## Numerical Solver for NH2018 S3

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## 1 Discretization of Equation (S3)

Equation (S3) in NH2018 reads

$$\frac{\partial}{\partial \mu} \left[ \frac{1}{2\Omega \mu} \frac{\partial}{\partial \mu} (u_{\text{REF}} \cos \phi) \right] + \frac{2\Omega a^2 H \mu}{R(1 - \mu^2)} e^{z/H} \frac{\partial}{\partial z} \left[ \frac{e^{(\kappa - 1)z/H}}{\frac{\partial \tilde{\theta}}{\partial z}} \frac{\partial}{\partial z} u_{\text{REF}} \cos \phi \right] = -a \frac{\partial}{\partial \mu} \left( \frac{q_{\text{REF}}}{2\Omega \mu} \right)$$
(1)

where  $\mu \equiv \sin \phi$ . With the substitutions  $\tilde{u} \equiv u_{\text{REF}} \cos \phi$ ,  $\tilde{q} \equiv \frac{q_{\text{REF}}}{2\Omega\mu}$  and static stability  $S \equiv \frac{\partial \tilde{\theta}}{\partial z}$ , equation (1) becomes

$$\frac{\partial}{\partial \mu} \left[ \frac{1}{2\Omega \mu} \frac{\partial}{\partial \mu} \tilde{u} \right] + \frac{2\Omega a^2 H \mu}{R(1 - \mu^2)} e^{z/H} \frac{\partial}{\partial z} \left[ \frac{e^{(\kappa - 1)z/H}}{S} \frac{\partial \tilde{u}}{\partial z} \right] = -a \frac{\partial \tilde{q}}{\partial \mu}$$
 (2)

Discretizing each term in equation (2) on the uniform  $\phi$  (index j) and z (index k) grids yields:

$$\begin{split} &\frac{\partial}{\partial \mu} \left[ \frac{1}{2\Omega \mu} \frac{\partial}{\partial \bar{\mu}} \tilde{u} \right] \\ &= \frac{1}{\cos \phi} \frac{\partial}{\partial \bar{\phi}} \left[ \frac{1}{2\Omega \sin \phi \cos \phi} \frac{\partial}{\partial \bar{\phi}} \tilde{u} \right] \\ &\approx \frac{1}{2\Omega (\Delta \phi)^2 \cos \phi_j} \left[ \frac{\tilde{u}_{j+1,k}}{\sin \phi_{j+1/2} \cos \phi_{j+1/2}} + \tilde{u}_{j,k} \left( \frac{1}{\sin \phi_{j+1/2} \cos \phi_{j+1/2}} + \frac{1}{\sin \phi_{j-1/2} \cos \phi_{j-1/2}} \right) + \frac{\tilde{u}_{j-1,k}}{\sin \phi_{j-1/2} \cos \phi_{j-1/2}} \right] \\ &\frac{2\Omega a^2 H \mu}{R(1 - \mu^2)} e^{z/H} \frac{\partial}{\partial z} \left[ \frac{e^{(\kappa - 1)z/H}}{\frac{\partial \bar{\phi}}{\partial z}} \frac{\partial}{\partial z} u_{\text{REF}} \cos \phi \right] \\ &= \frac{2\Omega a^2 H \sin \phi}{R \cos^2 \phi} e^{z/H} \frac{\partial}{\partial z} \left[ \frac{e^{(\kappa - 1)z/H}}{\frac{\partial \bar{\phi}}{\partial z}} \frac{\partial}{\partial z} u_{\text{REF}} \cos \phi \right] \\ &\approx \frac{2\Omega^2 a^2 H \sin \phi_j}{R \cos^2 \phi_j} \frac{e^{z_k/H}}{(\Delta z)^2} \left[ \tilde{u}_{j,k+1} \frac{e^{\frac{(\kappa - 1)z_{k+1/2}}{H}}}{S_{k+1/2}} - \tilde{u}_{j,k} \left( \frac{e^{\frac{(\kappa - 1)z_{k+1/2}}{H}}}{S_{k+1/2}} + \frac{e^{\frac{(\kappa - 1)z_{k-1/2}}{H}}}{S_{k-1/2}} \right) + \tilde{u}_{j,k-1} \frac{e^{\frac{(\kappa - 1)z_{k-1/2}}{H}}}{S_{k-1/2}} \right] \\ &- a \frac{\partial \tilde{q}}{\partial \mu} \\ &= - \frac{a}{\cos \phi} \frac{\partial \tilde{q}}{\partial \phi} \\ &\approx - \frac{a}{2\Delta \phi \cos \phi_j} \left( \tilde{q}_{j+1} - \tilde{q}_{j-1} \right) \end{aligned} \tag{5}$$

Substituting (3), (4) and (5) into (2) gives

$$a_{i,k}\tilde{u}_{i+1,k} + b_{i,k}\tilde{u}_{i-1,k} + c_{i,k}\tilde{u}_{i,k+1} + d_{i,k}\tilde{u}_{i,k-1} + e_{i,k}\tilde{u}_{i,k} = f_{i,k} \tag{6}$$

where

$$a_{j,k} = \frac{1}{\sin \phi_{j+1/2} \cos \phi_{j+1/2}} \tag{7}$$

$$b_{j,k} = \frac{1}{\sin \phi_{j-1/2} \cos \phi_{j-1/2}} \tag{8}$$

$$c_{j,k} = \frac{2\Omega^2 a^2 H \sin \phi_j}{R \cos \phi_j} e^{z_k/H} \frac{(\Delta \phi)^2}{(\Delta z)^2} \frac{e^{\frac{(\kappa - 1)z_{k+1/2}}{H}}}{S_{k+1/2}}$$
(9)

$$d_{j,k} = \frac{2\Omega^2 a^2 H \sin \phi_j}{R \cos \phi_j} e^{z_k/H} \frac{(\Delta \phi)^2}{(\Delta z)^2} \frac{e^{\frac{(\kappa - 1)z_{k-1/2}}{H}}}{S_{k-1/2}}$$
(10)

$$e_{j,k} = -(a_{j,k} + b_{j,k} + c_{j,k} + d_{j,k})$$
(11)

$$f_{j,k} = -\Omega a(\Delta \phi)(\tilde{q}_{j+1} - \tilde{q}_{j-1}) \tag{12}$$

Equation (6) is then solved using the successive over-relaxation algorithm.

## 2 The Successive Over-relaxation (SOR) Solver

The iterative procedure to solve equation (6) is

$$\tilde{u}_{j,k}^{new} = \tilde{u}_{j,k}^{old} - \omega \frac{\zeta_{j,k}}{e_{j,k}} \tag{13}$$

where  $\zeta_{j,k}$  is the residual

$$\zeta_{j,k} \equiv a_{j,k} \tilde{u}_{j+1,k} + b_{j,k} \tilde{u}_{j-1,k} + c_{j,k} \tilde{u}_{j,k+1} + d_{j,k} \tilde{u}_{j,k-1} + e_{j,k} \tilde{u}_{j,k} - f_{j,k}$$
(14)

and the over-relaxation parameter at the n-th step is given by

$$\omega_n = \begin{cases} \frac{1}{1 - \rho_{\text{jacobi}}^2/2} & \text{when } n = 1\\ \frac{1}{1 - \rho_{\text{jacobi}}^2 \omega_{n-1}/4} & \text{when } n > 1 \end{cases}$$

In NH2018,  $\rho_{\rm jacobi}^2$  used is 0.95. The condition to stop the iteration is when

$$\sum_{1}^{j_{\max}-1} \sum_{1}^{k_{\max}-1} \|\zeta_{j,k}\| < \epsilon \sum_{1}^{j_{\max}-1} \sum_{1}^{k_{\max}-1} \|f_{j,k}\|$$
(15)

where the tolerance parameter  $\epsilon = 10^{-5}$  in NH2018.

## 3 The boundary conditions

The boundary conditions for  $\tilde{u}$  are:

$$\tilde{u}(\phi = \frac{\pi}{2}, z = 0) = 0 \tag{16}$$

$$\tilde{u}(\phi = 0) = (\langle u \rangle + A) \tag{17}$$

(18)