

# Short summary of ATLAS+CMS Dark Matter Forum Recommendations For Signal MC

ATLAS+CMS Dark Matter Forum

May 5, 2015

This document outlines the choices made for the MC production and parameter scans for the simplified models recommended by the ATLAS/CMS Dark Matter Forum. This summary is aimed at defining the models and parameters for the upcoming signal Monte Carlo production by the two collaborations. Motivations for these choices, reinterpretation details, and a full bibliography will be provided in an upcoming document concluding the works of the Forum.

## *1 Prioritized lists of simplified models for MET+X searches*

### *1.1 Recommended models for all MET+X analyses*

For all MET+X analyses, simplified models should replace EFTs as the highest-priority dark matter interpretation. The consensus of the Forum is that all of these analyses should consider a set of simplified models involving a single, mediating particle with mass  $m_{med}$ . The two underlying choices for the model recommendations by the ATLAS/CMS DM Forum are:

1. the dark matter particle is a Dirac fermion of mass  $m_{DM}$ ;
2. only interactions consistent with Minimal Flavor Violation are allowed.

The highest priority models discussed within the Forum are:

- a. spin-1  $s$ -channel mediators with vector or axial-vector coupling to SM quarks  $g_q$  and vector or axial-vector coupling to a Dirac fermion WIMP  $g_{DM}$ ;
- b. spin-0  $s$ -channel mediators with scalar or pseudo-scalar coupling to SM quarks  $g_q$  and scalar or pseudo-scalar coupling to a Dirac fermion WIMP  $g_{DM}$ ;
- c. and colored spin-0  $t$ -channel mediators with coupling  $g$  at the vertex between quark, WIMP and mediator.

In all three cases, the coupling of the mediator to quarks  $g_q$  is identical for all three generations. In the case of the scalar and pseudo-scalar models, we assume a separate Yukawa coupling ( $y_q = m_q/v$ ) between the quarks and the mediators multiplying  $g_q$ ; the specific case of scalar mediator decaying into two top quarks is discussed in Section 1.3. The coupling between mediator and WIMP does not have any Yukawa structure.

There are thus four parameters in the first two models ( $m_{DM}$ ,  $m_{med}$ ,  $g_q$ , and  $g_{DM}$ ) and three parameters in the last model ( $m_{DM}$ ,  $m_{med}$ ,  $g$ ). We assume that no additional visible or invisible decays contribute to the width of the mediator. We provide formulae for calculation of the “minimal width” for any given choice of the four parameters. Early run 2 MET+X analyses are not sensitive to changes in the kinematic distributions that result from increasing the width up to  $\Gamma \approx m_{med}$ . The exceptional case of a narrow, off-shell mediator is addressed later.

All types of  $s$ -channel mediators are important, and ideally an analysis would have to include all cases. On the other hand, the vector and axial-vector models (including those with mixed couplings) produce nearly indistinguishable kinematic distributions for the MET+X analyses, and the scalar and pseudo-scalar models are also very similar. If one cannot generate all models, one could generate the full grid of mass points for only one of each type (for example, the pure axial vector and pure pseudo-scalar models, motivated respectively by complementarity with direct and indirect detection and galactic center excess interpretations).

Studies of the  $s$ -channel models include comparisons of the kinematic distributions across a scan of each of the four parameters. From these studies, we recommend generation of a minimal set of a total of approximately 30 points per simplified model for the monojet analysis, as shown in Table 1, combinations of parameter choices in order to capture the full variety of shapes possible in the models.

| $m_{DM}$ (GeV) | $m_{med}$ (GeV) |    |    |     |     |     |     |      |             |      |
|----------------|-----------------|----|----|-----|-----|-----|-----|------|-------------|------|
| 1              | 10              | 20 | 50 | 100 | 200 | 300 | 500 | 1000 | <b>2000</b> | 5000 |
| 10             | 10              | 15 | 50 | 100 |     |     |     |      |             | 5000 |
| 50             | 10              |    | 50 | 95  | 200 | 300 |     |      |             | 5000 |
| 150            | 10              |    |    |     | 200 | 295 | 500 |      |             | 5000 |
| 500            | 10              |    |    |     |     |     | 500 | 995  | <b>2000</b> | 5000 |
| 1000           | 10              |    |    |     |     |     |     | 1000 | <b>1995</b> | 5000 |

Table 1: Simplified model benchmarks for all  $s$ -channel simplified models (spin-1 and spin-0 mediators decaying to Dirac DM fermions taking the minimum width for  $g_q = g_{DM} = 1$ ). Points in **bold** are only generated for the vector/axial vector cases, while points in *italics* are generated for the monojet analysis but not for the search including heavy quarks. This table corresponds to 29 points for monojet vector/axial vector models, 26 points for monojet scalar/pseudoscalar models and 24 points for  $t\bar{t}$ +MET scalar/pseudoscalar models.

Namely the bulk of the grid scan for the  $s$ -channel models will cover from 10 GeV to 2 TeV (10 GeV to 1 TeV) in mediator mass for the vector and axial vector (scalar and pseudoscalar) models, and from 1 GeV to 1 TeV in WIMP mass. A high-mediator-mass point is

added for each of the WIMP masses in order to provide an EFT limit for theory reinterpretation and comparisons with previous results.

For very heavy mediators, changing the minimal width (by changing the coupling) also affects the signal kinematic distributions. It is unclear whether the MET+jet analysis would be able to observe these signals during Run 2; therefore we recommend a more limited scan which still allows further studies. For the highest mediator mass points for the vector and axial vector models, signals should be generated with couplings in a grid of  $g_{DM} = g_q = 1.0, 0.3, 0.5, 1.45$ , while for the scalar and pseudoscalar signals the recommended grid is  $g_{DM} = g_q = 1.0, 0.5, 1.0, 2.0, 3.0$ .

Results for other choices of parameters can be related to these through simple rescalings of the total rate, using cross-section formulae that will be provided by the forum.

The parameter space of the  $t$ -channel model has not yet been studied in the same detail as the  $s$ -channel models. **[Point being finalized: This model, and its parameter scan, is still under study.]** This model is parallel to, and partially motivated by, the squark of the MSSM. Since the  $t$ -channel model can produce signatures similar to those arising from squark production in the MSSM, investigations are needed of what ranges of  $g$ ,  $m_{med}$ , and  $m_{DM}$  are already excluded. Since the relevant coupling is arbitrary in the  $t$ -channel model, while it is weak-scale in the MSSM, additional production topologies need to be considered. The coordination of a generic  $t$ -channel scan with MSSM scans would need more thought.

#### 1.1.1 Technical implementation of the $s$ -channel models

There are several matrix element implementations of the  $s$ -channel vector and scalar mediated DM production. They are available in POWHEG, MadGraph and also MCFM. We recommend the POWHEG implementation of DM pair production +1 parton at NLO (the DMV model), and the associated theoretical uncertainties computed by POWHEG. POWHEG provides a set of weights in the LHE v3 format, which will be propagated through the ATLAS and CMS software to allow a determination of signal uncertainties without generating additional sets of samples. A prescription for determining these uncertainties from the weights is being provided through the forum by the POWHEG authors. For ATLAS and CMS, the necessary version of POWHEG is centrally available (V2.0 in both cases, revision 3049 for CMS and revision 3033 for ATLAS).

**[Point being finalized: The recommendations for assessing systematic uncertainties has to be approved by both collaborations. Once ready, ATLAS and CMS should discuss whether this can be harmonised.]**

*ATLAS-specific implementation details* The  $s$ -channel vector mediator model is included in this pilot MC15 request [ATL15b], and details about this request are in the monojet analysis twiki page [ATL15e]. The necessary lines for the evaluation of both scale and PDF uncertainties are available in the jobOptions according to instructions [ATL15f], and discussion is ongoing on how to propagate them to the software on the PAT-help hypernews (the thread can be found at this url <http://cern.ch/go/r8jk>). These weights must be available both in the ATHENA release both when the production is requested and at the analysis level, since it will not be possible to reweight these signal samples at truth-level for PDFs and scales as it was done in previous analyses. This requires the samples to be produced in a release that includes both Pythia8.2 (19.2.8.3 and newer) and LHE3 support. Once both are validated, samples will be produced including weights if they have the correct lines in the jobOptions, as it is the case for the pilot production above. If it is not possible to include the weights in time for the nominal MC15 production, then another request can be made that includes them at a later date.

#### 1.1.2 Technical implementation of the $t$ -channel model

The  $t$ -channel model is implemented in MadGraph. It is recommended to compute at LO and apply matching up to two additional partons beyond the Born level process. Theoretical uncertainties are to be calculated using the MadGraph SysCalc package, which performs scale and PDF variations in a similar manner as POWHEG. The relevant parameter card can be found on the Forum SVN repository [For15g].

*ATLAS-specific implementation details of  $t$ -channel model* This model has already been used as a benchmark for the ATLAS 8 TeV mono-Z search, and it is being investigated privately by both mono-jet and mono-W/Z analysers. Its ATLAS implementation using on-the-fly MadGraph will follow the pilot request for the monophoton D5 operator [ATL15b]. The treatment of the systematic uncertainties and their harmonization with CMS depends on the testing and in the inclusion of SysCalc [Sys15], a MadGraph module that is already available in MadGraphControl in ATHENA but has not been fully tested yet due to the lack of LHE3 support. As in the POWHEG case, Pythia 8.2 or higher is also necessary for LH3 weights to be handled correctly.

## 1.2 Specific models for analyses of MET+ $b$ quark(s)

### 1.2.1 Single $b$ +MET

We consider a simplified model for single- $b$ -jet signatures, with a colored scalar mediator coupling to Dark Matter and a  $b$ -quark. This model is described in Ref. [ABHL14]. The model is similar to the MSSM with a light bottom squark and neutralino, and is thus a flavor-specific example of a  $t$ -channel model. Preliminary studies show that generation of a grid of roughly 10 points in DM mass, 15 points in mediator mass and two coupling values is sufficient to cover the parameter space probed by the LHC with early searches.

This model is available in MadGraph, and the relevant card files are available in the Forum SVN repository [For15f].

### 1.2.2 $b\bar{b}$ +MET

As noted earlier, the  $t$ -channel model produces both mono-jet and di-jet +MET signatures after applying cuts. In the case when only a  $b$ -flavored mediator is accessible, both mono- $b$  and  $b\bar{b}$ +MET signatures are possible. **[Point being finalized: The details for the implementation of these models is still under investigation. If the studies are not available in time for the Forum conclusions, this will be included as a model to be investigated for future works.]**

## 1.3 Specific models for analyses of MET+ $top$ (s)

### 1.3.1 $t\bar{t}$ +MET

The pair production of top-flavored resonances that decay to a top quark and DM is the subject of intensive study by the experimental collaborations, and are not in the purview of this Forum.

Instead, we focus on the implications of the scalar/pseudo-scalar model introduced earlier, with the mediator produced in association with a  $t\bar{t}$  pair.

As mentioned in Sec. 1.1, we assume mass dependent Yukawa couplings to quarks for spin-0 scalar and pseudoscalar mediators. Couplings to top quarks will dominate, producing a tree-level signature of a  $t\bar{t}$  pair in association with missing transverse momentum from the two DM particles. The parameter scan for this model is based on three different kinematic regimes: the case of a very heavy mediator with  $m_{med} > 2m_{top}$ , the case of an on-shell mediator below the  $t\bar{t}$  threshold ( $m_{med} > m_{DM}$ ,  $m_{med} < 2m_{top}$ ) and the case of an off-shell mediator below the  $t\bar{t}$  threshold ( $m_{med} < 2m_{DM}$ ,  $m_{med} < 2m_{top}$ ). Large kinematic differences are found between on-shell and off-shell mediator regimes, while the total width changes slope depending on

whether one is above or below the  $t\bar{t}$  threshold. Scalar and pseudoscalar models show kinematic differences as the mediator mass is reduced, due to the difference in coupling of the top pair to the mediator. These considerations lead to the choice of the grid scan shown in Table 1, harmonised with the choices for the generic models and signatures, to be generated for both scalar and pseudoscalar cases separately.

The effect of varying the couplings (and therefore the mediator width) is small in the on-shell regime, but is significant in the case of very small or very large couplings and large mediator masses, with the latter effect due to parton distribution function suppression. LHC searches will be sensitive to those masses after roughly  $30 \text{ fb}^{-1}$  of data. A table of cross-sections will be provided, for several values of coupling strengths, to aid reinterpretation and validate the suggested scaling of the cross-sections that substitutes the scan in couplings for early searches.

*Technical implementation of models for  $t\bar{t}$ +MET* The implementation of this model is available in POWHEG, MadGraph and MCFM. In the case of searches with heavy quarks in the final state that are sensitive only to the tree-level mediator decays into DM particles in association with two top quarks, the MadGraph implementation is suggested for Run-2 searches. The relevant parameter cards are available in the Forum SVN repository [For15b].

*ATLAS-specific implementation details* On-the-fly MadGraph has already been requested for DC14 [ATL15d]. Instructions on how to run the private production for the tests performed within the Forum are also available on a Twiki [ATL15c]. The ATHENA implementation of this model using on-the-fly MadGraph will follow closely the MC15 pilot request for the monophoton D5 operator [ATL15b]. The implementation details necessary for the treatment of systematic uncertainties for MadGraph are described the  $t$ -channel model.

### 1.3.2 Single top+MET

Models of single top+MET production are different from the other models under consideration, since they truly involve a single top quark in the final state.

Two simplified models, described in Refs. [AFM11, BCDF15], are recommended for the signature of a single top quark+MET<sup>1</sup>:

- $s$ -channel production of a colored scalar resonance, decaying into a top quark and a fermion that can decay in two DM candidates;

<sup>1</sup> The non-resonant production contributes to both both MET+t and MET+ $t\bar{t}$  final states.

- $s$ - and  $t$ - channel non-resonant production of a top quark and a color singlet, vector mediator which then decays invisibly.

The parameters for the  $s$ -channel resonant production are:

- the couplings of the scalar resonance to down-type quarks, all considered equal and denoted with  $a_{res}^q$ ;
- the couplings of the DM particle to up-type quarks and to the mediator, denoted with  $a_{res}^{1/2}$ ;

The resonance width is computed according to the couplings, assuming no other decays in addition to those in DM and quarks. Fixing the branching fraction to DM particles to 100% would lead to a single value for the coupling of the resonance to DM particles.  $a_{res}^{1/2}$  may however become a parameter to be scanned, after further studies on whether changing the resonance width leads to significant changes in the kinematics of the model. This point will be settled after further studies.

The only relevant parameters considered for the  $s$ - and  $t$ -channel non-resonant production are the coupling of the vector mediator to up-type quarks, all considered equal and denoted with  $a_{non-res}$ . The width of the mediator does not affect the model kinematics. A proposal for the scan in the parameters of these models will be prepared after further study from the ATLAS and CMS analysers.

*Technical implementation of the single top+MET models* Card files for MadGraph are provided on the Forum SVN repository [For15e].

*ATLAS-specific implementation details* MadGraph samples have already been requested for DC14 [ATL15g]. The ATHENA implementation of this model using on-the-fly MadGraph will follow closely the MC15 pilot request for the monophoton D5 operator [ATL15b]. The implementation details necessary for the treatment of systematic uncertainties for MadGraph are described the  $t$ -channel model.

#### 1.4 Specific models for MET+EW boson searches

Searches with an electroweak boson (photon, Z, Higgs or W) in the final state should include the general models recommended in Section 1.1. In addition to these, we recommend models where an electroweak boson participates in the dark matter interaction itself, rather than appearing as final state radiation.

Though the requisite set simplified models have not yet been fully developed by the theoretical community, it is here especially that searches in MET+photon, W, Z, or H may play a more important

role than in the general models considered above. Until all necessary simplified models are developed, we must recommend use of EFT contact operators in some cases, provided the region of validity is carefully studied following the relevant recommendations in Section 2.

#### 1.4.1 *MET+photon,Z,W*

These searches should consider an effective  $VV\chi\chi$  vertex, with two electroweak bosons and two WIMPs, via an EFT contact operator. Such operators are available for scalar and fermion WIMPs and  $VV = ZZ, Z\gamma, WW, \gamma\gamma$  (so-called dimension 4, 5, and 7 operators [CNS<sup>+</sup><sub>13</sub>, CHH<sub>15</sub>]). We prioritize the Dirac fermion WIMP (dimension 7) operators, as those have been studied more widely and yield kinematic distributions that are distinct from any of the simplified models described above. Correlations between the different MET+photon, Z and W signatures are accounted for by the use of two model parameters encapsulating the various couplings ( $k_1$  and  $k_2$ ). Since only the cross section, and not the kinematic distributions, depend on the value of those parameters, we recommend the generation of these benchmarks using only a single value of these parameters. No difference in the kinematics is seen between the scalar and pseudoscalar models considered. This leads to the parameter scan for this benchmark to correspond to 8 different DM masses: 1, 10, 50, 100, 200, 400, 800 and 1300 GeV. The scale of the EFT for the generation is fixed at 3 TeV, and the couplings are set to  $k_1 = k_2 = 1$ .

*MET+W enhancements* The  $s$ -channel model mentioned in Subsection 1.1 has previously been simulated using different coupling parameters between the mediator and the up and down quarks, leading to an increase in cross-section for W boson production. It has recently been shown [BCD<sup>+</sup><sub>15</sub>] that this is due to a violation of gauge invariance. For this reason, we do not recommend this practice anymore, and restrict the grid scan to equal couplings between the quarks. Simplified models overcoming this issue are under development, beyond the timescale of the Forum. Meanwhile, the  $t$ -channel model of the previous subsection may be used to obtain different couplings and interference, but those studies are left for a longer timescale.

#### *Technical implementation of models with EW bosons in the final state*

These models are generated at leading order with MadGraph 2.2.2, using Pythia8 for the parton shower. Parameter cards can be found on the Forum SVN repository [For<sub>15a</sub>].



*ATLAS-specific implementation details* MadGraph samples have already been requested for DC14 [ATL15g]. The ATHENA implementation of this model using on-the-fly MadGraph will follow closely the MC15 pilot request for the monophoton D5 operator [ATL15b]. The implementation details necessary for the treatment of systematic uncertainties for MadGraph are described the  $t$ -channel model.

#### 1.4.2 MET+Higgs

**[Point being finalized: These models, and their parameter scan, will be discussed at the next DM Forum meeting on Thursday May 4th.]**

These searches should also consider models with an effective  $VV\chi\chi$  vertex, with two Higgs bosons and two WIMPs coupled via an EFT contact operator. Such operators are available for scalar and fermion WIMPs at dimension 5 [CDM<sup>+</sup>14]. The recommended parameter scan and implementation follows that of the  $VV = ZZ, Z\gamma, WW, \gamma\gamma$  cases, and the model can be found on the Forum SVN repository [For15d].

Two benchmark simplified models [CDM<sup>+</sup>14] are recommended for MET+Higgs searches:

- A model where a vector mediator ( $Z'$ ) is exchanged in the  $s$ -channel, radiates a Higgs or a  $Z$  boson and decays into two DM particles.
- A model where a scalar mediator couples to the SM only through the SM Higgs and decays to two DM particles.

Preliminary studies of these models show that the parameter scan for these models (on DM mass, mediator mass, mediator couplings and width, mixing of the mediator with the  $Z$  boson) can be reduced to a scan in DM mass and mediator mass that follows Table 1 while keeping the other model parameters fixed to the following values:

- coupling between mediator and SM Higgs boson  $g_{hZ'Z'}/m_{med} = 1$
- mediator-DM coupling  $g_{DM} = 1$
- mediator-quark  $g_q = 1/3$
- mixing angle between the new baryonic Higgs in this model and the SM Higgs  $\sin(\theta) = 0.3$

for the vector mediator model, and the following values:

- Yukawa coupling to DM  $y_{DM} = 1$
- new physics coupling between scalar and SM Higgs  $b = 3$
- mixing angle  $\sin(\theta) = 0.3$

for the scalar model.

### Technical implementation of models with Higgs bosons in the final state

These models are generated at leading order with MadGraph 2.2.2, using Pythia8 for the parton shower. The MadGraph implementations of those models can be found on the Forum SVN repository [For15c].

*ATLAS-specific implementation details* MadGraph samples have already been requested for DC14 [ATL15a]. The ATHENA implementation of this model using on-the-fly MadGraph will follow closely the MC15 pilot request for the monophoton D5 operator [ATL15b]. The implementation details necessary for the treatment of systematic uncertainties for MadGraph are described the  $t$ -channel model.

### 1.5 Prioritization of parameter scans and external constraints

The prioritization of parameter scans according to existing constraints has been widely discussed. It is generally difficult to map collider results to results of other searches for dark matter, such as direct detection cross-section. This mapping is highly model-dependent, and in order to do this with certainty one must have a fully specified theory of dark matter interactions.

The problems that arise when mapping collider results on to the results of other experiments, such as the direct detection scattering cross-section, arise with equal difficulty in the reverse direction, when attempting to apply results from non-collider searches to collider models. One should consider non-collider and relic density constraints only with a complete theory in mind; for this reason the experiments should provide information on the widest sensible range of parameters so that others who are interested in specific models can apply the constraints within their model of choice.

Another type of constraint 'external' to the MET+X analyses comes from the LHC experiments themselves. Resonant signals of dark matter mediators are a key difference between simplified models and the EFT scenario. Direct searches for these mediators, for example in di-jet,  $t\bar{t}$ , and dilepton final states, are thus a crucial part of the LHC search program. These are discussed in the forum report but specific recommendations for them are beyond the scope of the document. It is tempting to apply the constraints from these mediator searches to reduce the parameter space studied in the MET+X searches, but one should remember that the proposed simplified models are intended as representative benchmarks. There is a wide consensus amongst the participants of the forum that these constraints cannot be applied to the parameter scans in any rigorous quantitative way, without a more complete theory of dark matter. However, even if the parame-

ter scans and the searches are not optimized with those constraints in mind, we intend to make all information available to the community to exploit the unique sensitivity of colliders to all possible Dark Matter signatures.

## 2 Recommendation for contact operator interpretation

Except in the cases explicitly mentioned above, EFT contact operators should no longer be a target for optimization: in general there is a poor relationship between the kinematic distributions predicted by the EFTs and the types of signals to which Run 2 analyses will be sensitive. On the other hand, many theorists participating in the forum insist that the information conveyed by EFT limits is valuable. For the V/A and S/P simplified models, the highest mediator mass has been chosen to provide an equivalent of the D5/D8 and D1<sup>2</sup> EFT interpretations, which are the limiting-cases of these models. We recommend to cross-check the kinematics of the highest-mass simplified model point with the corresponding contact interaction benchmark limit, after having ensured the robustness of the contact interaction benchmark point tested (e.g. with the ratio of valid events  $R$  defined in [Col15]).

<sup>2</sup> Following the notation of Ref. [GIR<sup>+</sup>10]

In case limits on other EFT operators are requested by theorists, and they cannot be obtained by suitable combinations of simplified models, the forum proposes two truncation procedures to ensure the EFT prediction is defensible:

- for cut and count analyses (all of the searches done so far), a very conservative truncation for EFT limits from Ref. [RWZ15] is proposed as a modification of the approach ATLAS has used for 8 TeV papers [Col15]
- for shape-based analyses, the references above may be used to obtain the criteria that have to be applied to discard generated events, rather than rescaling the limit as done for cut and count analyses.

**[Point being finalized: when the completion involves loops, the kinematics may vary significantly depending on the scale of the particles in the loops, applying a truncation may not be as straightforward as when we do have a simpler completion.]**

With Run 1 analyses, both experiments are very familiar with simulation of the EFT contact operators. The forum has focused on validity procedures that may be applied at analysis-level after sample generation. Thus, so long as samples of contact operators are actually produced where required, the forum recommendations here have no direct bearing on MC requests.

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