

《物理与人工智能》

AI与量子多体物理

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What is emergence?

- Emergence occurs when a complex entity has properties or behaviors that its parts do not have on their own, and emerge only when they interact in a wider whole.

4 August 1972, Volume 177, Number 4047

SCIENCE

More Is Different

P. W. Anderson

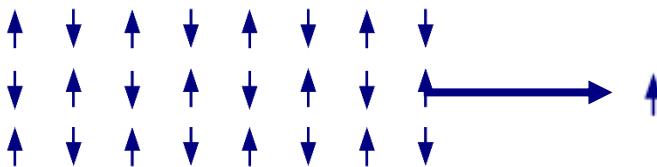
Universal macroscopic
behaviors

✓ Image
recognition:
flower



Numerous microscopic
degrees of freedom

✓ Many-body
physics:
antiferromagnet



a single
pixel

✓ And more:
wet

a single
spin

macroeconomy

water molecule

individual



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What are phases?



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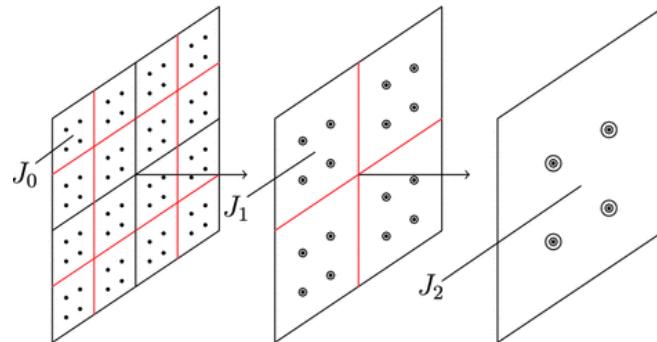
- Renormalization group – a pathway towards understanding collective behaviors



Kenneth G. Wilson
1982 Nobel Prize
in Physics

$$H = -J \sum_{\langle ij \rangle} \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j$$

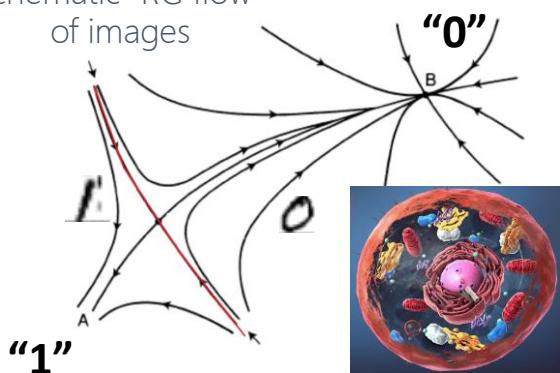
$J_\infty \rightarrow 0$ J_c $J_\infty \rightarrow \infty$
paramagnetic ferromagnetic
phase phase



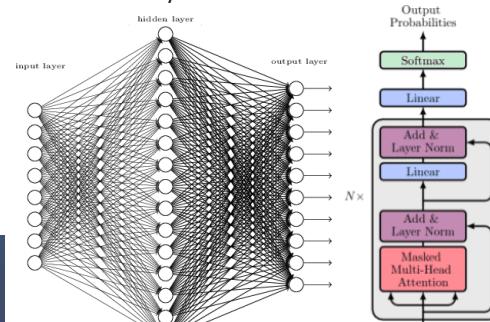
universal properties irrespective of microscopic details

- Image recognition, AI, even life, are collective behaviors

schematic "RG flow"
of images



No two cells are
exactly the same!



John Hopfield & Geoffrey Hinton
2024 Nobel Prize in Physics



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AI for Quantum Phase Recognition

Review: discriminative AI



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AI for Quantum Phases Recognition



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not only but also

↓

1:	1	1	1	1	...
2:	2	2	2	2	...
3:	3	3	3	3	...

The Fock space is exponentially large $2^{30} \sim 100000000$, $2^{1000} \sim 100 \dots 00$ with >300 zeros!



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AI for Quantum Phases Recognition

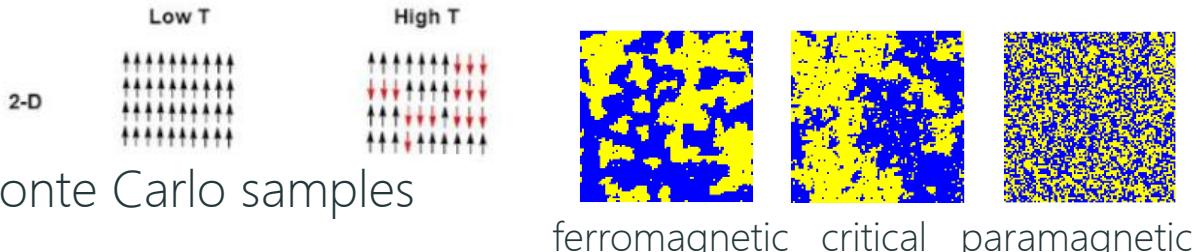


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- The Ising model:

$$H(\sigma) = -J \sum_{\langle i j \rangle} \sigma_i \sigma_j$$

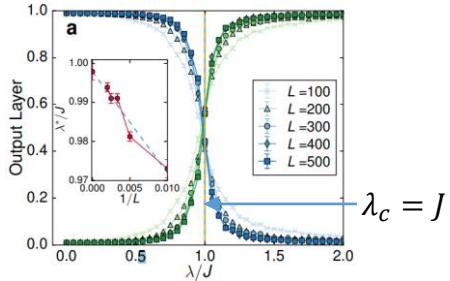
two phases: ferromagnetic (ordered) vs paramagnetic (disordered)



- The Aubry-Andre model (quasicrystal):

$$H = -J \sum_i (c_i^\dagger c_{i+1} + \text{h.c.}) + 2\lambda \sum_i \cos(2\pi\phi_i) c_i^\dagger c_i$$

Numerical data: LDOS
delocalized vs localized



Juan Carrasquilla, Roger G. Melko, 2017.



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AI for Quantum Phases Recognition



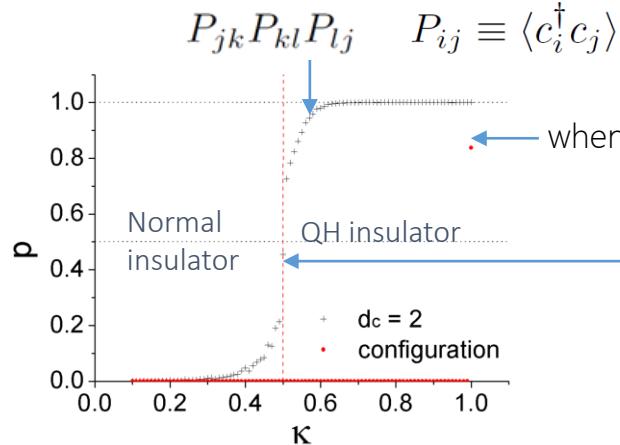
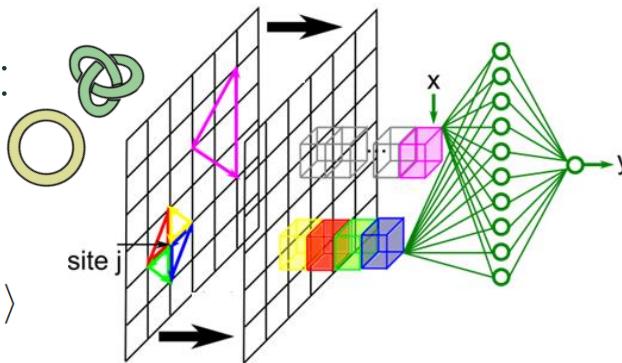
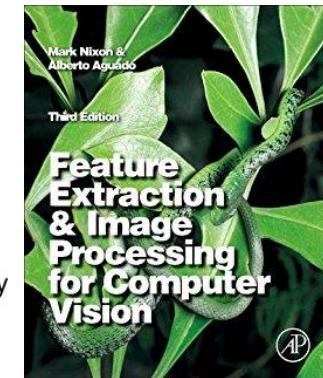
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- Quantum and topological phases: the compatibility issue
a feature selection layer to bridge between

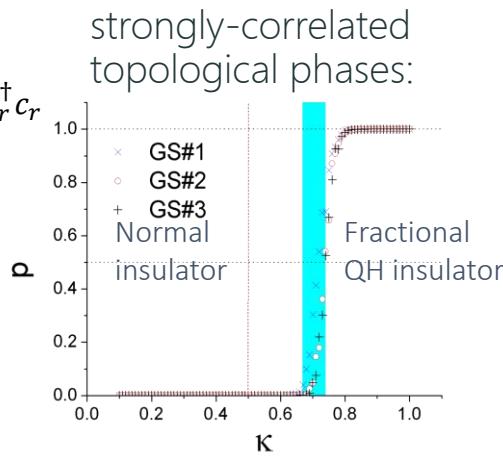
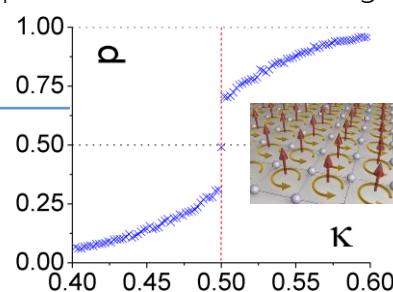


- E.g., topological phases:

relevant operators in
the Kubo formula for σ_{xy}



when operators are irrelevant, e.g., $n_r = c_r^\dagger c_r$



Yi Zhang, E.-A. Kim, 2017.



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AI for Quantum Phases Recognition



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- From quantum many-body models, for quantum many-body models:

- Observation 1: various model parameters

$$H = \sum_{rr'} t_{rr'} c_r^\dagger c_{r'} + \sum_r \mu_r c_r^\dagger c_r + \sum_r U_r c_r^\dagger c_r n_r^f + \dots$$

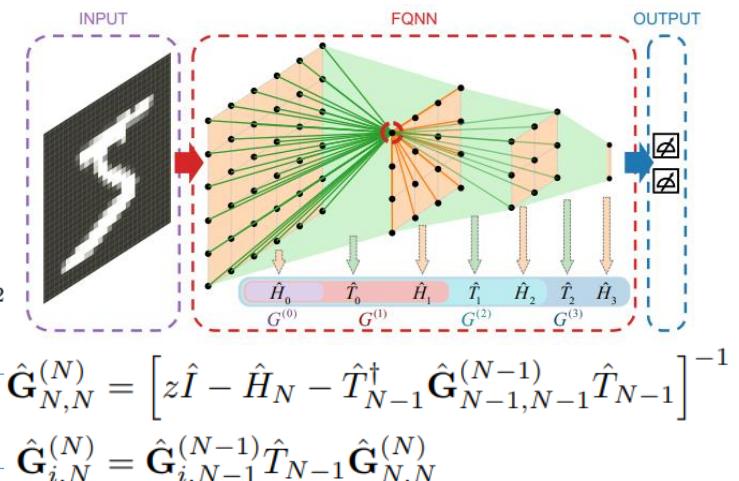
- Observation 2: nonlinear physical properties

LDOS: $y_k = -\frac{1}{\pi} \Im \left[\hat{\mathbf{G}}_{L,L}^{(L)} \right]_{k,k}$ Conductance $y = \sum_m \left| \left[\hat{\mathbf{G}}_{0,L}^{(L)} \right]_{m,1} \right|^2$

$$\hat{\mathbf{G}}_{N,N}^{(N)} = \left[z\hat{I} - \hat{H}_N - \hat{T}_{N-1}^\dagger \hat{\mathbf{G}}_{N-1,N-1}^{(N-1)} \hat{T}_{N-1} \right]^{-1}$$

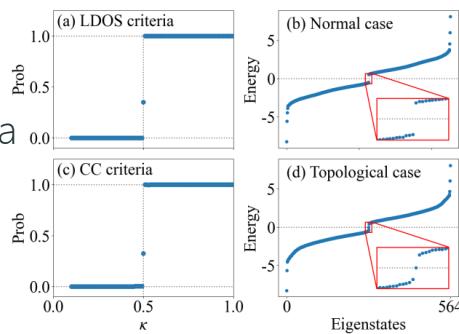
$$\hat{\mathbf{G}}_{i,N}^{(N)} = \hat{\mathbf{G}}_{i,N-1}^{(N-1)} \hat{T}_{N-1} \hat{\mathbf{G}}_{N,N}^{(N)}$$

- Observation 3: efficient recursive methods

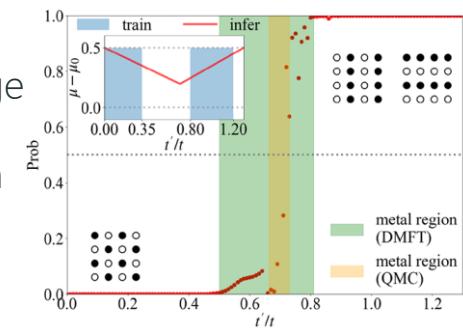


- Direct quantum and topological phase recognitions without presumptions:

map and recognize topological phases via collective properties of FQNN models:



recognize emergent charge density wave phases – (un)controlled estimates in FQNNs (original) models:



Pei-Lin Zheng, et al. 2023, 2024.



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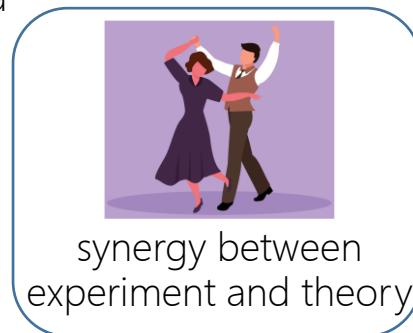
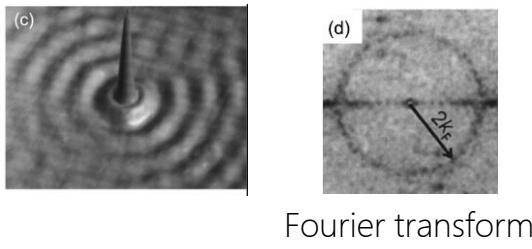
AI for Quantum Phases Recognition



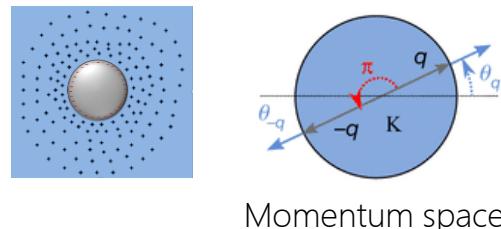
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- Interpretation of big, complex experimental data

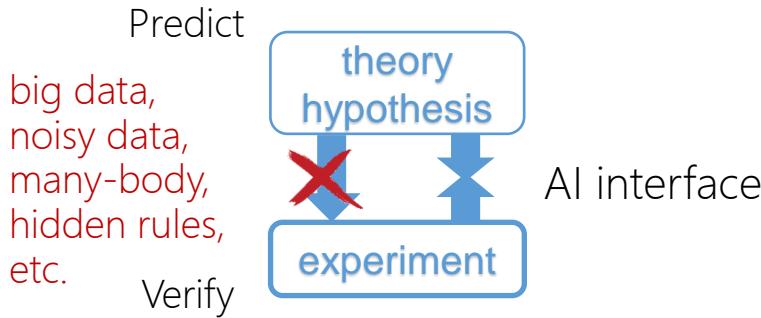
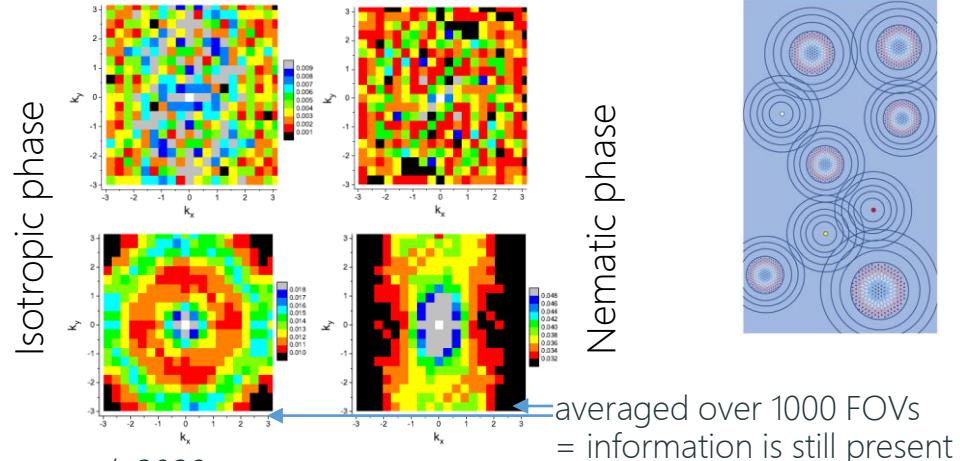
Quasi-particle interference pattern as a Fermi-surface probe of electron liquids



Friedel oscillations as a screening and back-scattering process around a local impurity



- However, our theoretical capacity largely lags behind our real-world complexity:



Yi Zhang, et al. 2020.



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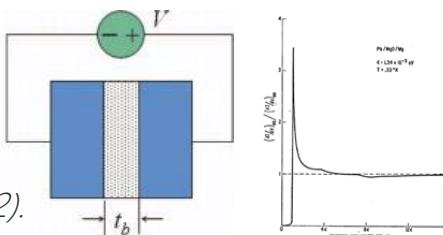
AI for Quantum Phases Recognition



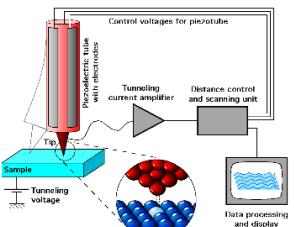
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- How experiments have evolved: from

Giaever et al.,
Phys. Rev. **126**, 941 (1962).

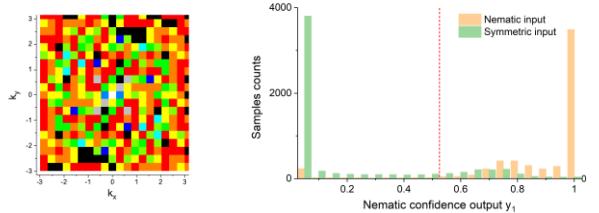


to

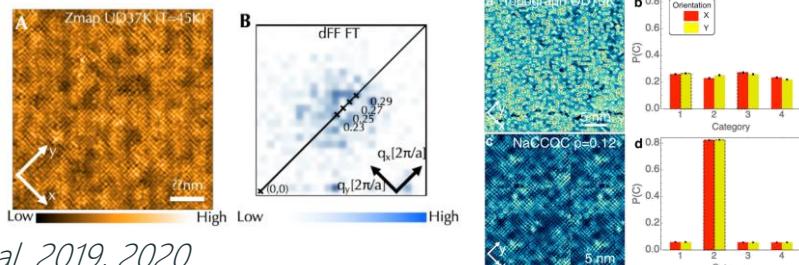


- Idea: Train with the big noisy data, trained for the big noisy data

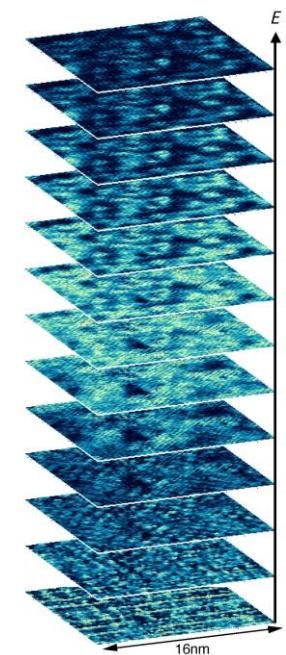
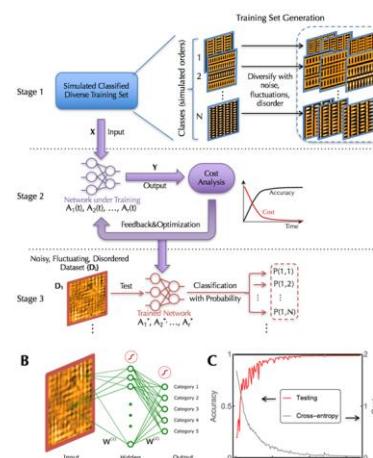
- Recognition of nematic phases from STM data



- Recognition of CDW phases from STM data



Yi Zhang, et al. 2019, 2020.



generalizable to other big noisy experimental data, e.g. neutron scattering

Anjana Samarakoon, et al. 2020.



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AI for Quantum Monte Carlo Methods

Review: generative AI



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MC and Quantum MC Methods

- Statistical mechanics: $W(\vec{x}) = \frac{1}{Z} \exp\left[-\frac{E(\vec{x})}{k_B T}\right]$, $Z = \sum_{\vec{x}} \exp\left[-\frac{E(\vec{x})}{k_B T}\right]$

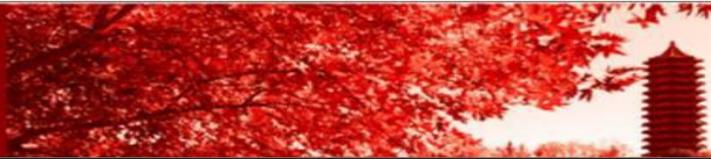
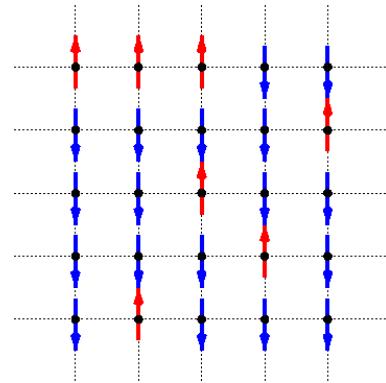
- Example: Ising model $E(\vec{x}) = -\sum_{\langle ij \rangle} J x_i x_j$, $x_i = \pm 1$

- The Metropolis Algorithm: (also used in simulated annealing)

1. Generate a random initial state $\vec{x}_{t=0}$ with energy $E(\vec{x}_{t=0})$;
2. Flip a random spin $x_i \rightarrow -x_i$ and calculate the energy $E(\vec{x}_?)$ of this trial state $\vec{x}_?$;
3. Calculate the difference in energy generated by the spin flip, $\Delta E = E(\vec{x}_?) - E(\vec{x}_t)$;
 - If $\Delta E \leq 0$ (the trial spin state is energetically favorable), accept the spin flip;
 - If $\Delta E > 0$, accept the spin flip with probability $p = \exp(-\Delta E / k_B T)$;
4. Measure the target physical quantities, e.g., energy, magnetization, etc.
5. Repeat steps (2) to (4) until sufficient number N of uncorrelated samples are obtained.

The target probabilities are guaranteed by detailed balance:

$$\frac{W(A)}{W(B)} = \frac{P(B \rightarrow A)}{P(A \rightarrow B)} = \exp\left(-\frac{E_A - E_B}{k_B T}\right)$$

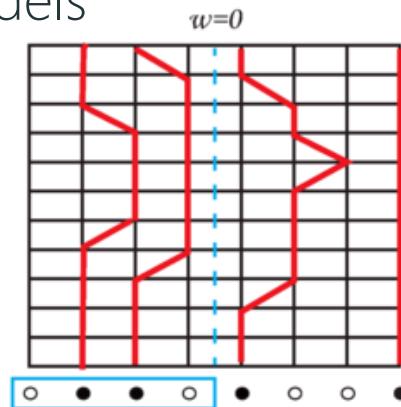


MC and Quantum MC Methods

- Also applicable to certain quantum many-body models

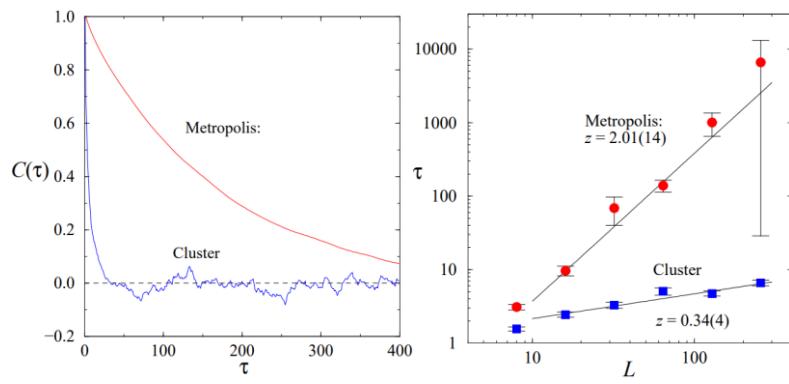
- Auxiliary-Field Quantum Monte Carlo
- Path Integral Monte Carlo
- Determinantal Monte Carlo
- Stochastic Series Expansion Quantum Monte Carlo, etc.

commonly sampling configurations in $(d+1)$ -dims space-(imaginary)-time



- However, local minima cause critical slowing down – cluster update:

- Choose a random site x_i .
- Add neighbor site $x_j = x_i$ into the cluster with probability $p = 1 - e^{-2\beta J}$.
- Grow the cluster until all neighbors are considered. Flip cluster.



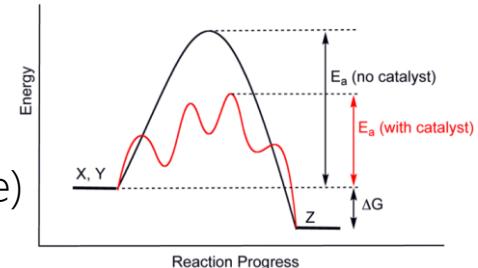
AI for Quantum MC Methods



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- Pros and cons of cluster update:

- probability $W(\vec{x})$ ensured via detailed balance
- global updates with high efficiency (100% acceptance rate)
- yet, heavily reliant on the model



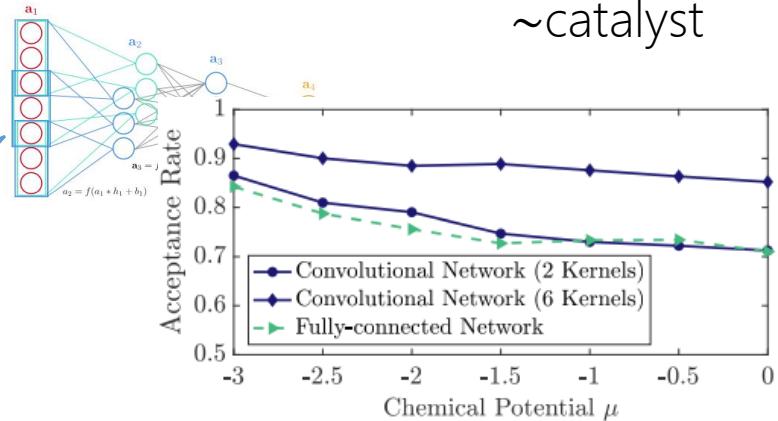
We cannot flip a random cluster with detailed-balance probability, which equals exponentially small acceptance rate! → globally distinctive states with similar weights
~catalyst

- Idea: fitting $W(\vec{x})$ with an AI model:

then accept cluster with acceptance rate:

$$\alpha(A \rightarrow B) = \min\{1, e^{-\beta[(E_B - E_B^{\text{eff}}) - (E_A - E_A^{\text{eff}})]}\}$$

after which the MC is exact.



Huitao Shen, Junwei Liu, and Liang Fu, 2018.



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Generative models

- Graph models: probability distribution with statistical mechanics

$$\mu(v) = \frac{1}{Z} \exp \left\{ \sum_i \theta_i v_i + \sum_{(i,j) \in E} \theta_{ij} v_i v_j \right\}$$

- Restricted Boltzmann Machine

"Restricted"

Similar to and trainable as ANN:

- A *binary* graphic model with no intra-layer connections

- The configuration probability follows Boltzmann distribution

$$P(\mathbf{X}, \mathbf{H}) = \frac{1}{Z} \exp(-E(\mathbf{X}, \mathbf{H}))$$

$$E(\mathbf{X}, \mathbf{H}) = -\mathbf{X}^T \mathbf{b} - \mathbf{c}^T \mathbf{H} - \mathbf{X}^T \mathbf{W} \mathbf{H}$$

$$\mathbf{H} = (H_1, \dots, H_J)^T$$

$$\mathbf{X} = (X_1, \dots, X_I)^T$$

- \mathbf{W} , \mathbf{b} , and \mathbf{c} as model parameters, after training:

maximize the likelihood of given data
or fit to a given distribution



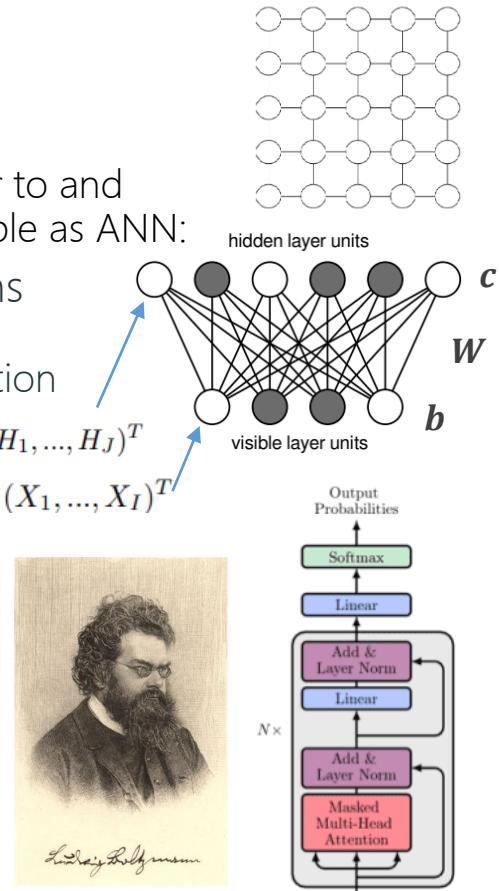
Generating handwritten digits:



- Generative Pre-trained Transformer (GPT)



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AI for Quantum MC Methods



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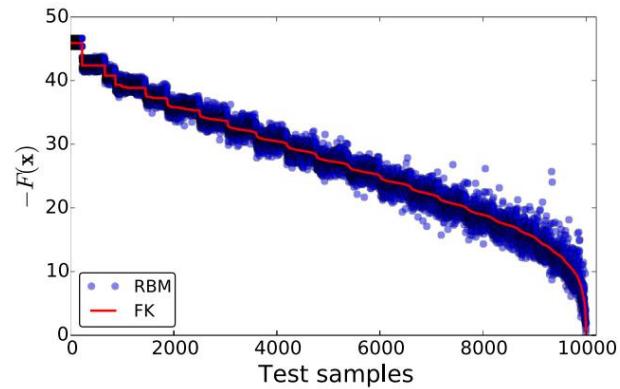
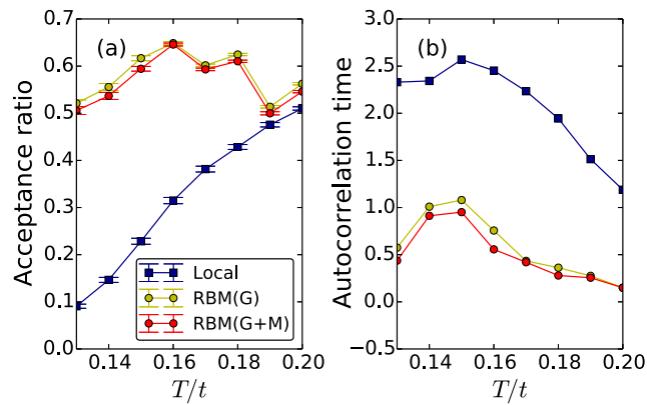
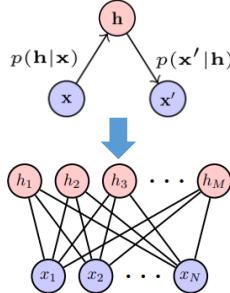
- Example: the Falicov-Kimball model on 2D square lattice

$$\hat{H}_{\text{FK}} = \sum_{i,j} \hat{c}_i^\dagger \mathcal{K}_{ij} \hat{c}_j + U \sum_{i=1}^N \left(\hat{n}_i - \frac{1}{2} \right) \left(x_i - \frac{1}{2} \right)$$
$$p_{\text{FK}}(\mathbf{x}) = e^{-F_{\text{FK}}(\mathbf{x})}/Z_{\text{FK}}$$
$$\begin{cases} x_i \in \{0,1\} \\ \mathcal{K}_{ij} = -t \quad U/t = 4 \\ \beta = 1/T \end{cases}$$

- The trained RBM successfully captures the probability distribution:

compensate with: $A(\mathbf{x} \rightarrow \mathbf{x}') = \min \left[1, \frac{p(\mathbf{x})}{p(\mathbf{x}')} \cdot \frac{p_{\text{FK}}(\mathbf{x}')}{p_{\text{FK}}(\mathbf{x})} \right]$

- Nonlocal updates from hidden variables:



drastically improved
acceptance rate and
autocorrelation time

Li Huang and Lei Wang, 2017.



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AI for Quantum Control and Optimization

Review: reinforcement learning



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Quantum Processes – Quantum Compiling



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- ◆ Classical computer:

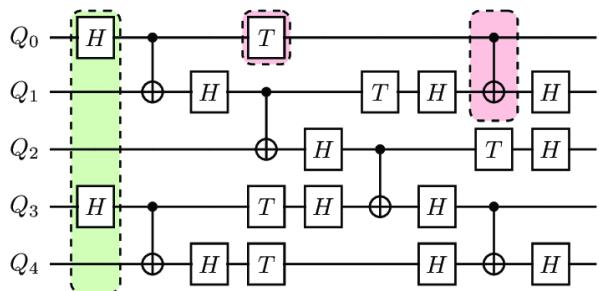


Logical gates:

AND		<table border="1"> <thead> <tr> <th>IN 1</th> <th>IN 2</th> <th>OUT</th> </tr> </thead> <tbody> <tr> <td>0</td> <td>0</td> <td>0</td> </tr> <tr> <td>0</td> <td>1</td> <td>0</td> </tr> <tr> <td>1</td> <td>0</td> <td>0</td> </tr> <tr> <td>1</td> <td>1</td> <td>1</td> </tr> </tbody> </table>	IN 1	IN 2	OUT	0	0	0	0	1	0	1	0	0	1	1	1
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- ◆ Quantum computer:



Fundamental quantum gates:

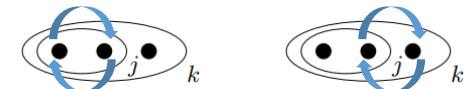
braiding of Fibonacci anyons

$$\text{Hadamard (H)} \quad \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

$$\pi/8 (\mathbf{T}) \quad \begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{bmatrix}$$

$$\sigma_1 = \begin{pmatrix} \eta^{-4} & 0 \\ 0 & \eta^3 \end{pmatrix} \quad \sigma_2 = \begin{pmatrix} -\phi^{-1}\eta^{-1} & \phi^{-\frac{1}{2}}\eta^{-3} \\ \phi^{-\frac{1}{2}}\eta^{-3} & -\phi^{-1} \end{pmatrix}$$

$$\eta = e^{i\pi/5} \quad \phi = (\sqrt{5} + 1)/2$$



Chetan Nayak, et al., 2008.

- Goal: find *fast* a *short* sequence $U \approx U_1^{n_1} U_2^{n_2} U_1^{n_3} U_2^{n_4} \dots$ close to U_{tar}
 - ◆ brute-force: good length complexity but bad time complexity
 - ◆ Solovay-Kitaev (recursive): good time complexity but bad length complexity



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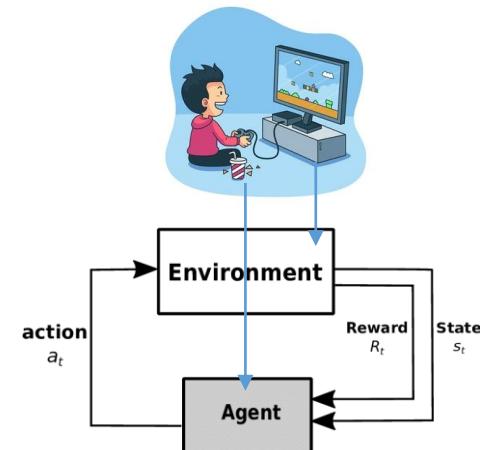
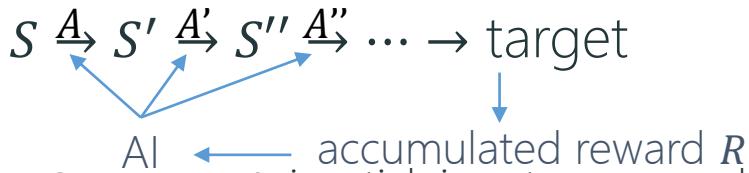


Reinforcement learning

- Reinforcement learning: An agent that interacts with an environment and maximizes reward (minimizes penalty)



S : current state; A : action upon state; R : reward



- Video games: S : screen; A : joystick input; R : score, life, cleared levels ...
- Chess, Go: S : current board configuration; A : next move; R : win ...
- Rubik's cube: S : current colorings; A : next twist; R : (minus) steps taken ...

Training the AI model self-consistently with the Bellman equation:

$$Q^{new}(s_t, a_t) \leftarrow \underbrace{Q(s_t, a_t)}_{\text{old value}} + \underbrace{\alpha}_{\text{learning rate}} \cdot \underbrace{\left(\underbrace{r_t}_{\text{reward}} + \underbrace{\gamma}_{\text{discount factor}} \cdot \underbrace{\max_a Q(s_{t+1}, a)}_{\text{estimate of optimal future value}} - \underbrace{Q(s_t, a_t)}_{\text{old value}} \right)}_{\text{temporal difference}}$$

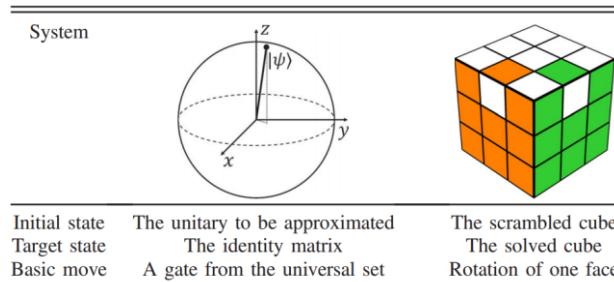
new value (temporal difference target)



AI for Quantum Control and Optimization

- Comparison between Rubik's cube and quantum compiling:

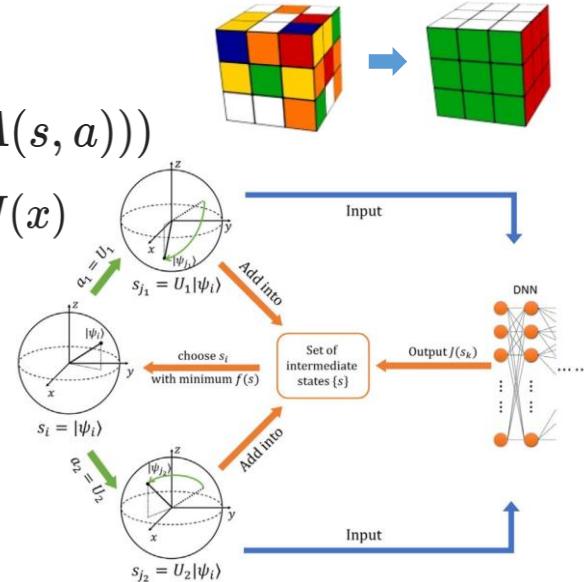
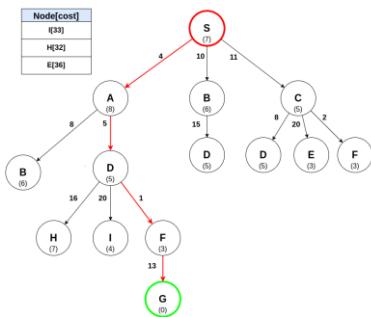
- ◆ S : current configuration / unitary U
- ◆ A : applied rotation / elementary gate U_i
- ◆ R : expected distance towards solution



- Combine the cost(-to-go) function:

$$\text{Q-learning: } JI(s) = \min_a (1 + J(A(s, a)))$$

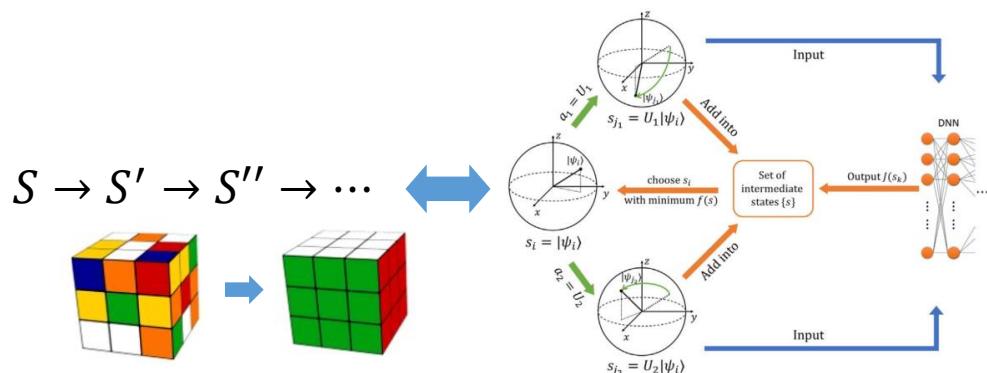
with the weighted A* search: $f(x) = \lambda g(x) + J(x)$



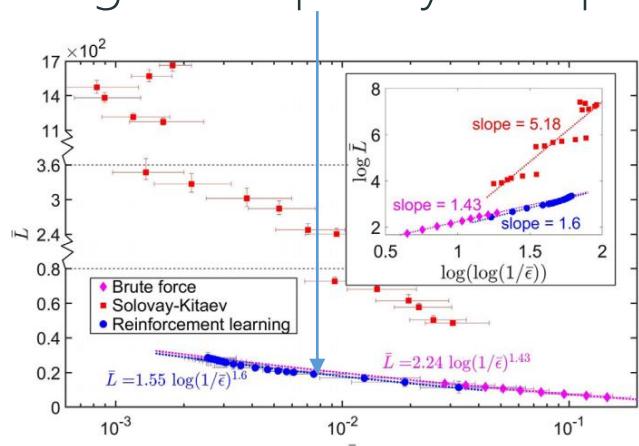
AI for Quantum Control and Optimization



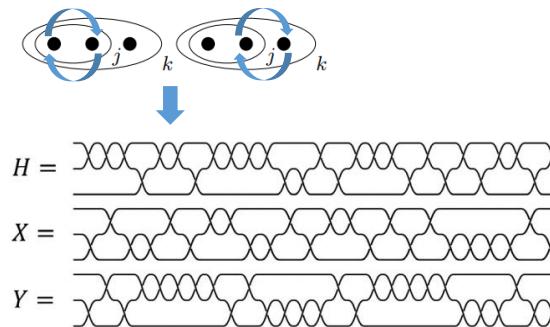
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- Time complexity: comparable to the SK recursion, very efficient
- Length complexity: comparable to brute force



Typical target-unitary examples :



better than
 $O(10^{-3})$ precision

A novel *good-enough* solver

Yuan-Hang Zhang, Pei-Lin Zheng, Yi Zhang, and Dong-Ling Deng, 2020.



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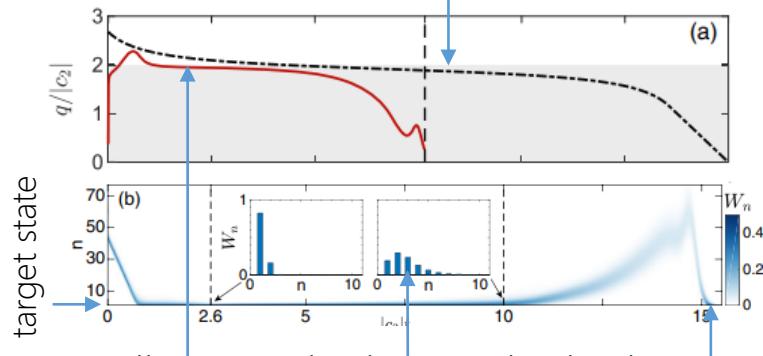


AI for Quantum Control and Optimization

- Quantum process – quantum state preparation, e.g., the Dicke state

$$H = \frac{c_2}{2N} \mathbf{L}^2 - q(t) N_0 \quad c_2 < 0 \quad |\psi_{\text{Dicke}}^{(0)}\rangle \equiv |N, L_z = 0\rangle$$

- Conventionally, adiabatic evolution, slowly turn off $q(t)$ to keep at ground state



- AI allows to think outside the box:

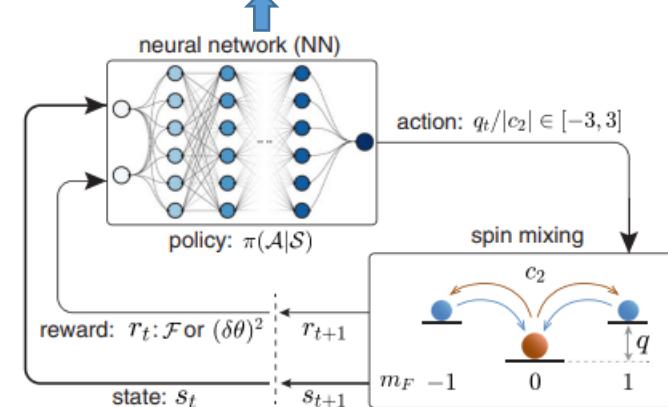
A faster process is obtained via reinforcement learning

$$s_t: \rho_0 = \langle N_0 \rangle / N, \langle \delta N_0^2 \rangle / N^2, |\langle a_{+1}^\dagger a_{-1}^\dagger a_0^2 \rangle| / N^2 \\ \theta_s = \arg \langle a_{+1}^\dagger a_{-1}^\dagger a_0^2 \rangle$$

Excited states are generated in the meantime – no adiabaticity

Nevertheless, final state large overlap with target $\mathcal{F} = |\langle \psi(t) | \psi_{\text{Dicke}}^{(0)} \rangle|^2$

Shuai-Feng Guo, et al. (2021)



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Summary



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- AI for quantum phases: numerical and experimental data and models
- AI for quantum methods: synergy and catalyst for algorithmic efficiency
- AI for quantum control: quantum compiling and state preparation
- Discussions:
 - ◆ No black magic: performance bounded from above by the quality of the samples.
 - ◆ Even for the best case scenario, AI methods are approximate.
 - ◆ Use the knowledge and intuition to improve, every bit helps!
 - ◆ Reverse thinking and consider AI for reverse thinking
 - ◆ Sometimes, trying an idea out is the best way to verify its practicality.
- We are still at an early stage of AI for Physics and Science.

*Know your target
and limitations!*



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