

Lecture 17

Rigid body in space: Stability; Poinsot construction; Effect of energy dissipation.

6 - 19 October, 2021

Stability

Stability of rigid body ($I_1 \neq I_2 \neq I_3$) in space in *pure* spin about a principal axis of inertia.

Definition: System is *stable* if *small perturbations* of the motion remain *small* for all time.

I. BFCS is principal CS $\{\mathcal{E}, G, \hat{\mathbf{e}}_i\}$ of \mathbf{I}^G .

II. *Base state:* Body spins about $\hat{\mathbf{e}}_3$ at Ω_0 .

III. *Small perturbations at $t = 0$:*

$$\Omega_0 \hat{\mathbf{e}}_3 \xrightarrow{\text{pert.}} \omega_1^0 \hat{\mathbf{e}}_1 + \omega_2^0 \hat{\mathbf{e}}_2 + (\Omega_0 + \omega_3^0) \hat{\mathbf{e}}_3, \quad |\omega_i^0| \ll |\Omega_0|$$

$$\text{IV. Linearized AMB}_{/G}: \begin{cases} 0 = I_1 \dot{\omega}_1 - (I_2 - I_3) \omega_2 \Omega_0 \\ 0 = I_2 \dot{\omega}_2 - (I_3 - I_1) \omega_1 \Omega_0 \\ 0 = I_3 \dot{\omega}_3 \end{cases}$$

V. Spin about major / minor axis is *stable*:

$$\omega_1 = \omega_{01} \cos(\nu t + \phi), \quad \omega_2 = \omega_{02} \sin(\nu t + \phi)$$

VI. Spin about intermediate axis is *unstable*:

$$\omega_{1,2} = \omega_{01,02}^+ \exp(\nu t) + \omega_{01,02}^- \exp(-\nu t)$$

Application



- I. Perturbation to pure spin about axes of *maximum* and *minimum* inertia remain small because these axes are **stable** axes of rotation.
- II. Pure spin about the *intermediate* axis is **unstable**, hence small initial perturbations magnify with time leading to tumbling motion.

Poinsot construction

I. Rigid body in space with $I_1 > I_2 > I_3$:

1. $\mathbf{h}^G = H \hat{\mathbf{E}}_3 = \text{const.}; E_k = T = \text{const.}$

2. BFCs $\{\mathcal{E}, G, \hat{\mathbf{e}}_i\}$ is principal CS of \mathbf{I}^G

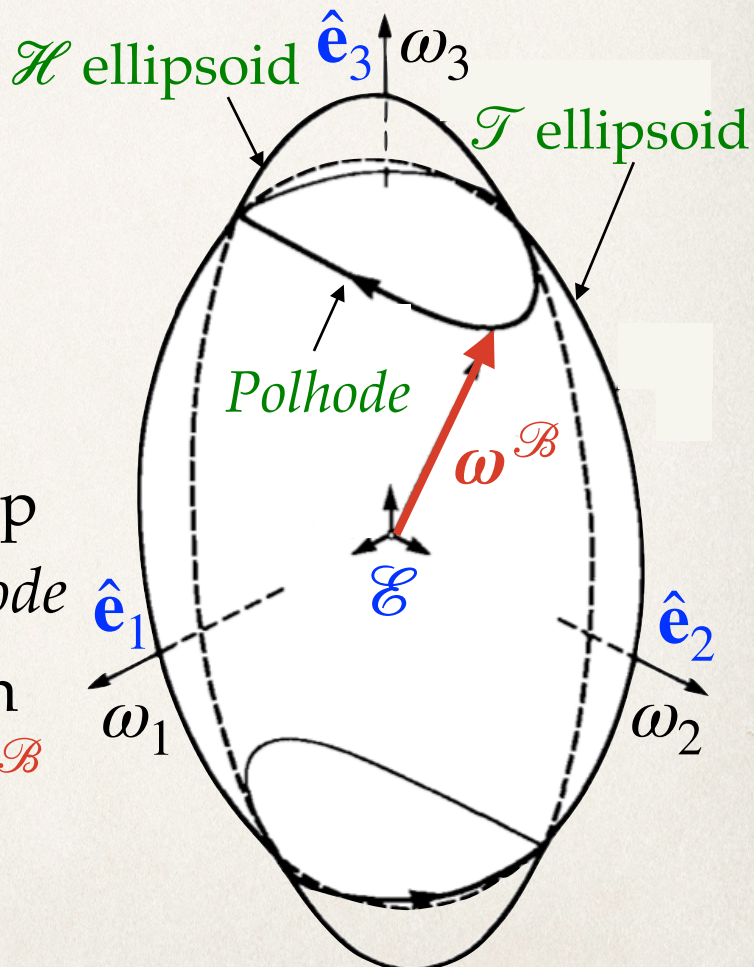
II. \mathcal{H} ellipsoid: $I_1^2 \omega_1^2 + I_2^2 \omega_2^2 + I_3^2 \omega_3^2 = H^2$.

III. \mathcal{T} ellipsoid: $I_1 \omega_1^2 + I_2 \omega_2^2 + I_3 \omega_3^2 = 2T$.

IV. Construct \mathcal{H} and \mathcal{T} ellipsoids in \mathcal{E} .

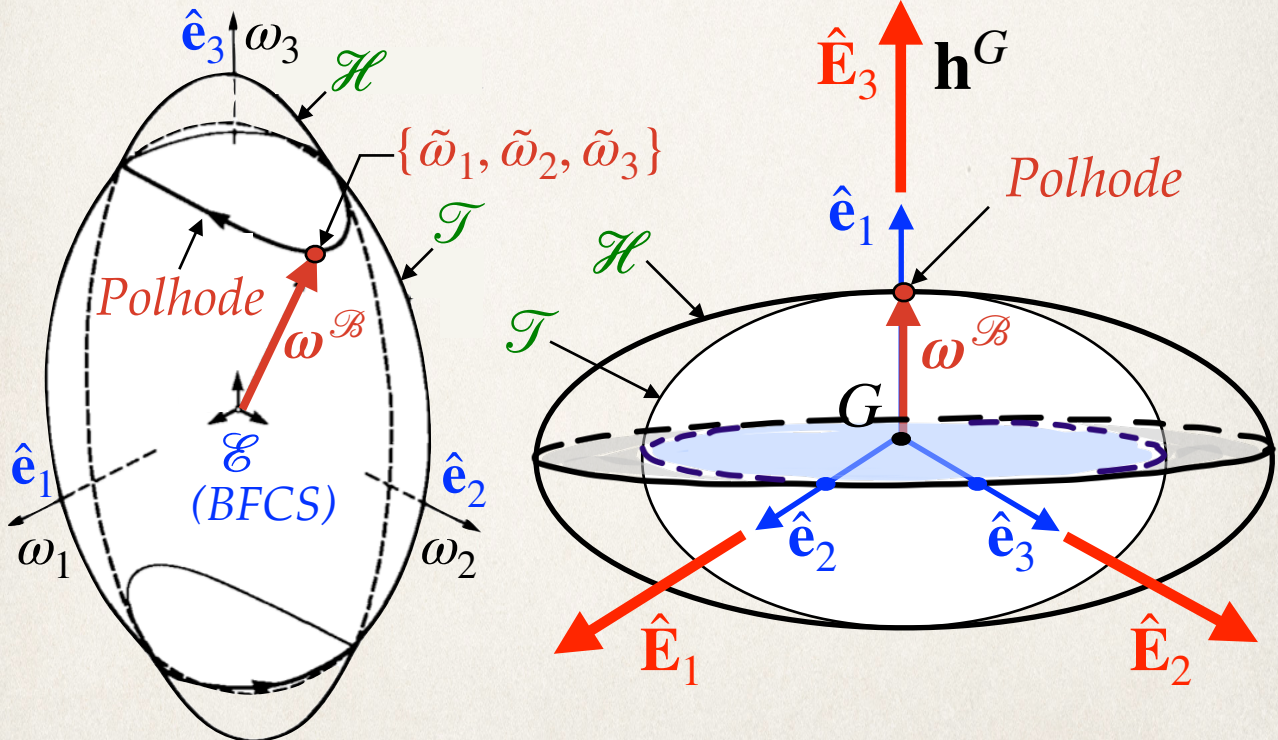
V. \mathcal{H} and \mathcal{T} ellipsoids will intersect along curves called *polhodes*.

1. $\boldsymbol{\omega}^{\mathcal{B}}$ vector's tip follows a *polhode*
2. $\boldsymbol{\omega}^{\mathcal{B}}$ changes in general $\Rightarrow \boldsymbol{\alpha}^{\mathcal{B}} \Rightarrow \text{Tumbling!}$



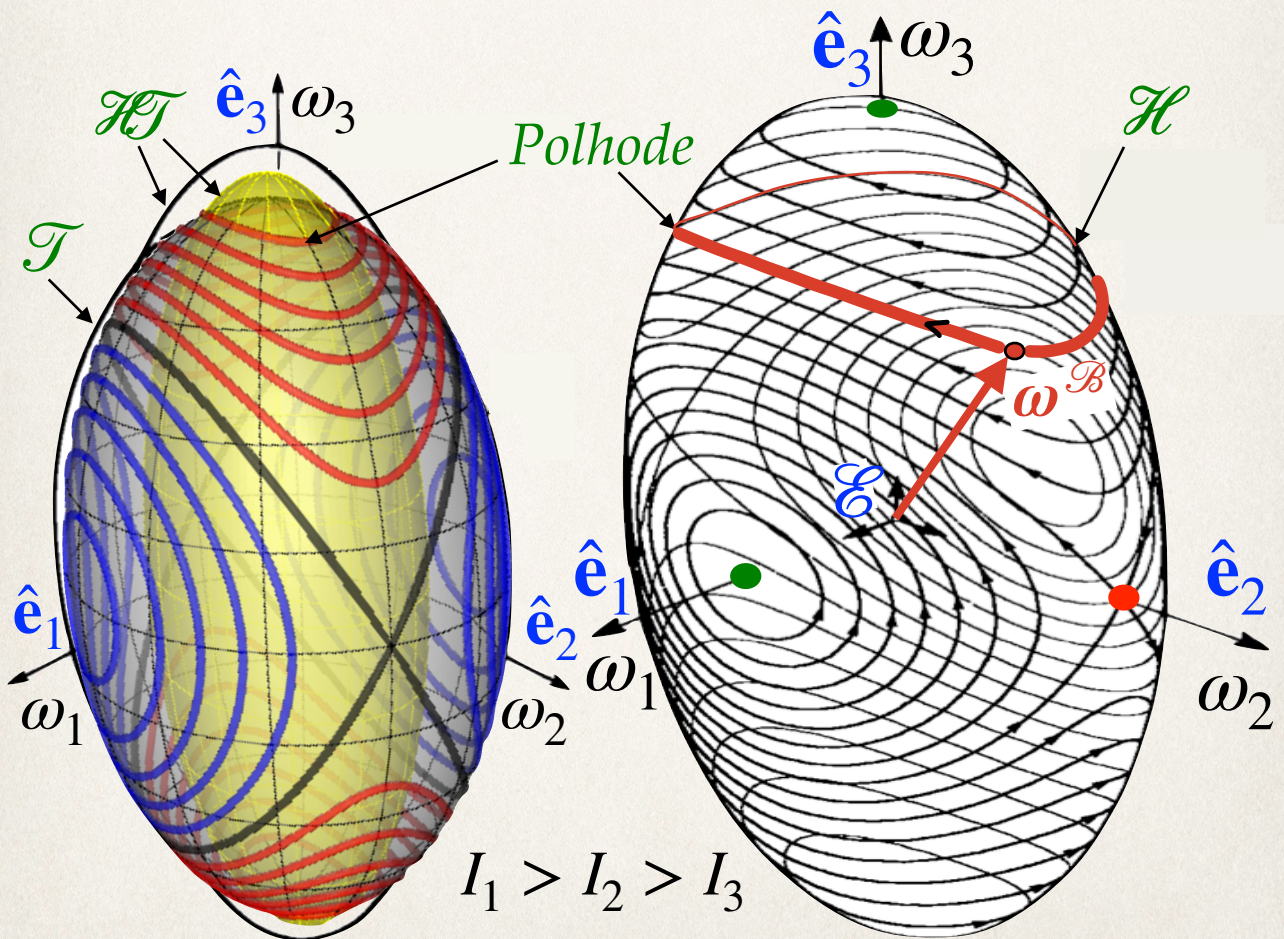
Poinsot construction

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- Diagram illustrating a crystal lattice point G (black dot) with three red arrows representing unit vectors \hat{E}_1 , \hat{E}_2 , and \hat{E}_3 . A black vector \mathbf{h}^G points upwards from G . A red arc labeled ϵ_0 is shown between the vertical axis and the \hat{E}_2 vector.



Poinsot construction

- I. On \mathcal{H} : Family of *polhodes* generated by T .
- II. *Polhode* fixed by $|\mathbf{h}^G| = H$ and $E_k = T$.
- III. *Polhodes* can not intersect.
- IV. For a given rigid body H and T are known \implies *Polhode* fixed
 1. $\implies \boldsymbol{\omega}^{\mathcal{B}}$ can only move on this *polhode*.



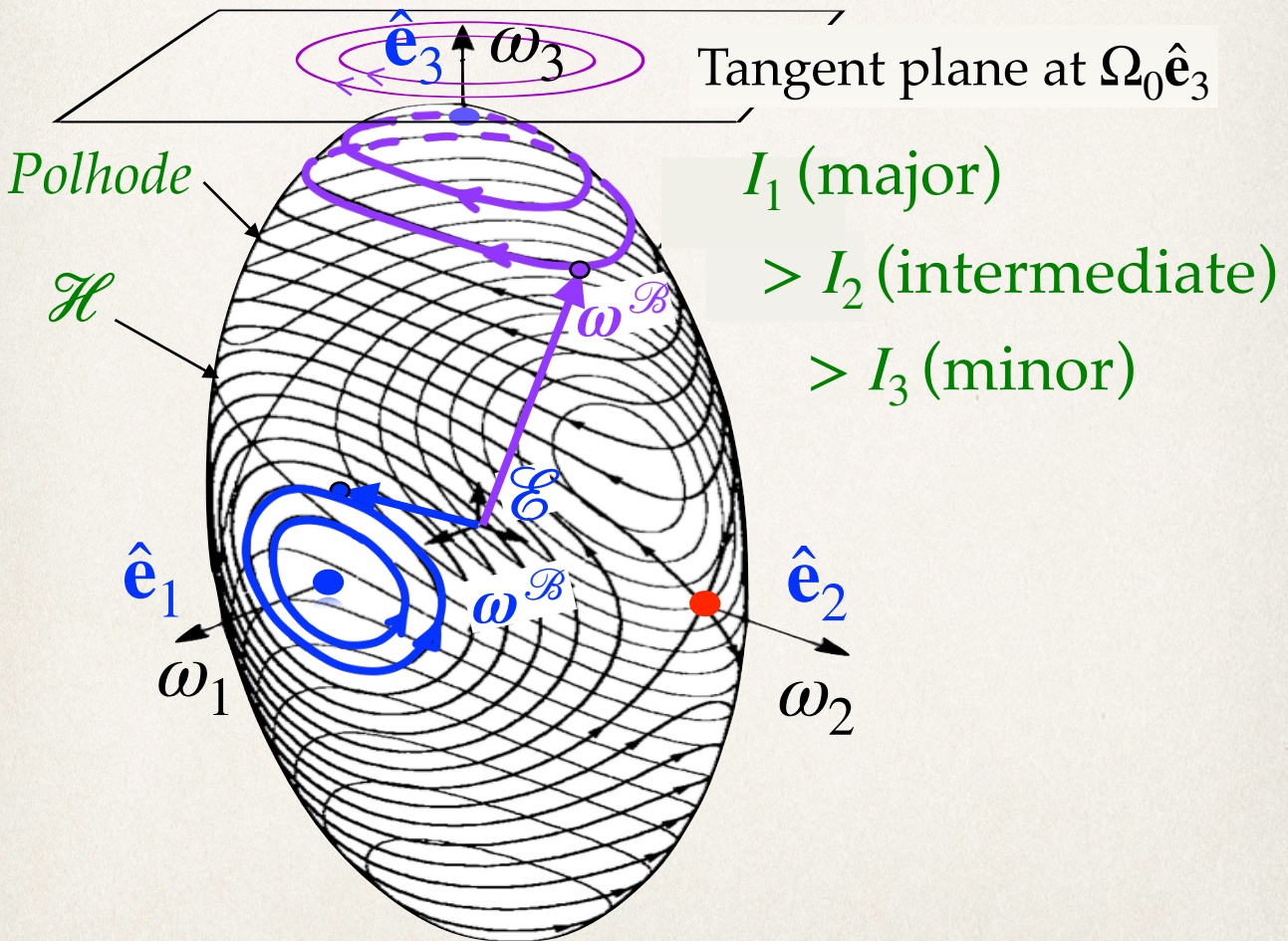
Perturbed motion

1. *Base state*: Pure spin about $\hat{\mathbf{e}}_1$ or $\hat{\mathbf{e}}_3 \Rightarrow$ **Stable**.
2. **Perturbed motion**. $\omega^{\mathcal{B}}$ moves on *polhodes* that are *centered around base state*.
3. *Small perturbations*: *Polhodes* \approx *Ellipses*.

i. **Example**. *Base state* is spin at Ω_0 about $\hat{\mathbf{e}}_3$.

Perturbed motion: $\omega_3 = \Omega_0$,

$$\omega_1 = \omega_{01} \cos(\nu t + \phi), \quad \omega_2 = \omega_{02} \sin(\nu t + \phi),$$

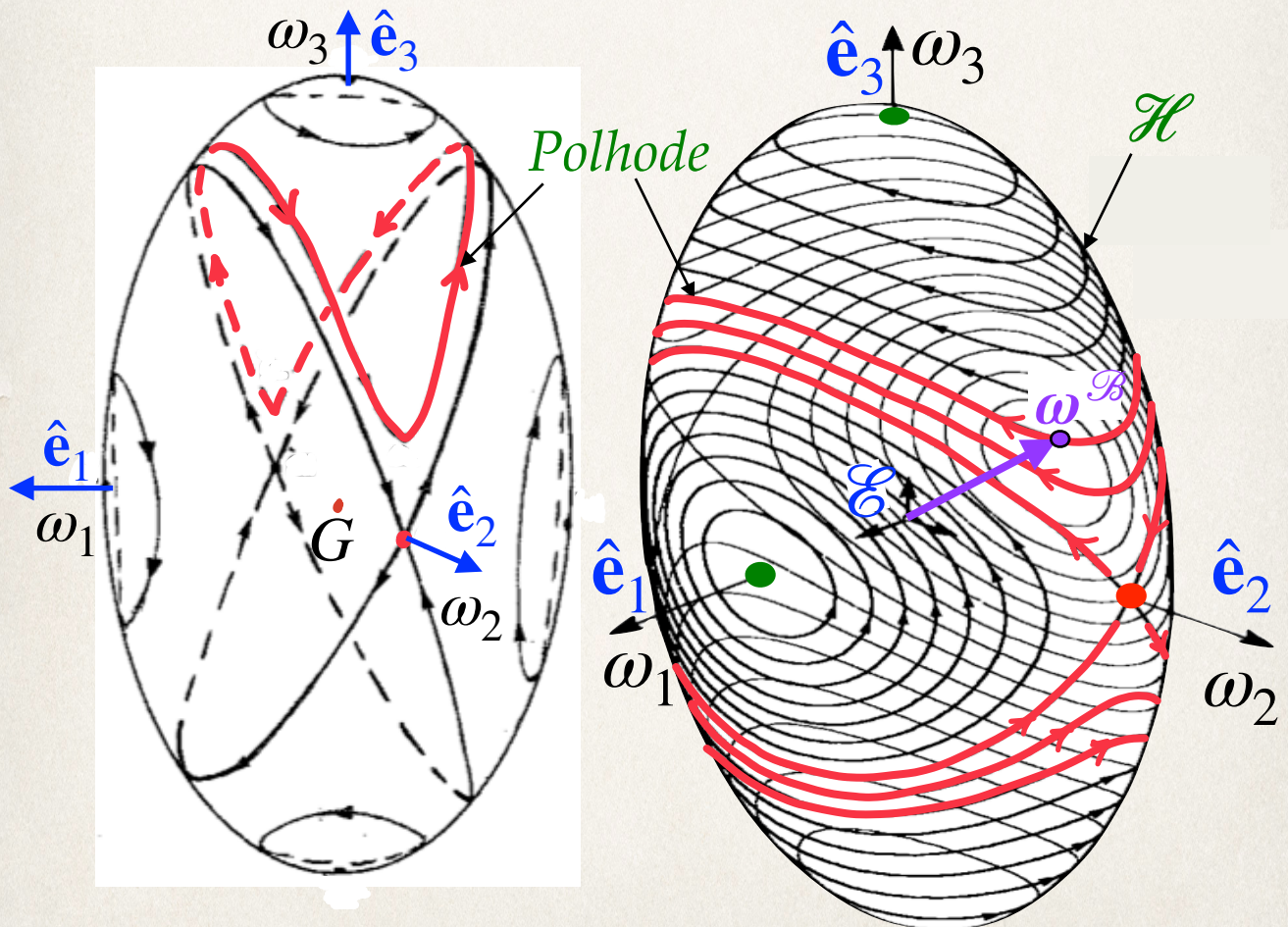


Perturbed motion

1. *Base state*: Pure spin about $\hat{\mathbf{e}}_2 \Rightarrow$ **Unstable**.
2. *Perturbed motion* deviates from pure spin.
3. *Polhodes* near $\hat{\mathbf{e}}_2$ go far away from $\hat{\mathbf{e}}_2$. Thus, $\omega^{\mathcal{B}}$ different from pure spin \Rightarrow **Tumbling!**
4. *Base state* is a **Saddle Point**. *Perturbed motion*:

$$\omega_{1,3} = \omega_{01,03}^+ \exp(\nu t) + \omega_{01,03}^- \exp(-\nu t)$$

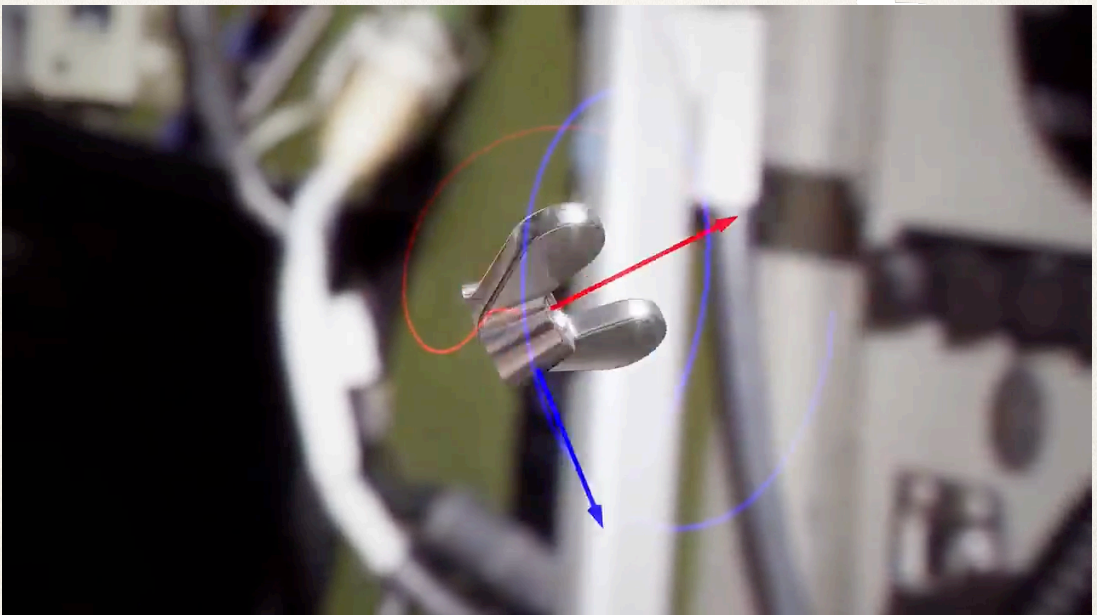
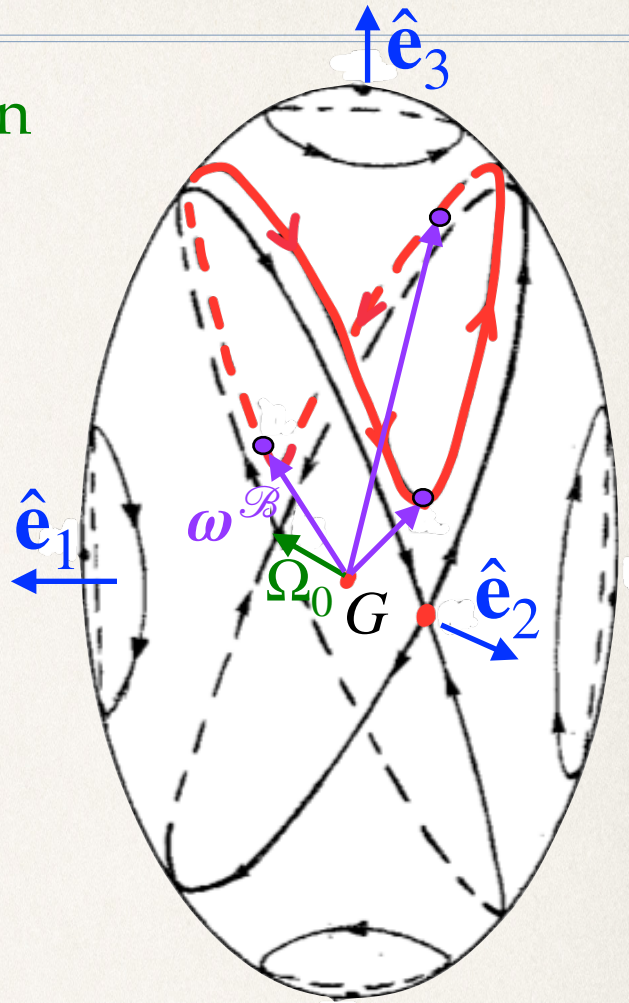
$$I_1 (\text{major}) > I_2 (\text{intermediate}) > I_3 (\text{minor})$$



Application

I. *Base state: Pure spin*
 Ω_0 along intermediate axis $-\hat{\mathbf{e}}_2$.

II. When *perturbed*,
 $\omega^{\mathcal{B}}$ can traverse a
polhode in which
 $\omega^{\mathcal{B}}$ points nearly
along $-\hat{\mathbf{e}}_2$ to
aligning along $\hat{\mathbf{e}}_2$!



Energy dissipation

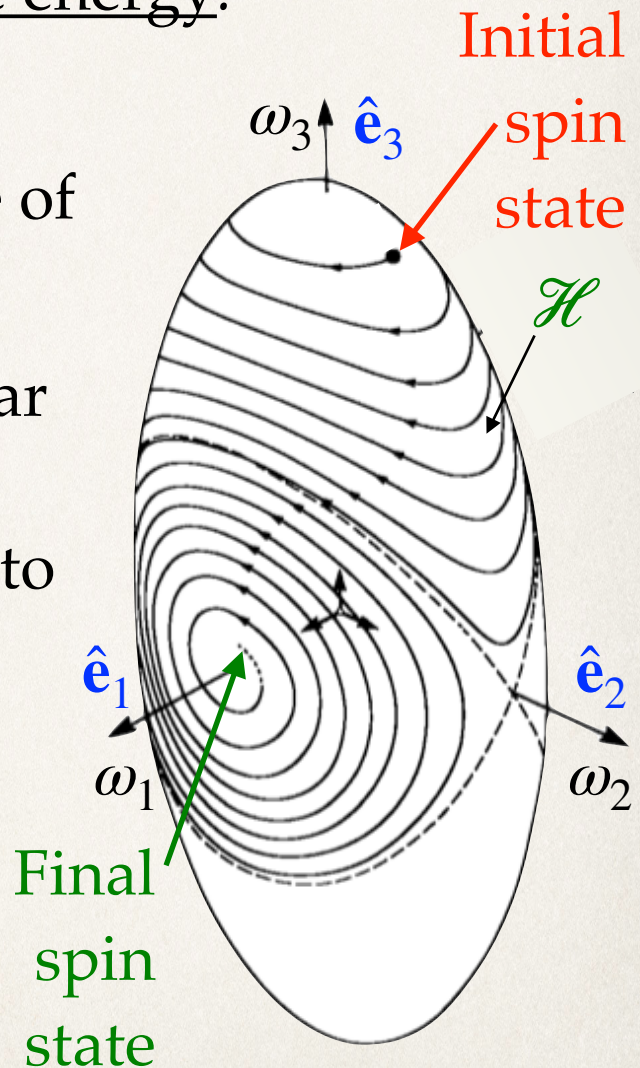
- I. Consider rigid body ($I_1 > I_2 > I_3$).
- II. For fixed H , *pure* rotation about
 1. $\hat{\mathbf{e}}_1$ has *minimum* KE: $T_{\min} = H^2/I_1$;
 2. $\hat{\mathbf{e}}_3$ has *minimum* KE: $T_{\max} = H^2/I_3$.

III. All bodies dissipate energy.

IV. **Energy dissipation**
drives body to state of
minimum energy.

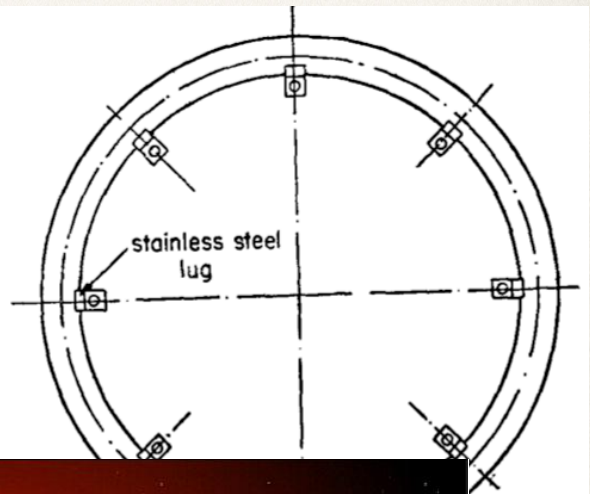
V. Body rotating in near
pure spin about $\hat{\mathbf{e}}_3$
(high KE) is driven to
rotate near $\hat{\mathbf{e}}_1$ (low
KE).

VI. This is called
nutational damping.



Application

- I. **Asteroids.** Explains why many are in pure spin about axis of maximum inertia.
- II. **Satellite nutation dampers.** To keep orientation of a satellite fixed.
 1. *Challenging engineering:* Want maximum damping with space and weight constraints.



The stabilisation system

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Abstract. The attitude stabilisation of *Aryabhata* was accomplished by spinning it about its axis of maximum moment of inertia. The spin stabilisation ensures satisfactory thermal control, uniform power generation through the body mounted solar panels and the scan capability for the scientific payloads. To bring down the nutation of the spinning spacecraft to a value well within the specified limits, a fluid-in-tube damper was also provided.

Proc. Indian Acad. Sci. Vol. C 1, No. 2, September 1978, pp. 135-143.



THERE ARE SATELLITES WATCHING ME.
ALL THE TIME. ALL MACHINES ARE REPORTING
TO THE WATCHERS. THE WATCHERS ARE
ALWAYS WATCHING!

RIGHT. OF COURSE. BUT IS
ANYTHING BOTHERING YOU?



It's getting harder and harder to diagnose paranoia.