

# The Quantum Monte Carlo Simulation of Transverse Field Ising Model

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January 5, 2021

# Outline

**1** Introduction to Transverse Field Ising

2 Worm Algorithm

3 C++ Implement

4 Numerical Results

# Classical Ising Model

Hamiltonian:

$$\mathcal{H} = -t \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z - h_z \sum_i \sigma_i^z$$

Thermal Fluctuation  $\implies$  Thermal(Classical) Phase Transition

Low T: Ordered state

High T: Paramagnetic state

| Exponent |             | Definition                    | Ising Value |
|----------|-------------|-------------------------------|-------------|
| $\alpha$ | $C$         | $\propto (T - T_c)^{-\alpha}$ | 0           |
| $\beta$  | $M$         | $\propto (T_c - T)^\beta$     | 1/8         |
| $\gamma$ | $\chi$      | $\propto (T - T_c)^{-\gamma}$ | 7/4         |
| $\delta$ | $M$         | $\propto h^{1/\delta}$        | 15          |
| $\nu$    | $\xi$       | $\propto (T - T_c)^{-\nu}$    | 1           |
| $\eta$   | $\Gamma(n)$ | $\propto  n ^{2-d-\eta}$      | 1/4         |

Critical exponents of Classical Ising Model

# Quantum Ising Model

Hamiltonian of Transverse Field Ising model (TFI):

$$\mathcal{H} = -t \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z - h \sum_i \sigma_i^x$$

$[\sigma_i^z, \sigma_i^x] \neq 0 \implies$  Quantum Fluctuation!

Quantum Phase Transition:

$h \ll t$ : Spins aligned in z-direction.

$h \gg t$ : Spins aligned in x-direction.

# Theoretical Calculation

Jordan-Wigner transformation:

$$\sigma_i^x = 1 - 2c_i^\dagger c_i$$

$$\sigma_i^z = -\prod_{j<i} (1 - 2c_j^\dagger c_j) (c_i + c_i^\dagger)$$

$c_i$ s are fermion operators.

Hamiltonian becomes:

$$\mathcal{H} = -t \sum_i (c_i^\dagger c_{i+1} + c_{i+1}^\dagger c_i + c_i^\dagger c_{i+1}^\dagger + c_{i+1} c_i - 2g c_i^\dagger c_i + g)$$

with  $g = h/t$

# Theoretical Calculation

Solve it by Bogoliubov transformation:

$$\gamma_k = u_k c_k - v_k c_{-k}^\dagger$$

$$\gamma_k^\dagger = u_k c_k^\dagger - v_k c_{-k}$$

with  $c_k$  Fourier transformation of  $c_i$ . And

$$u_k = \cos\left(\frac{\theta_k}{2}\right), v_k = \sin\left(\frac{\theta_k}{2}\right)$$

$$\tan(\theta_k) = \frac{\sin(ka)}{g - \cos(ka)}$$

# Theoretical Calculation

Switch to Majorana representation (continuous limit):

$$\psi(x_i) = -i(c_i - c_i^\dagger) = \frac{1}{2N} \sum_k (u_1(k) \gamma_k^\dagger e^{ikx_i} + h.c.)$$

$$\bar{\psi}(x_i) = c_i + c_i^\dagger = \frac{1}{2N} \sum_k (u_2(k) \gamma_k^\dagger e^{ikx_i} + h.c.)$$

$$\psi(x_i, t) = e^{i\mathcal{H}t} \psi(x_i) e^{-i\mathcal{H}t}, \bar{\psi}(x_i, t) = e^{i\mathcal{H}t} \bar{\psi}(x_i) e^{-i\mathcal{H}t}$$

$\implies$  Solution of 1+1D Dirac equation



# Theoretical Calculation

The action:

$$S = \frac{i}{2} \int \psi(\partial_0 - \partial_1)\psi + \bar{\psi}(\partial_0 + \partial_1)\bar{\psi} + (1 - g)\psi\bar{\psi}$$

$g = 1 \implies$  massless, scale invariance, i.e.

$$S' = S \text{ when } x \rightarrow \lambda x$$

$\implies$  Critical Point at  $h = t!!$

The two point correlation functions at criticality:

$$\langle \sigma_i^z \sigma_{i+n}^z \rangle \propto \frac{1}{|n|^{1/4}}$$

$$\langle \sigma_i^x \sigma_{i+n}^x \rangle \propto \frac{1}{|n|^2}$$

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# Holstein-Primakoff Transformation

TFI:

$$\mathcal{H} = -t \sum_{\langle ij \rangle} \sigma_i^x \sigma_j^x - h \sum_i \sigma_i^z = K + U$$

Hopping and pairing terms:

$$K = K_1 + K_2 = -t \sum_{\langle ij \rangle} (\sigma_i^+ \sigma_j^- + \text{h.c.}) - t \sum_{\langle ij \rangle} (\sigma_i^+ \sigma_j^+ + \text{h.c.})$$

HP transformation:  $b_i(b_i^\dagger) = \sigma_i^+(\sigma_i^-)$ ,  $n_i = b_i^\dagger b_i = (\sigma_i^z + 1)/2$

$$\Rightarrow \mathcal{H} = -t \sum_{\langle ij \rangle} (b_i^\dagger b_j + \text{h.c.}) - t \sum_{\langle ij \rangle} (b_i^\dagger b_j^\dagger + \text{h.c.}) - \mu \sum_i n_i$$

where  $\mu = 2h$ .

# $\mathcal{Z}$ Configurations

$$\begin{aligned}
 \mathcal{Z} &= \text{Tr} \left( e^{-\beta \mathcal{H}} \right) = \sum_{\alpha_0} \langle \alpha_0 | e^{-\beta \mathcal{H}} | \alpha_0 \rangle \\
 &= \lim_{d\tau = \frac{\beta}{n}, n \rightarrow \infty} \sum_{\{\alpha_i\}} \langle \alpha_0 | e^{-\mathcal{H}d\tau} | \alpha_{n-1} \rangle \cdots \langle \alpha_i | e^{-\mathcal{H}d\tau} | \alpha_0 \rangle \\
 &= \sum_{\alpha_0} \sum_{\mathcal{N}} \int_0^\beta \int_{\tau_1}^\beta \cdots \int_{\tau_{\mathcal{N}-1}}^\beta \prod_{k=1}^{\mathcal{N}} d\tau_k F(t, h) \\
 F(t, h) &= t^{\mathcal{N}_h + \mathcal{N}_p} e^{-\int_0^\beta U(\tau) d\tau}
 \end{aligned}$$

$\mathcal{N}_h$  and  $\mathcal{N}_p$ : number of hopping and pairing kinks.  $\mathcal{N}_h + \mathcal{N}_p = \mathcal{N}$ .

$\mathcal{Z}$  configuration: several loops of spin-up (and kinks) in the  $(d+1)$ D space-time.

Statistical weight of a  $\mathcal{Z}$  configuration:

$$W_{\mathcal{Z}}(t, h) = \prod_{k=1}^{\mathcal{N}} d\tau F(t, h)$$

# $\mathcal{G}$ Configurations

$$\mathcal{G}(\mathbf{x}_I, \tau_I; \mathbf{x}_M, \tau_M) = \text{Tr} \left[ T_\tau \left( \sigma_I^x(\tau_I) \sigma_M^x(\tau_M) e^{-\beta \mathcal{H}} \right) \right]$$

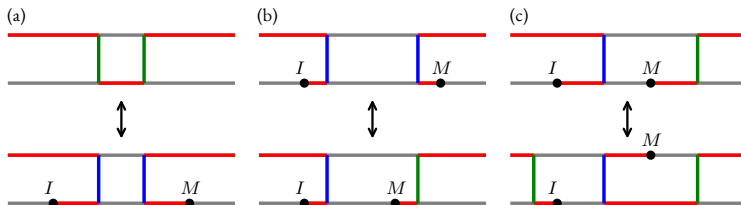
$\mathcal{G}$  configuration:  $\mathcal{Z}$  configuration + an open path of spin-up with two ending points “Ira” ( $I$ ) and “Masha” ( $M$ ).

Statistical weight of a  $\mathcal{G}$  configuration:

$$W_{\mathcal{G}} = \frac{d\tau_I d\tau_M}{\omega_G} \prod_{k=1}^{\mathcal{N}} d\tau_k F(t, h)$$

- Key idea: move the defect  $M$  of the open path to produce different  $\mathcal{G}$  and  $\mathcal{Z}$  configurations for sampling. (Like a worm wriggling!)

# Updates



(a) Create/annihilate defects  $I$  and  $M$ .

$$\frac{d\tau_I}{\beta N} \frac{d\tau_M}{\tau_a} \cdot W_Z \cdot \mathcal{P}_{\text{crea}} = \mathcal{A}_a \cdot W_G \cdot \mathcal{P}_{\text{anni}} \Rightarrow \begin{cases} P_{\text{crea}} = \min \left[ 1, \mathcal{A}_a \tau_a \frac{\beta N}{\omega_G} \frac{F_{\text{new}}}{F_{\text{old}}} \right] \\ P_{\text{anni}} = \min \left[ 1, \frac{1}{\mathcal{A}_a \tau_a} \frac{\omega_G}{\beta N} \frac{F_{\text{new}}}{F_{\text{old}}} \right] \end{cases}$$

(b) Move imaginary time of defect  $M$ .

$$P_{\text{move}} = \min \left\{ 1, \frac{F_{\text{new}}}{F_{\text{old}}} \right\}$$

(c) Insert/delete a kink.

$$\mathcal{A}_c \cdot \frac{1}{z_d} \cdot \frac{d\tau}{\tau_c} \cdot W \cdot \mathcal{P}_{\text{inse}} = \mathcal{A}_c \cdot \frac{1}{z_d} \cdot \frac{1}{n_k} \cdot W_+ \cdot \mathcal{P}_{\text{dele}} \Rightarrow \begin{cases} P_{\text{inse}} = \min \left[ 1, \frac{\tau_c}{n_k + 1} \frac{F_{\text{new}}}{F_{\text{old}}} \right] \\ P_{\text{dele}} = \min \left[ 1, \frac{n_k}{\tau_c} \frac{F_{\text{new}}}{F_{\text{old}}} \right] \end{cases}$$

$\tau_a, \tau_b, \tau_c$ : ranges for choosing random time displacement.

*A priori* probabilities:  $\mathcal{A}_a + \mathcal{A}_b + 2\mathcal{A}_c = 1$ .

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# The Grid

- ▶ Use one bit to represent one spin
- ▶ Use one bit to represent one kink between two world lines
- ▶ Use bitwise **XOR** with mask to flip spins and calculate the correlation
- ▶ Use **OR** to get the state of one spin
- ▶ Use bitwise **OR** to get the positions without kinks
- ▶ Use **\_\_popcnt** to count spin-up number
- ▶ Use **\_tzcnt\_u32** to find the first set bit to find the starting point of one loop in the grid

# Warping Number of Loops

- ▶ Depth first search
- ▶ Must traverse the whole grid to find all loops and calculate the warpping number of each loop
- ▶ This doesn't affect the performance since its time consumption is about 1% of creating and annihilating once
- ▶ Use a direction to make sure that the current point doesn't goes it way back
- ▶ Mark the spins “walked” to avoid traversing one loop twice

# Miscellaneous

- ▶ Use OpenGL to implement visualization
- ▶ Reached 1000 times creating and annihilating per second under grid size  $N=64$  with single thread
- ▶ For  $R_{\downarrow}$ , we can simply change the sign of  $h$

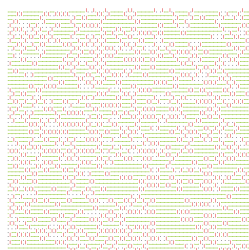
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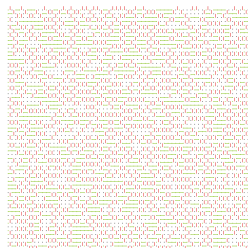
# Visualization



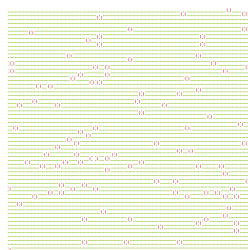
(a)  $h = 0.01$ .



(c)  $h = 1.0$ .



(b)  $h = 0.5$ .



(d)  $h = 1.5$ .

# Calculated Quantities

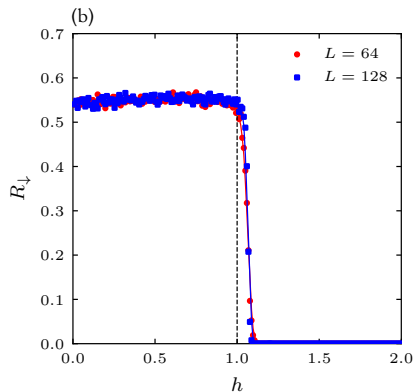
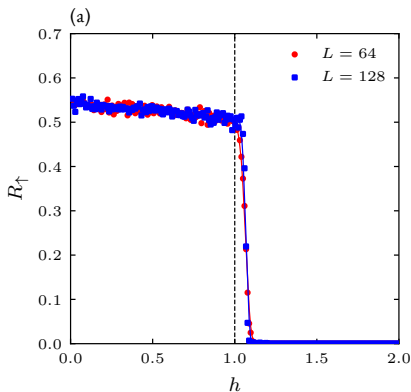
- ▶ Topological quantities: winding number  $W^l \Rightarrow$  total winding number  $W \Rightarrow$  wrapping probability  $R$ .
- ▶ Physical quantities: correlation functions.

$$\overline{\sigma_i^z} \equiv \frac{1}{L^2} \sum_{j,k} \sigma_{j,k}^z$$

$$\overline{\sigma_i^z \sigma_{i+n}^z} \equiv \frac{1}{L^2} \sum_{j,k} \sigma_{j,k}^z \sigma_{j+n,k}^z$$

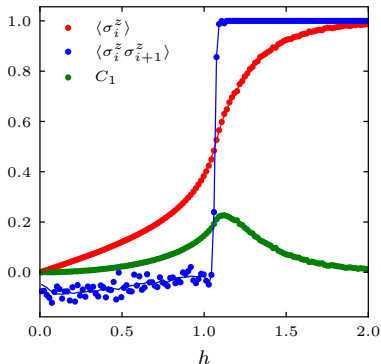
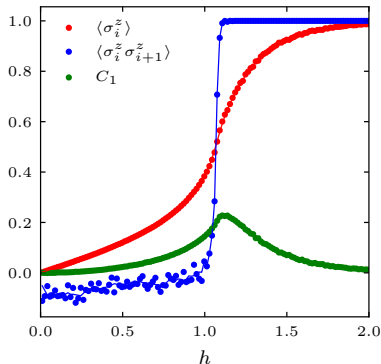
$$C_n \equiv \overline{\sigma_i^z \sigma_{i+n}^z} - \overline{\sigma_i^z}^2$$

# Wrapping Probabilities



Wrapping probabilities  $R_{\uparrow}$  and  $R_{\downarrow}$  versus the transverse field  $h$ . (a) Spin-up. (b) Spin-down.

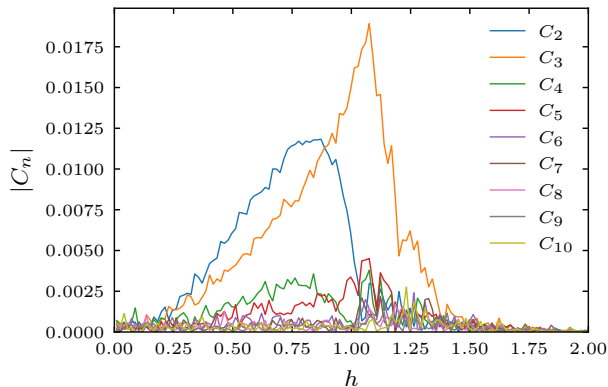
# Correlation Functions



Expectation values of spin states  $\overline{\sigma_i^z}$ , product of NN spins  $\overline{\sigma_i^z \sigma_{i+1}^z}$  and correlation function  $C_1$  versus transverse field  $h$ . (a)  $L = 64$ . (b)  $L = 128$ .



# Correlation Functions



Correlation functions  $|C_n|$  versus the transverse field  $h$ .

*Thank you for listening!*