

# The Physics of Simple Waves

## Shock Waves

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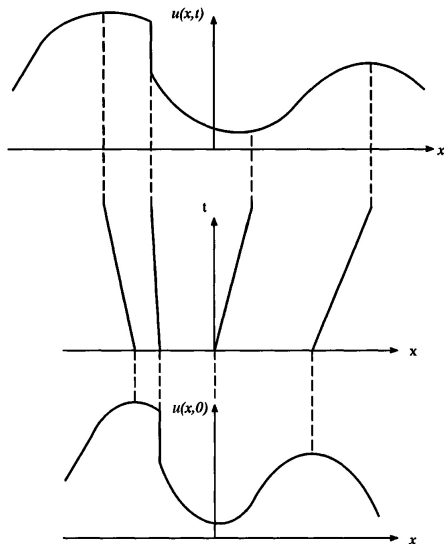
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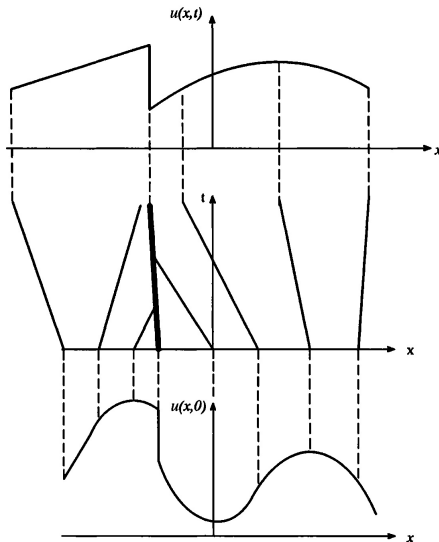
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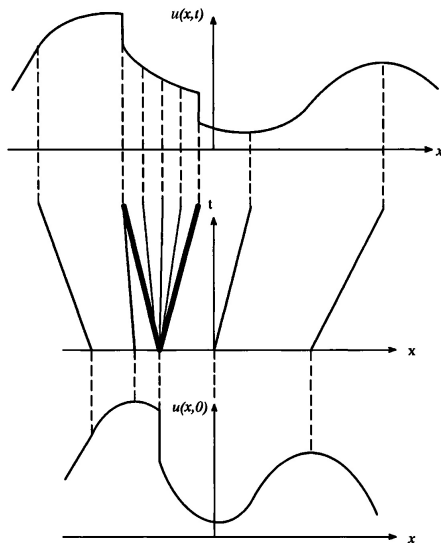
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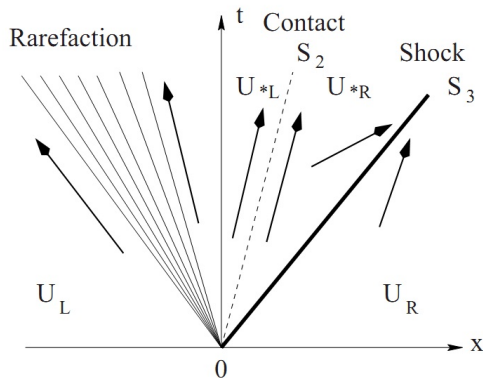
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- Introduction to Simple Waves



- Rankine Hugoniot Relation help us to describe the behavior of shockwaves traveling normal to the prevailing flow.

$$\vec{f}_R - \vec{f}_L = S(\vec{u}_R - \vec{u}_L) \quad (1)$$

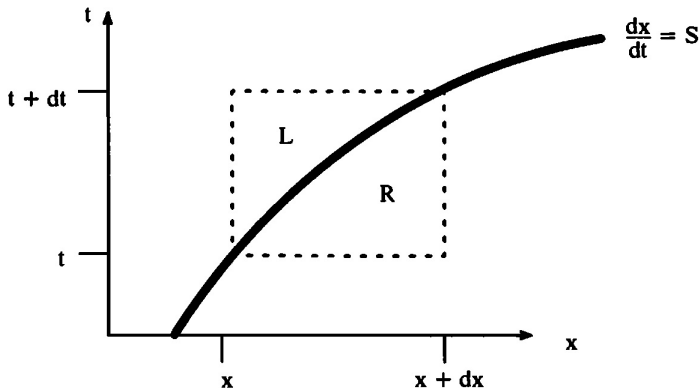
- Vectors of conserved quantities:

$$\vec{u} = \begin{bmatrix} \rho \\ \rho u \\ \rho \mathbf{e}_T \end{bmatrix} = \begin{bmatrix} u_1 \\ u_2 \\ u_3 \end{bmatrix} \quad (2)$$

$$\vec{f} = \begin{bmatrix} \rho u \\ \rho u^2 + p \\ (\rho \mathbf{e}_T + p)u \end{bmatrix} = \begin{bmatrix} f_1 \\ f_2 \\ f_3 \end{bmatrix} \quad (3)$$



- Graphically:



**Figure 2.2** The derivation of the Rankine–Hugoniot relations.

- Suppose that a shock that separates two regions of uniform flow  $u_L$  and  $u_R$  is traveling at a constant speed.

$$\rho_L(u_L - S) = \rho_R(u_R - S) \quad (4)$$

$$\rho_L(u_L - S)^2 + p_L = \rho_R(u_R - S)^2 + p_R \quad (5)$$

$$\rho_L(u_L - S)(E_L) = \rho_R(u_R - S)(E_R)$$

- we notice that the energy relation can be simplified by using relation (4) and:

$$E = \frac{1}{2}u^2 + h$$

- Which would lead us to a preferred form of the energy description:

$$E_L = E_R$$

$$\frac{1}{2}(u_L - S)^2 + h_L = \frac{1}{2}(u_R - S)^2 + h_R$$

- Suppose that a shock that separates two regions of uniform flow  $u_L$  and  $u_R$  is traveling at a constant speed  $S$ .

$$\begin{aligned}\rho_L(u_L - S) &= \rho_R(u_R - S) \\ \rho_L(u_L - S)^2 + p_L &= \rho_R(u_R - S)^2 + p_R \\ \frac{1}{2}(u_L - S)^2 + h_L &= \frac{1}{2}(u_R - S)^2 + h_R\end{aligned}\tag{6}$$

- We are basically targeting three main relations to describe the physics of the shock wave:
  - Energy Relation:  $\rho_R$  and  $\rho_L$
  - Flux Relation:  $u_R$  and  $u_L$
  - Pressure Relation:  $P_R$  and  $P_L$
- To do so, we start with Equations that we derive from the Hugoniot Relation, to formulate a way to describe this relations.

- Combining the eqns (4) & (5)

$$\rho_L \left( \frac{\rho_R}{\rho_L} (u_R - S)^2 \right) = \rho_R (u_R - S)^2 + P_R - P_L$$

- we can develop the following two relations:

$$(u_R - S)^2 = \frac{\rho_L}{\rho_R} \left( \frac{P_R - P_L}{\rho_R - \rho_L} \right) \quad (7)$$

$$(u_L - S)^2 = \frac{\rho_R}{\rho_L} \left( \frac{P_R - P_L}{\rho_R - \rho_L} \right) \quad (8)$$

- Substituting (7) and (8) in (6):

$$\frac{1}{2}(u_L - S)^2 + h_L = \frac{1}{2}(u_R - S)^2 + h_R$$

- we can now get a relation of enthalpies:

$$\frac{1}{2}(P_L - P_R) \left( \frac{1}{\rho_L} + \frac{1}{\rho_R} \right) = h_R - h_L \quad (9)$$

- Now using the general formulation of enthalpies:

$$h = e + \frac{P}{\rho}$$

- where  $e$  is internal energy,
- By performing some algebraic manipulations we can obtain:

$$\frac{1}{2} (P_L + P_R) \left( \frac{1}{\rho_L} - \frac{1}{\rho_R} \right) = e_R - e_L \quad (10)$$

- we have obtain the first of the relations we need.



- but we know that:

$$e = c_v T, \quad c_v = \frac{R}{\gamma - 1}, \quad P = \rho R T$$

- Therefore we can describe internal energy as:

$$e = \frac{P}{(\gamma - 1)\rho}$$

- Substituting the new formulation of internal energy in (10)

$$\frac{1}{2} (P_L + P_R) \left( \frac{1}{\rho_L} - \frac{1}{\rho_R} \right) = e_R - e_L$$

- and performing some algebraic manipulations we can obtain:

$$\frac{\rho_L}{\rho_R} = \frac{\left( \frac{P_L}{P_R} \right) + \left( \frac{\gamma-1}{\gamma+1} \right)}{\left( \frac{\gamma-1}{\gamma+1} \right) \left( \frac{P_L}{P_R} \right) + 1} \quad (11)$$

- Wich establish a very useful relation between the density ratio  $\rho_L/\rho_R$

- We wish to derivate a relation  $P_L$  &  $P_R$ ; but first, it is convenient to developpe a relation  $a_L/a_R$  using the  $\rho_L/\rho_R$  ratio that we already now.
- using the general fomulation of sound speed "a" for a perfect gas:

$$a = \gamma RT = \frac{\gamma P}{\rho}$$

$$\frac{a_R^2}{a_L^2} = \frac{P_R}{P_L} \left( \frac{\rho_L}{\rho_R} \right)$$

- and performing some algebraic manipulations we can obtain:

$$\frac{a_R^2}{a_L^2} = \frac{P_R}{P_L} \left( \frac{\left(\frac{P_L}{P_R}\right) + \left(\frac{\gamma-1}{\gamma+1}\right)}{\left(\frac{\gamma-1}{\gamma+1}\right) \left(\frac{P_L}{P_R}\right) + 1} \right) \quad (12)$$

- We now introduce Mach Numbers

$$M_R = u_R / a_R \quad M_{shock} = S / a_R$$

- Which would leads to:

$$M_R - M_{shock} = \frac{u_R - S}{a_R} \quad (13)$$

- By manipulating eqns: (7), (11) and (13):

$$(u_R - S)^2 = \frac{\rho_L}{\rho_R} \left( \frac{P_R - P_L}{\rho_R - \rho_L} \right)$$

$$\frac{\rho_L}{\rho_R} = \frac{\left( \frac{P_L}{P_R} \right) + \left( \frac{\gamma-1}{\gamma+1} \right)}{\left( \frac{\gamma-1}{\gamma+1} \right) \left( \frac{P_L}{P_R} \right) + 1}$$

$$M_R - M_{shock} = \frac{u_R - S}{a_R}$$

- theory leads us to the density and pressure ratios across the shock as functions of the relative Mach Number  $M_R - M_{Shock}$ , namely

$$\frac{\rho_L}{\rho_R} = \frac{(\gamma + 1)(M_R - M_{shock})^2}{(\gamma - 1)(M_R - M_{shock})^2 + 2} \quad (14)$$

$$\frac{P_L}{P_R} = \frac{2\gamma(M_R - M_{shock})^2 - (\gamma - 1)}{\gamma + 1} \quad (15)$$

- from (15) we can notice the following relation:

$$M_R - M_{shock} = -\sqrt{\left(\frac{\gamma+1}{2\gamma}\right) \left(\frac{P_L}{P_R}\right) + \left(\frac{\gamma-1}{2\gamma}\right)} \quad (16)$$

- which leads to an expression for the shock speed as a function of the pressure ratio across the shock, namely

$$S = u_R + a_R \sqrt{\left(\frac{\gamma+1}{2\gamma}\right) \left(\frac{P_L}{P_R}\right) + \left(\frac{\gamma-1}{2\gamma}\right)} \quad (17)$$

- Is important to notice that at  $P_L/P_R$  approaches unity it approaches the characteristic speed  $\lambda_+ = u_R + a_R$  as expected.

- we can relate the  $u_L$  and  $u_R$  by using (4)

$$\rho_L(u_L - S) = \rho_R(u_R - S) \quad (18)$$

- which leads to an expression for  $u_L$ , namely

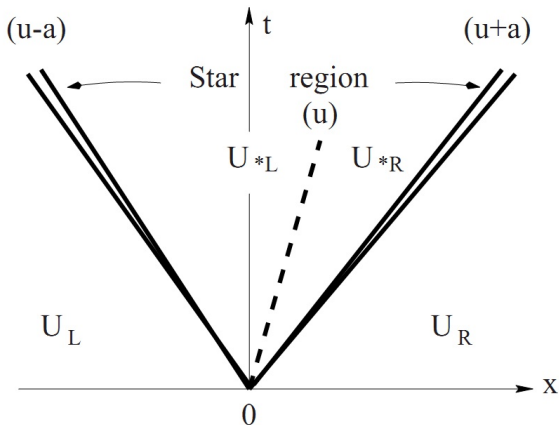
$$u_L = \left(1 - \frac{\rho_R}{\rho_L}\right) S + u_R \left(\frac{\rho_R}{\rho_L}\right) \quad (19)$$

- which leads to today's last relation:

$$u_L = u_R + \frac{a_R}{\gamma} \frac{\frac{P_L}{P_R} - 1}{\sqrt{\left(\frac{\gamma+1}{2\gamma}\right) \left(\frac{P_L}{P_R}\right) + \left(\frac{\gamma-1}{2\gamma}\right)}} \quad (20)$$



- Shockwaves to the left or to the right?



- Some literature like to use the following notation for the Right running Shockwave:

$$\frac{\rho_*}{\rho_R} = \frac{\left(\frac{P_*}{P_R}\right) + \left(\frac{\gamma-1}{\gamma+1}\right)}{\left(\frac{\gamma-1}{\gamma+1}\right) \left(\frac{P_*}{P_R}\right) + 1}$$

$$\frac{\rho_*}{\rho_R} = \frac{(\gamma+1)(M_R - M_{shock})^2}{(\gamma-1)(M_R - M_{shock})^2 + 2}$$

$$\frac{P_*}{P_R} = \frac{2\gamma(M_R - M_{shock})^2 - (\gamma-1)}{\gamma+1}$$

$$u_* = u_R + \frac{a_R}{\gamma} \frac{\frac{P_*}{P_R} - 1}{\sqrt{\left(\frac{\gamma+1}{2\gamma}\right) \left(\frac{P_*}{P_R}\right) + \left(\frac{\gamma-1}{2\gamma}\right)}}$$

- For the left running Shockwave, the equations are totally analogous:

$$\frac{\rho_*}{\rho_L} = \frac{\left(\frac{P_*}{P_L}\right) + \left(\frac{\gamma-1}{\gamma+1}\right)}{\left(\frac{\gamma-1}{\gamma+1}\right) \left(\frac{P_*}{P_L}\right) + 1}$$

$$\frac{\rho_*}{\rho_L} = \frac{(\gamma+1)(M_R - M_{shock})^2}{(\gamma-1)(M_R - M_{shock})^2 + 2}$$

$$\frac{P_*}{P_L} = \frac{2\gamma(M_L - M_{shock})^2 - (\gamma-1)}{\gamma+1}$$

$$u_* = u_L + \frac{a_L}{\gamma} \frac{\frac{P_*}{P_L} - 1}{\sqrt{\left(\frac{\gamma+1}{2\gamma}\right) \left(\frac{P_*}{P_L}\right) + \left(\frac{\gamma-1}{2\gamma}\right)}}$$

- Equations (14), (15) and (20) define a shock for given initial conditions  $(\rho_R, u_R, P_R)^T$  ahead of the shock and a chosen mach number  $M_s$  or equivalently a shock Speed  $S$ .

$$\frac{\rho_L}{\rho_R} = \frac{(\gamma + 1)(M_R - M_{shock})^2}{(\gamma - 1)(M_R - M_{shock})^2 + 2}$$

$$\frac{P_L}{P_R} = \frac{2\gamma(M_R - M_{shock})^2 - (\gamma - 1)}{\gamma + 1}$$

$$u_L = u_R + \frac{a_R}{\gamma} \frac{\frac{P_L}{P_R} - 1}{\sqrt{\left(\frac{\gamma+1}{2\gamma}\right) \left(\frac{P_L}{P_R}\right) + \left(\frac{\gamma-1}{2\gamma}\right)}}$$