

## 1 Slater-type orbitals

A real Slater-type orbital (STO) is defined as follows:

$$\chi_{nlm}(\zeta, \mathbf{r}) = N_n(\zeta) R_n(r, \zeta) S_{lm}(\theta, \phi) \quad (1)$$

where

$$N_n(\zeta) = \frac{(2\zeta)^{n+\frac{1}{2}}}{\sqrt{(2n)!}} \quad (2)$$

is the normalization factor,

$$R_n(r, \zeta) = r^{n-1} e^{-\zeta r} \quad (3)$$

is the radial part and  $S_{lm}(\theta, \phi)$  is a real spherical harmonic.

## 2 Prolate spheroidal coordinates

For two-center integrals, it is more convenient to align the coordinate system of the two centers along the  $z$ -axis and use prolate spheroidal coordinates. In order to switch from spherical coordinates to the new coordinate system, we use the following identities:

$$\begin{aligned} r_a &= \frac{R(\mu + \nu)}{2}; & \cos \theta_a &= \frac{1 + \mu\nu}{\mu + \nu}; & \sin \theta_a &= \frac{\sqrt{(\mu^2 - 1)(1 - \nu^2)}}{\mu + \nu}; \\ r_b &= \frac{R(\mu - \nu)}{2}; & \cos \theta_b &= \frac{1 - \mu\nu}{\mu - \nu}; & \sin \theta_b &= \frac{\sqrt{(\mu^2 - 1)(1 - \nu^2)}}{\mu - \nu} \end{aligned} \quad (4)$$

Inside the integrals this change gives

$$dV = \frac{R^3}{8} (\mu + \nu)(\mu - \nu) d\mu d\nu d\phi \quad (5)$$

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## 3 Overlap integrals

The one-center overlap integrals

$$S_{nlm, n'l'm'}(\zeta, \zeta') = \int \chi_{nlm}^*(\zeta, \mathbf{r}) \chi_{n'l'm'}(\zeta', \mathbf{r}) dV \quad (6)$$

is the most straightforward molecular integrals and is readily evaluated in spherical coordinates. The integration of the radial part yields

$$\int_0^\infty R_n(\zeta, r) R_{n'}(\zeta', r) r^2 dr = \frac{(n + n')!}{(\zeta + \zeta')^{n+n'+1}} \quad (7)$$

and using the orthonormality condition of the spherical harmonics, the analytical form of one-center overlap is

$$S_{nlm, n'l'm'}(\zeta, \zeta') = N_n(\zeta) N_{n'}(\zeta') \frac{(n + n')!}{(\zeta + \zeta')^{n+n'+1}} \delta_{ll'} \delta_{mm'} \quad (8)$$