

Research Summary

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Outline

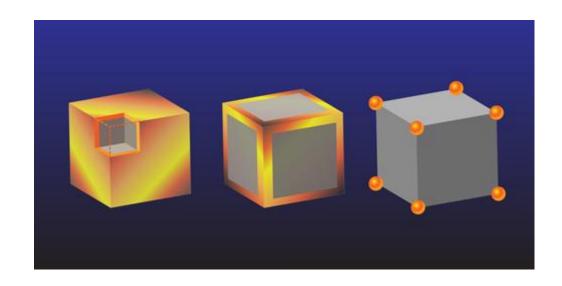
➤ Higher-order Topology in Graphyne Families **Binbin Liu** et al., Phys. Rev. B 106, 035153 (2022).

Xu-Tao Zeng, **Binbin Liu**, et al., arXiv: 2302.13090. (Submitted to PRB.)

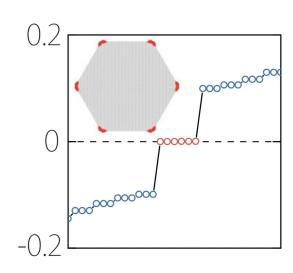
- ➤ Non-centered Inversion Symmetry in Momentum Space **Binbin Liu†**, et al. (To be submitted to PRL.)
- Large Bilinear Magnetoresistance (BMR) from Rashba Spin-Splitting on the Surface of a Topological Insulator Wang Yang*, Binbin Liu*, et al., Phys. Rev. B 106, L241401 (2022).

Higher-order Topology in Graphyne Families

Introduction



3D Higher-order TI



2D Second-order TI

Real Chern number (PT)

$$(-1)^{\nu_R} = \prod_{i=1}^4 (-1)^{\lfloor (n_-^{\Gamma_i}/2) \rfloor}$$

Real Chern insulator / Stiefel - Whitney insulator

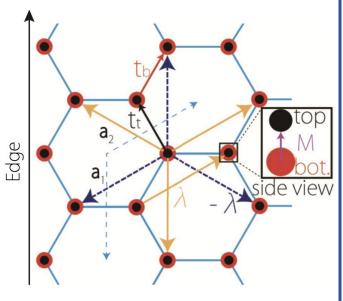
Fractional Charge (Crystalline)

$$Q_{corner}^{(2)} = \frac{e}{4} \left[-n_{C_2}^X - n_{C_2}^Y + \left(n_{C_2}^M + n_{C_2}^\Gamma \right) \right] \text{ mod } e$$

Quadrupole insulator / TCI

$$\mathcal{C}_{6z} \;\; Q_{corner}^{(2)} = e rac{
u_R}{2}.$$

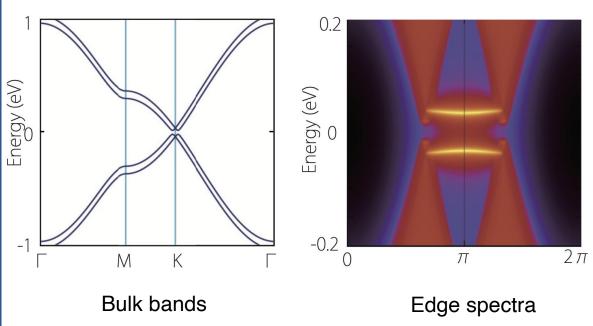
SOTI models



2D model

$$\mathcal{C}_{2z}\mathcal{T}$$
 \mathcal{C}_{6z} and \mathcal{C}_{2y} $u_R = 1$.

$$H^{4BTB}(\mathbf{k}) = \begin{bmatrix} G_t & M\sigma_0 + iS \\ M\sigma_0 - iS & G_b \end{bmatrix}$$

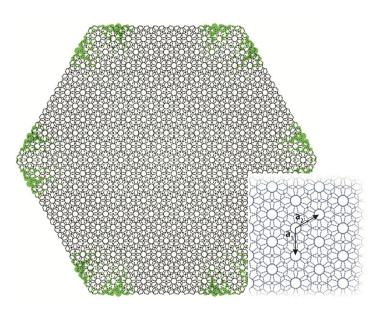


z
 3D model $\mathcal{H}_{
m 3D}=\mathcal{H}_{
m 2D}+2t_z\cos k_z\sigma_z au_0$

Can be readily realized in **Graphyne families!**

SOTI in Graphyne Families

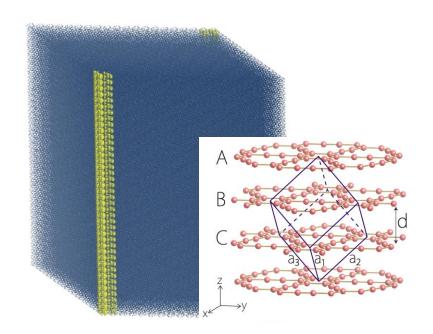
Moiré α-Graphyne (tbGPY)



General way to SOTI phase in moiré materials

B. Liu et al., Phys. Rev. B 106, 035153 (2022).

Bulk γ-Graphyne Nat. Synth. 1, 449 (2022).



The first 3D RCI material

X.-T. Zeng, **B. Liu,** et al., arXiv: 2302.13090.

2D SOTI & RCI: Graphdiyn(PRL 2019), tbGPY (PRB 2022)

SO real nodal-line semimtal: 3D graphdiyn (PRL 2022)

3D SOTI & RCI: bulk γ-graphyne (arXiv 2023)

Non-centered Inversion Symmetry in Momentum Space

Projective Symmetry

Projective symmetry argebra.

$$(PL_i)^2 = \theta$$

$$\theta = \pm 1$$

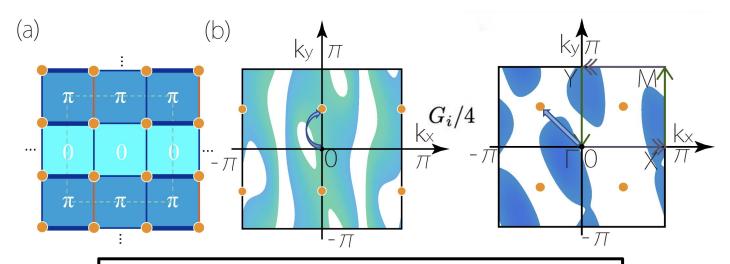
Cohomology invariant

$$PL_iP = -L_i^{-1} = e^{i(k_i + G_i/2)a_i}$$

$$\mathcal{P}_i = P\mathcal{L}_{G_i/2}$$

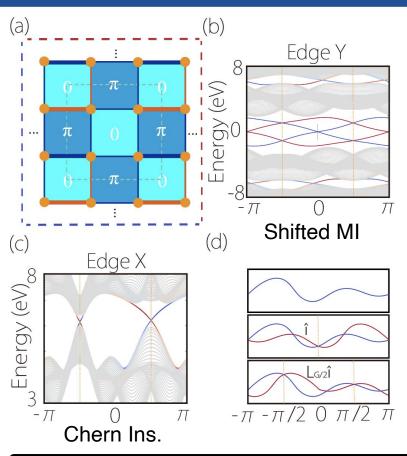
Half reciprocal lattice translation

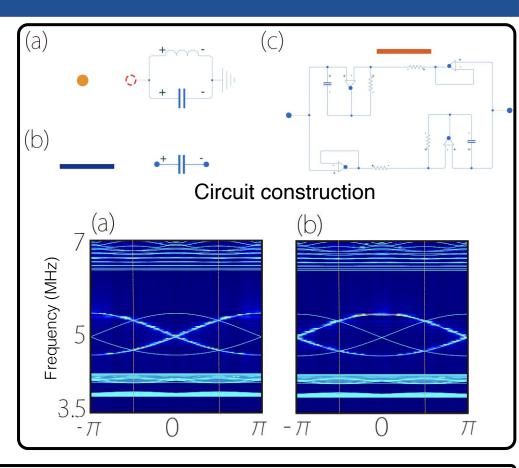
$$UH_i(-\mathbf{k}_{rest}, -k_i + G_i/2)U^{\dagger} = H_i(\mathbf{k}_{rest}, k_i)$$

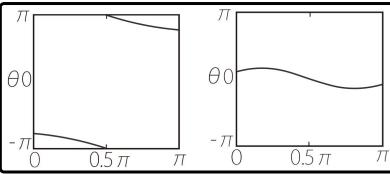


Inversion centers: *Quater* reciprocal lattice translation

Topological States & Simulations







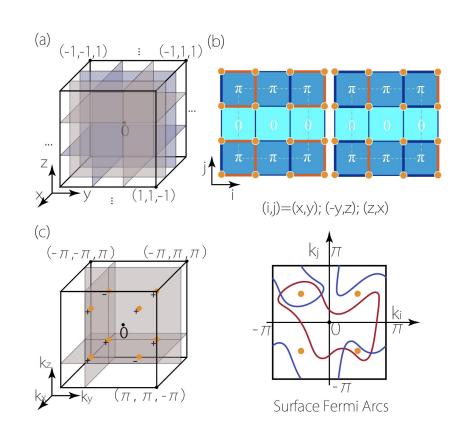
$$u=rac{1}{2\pi}[\gamma(0)+\gamma(\pi)]\mod 2$$

$$=\zeta\mod 2$$

$$(-1)^{
u}=\prod_{j=a,b}\prod_{i}^{N}\lambda_{i}(\Gamma_{j})$$
 Topological invariant

Extending to 3D

POSTER



$$\mathcal{P}_{xyz}^{3D} = U_{3D} \mathcal{L}_{G_x/2} \mathcal{L}_{G_y/2} \mathcal{L}_{G_z/2} \hat{I}_{xyz}$$

To be submitted to PRL.

Non-centered inversion symmetry in momentum space Bin-Bin Liu†, Xian-Lei Sheng†, Y. X. Zhao†, and Shengyuan A. Yang School of Physics, Beihang University

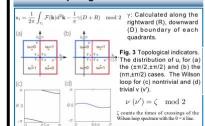
Introduction

Symmetry is of paramount import in physics. Spacial inversion symmetry is one of the most important symmetries which defines the parity. It inverses the real space position as well as the momentum of a physical system relative to the original point in the real and momentum spaces. In this letter, we discover a new kind of inversion symmetry which inverses the real space position relative to the origin but is non-centered in the momentum space We derive this non-centered momentum with projectively enriched symmetry algebras. We demonstrate the idea in Hamiltonians on flux lattices with Z2 gauge field, which naturally realize the noncentered momentum inversion symmetries. Intriguingly, we discover unconventional topological indicators emergent from the gauge fields, which are protected by the momentum-non-centered symmetries. The topological numbers can be read off from Wilson loops evolving along both reciprocal lattice vectors. We find that non-trivial topological indicators defined along both vectors signal a novel topological insulator, with both of its x and y edges (four edges) hosting edge states linked by the non-centered inversion symmetry in each pairs, drastically distinct from previous topological insulators with nomal inversion symmetry. We provide circuits simulations which agree with the theoretical prediction. Furthermore, we construct a 3D topological insulator with surface states satisfying the momentum-non-centered inversion symmetry Our theory opens new routes for exploring physical and topological consequences associated with the momentum-non-centered inversion symmetries which can be realized and tested in circuits and other artificial systems

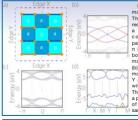
Projective Symmetry

Normal Inversion symmetry (NIS): $P = U\hat{I}$ $UH(-k_x, -k_y)U^{\dagger} = H(k_x, k_y).$ Momentum-non-centered inversion symmetry (MNCIS) from projective symmetry $\mathcal{P}\mathcal{L}_{v}\mathcal{P}^{\dagger} = -\mathcal{L}_{v}^{-1} = e^{(ik_{y}+iG_{y}/2)b}$ $P = U \mathcal{L}_{G_y/2} \hat{I}$ or $P = U \mathcal{L}_{G_x/2} \mathcal{L}_{G_y/2} \hat{I}$ $UH(-k_x + \pi, -k_y + \pi)U^{\dagger} = H(k_x, k_y)$ -п -п/2 О п/2 п

Fig. 1 Bulk states with MNCIS. Fig. 2 Edge states with NIS and MNCIS Topological invariants



Topological insulators with MNCIS

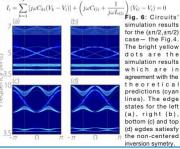


The lattice model. The red bond is attached with a π flux, with corresponding flux pattern indicated by 0 or m in each plaque. The magnitude of the bond. (b) (c) The edge states in momentum space along Y and X edges, agree with Fig. 2 (bottom). (d The band structure along of the BZ, which is the

Circuit simulations



Kirchhoff's current law:



for the $(\pm \pi/2, \pm \pi/2)$ case- the Fig.4. The bright yellow dots are the simulation results which are in agreement with the theoretica predictions (cvan lines) The edge states for the left (a), right (b) bottom (c) and top (d) egdes satiesfy the non-centered inversion symetry.

Conclusions

Discovered a novel non-centered inversion symmetry nomentum space from projective symmetry algebras.

Discovered novel topological insulator states with the MNCIS.

Formulated a topological indicator for the states and derived several equivalent ways to calculate it: including the parity counting from the topological quantum chemistry, the integration form from K-theory, and the Wilson loop method.

Proposed circuits setup whoes circuits' Laplacian from Kirchhoff's current law to the Hamiltonian. Simulated the circuits

Detected the oddity of Chern numbe using the topological

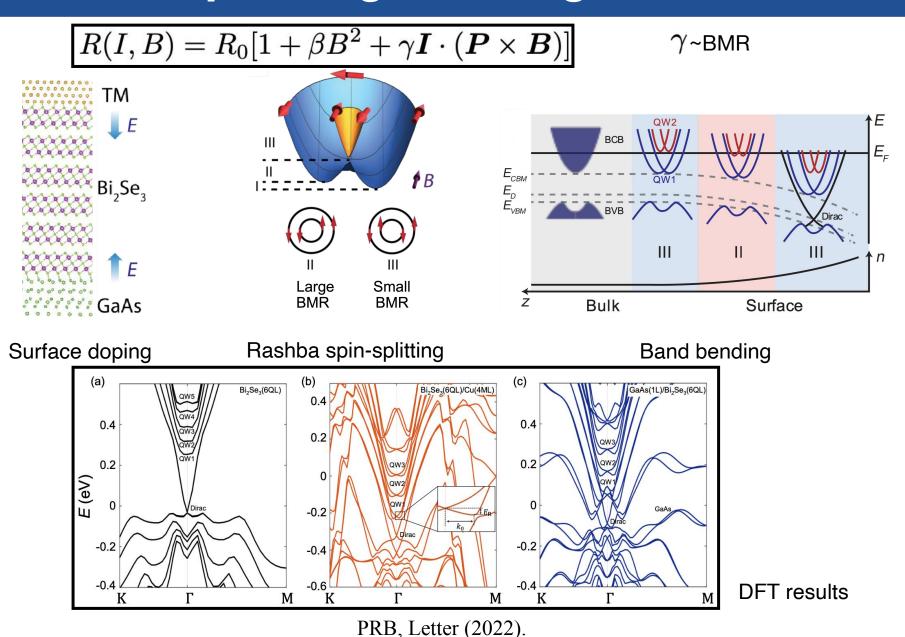
Constructed a 3d topological indicator protected by the non centered inversion symmetry.

References

[1] Z. Y. Chen, S. A. Yang, and Y. X. Zhao, Nature Communications 13, 2215 (2022). [2] Y. X. Zhao, Y.-X. Huang, and S. A. Yang, Phys. Rev. B 102, 161117 (2020).

Large BMR in Rashba States on the Surface of a TI

Explaining the Large BMR



Outlooks

nonHermitian **Tensor networks** Complex systems Application Circuit Hyperbolic fluid metamaterial resi Correlated electrons Topological everything material

Outlooks

Scientific writing Advising Idea contributing Fast learning



Thank you!

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