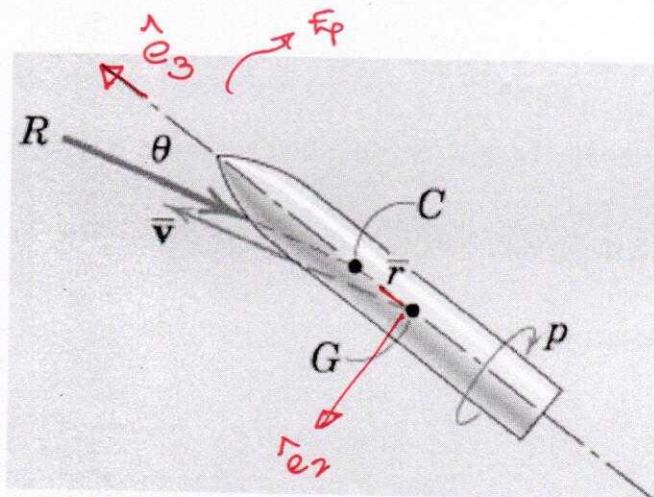


Problem 2

To find spin 'p' such that $\dot{\theta} = 0$.



Defining a BFCS $\{E, G, \hat{e}_i\}$

such that \hat{e}_3 is along symmetry axis. Assuming

that at this instant E_p is aligned with E_0 , the observer CS $\{E, G, \hat{e}_i\}$

KINEMATICS

Choosing the 3-1-3 Euler angle sequence, the kinematics of the motion of the rocket can be defined as per equations developed in Lecture 18, viz.

if $\underline{\omega}^B = \dot{\phi} \hat{E}_3 + \dot{\theta} \hat{e}_1' + \dot{\psi} \hat{e}_3$ which can then be written as

$$\omega_1 = \dot{\phi} \sin \theta \sin \psi + \dot{\theta} \cos \psi$$

$$\omega_2 = \dot{\phi} \sin \theta \cos \psi - \dot{\theta} \sin \psi$$

$$\omega_3 = \dot{\phi} \cos \theta + \dot{\psi}$$

and the angular accelerations are

$$\alpha_1 = \frac{d}{dt} (\omega_1) = \frac{d}{dt} (\dot{\phi} \sin \theta \sin \psi + \dot{\theta} \cos \psi)$$

$$\alpha_1 = \ddot{\phi} \sin \theta \sin \psi + \dot{\phi} \dot{\theta} \cos \theta \sin \psi + \dot{\phi} \sin \theta \dot{\psi} \cos \psi + \ddot{\theta} \cos \psi + \dot{\theta} \dot{\psi} (-\sin \psi)$$

$$\alpha_2 = \ddot{\phi} \sin \theta \cos \psi + \dot{\phi} \dot{\theta} \cos \theta \cos \psi - \dot{\phi} \dot{\psi} \sin \theta \sin \psi - \ddot{\theta} \sin \psi - \dot{\theta} \dot{\psi} \cos \psi$$

$$\alpha_3 = \ddot{\phi} \cos \theta - \dot{\phi} \dot{\theta} \sin \theta + \ddot{\psi}$$

(2)

Since we wish to find the stabilized solution,
 $\ddot{\phi} = 0$, $\ddot{\psi} = 0$, $\ddot{\theta} = 0$, given that $\dot{\theta} = 0$.

Moreover since ψ is the spin angle (about \hat{e}_3) the symmetry axis, $\psi = 0$ when BFCs and observer CS are aligned (or ψ does not affect the solution).

Applying the above values, we get,

$$\omega_1 = 0 \quad (\because \sin\psi = 0, \dot{\theta} = 0)$$

$$\omega_2 = \dot{\phi} \sin\theta \quad (\because \cos\psi = 1, \dot{\theta} = 0)$$

$$\omega_3 = \dot{\phi} \cos\theta + \dot{\psi}$$

and

$$\dot{\omega}_1 = \alpha_1 = \dot{\phi} \dot{\psi} \sin\theta$$

$$\dot{\omega}_2 = \alpha_2 = 0$$

$$\dot{\omega}_3 = \alpha_3 = 0$$

Now let's look at the kinetics.

Using the Euler's equations

$$M_1 = I_1 \dot{\omega}_1 - (I_2 - I_3) \omega_2 \omega_3$$

$$M_2 = I_2 \dot{\omega}_2 - (I_3 - I_1) \omega_3 \omega_1$$

$$M_3 = I_3 \dot{\omega}_3 - (I_1 - I_2) \omega_1 \omega_2$$

The applied moments are

$$M_1 = R \bar{H} \sin \alpha$$

$$M_2 = 0$$

$$M_3 = 0$$

Thus the Euler eqns give (recognizing $I_1 = I_2 = I_0$
& $I_3 = I$)

M_1 eqn $R \bar{H} \sin \alpha = I_0 \dot{\phi} \dot{\psi} \sin \alpha - (I_0 - I) \dot{\phi} \sin \alpha (\dot{\phi} \cos \alpha + \dot{\psi})$

M_2 Eqn $0 = I_0 (0) - (I - I_0) (\dot{\phi} \cos \alpha + \dot{\psi}) (0) = 0$
 \Rightarrow identically satisfied ✓

M_3 eqn $0 = I (0) - (I_0 - I_0) (0) (\dot{\phi} \sin \alpha) = 0$
 \Rightarrow identically satisfied. ✓

The M_1 eqn gives.

$$R \bar{H} \sin \alpha = I_0 \dot{\phi} \dot{\psi} \sin \alpha - (I_0 - I) (\dot{\phi}^2 \sin \alpha \cos \alpha + \dot{\phi} \dot{\psi} \sin \alpha)$$

$$\Rightarrow R \bar{H} = I_0 \dot{\phi} \dot{\psi} + (I - I_0) \dot{\phi}^2 \cos \alpha + (I - I_0) \dot{\phi} \dot{\psi}$$

$$= I \dot{\phi} \dot{\psi} + (I - I_0) \dot{\phi}^2 \cos \alpha$$

$$\Rightarrow [(I - I_0) \cos \alpha] \dot{\phi}^2 + [I \dot{\psi}] \dot{\phi} - R \bar{H} = 0$$

Note. $\dot{\psi}$ is the spin rate 'p' & $\dot{\phi}$ = precession rate = $\dot{\phi}$

$$\Rightarrow [(I - I_0) \cos \alpha] \dot{\phi}^2 + [I p] \dot{\phi} - R \bar{H} = 0$$

(4)

We thus have a quadratic equation for the precession rate in terms of the spin rate, 'p'

$$\dot{\phi} = \frac{-I_p \pm \sqrt{(I_p)^2 - 4[(I - I_0) \cos \alpha] [-R \hbar]}}{2(I - I_0) \cos \alpha}$$

For a feasible solution, term within the square root should be ≥ 0 , i.e.,

$$(I_p)^2 - 4[(I - I_0) \cos \alpha] [-R \hbar] \geq 0$$

Noting $I_0 > I$ (for the socket), the condition for 'p' is

$$p \geq \frac{2}{I} \sqrt{(I_0 - I) R \hbar \cos \alpha}$$