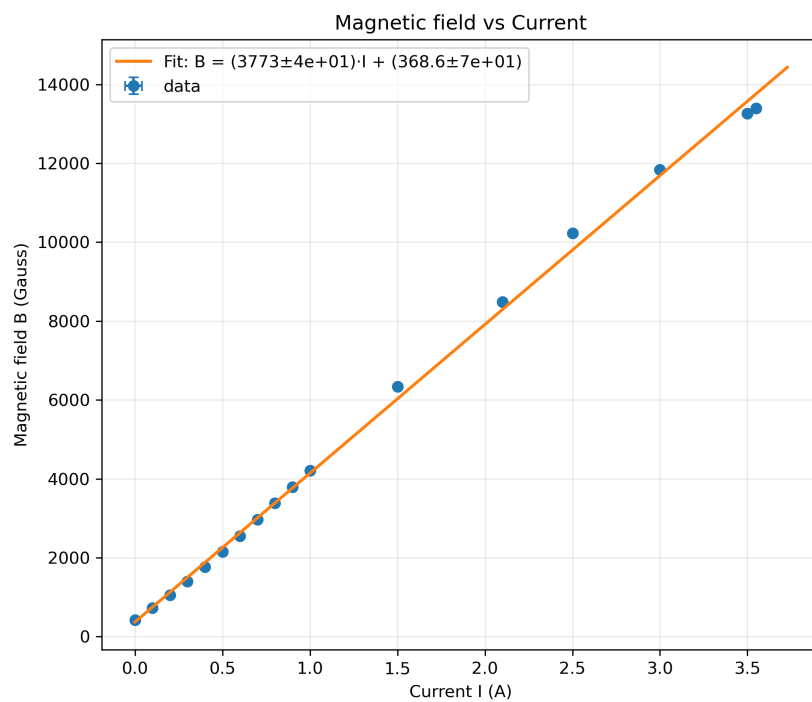
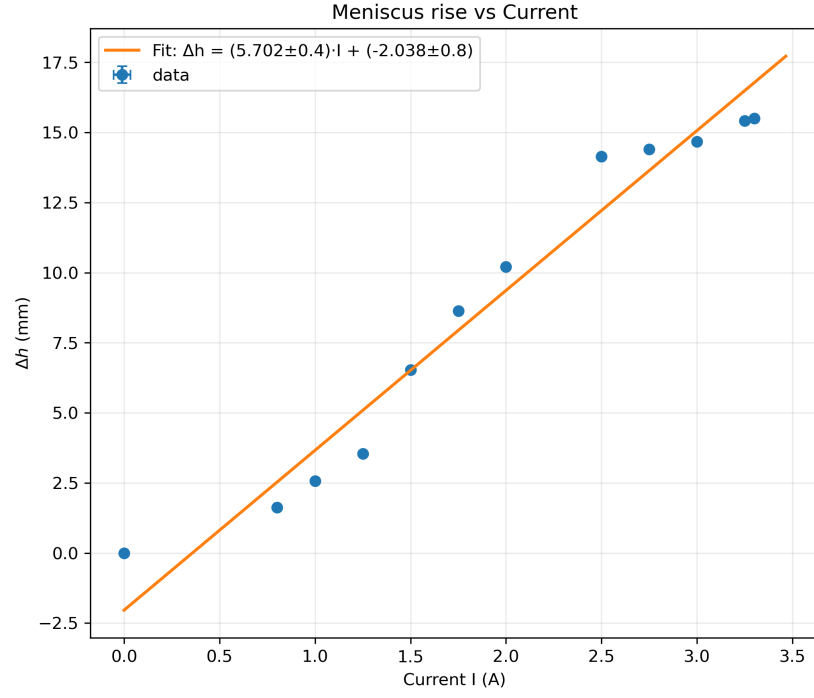
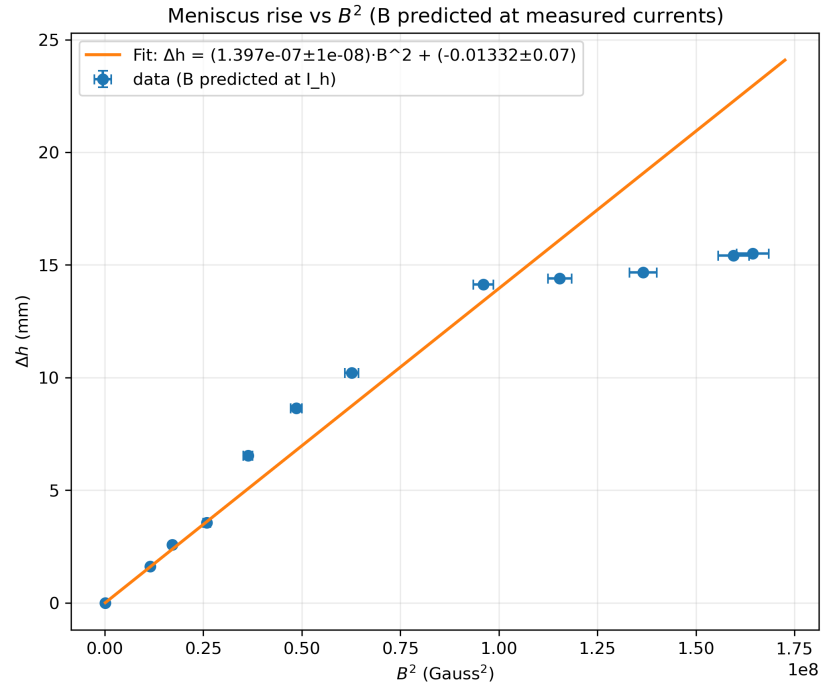
B. Plots

FIG. 1: Magnetic field B vs current I with linear fit.

FIG. 2: Meniscus rise Δh vs current I .FIG. 3: Meniscus rise Δh vs squared magnetic field B^2 .

IV. RESULT AND CONCLUSION

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REFERENCES