

Collaboration between INSP (Paris), CNEA (Bariloche), IFR (Rosario)

- ETACHA was developed following experiments at the ALICE accelerator in Orsay to predict (and optimize) the production of highly stripped ions (bare and hydrogenlike).
- Our first model considered ions with only up to 10 e⁻ (n=1 & 2 shells).
- In 1996, a new version for PC was elaborated, with several improvements such as an extension to up to 28 electrons (n=1,2,3 shells). This allows to consider ions such as 30 MeV/n ⁵⁶Pb as delivered by GANIL
- End of 2001, the code was further extended, with up to 60 e⁻, but much more remains to do to account for many electron ions ...

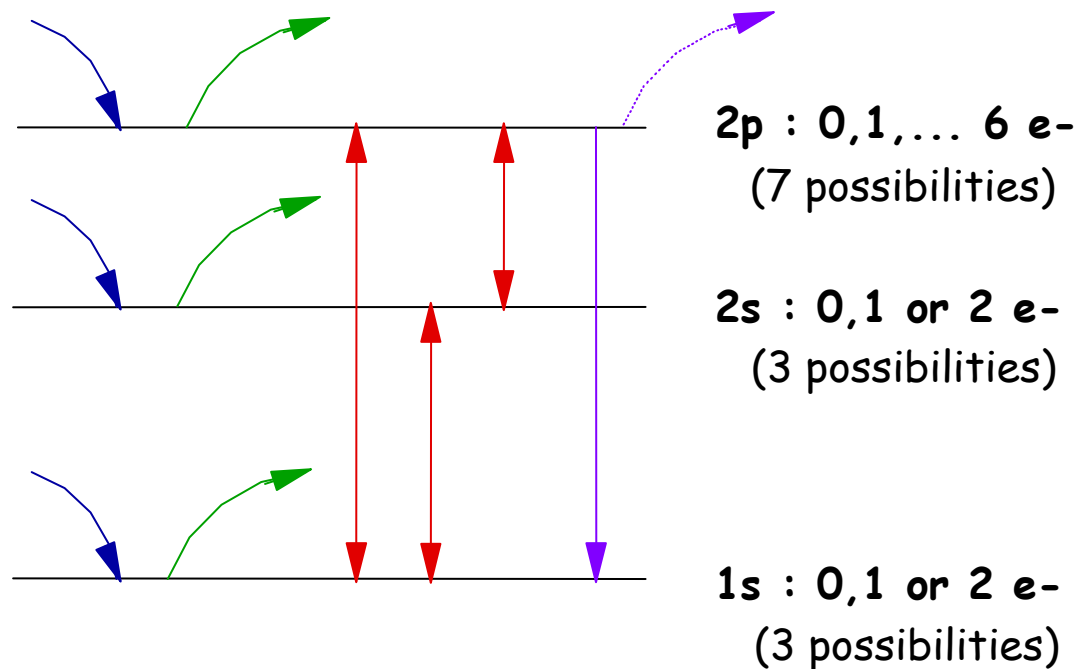
In any case, the starting point is a code that was well suited to highly charged ions

Basic considerations in the model

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Atomic processes and charge states :

- each charge state results from the combination of various configurations with the same electron number
- the population of each configuration results from the competition between various atomic processes :



- capture
- ionisation
- excitation
- radiative (and Auger) processes

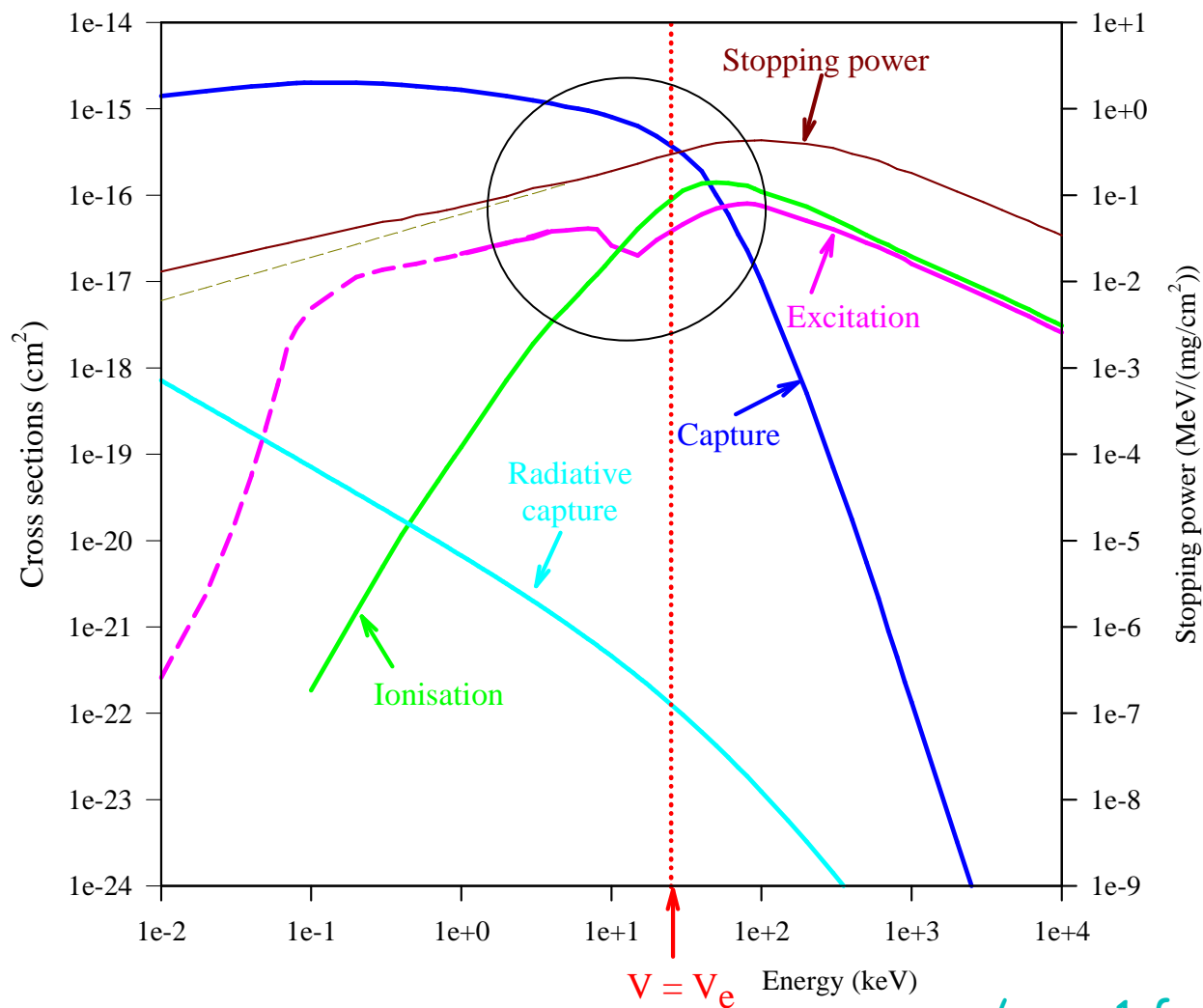
Atomic processes and charge states :

- each charge state results from the combination of various configurations with the same electron number
- 2 shell model :
(10 electrons) $Y(n1s).Y(n2s).Y(n2p) \rightarrow 3+3+7=13$ configurations or
 $Y(n1s,n2s,n2p) \rightarrow 3 \times 3 \times 7 = 63$ correlated configurations
(better results)
- 3 shell model :
(28 electrons) $Y(n1s,n2s,n2p,n3s,n3p,n3d) \rightarrow 3 \times 3 \times 7 \times 3 \times 7 \times 11 = 14553$ (!) correlated configurations or
 $Y(n1s,n2s,n2p,n3) \rightarrow 3 \times 3 \times 7 \times 19 = 1197$ partially correlated +
 $Y(n3s), Y(n3p), Y(n3d)$ calculated separately $\rightarrow 1218$ equations or
 $Y(n1s,n2s,n2p)$ then $Y(n12,n3) \rightarrow 63+21+11 \times 19 = 293$ equations
- 4 shell model :
(60 electrons) $Y(n1s,n2s,n2p,n3,n4) \rightarrow 3 \times 3 \times 7 \times 19 \times 33 = 39501$ (!) configurations or
 $Y(n1s,n2s,n2p)$ then $Y(n12,n3)$ then $Y(n123,n4) \rightarrow$
 $293+33+29 \times 33 = 293+990 = 1283$ coupled differential equations :

$$\frac{dY_i(x)}{dx} = \sum_j Y_j(x) \sigma_{ji} - Y_i(x) \sum_j \sigma_{ij}$$

Cross sections (p → H system)

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$$\bar{n}_{e-} = n_{\max} \frac{\sigma_{\text{gain}}}{\sigma_{\text{loss}} + \sigma_{\text{gain}}}$$

$v/v_{1s} = 1$ for Ar at 8 MeV/u

$v/v_{1s} = 1$ for Kr at 34 MeV/u

$v/v_{1s} = 1$ for Pb at 230 MeV/u

➤ ETACHA has proved to be accurate in the high velocity regime, ie, for ions such as

- C ions, $E \sim 1 \text{ MeV/u}$
- Ar ions, $E \sim 8 \text{ MeV/u}$
- Kr ions, $E \sim 30 \text{ MeV/u}$
- Pb ions, $E \sim 200 \text{ MeV/u}$

$$K = \frac{Z_c}{Z_p} \times \frac{V_e}{V_p} \ll 1$$

➤ Challenging systems need to be considered in the context of S3 :

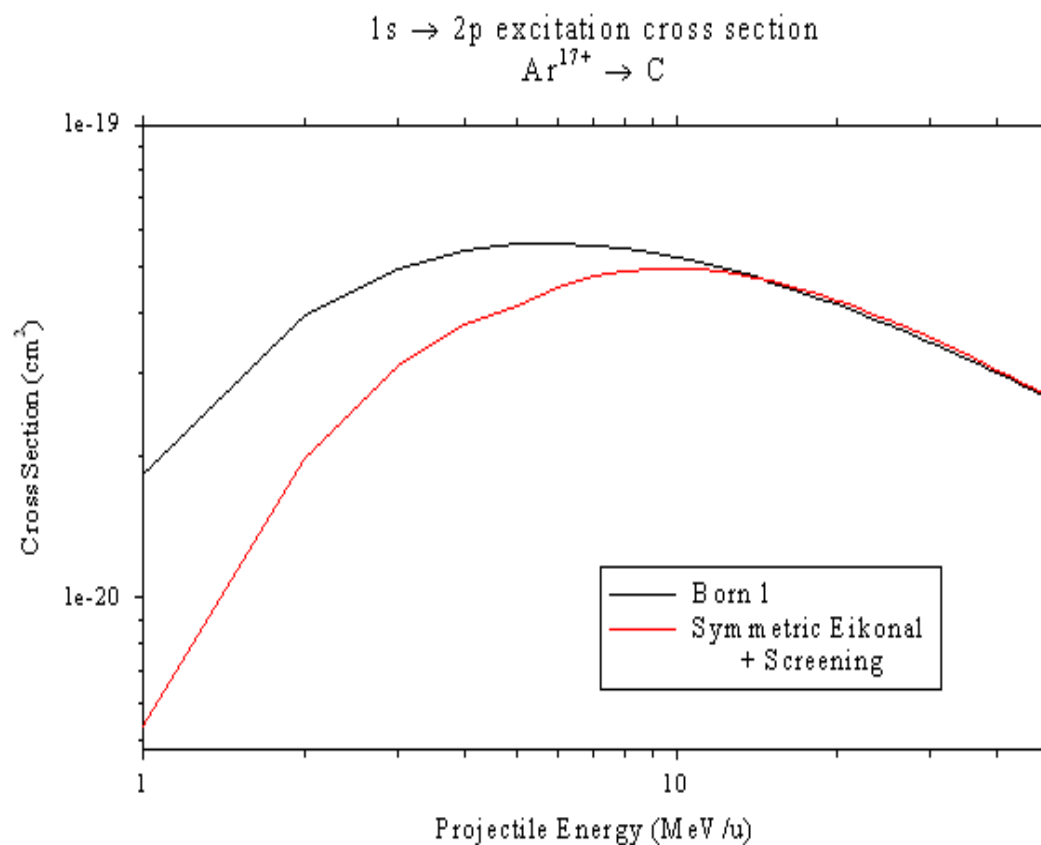
- (Exotic) Sn ions, $E \sim 1 \text{ MeV/u}$
- SHE ions, $E \sim 3 - 0.1 \text{ MeV/u}$

and probably many others ...

Cross sections for strong perturbation

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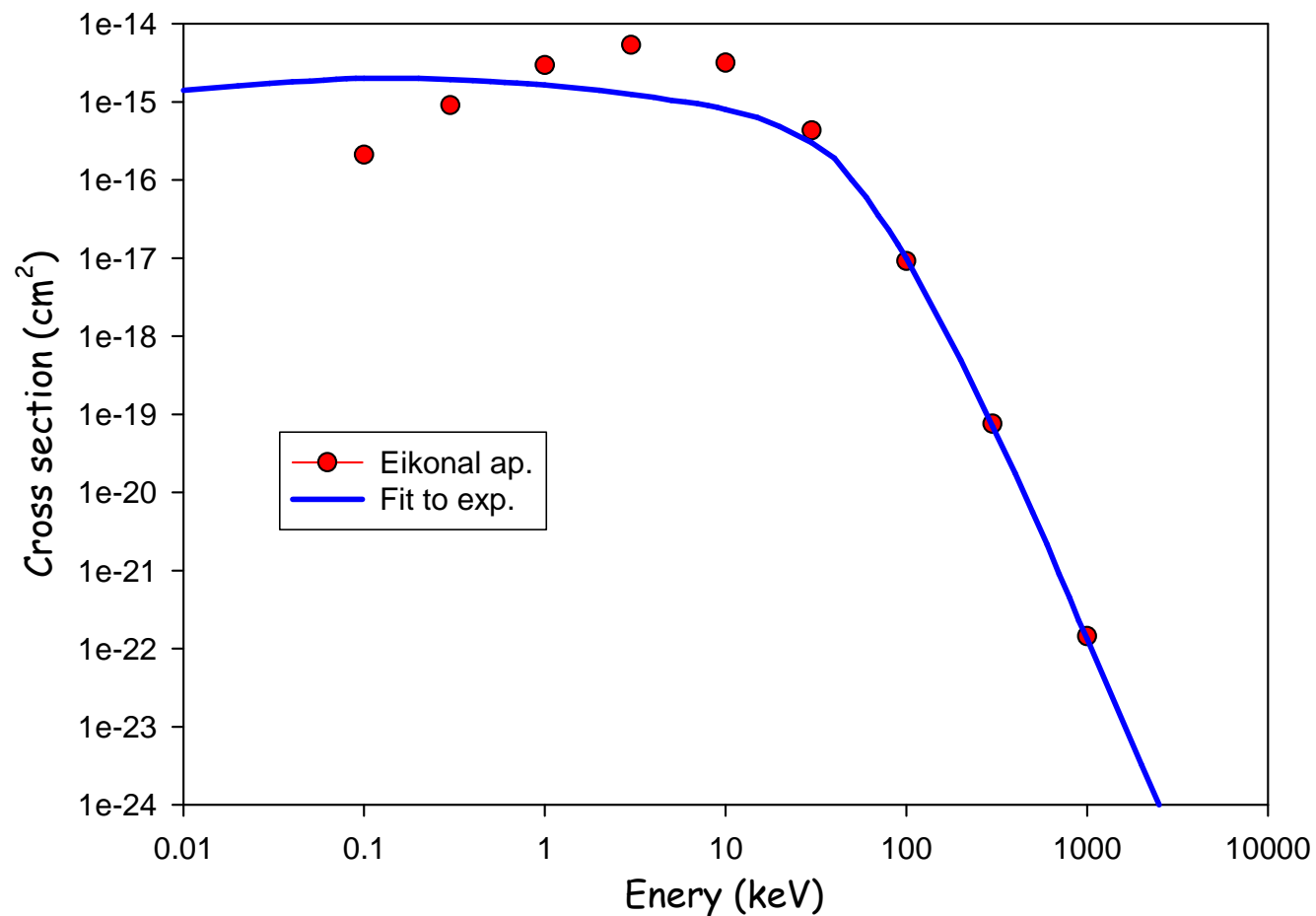
At low velocities, excitation and ionization cross sections strongly deviate from (simple) Born 1 theories



Cross sections : capture at low velocity

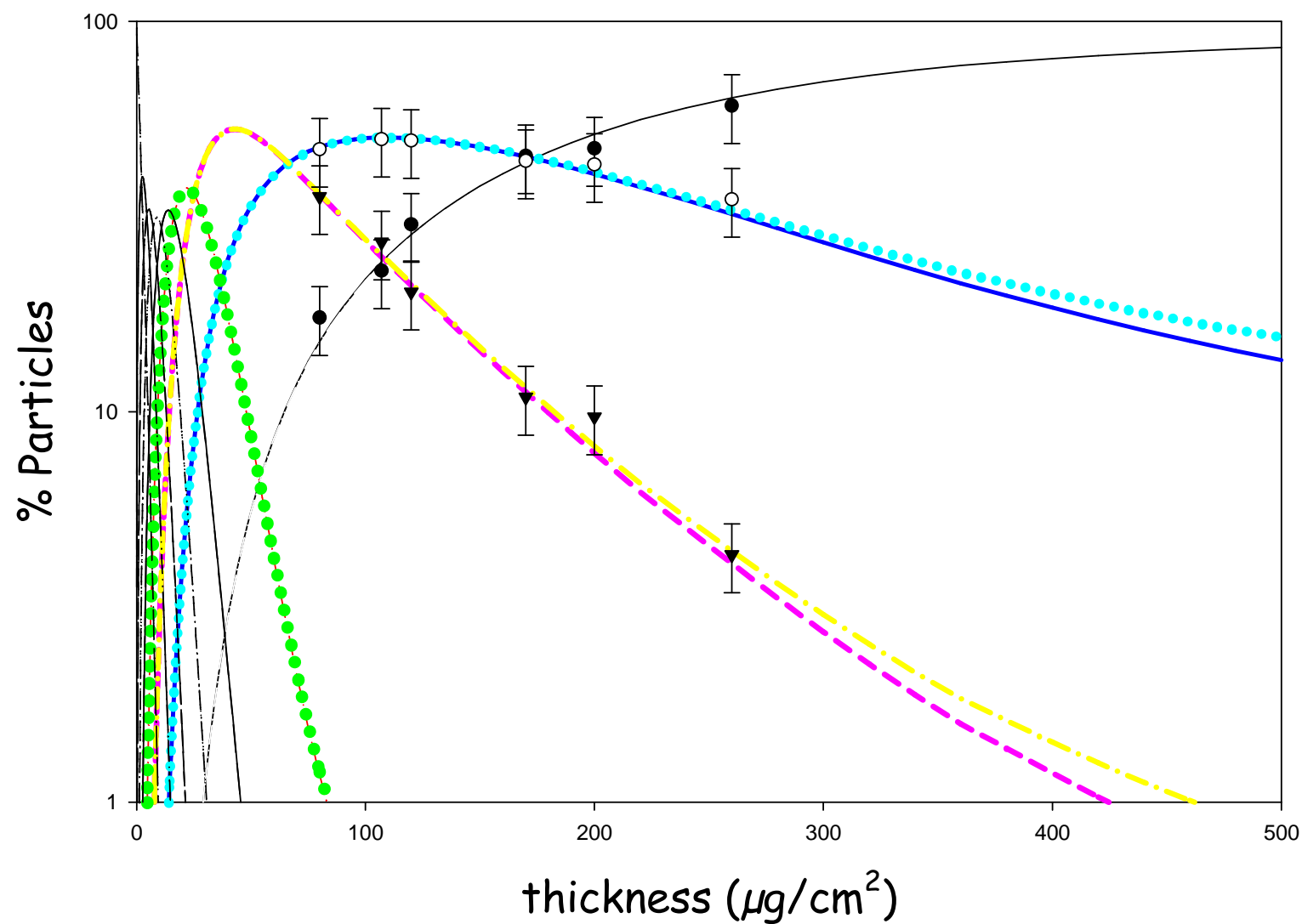
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Capture cross sections (p \rightarrow H)



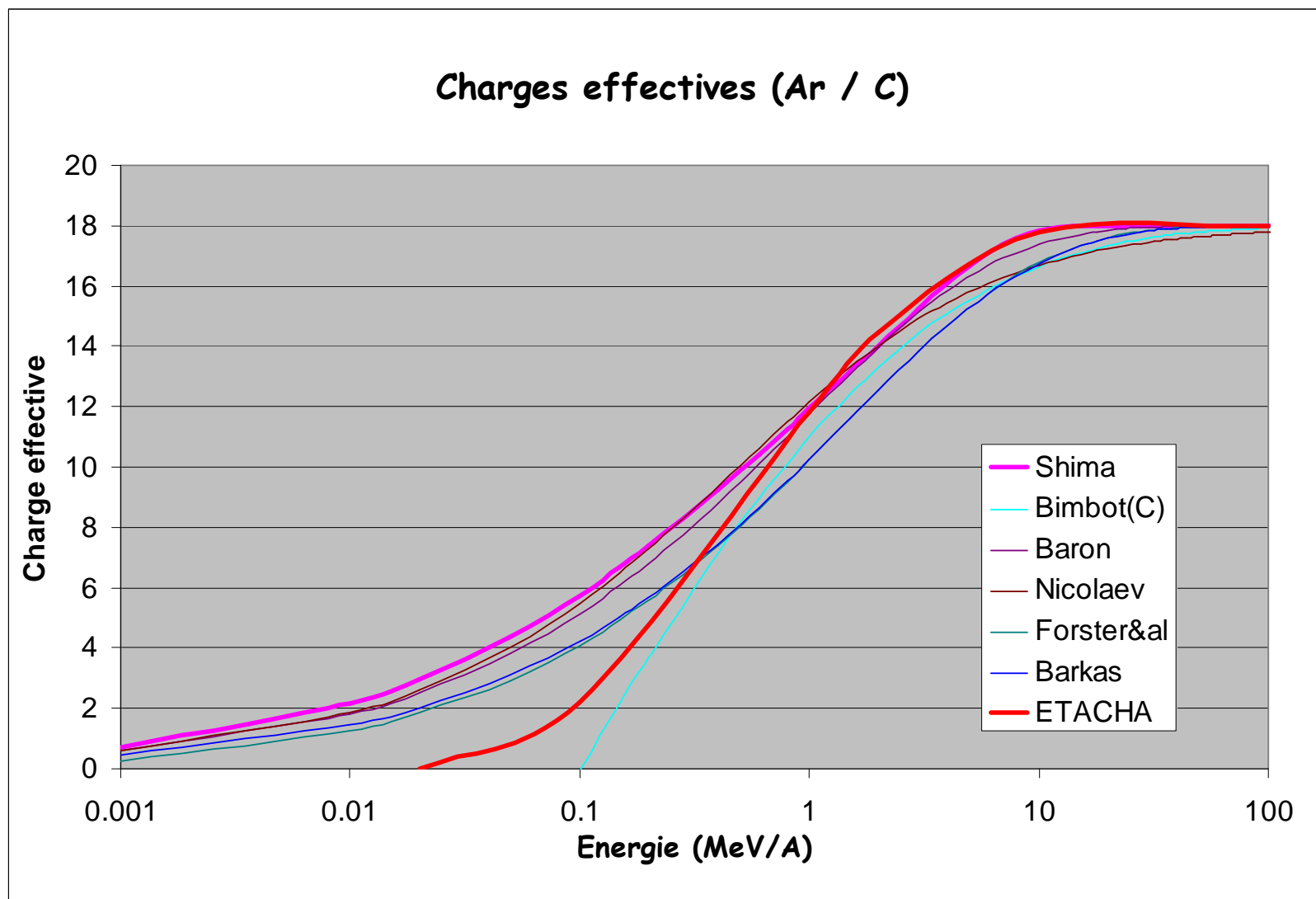
The CDW approximation appears to diverge at low velocity

Ar(+10)-C 13.6 MeV SEIK & CDW



Some results

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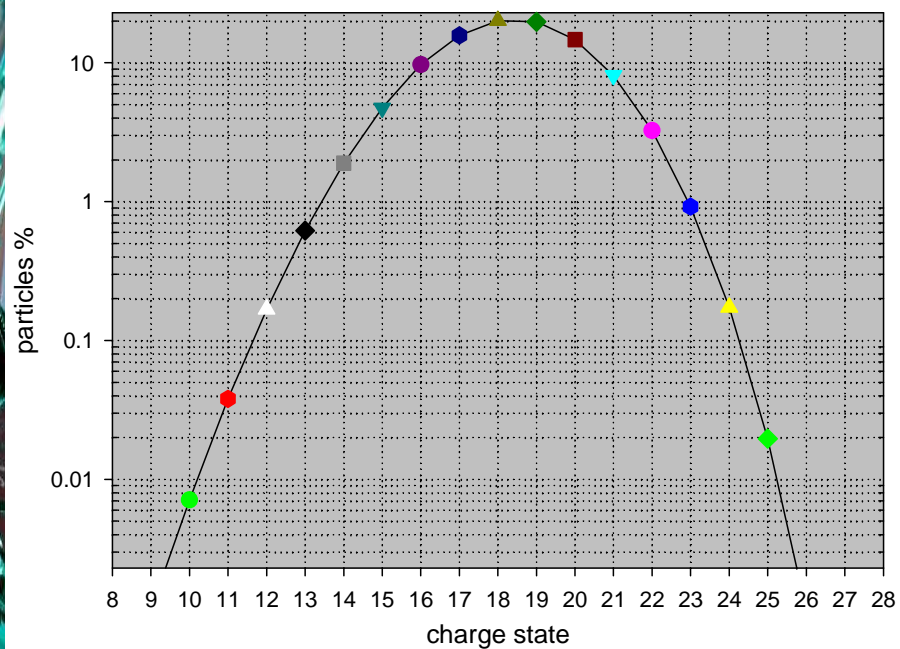


Some preliminary results

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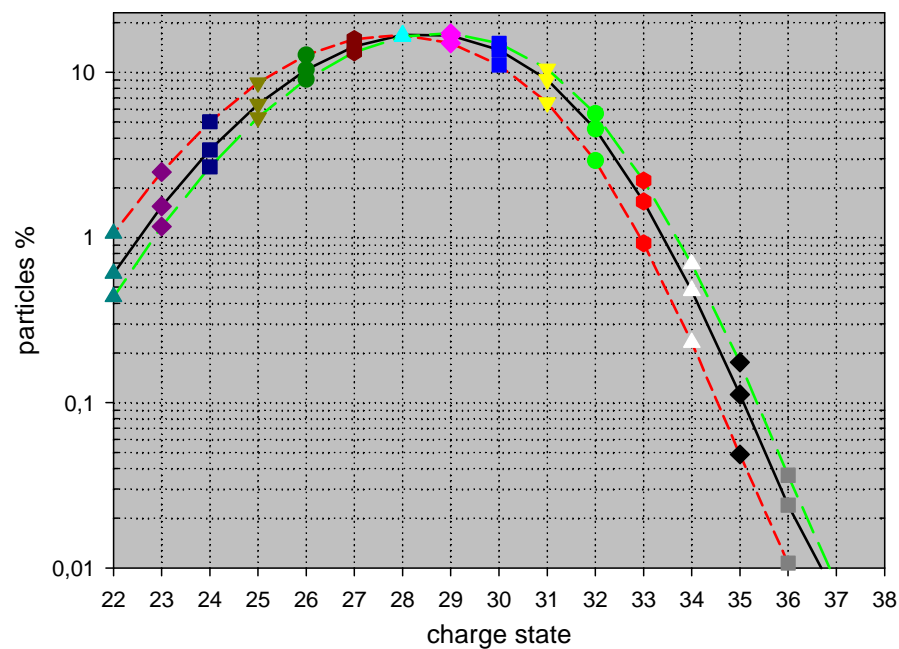


Charge state distribution $^{58}\text{Ni}^{19+}$ on ^{46}Ti @ 3.15 MeV/u
for 500 $\mu\text{g}/\text{cm}^2$ of Ti



$$\langle Q \rangle_{\text{Shima}} = 23$$

Charge state distribution $^{100}\text{Sn}^{50+}$ on ^{46}Ti @ 0.981 MeV/u
for 100, 250 and 500 $\mu\text{g}/\text{cm}^2$ of Ti



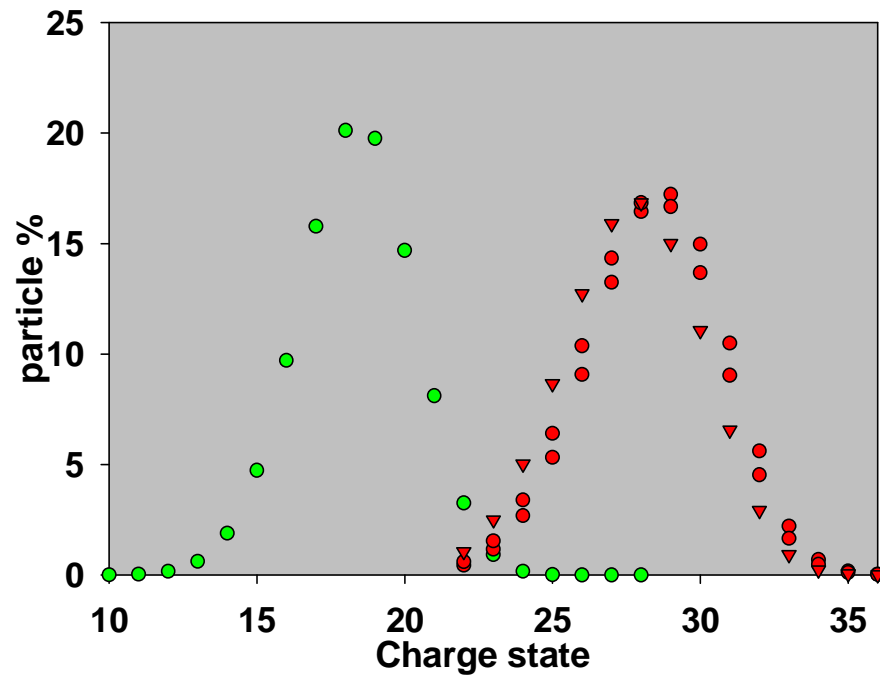
$$\langle Q \rangle_{\text{Shima}} = 32$$

Some preliminary results

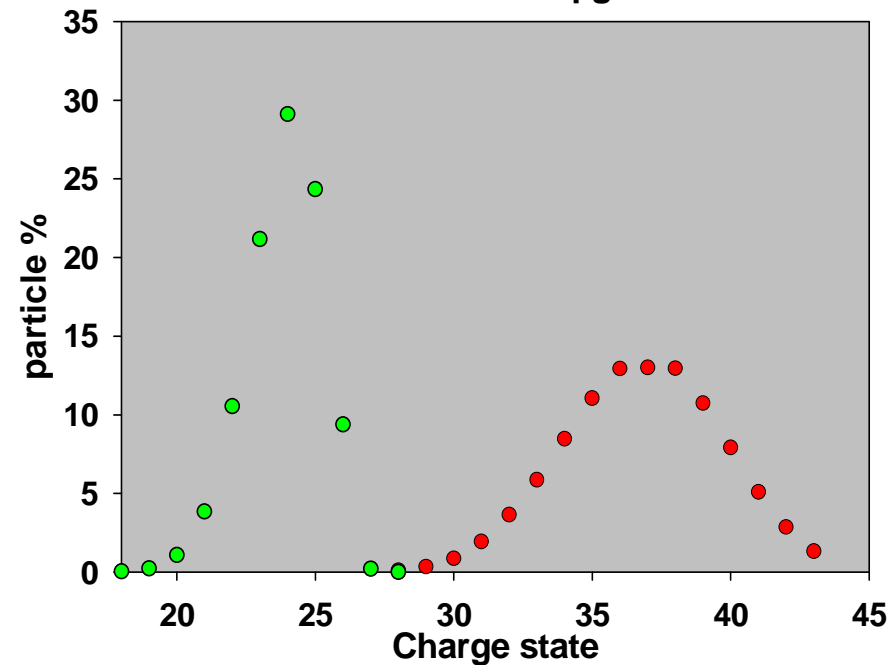
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Ni and Sn at the exit of Ti



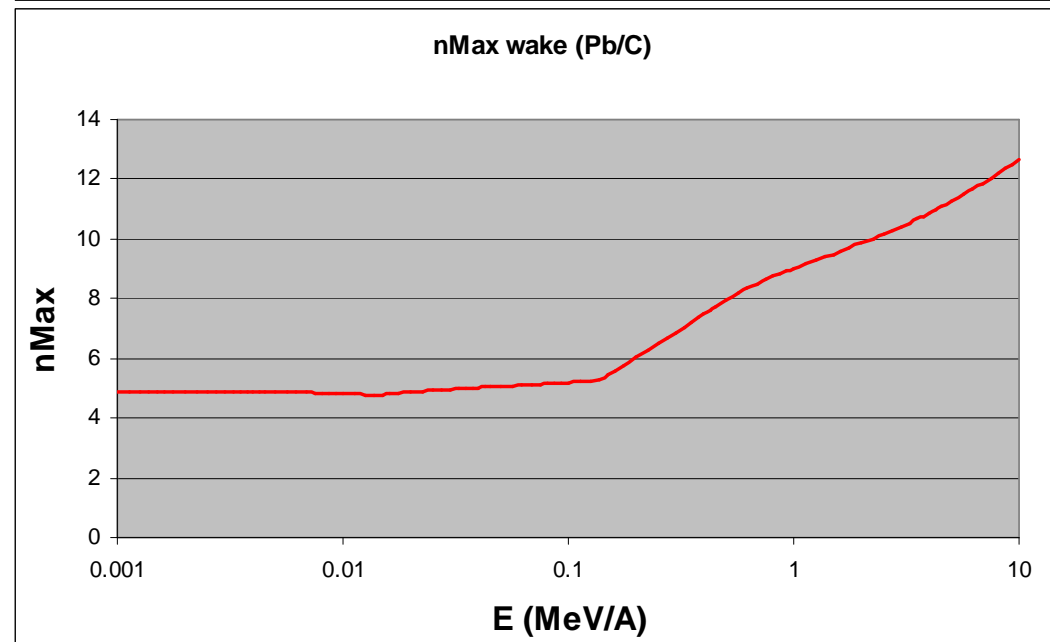
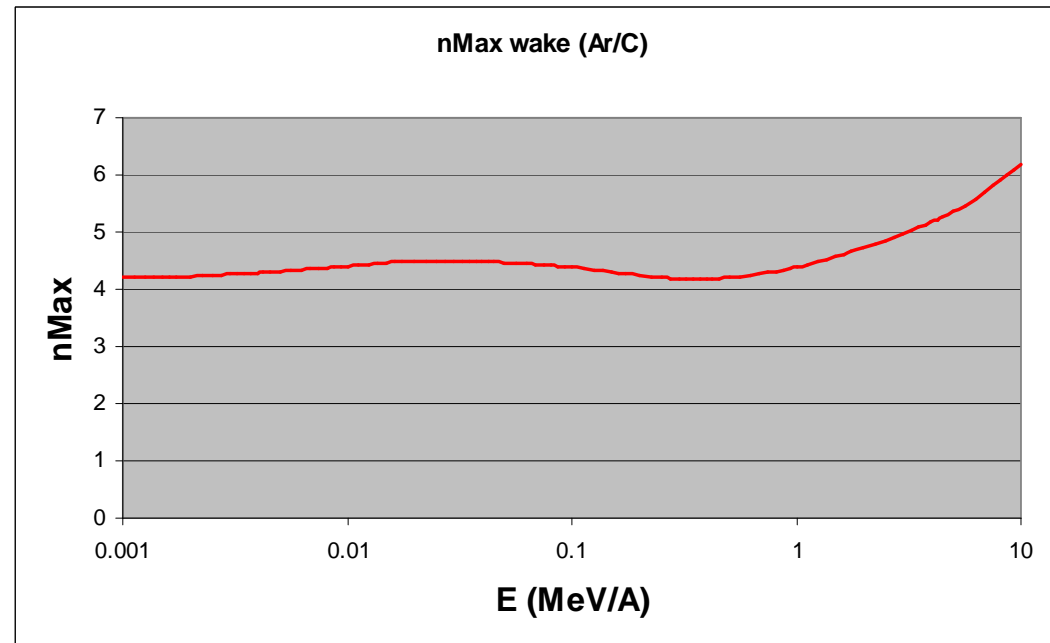
Ni and Sn after 80 $\mu\text{g}/\text{cm}^2$ C



- Cross sections are supposed to be (independent electron model) :
 - proportional to the number of electrons in the initial state
 - Proportional to the number of vacancies in the final state
- This is probably a good enough approximation for the present purpose, provided not too much electrons or holes are involved. Checks need to be done.
- Cross sections, however, also depend on the effective (mean) charge of the projectile. ETACHA re-calculates cross sections each time the mean charge state (or projectile energy) has changed by more than a few %.
Screening formula are used, that are valid as long as there is not too large a departure from hydrogenlike states : for many electron ions, shell effects may need to be considered.
- For ions with more than 60 electrons, new “tricks” need to be found
...

Dynamical screening in solid targets

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- Cross sections for low velocity / strong perturbation are on the way to be under control
- The system of rate equations has to be revisited if calculations are to be done for more than 60 electron - ions
- Screening and shell effects need careful examination.

