

Miscellanea

$$\left| \frac{\partial(x, y, z)}{\partial(r, \theta, \phi)} \right| = r^2 \sin \theta$$
$$\left| \frac{\partial(x, y)}{\partial(r, \theta)} \right| = \left| \frac{\partial(x, y, z)}{\partial(r, \phi, z)} \right| = r$$
$$\int \frac{dx}{x} = \ln |x| + C$$

$$\sin \theta \xrightarrow{\frac{d}{d\theta}} \cos \theta \rightarrow -\sin \theta$$

Trigonometry

$$\frac{\sin a}{A} = \frac{\sin b}{B} = \frac{\sin c}{C} \quad (\text{law of sines})$$

$$C = \sqrt{A^2 + B^2 - 2AB \cos(c)} \quad (\text{law of cosines})$$

Identities

$$\sin^2 \theta + \cos^2 \theta = 1$$

$$\sin(\theta \pm \frac{\pi}{2}) = \pm \cos \theta$$

$$\sin(\alpha \pm \beta) = \sin \alpha \cos \beta \pm \cos \alpha \sin \beta$$

$$\cos(\alpha \pm \beta) = \cos \alpha \cos \beta \mp \sin \alpha \sin \beta$$

$$\sin^2 \theta = \frac{1 - \cos(2\theta)}{2}$$

$$\cos^2 \theta = \frac{1 + \cos(2\theta)}{2}$$

$$\sin(2\theta) = 2 \sin \theta \cos \theta$$

$$\begin{aligned} \cos(2\theta) &= \cos^2 \theta - \sin^2 \theta \\ &= 2 \cos^2 \theta - 1 \\ &= 1 - 2 \sin^2 \theta \end{aligned}$$

Derivatives

$$\frac{d(uv)}{dx} = u'v + uv'$$

$$\frac{d(u/v)}{dx} = \frac{u'v - uv'}{v^2}$$

$$\frac{d(u(v))}{dx} = u'(v)v'$$

Kinematics

$$\vec{v}(t) = \vec{v}_0 + \vec{a}t \quad (\text{velocity})$$

$$\vec{\omega}(t) = \vec{\omega}_0 + \vec{\alpha}t \quad (\text{rot. form})$$

$$\vec{x}(t) = \vec{x}_0 + \vec{v}t + \frac{1}{2}\vec{a}t^2 \quad (\text{displacement})$$

$$\vec{\theta}(t) = \vec{\theta}_0 + \vec{\omega}t + \frac{1}{2}\vec{\alpha}t^2 \quad (\text{rot. form})$$

$$v_f = \sqrt{v_i^2 + 2a\Delta x} \quad (v_f)$$

Momentum

$$\vec{p} = m\vec{v} \quad (\text{momentum})$$

$$\vec{\ell} = I\vec{\omega} = \vec{r} \times \vec{p} \quad (\text{rot. form})$$

$$\dot{\vec{P}} = \vec{F}_{\text{ext}} \quad (\text{3rd law})$$

$$\dot{\vec{\ell}} = \vec{\tau} = I\vec{\alpha} = \vec{r} \times \vec{F} \quad (\text{rot. form})$$

$$\begin{aligned} \Delta \vec{P} &= \vec{P}_f - \vec{P}_i = \vec{J} \\ &= \int \vec{F} dt \\ &= \vec{F}_{\text{avg}} \Delta t \quad (\text{impulse}) \end{aligned}$$

Total Angular Momentum

$$\begin{aligned} \vec{L} &= \sum_{i=1}^N \vec{\ell}_i \\ &= \vec{R}_{\text{CM}} \times \vec{P}_{\text{CM}} \quad (\vec{L}_{\text{orb}} \text{ of CM about axis}) \\ &+ \sum_{i=1}^N r'_i \times \vec{p}'_i \quad (\vec{L}_{\text{spin}} \text{ of body about CM}) \end{aligned}$$

Collisions

An *elastic* collision conserves kinetic energy.

$$m_1 \vec{v}_1 + m_2 \vec{v}_2 = m_1 \vec{v}_1' + m_2 \vec{v}_2'$$

An *inelastic* collision does not.

$$m_1 \vec{v}_1 + m_2 \vec{v}_2 = m \vec{v}'$$

Changing Mass

Rockets

$$\vec{F}_{\text{th}} = -\dot{m} \vec{v}_{\text{ex}} \quad (\text{thrust})$$

$$\vec{v}(t) = \vec{v}_0 + \vec{v}_{\text{ex}} \ln \left(\frac{m_0}{m(t)} \right)_{(\Delta V)}$$

Center of Mass

$$\vec{R} = \frac{1}{M} \sum_{i=1}^N m_i \vec{r}_i \quad (\text{CM summation})$$

$$M = \int \rho dV \quad (\text{total mass})$$

$$\vec{R} = \frac{1}{M} \int \vec{r} dm \quad (\text{CM integral})$$

dm can usually be rewritten in terms of dV

$$\vec{R} = \frac{1}{M} \int \vec{r} \rho(\vec{r}) dV$$

Look for symmetries by which you can substitute \vec{R} and \vec{r} for scalars.

Moment of Inertia

$$\sum_{i=1}^N = m_i r_{\perp}^2 \quad (\text{summation})$$

$$I = \int r_{\perp}^2 dm \quad (\text{integral})$$

Parallel Axis Theorem

$$I = I_{\text{CM}} + Md^2$$

where $I \parallel I_{\text{CM}}$ and d is the distance between I & I_{CM}

Perpendicular Axis Theorem

For *planar lamina* (flat, plate-like objects)

$$I_z = I_x + I_y$$

Work & Energy

$$\begin{aligned} W &= \int_C \vec{F} \cdot d\vec{r} \\ &= \int_a^b \vec{F}(\vec{r}(t)) \cdot \vec{r}'(t) dt \\ &= \int_{x,y} \langle F_x, F_y \rangle \cdot \langle dx, dy \rangle \quad (\text{work integral}) \end{aligned}$$

$$W = \Delta T \quad (\text{Work-Energy Theorem})$$

$$\vec{F} = -\nabla U \quad (\text{potential from force})$$

Conservative Forces

A *conservative force* is one with a corresponding potential energy. This means:

- the force depends *only* on position
- the work done by that force is the same for *any* path $\implies \nabla \times \vec{F} = 0$

Lagrangian Mechanics

$$\mathcal{L} = T - U \quad (\text{Lagrangian})$$

$$\frac{d}{dt} \underbrace{\frac{\partial \mathcal{L}}{\partial \dot{q}_i}}_{\text{gnr. } p} = \underbrace{\frac{\partial \mathcal{L}}{\partial q_i}}_{\text{gnr. } F} \quad (\mathcal{L} \text{ eqn.})$$