Computer modeling of physical phenomena: path integrals and quantum Monte Carlo

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Density matrix – central object of quantum statistics

Path integrals (by Feynman)

References

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Density matrix – central object of quantum statistics

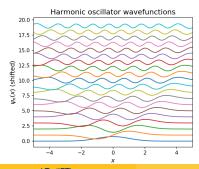
Path integrals (by Feynman)

Quantum harmonic oscillator

quantum particle in a potential

$$H = \frac{1}{2m}p^2 + \frac{1}{2}m\omega^2x^2, \quad p = -i\hbar\partial_x$$

- $\hbar = m = \omega = 1$ in our numerics
- wave functions $\psi_n(x)$ and eigenvalues $E_n = n + \frac{1}{2}$



 $\begin{array}{ll} \textbf{procedure harmonic-wavefunction} \\ \textbf{input} \ x \end{array}$

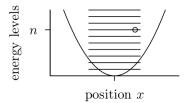
$$\psi_{-1}^{\text{h.o.}}(x) \leftarrow 0$$
 (unphysical, starts recursion)
 $\psi_{0}^{\text{h.o.}}(x) \leftarrow \pi^{-1/4} \exp\left(-x^2/2\right)$ (ground state)

for
$$n = 1, 2, \dots$$
 do

$$\begin{cases} \psi_n^{\text{h.o.}}(x) \leftarrow \sqrt{\frac{2}{n}} x \psi_{n-1}^{\text{h.o.}}(x) - \sqrt{\frac{n-1}{n}} \psi_{n-2}^{\text{h.o.}}(x) \\ \text{output } \{\psi_n^{\text{h.o.}}(x), \psi_1^{\text{h.o.}}(x), \dots \} \end{cases}$$

J.T. (IFT) – Path –

Thermal occupation



• thermodynamics – partition function:

$$\mathcal{Z}(\beta) = \sum_{n} e^{-\beta E_n} = \ldots = \frac{1}{2 \sinh(\beta/2)}$$

• particle at level "n" at position x, probability:

$$|\psi_n(x)|^2 \times e^{-\beta E_n}/\mathcal{Z}$$

Density matrix

thermal occupation (prob. of being at x):

$$\pi(x) = \frac{1}{\mathcal{Z}} \sum_{n} e^{-\beta E_n} |\psi_n(x)|^2$$

• general object, encapsulates knowledge of ψ_n , E_n

$$\rho(\mathbf{X}, \mathbf{X}', \beta) = \sum_{n} \psi_{n}^{*}(\mathbf{X}) \psi_{n}(\mathbf{X}') e^{-\beta \mathsf{E}_{n}}$$

diagonal part of density matrix

$$\mathcal{Z}(eta) = \int dx \,
ho(x, x, eta), \ \ \pi(x) =
ho(x, x, eta)/\mathcal{Z}$$

Free density matrix

• kinetic Hamiltonian $H_0 = \frac{1}{2m}p^2$; plane waves

$$\langle x|k\rangle = \frac{1}{\sqrt{2\pi}}e^{ikx} \ (p=\hbar k)$$

introducing resolution of unity

$$\rho_0(x,x',\beta) = \langle x|e^{-\beta H_0}|x'\rangle = \int \frac{dk}{2\pi} e^{ik(x-x')} e^{-\beta \frac{\hbar^2 k^2}{2m}}$$

performing the Gaussian integral we finaly get

$$\rho_0(x,x',\beta) = \sqrt{\frac{m}{2\pi\hbar^2\beta}} e^{-\frac{m}{2\hbar^2\beta}(x-x')^2}$$

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Path integrals (by Feynman)

Imaginary time

 note: there is a close analogy (formal mapping) between the density matrix and the unitary evolution (propagator):

$$\langle x|e^{-\beta H}|x'\rangle \iff \langle x|e^{-itH}|x'\rangle$$

- imaginary time $\beta \longleftrightarrow it$ real time
- convolution property (combined evolution in real time):

$$\int dx' \, \rho(x,x',\beta_1) \rho(x',x'',\beta_2) = \ldots = \rho(x,x'',\beta_1+\beta_2)$$

- Path -10/16

Path integral

applying repeatedly the convolution:

$$\rho(x_0,x_N,\beta) = \int \ldots \int dx_1 \ldots dx_{N-1} \, \rho(x_0,x_1,\beta/N) \ldots \rho(x_{N-1},x_N,\beta/N)$$

- path is a sequence $\{x_0, \dots, x_N\}$, we imagine x_k at time $k\Delta \tau$ $(\Delta \tau = \beta/N)$; path in d+1 dimensions
- density functions and partition function are represented as integrals over paths $\{x_0, \dots, x_N\}$ (formally also $N \to \infty$)

Trotter decomposition

• we use a (sec. order) Suzuki-Trotter decomposition

$$e^{-\Delta \tau (H_0+V)} = e^{-\frac{1}{2}\Delta \tau V} e^{-\Delta \tau H_0} e^{-\frac{1}{2}\Delta \tau V}$$

so

$$\rho(\mathbf{X}, \mathbf{X}', \Delta \tau) = \mathbf{e}^{-\frac{1}{2}\Delta \tau V(\mathbf{X})} \rho_0(\mathbf{X}, \mathbf{X}', \Delta \tau) \mathbf{e}^{-\frac{1}{2}\Delta \tau V(\mathbf{X}')}$$

• we know explicit formula for the free density matrix $\rho_0(x, x', \Delta \tau)$

Expansion for $\pi(x)$

• collecting all terms in the path integral (with x_0 fixed and p.b.c.)

$$\pi(x_0) = \frac{1}{Z} \int \mathcal{D}\{x_i\} e^{-\Delta \tau V(x_0)} \rho_0(x_0, x_1, \Delta \tau) \dots e^{-\Delta \tau V(x_{N-1})} \rho_0(x_{N-1}, x_0, \Delta \tau)$$

form of this path integral with "Euclidean" action S:

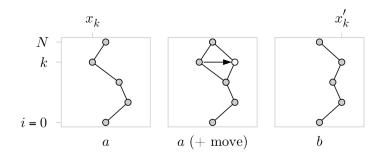
$$\pi(x_0) = \frac{1}{\mathcal{Z}} \int \mathcal{D}\{x_i\} e^{-S(\{x_i\})},$$

where
$$S(\{x_i\}) = \sum_{i=2\Delta\tau} \frac{1}{2\Delta\tau} (x_i - x_{i+1})^2 + \sum_{i=2\Delta\tau} \Delta\tau V(x_i)$$

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Path integrals (by Feynman)

Path sampling algorithm



- choose random element k, accept the move $x_k \to x_k + \delta$ with Metropolis algorithm
- only the segments $\{x_{k-1}, x_k\}$ and $\{x_k, x_{k+1}\}$ are involved

Pseudo-code

```
procedure naive-harmonic-path
input \{x_0, ..., x_{N-1}\}
\Delta_{\tau} \leftarrow \beta/N
k \leftarrow \operatorname{nran}(0, N-1)
k_{\pm} \leftarrow k \pm 1 \text{ modulo N}
x'_{k} \leftarrow x_{k} + \operatorname{ran}(-\delta, \delta)
\pi_a \leftarrow \rho^{\text{free}}\left(x_{k_-}, x_k, \Delta_\tau\right) \rho^{\text{free}}\left(x_k, x_{k_+}, \Delta_\tau\right) \exp\left(-\frac{1}{2}\Delta_\tau x_k^2\right) \\ \pi_b \leftarrow \rho^{\text{free}}\left(x_{k_-}, x_k', \Delta_\tau\right) \rho^{\text{free}}\left(x_k', x_{k_+}, \Delta_\tau\right) \exp\left(-\frac{1}{2}\Delta_\tau x_k^2\right) 
\Upsilon \leftarrow \pi_b/\pi_a
if (\operatorname{ran}(0,1) < \Upsilon)x_k \leftarrow x_k'
output \{x_0, ..., x_{N-1}\}
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