

# viscosity.h Documentation

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## Section 1

### Introduction and Background

**viscosity.h** is constructed exclusively for NS solver 'centered.h' (see corresponding doc for more information.), whose purpose is to solve viscos equation implicitly using poisson equation solver built in 'poisson.h'. The governing equation is

$$\rho_{n+\frac{1}{2}} \left[ \frac{\mathbf{u}^* - \mathbf{u}'}{\Delta t} \right] = \nabla \cdot [2\mu_{n+\frac{1}{2}} \mathbf{D}^*] \quad (1)$$

where  $\mathbf{D}^* = [\nabla \mathbf{u}^* + (\nabla \mathbf{u}^*)^T]/2$ ,  $\mathbf{u}'$  is known variable and  $\mathbf{u}^*$  is the desired output.

Consider integral form of equation 1

$$\int_{\Omega} \rho_{n+\frac{1}{2}} \left[ \frac{\mathbf{u}^* - \mathbf{u}'}{\Delta t} \right] dV = \int_{\Omega} \nabla \cdot [2\mu_{n+\frac{1}{2}} \mathbf{D}^*] dV = \int_{\partial\Omega} [2\mu_{n+\frac{1}{2}} \mathbf{D}^*] \cdot \mathbf{n} dS \quad (2)$$

The discrete form of this equation reads (component form is presented instead of tensor for sake of convenience).

$$\Delta \rho_{n+\frac{1}{2}} \left[ \frac{u_i^* - u_i'}{\Delta t} \right] = \sum_{f=1}^6 (n_f \mu_{n+\frac{1}{2}} \left( \frac{\partial u_j^*}{\partial x_i} + \frac{\partial u_i^*}{\partial x_j} \right))_f \quad i = x, y, z \quad j = normal(f) \quad (3)$$

where  $f$  in summation and subscript represents surfaces of single cell,  $j$  represents coordinate component of face normal disregarding negative or positive direction,  $n_f$  indicates whether the face normal is consistent with positive direction of coordinate, taking 2D cell as an example

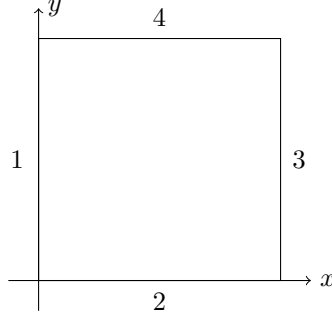


Figure 1: 2D cell example.

Unfold *R.H.S.* of equation 3 based on 2D cell depicted by figure 1 yields

$$R.H.S = [\mu_3(\frac{\partial u_x^*}{\partial x_i} + \frac{\partial u_i^*}{\partial x})_3 - \mu_1(\frac{\partial u_x^*}{\partial x_i} + \frac{\partial u_i^*}{\partial x})_1 + \mu_4(\frac{\partial u_y^*}{\partial x_i} + \frac{\partial u_i^*}{\partial y})_4 - \mu_2(\frac{\partial u_y^*}{\partial x_i} + \frac{\partial u_i^*}{\partial y})_2] \quad i = x, y \quad (4)$$

Constructing derivative is simple for aligned direction but cumbersome for splitting direction.

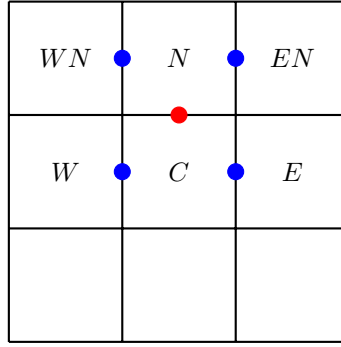


Figure 2: Sketch for derivative calculating.

Still take 2D cell as an example, as shown in figure 2, aligned direction derivative *e.g.*  $(\frac{\partial u_y^*}{\partial y})_n$  (lower case indicates corresponding cell face) can be obtained by simply central difference scheme

$$(\frac{\partial u_y^*}{\partial x_y})_n = \frac{u_N^* - u_C^*}{\Delta} \quad (5)$$

However, for splitting direction derivative *e.g.*  $(\frac{\partial u_y^*}{\partial x})_n$  highlighted by red point above, there is no direct method to calculate the desired but average the ambient derivatives (highlighted by blue point), which are obtained following the direct method.

$$(\frac{\partial u_y^*}{\partial x})_n = \frac{(\frac{\partial u_y^*}{\partial x})_{wn} + (\frac{\partial u_y^*}{\partial x})_{en} + (\frac{\partial u_y^*}{\partial x})_w + (\frac{\partial u_y^*}{\partial x})_e}{4} \quad (6)$$

Now consider constructing  $x$  component of equation 3 for cell locates  $(0,0)$ . We herein denote  $x, y$  components of  $\mathbf{u}$  as  $u, v$  but not  $u_x, u_y$  for sake of clarity. The *L.H.S* can be directly expressed as

$$L.H.S = \Delta \rho_{n+\frac{1}{2}} [\frac{u_{0,0}^* - u'_{0,0}}{\Delta t}] \quad (7)$$

Then the first two terms of equation 4 which can be obtained directly are

$$R.H.SF2 = 2\mu_{(\frac{1}{2},0)}\left(\frac{u_{1,0}^* - u_{0,0}^*}{\Delta}\right) - 2\mu_{(-\frac{1}{2},0)}\left(\frac{u_{0,0}^* - u_{-1,0}^*}{\Delta}\right) \quad (8)$$

Finally the last two terms that is calculated by averaging yields

$$\begin{aligned} R.H.SL2 = & \mu_{(0,\frac{1}{2})}\left[\frac{u_{(0,1)}^* - u_{(0,0)}^*}{\Delta} + \frac{v_{(1,1)}^* - v_{(-1,1)}^* + v_{(1,0)}^* - v_{(-1,0)}^*}{4\Delta}\right] \\ & - \mu_{(0,-\frac{1}{2})}\left[\frac{u_{(0,0)}^* - u_{(0,-1)}^*}{\Delta} + \frac{v_{(1,0)}^* - v_{(-1,0)}^* + v_{(1,-1)}^* - v_{(-1,-1)}^*}{4\Delta}\right] \end{aligned} \quad (9)$$

Rearrange the equations we have

$$u_{(0,0)}^* = \frac{\frac{\Delta t}{\rho}(2\mu_{(\frac{1}{2},0)}u_{(1,0)}^* + 2\mu_{(-\frac{1}{2},0)}u_{(-1,0)}^* + \mathcal{A} - \mathcal{B}) + \Delta^2 u'_{(0,0)}}{\Delta^2 + \frac{\Delta t}{\rho}(2\mu_{(\frac{1}{2},0)} + 2\mu_{(-\frac{1}{2},0)} + \mu_{(0,\frac{1}{2})} + \mu_{(0,-\frac{1}{2})})} \quad (10)$$

where

$$\mathcal{A} = \mu_{(0,\frac{1}{2})}(u_{(0,1)}^* + \frac{v_{(1,1)}^* + v_{(1,0)}^* - v_{(-1,1)}^* - v_{(-1,0)}^*}{4}) \quad (11)$$

$$\mathcal{B} = \mu_{(0,-\frac{1}{2})}(-u_{(0,-1)}^* + \frac{v_{(1,-1)}^* + v_{(1,0)}^* - v_{(-1,-1)}^* - v_{(-1,0)}^*}{4}) \quad (12)$$

Now, we obtain the implicit discrete linear expressions for desired valuable  $\mathbf{u}^*$ . Moreover, if we modify the governing equation into a more general form

$$\mathbf{u}' = \mathbf{u}^* - \frac{\Delta t \nabla \cdot [2\mu_{n+\frac{1}{2}} \mathbf{D}^*]}{\rho_{n+\frac{1}{2}}} = \mathcal{D}(\mathbf{u}^*) \quad (13)$$

here  $\mathcal{D}$  is a linear operator that satisfies

$$RES^k = \mathbf{u}' - \mathcal{D}(\mathbf{u}^{*,k}) \quad (14)$$

$$\mathcal{D}(\delta \mathbf{u}^{*,k}) = RES^k \quad (15)$$

$$\mathbf{u}^{*,k+1} = \mathbf{u}^{*,k} + \delta \mathbf{u}^{*,k} \quad (16)$$

We denote by  $k$  the  $k$ th iterate approximating result of corresponding value. Then the multicycle solver constructed in 'poisson.h' can be applied to solve viscosity equation with specific relaxation and residual functions, which are built in this headfile.

## Section 2

## Functions

### 2.1 Define

#### 2.1.1 Parameters

Name	Data type	Status	Option	Representation (before/after)
<i>mu</i>	face vector	unchanged	compulsory	$\mu^{n+\frac{1}{2}}$
<i>rho</i>	scalar	unchanged	compulsory	$\rho^{n+\frac{1}{2}}$
<i>dt</i>	double	unchanged	compulsory	$\Delta t$

#### 2.1.2 Program Workflow

##### Pre-Define

Define data structure Viscosity called in the total solver

$$\textcolor{red}{mu} = \mu^{n+\frac{1}{2}} \quad \textcolor{red}{rho} = \rho^{n+\frac{1}{2}}$$

$$\textcolor{red}{dt} = \Delta t$$

##### macro

macro defined here to set factor *lambda* especially for 'axi.h'. Otherwise, for cases in Cartesian mesh, the factor is set to be unity.

```

1  #include "poisson.h"
2
3  struct Viscosity {
4      face vector mu;
5      scalar rho;
6      double dt;
7  };
8
9  #if AXI
10 # define lambda ((coord){1., 1. + dt/rho[]*(mu.x[] + mu.x[1] + \
11 mu.y[] + mu.y[0,1])/2./sq(y)})
12 #else // not AXI
13 # if dimension == 1
14 #   define lambda ((coord){1.})
15 # elif dimension == 2
16 #   define lambda ((coord){1.,1.})

```

```

17 # elif dimension == 3
18 #   define lambda ((coord){1.,1.,1.})
19 #endif
20 #endif

```

## 2.2 relax\_viscosity

### 2.2.1 Parameters

Name	Data type	Status	Option	Representation (before/after)
<i>a</i>	scalar*	update	compulsory	$\delta \mathbf{u}^{*,k} / \delta \mathbf{u}^{*,k+1}$
<i>b</i>	scalar*	unchanged	compulsory	$RES$
<i>dt</i>	double	unchanged	compulsory	$\Delta t$
<i>l</i>	int	unchanged	compulsory	mesh level
<i>data</i>	struct Vviscosity	unchanged	compulsory	$\mu^{n+\frac{1}{2}}, \rho^{n+\frac{1}{2}}, \Delta t$

### 2.2.2 Worth Mentioning Details

The purpose of this function is to implement equation 10 and push forward iteration of G-S (or Jacobi) method for a single round on level  $l$  mesh. Note the object processed is  $\delta \mathbf{u}^*$  defined previously instead of  $\mathbf{u}^*$ . Therefore, the equation it solves is equation 15 with given residual  $RES$  computed in another function.

The Basilisk self-defined macro `foreach_level_or_leaf(1)` takes  $l$  as its parameter and traverses each cell on level  $l$  or leaf cell (the cell that cannot be divided) whose level is less than  $l$ . Interested readers are referred to 'poisson.h Documentation' or Van Hooft *et.al.*[1] for more information about mesh hierarchy in Basilisk.

### 2.2.3 Program Workflow

**Starting Point****input:**

$$a = \delta \mathbf{u}^* \quad b = RES$$

$$l = meshlevel$$

$$data = \mu^{n+\frac{1}{2}}, \rho^{n+\frac{1}{2}}, \Delta t$$

**data assignment** $\mu, \rho, \Delta t$  is extracted from *data**u* and *r* serve as pointer and point to  $\delta \mathbf{u}^*$ , *RES* respectively**iteration type**

The algorithm of iteration is selected by macro, if JACOBI

is false, then a pointer *w* is created and points to *u* (*a*).Then all modification hereinafter to *w* is direct exert on *u*.

```

1 static void relax_viscosity (scalar * a, scalar * b, int l, void * data)
2 {
3     struct Viscosity * p = (struct Viscosity *) data;
4     (const) face vector mu = p->mu;
5     (const) scalar rho = p->rho;
6     double dt = p->dt;
7     vector u = vector(a[0]), r = vector(b[0]);
8
9     #if JACOBI
10     vector w[];
11 #else
12     vector w = u;
13 #endif

```

**Relaxation Compute**

implement of equation 10 for different dimension.

$$w = \delta \mathbf{u}_{TBD}^{*,k+1}$$

```

1 foreach_level_or_leaf (1) {
2     foreach_dimension()
3     w.x[] = (dt/rho[]*(2.*mu.x[1]*u.x[1] + 2.*mu.x[]*u.x[-1]
4 #if dimension > 1
5     + mu.y[0,1]*(u.x[0,1] +
6     (u.y[1,0] + u.y[1,1])/4. -
7     (u.y[-1,0] + u.y[-1,1])/4.)
8     - mu.y[]*(- u.x[0,-1] +
9     (u.y[1,-1] + u.y[1,0])/4. -

```

```

10      (u.y[-1,-1] + u.y[-1,0])/4.)
11  #endif
12  #if dimension > 2
13      + mu.z[0,0,1]*(u.x[0,0,1] +
14      (u.z[1,0,0] + u.z[1,0,1])/4. -
15      (u.z[-1,0,0] + u.z[-1,0,1])/4.)
16      - mu.z[]*(- u.x[0,0,-1] +
17      (u.z[1,0,-1] + u.z[1,0,0])/4. -
18      (u.z[-1,0,-1] + u.z[-1,0,0])/4.)
19  #endif
20      ) + r.x[]*sq(Delta))/
21  (sq(Delta)*lambda.x + dt/rho[]*(2.*mu.x[1] + 2.*mu.x[]
22  #if dimension > 1
23      + mu.y[0,1] + mu.y[]
24  #endif
25  #if dimension > 2
26      + mu.z[0,0,1] + mu.z[]
27  #endif
28      ));
29  }

```



### Update

If *JACOBI* == *TRUE*  
 $\mathbf{u} = \delta \mathbf{u}^{*,k+1} = \frac{1}{3} \delta \mathbf{u}^{*,k} + \frac{2}{3} \delta \mathbf{u}_{TBD}^{*,k+1}$   
 else (*G* - *S* iteration)  
 $\mathbf{u} = \delta \mathbf{u}^{*,k+1} = \delta \mathbf{u}_{TBD}^{*,k+1}$   
*a* is modified along with *u*

```

1  #if JACOBI
2      foreach_level_or_leaf (1)
3          foreach_dimension()
4              u.x[] = (u.x[] + 2.*w.x[])/3.;
5  #endif
6
7  #if TRASH
8      vector u1[];
9      foreach_level_or_leaf (1)
10          foreach_dimension()
11              u1.x[] = u.x[];
12      trash ({u});
13      foreach_level_or_leaf (1)

```

```

14     foreach_dimension()
15         u.x[] = u1.x[];
16 #endif
17 }

```

## 2.3 residual\_viscosity

### 2.3.1 Parameters

Name	Data type	Status	Option	Representation (before/after)
<i>a</i>	scalar*	unchanged	compulsory	$\mathbf{u}^{*,k}$
<i>b</i>	scalar*	unchanged	compulsory	$\mathbf{u}'$
<i>resl</i>	scalar*	<b>output</b>	compulsory	$RES^k = \mathbf{u}' - \mathcal{D}(\mathbf{u}^{*,k})$
<i>data</i>	struct Vsicosity	unchanged	compulsory	$\mu^{n+\frac{1}{2}}, \rho^{n+\frac{1}{2}}, \Delta t$

### 2.3.2 Worth Mentioning Details

This function updates residual by solving equation 13 and returns the maximum component. Two computational methods are provided for tree and non-tree grid. The first method for tree grid costs more computational resource compared with the second one due to extra traversal or cell face but preserves 2nd order on tree grid. While the second method is of 2nd order with caesian but 1st with tree grid.

Consider quadtree for 2D cases shown in figure 3 where coarser grid is of level  $l$  and finer grid is of level  $l+1$ . ■ represents example targert cell whose residual will be compute to display the difference between two methods. ● represents active point located at cell center where stores value for each cell. ● and ● are ghost cells which do not exist initially. ● is calculated by bilinear interpolation and is called as adjacent cell for finer cell while ● is obtained by averaging and is called from coarser cell. Moreover, cell faces highlighted by letters is another issue that worth mentioning. As can be seen from figure, target cell has five marked faces instead of four that face  $A$  is divided into subfaces  $A_1$  and  $A_2$  due to tree grid. However when called face value from centered cell perspective *i.e.* invoke value on face  $A$  from target cell, the value returned is the average of  $A_1$  and  $A_2$ .



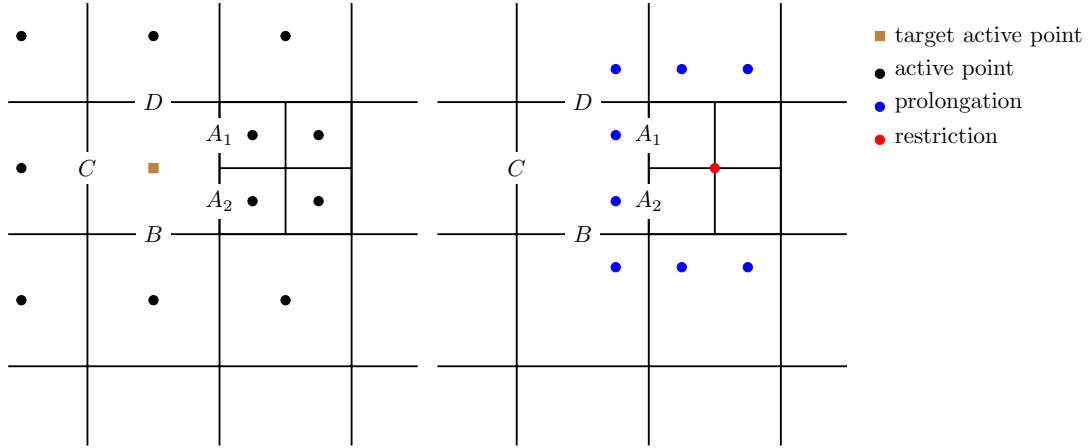


Figure 3: Sketch for 2D quadtree sample, active point and those ghost cells are depicted separately for clarity.

The first method which conserves 2nd order on tree grid calculates and stores value at each face. Therefore,  $\bullet$  as well as  $\bullet$  at level  $l + 1$  is used to compute desired value at face  $A_1$  and  $A_2$ . Then value on target cell is obtained by linear combination of its face value. In conclusion, all marked points ( $\bullet$  is used when compute face value on  $B, D$ ) in figure 3 participate in determining target value.

However, second method employs direct way to calculate target value from cell centered stand. Therefore, target value is only decided by  $\bullet$  and  $\bullet$  on level  $l$ . Thus, applying second method to tree grid will lead to accuracy decrease. Otherwise, both methods act the same way on cartesian and the second method runs even faster.

### 2.3.3 Program Workflow

#### Starting Point

input:

$$a = \mathbf{u}^{*,k} \quad b = \mathbf{u}'$$

$$resl = RES^k$$

$$data = \mu^{n+\frac{1}{2}}, \rho^{n+\frac{1}{2}}, \Delta t$$

data assignment

$\mu, \rho, \Delta t$  is extracted from  $data$

$u, r$  and  $res$  serve as pointer and point to  $\delta \mathbf{u}^{*,k}, \mathbf{u}', RES^k$  respectively.

```

1 static double residual_viscosity (scalar * a, scalar * b, scalar * resl,
2     void * data)
3 {
4     struct Viscosity * p = (struct Viscosity *) data;
5     (const) face vector mu = p->mu;
6     (const) scalar rho = p->rho;

```

```

7  double dt = p->dt;
8  vector u = vector(a[0]), r = vector(b[0]), res = vector(resl[0]);
9  double maxres = 0.;

```



### Tree Grid Update

#### ghost cell:

Set prolongation ghost cells for all components manually.

#### tree-flux compute

Compute flux *tau* on each face for three dimension based on equation 4.

#### tree-flux assmeble

Traverse all cell and asseemble  $\mathbf{D}^*$  as *d* from cell centered perspective.

#### residual compute

Then the residual is updated by equation 13 and stored in *res*. The maximum of all components is returned through double data *maxres*.

```

1  #if TREE
2  boundary ({u});
3
4  foreach_dimension() {
5      face vector taux[];
6      foreach_face(x)
7          taux.x[] = 2.*mu.x[]*(u.x[] - u.x[-1])/Delta;
8      #if dimension > 1
9          foreach_face(y)
10             taux.y[] = mu.y[]*(u.x[] - u.x[0,-1] +
11                (u.y[1,-1] + u.y[1,0])/4. -
12                (u.y[-1,-1] + u.y[-1,0])/4.)/Delta;
13      #endif
14      #if dimension > 2
15          foreach_face(z)
16             taux.z[] = mu.z[]*(u.x[] - u.x[0,0,-1] +
17                (u.z[1,0,-1] + u.z[1,0,0])/4. -
18                (u.z[-1,0,-1] + u.z[-1,0,0])/4.)/Delta;
19      #endif
20      foreach (reduction(max:maxres)) {
21          double d = 0.;
22          foreach_dimension()
23              d += taux.x[1] - taux.x[];
24          res.x[] = r.x[] - lambda.x*u.x[] + dt/rho[]*d/Delta;
25          if (fabs (res.x[]) > maxres)
26              maxres = fabs (res.x[]);

```

```

27     }
28 }

```



### NonTree Grid Update residual compute

Then the residual is updated by equation 13 and stored in *res*. The maximum of all components is returned through double data *maxres*.

```

1  #else
2  foreach (reduction(max:maxres))
3      foreach_dimension() {
4          res.x[] = r.x[] - lambda.x*u.x[] +
5              dt/rho[]*(2.*mu.x[1,0]*(u.x[1] - u.x[])
6              - 2.*mu.x[]*(u.x[] - u.x[-1]))
7          #if dimension > 1
8              + mu.y[0,1]*(u.x[0,1] - u.x[] +
9              (u.y[1,0] + u.y[1,1])/4. -
10             (u.y[-1,0] + u.y[-1,1])/4.)
11             - mu.y[]*(u.x[] - u.x[0,-1] +
12             (u.y[1,-1] + u.y[1,0])/4. -
13             (u.y[-1,-1] + u.y[-1,0])/4.)
14          #endif
15          #if dimension > 2
16              + mu.z[0,0,1]*(u.x[0,0,1] - u.x[] +
17              (u.z[1,0,0] + u.z[1,0,1])/4. -
18              (u.z[-1,0,0] + u.z[-1,0,1])/4.)
19              - mu.z[]*(u.x[] - u.x[0,0,-1] +
20              (u.z[1,0,-1] + u.z[1,0,0])/4. -
21              (u.z[-1,0,-1] + u.z[-1,0,0])/4.)
22          #endif
23              )/sq(Delta);
24          if (fabs (res.x[]) > maxres)
25              maxres = fabs (res.x[]);
26      }
27  #endif
28  return maxres;
29 }
30
31 #undef lambda

```

## 2.4 viscosity

### 2.4.1 Parameters

Name	Data type	Status	Option/Default	Representation (before/after)
<i>u</i>	vector	update	compulsory	$\mathbf{u}'/\mathbf{u}^*$
<i>mu</i>	face vector	unchanged	compulsory	$\mu^{n+\frac{1}{2}}$
<i>rho</i>	scalar	unchanged	compulsory	$\rho^{n+\frac{1}{2}}$
<i>dt</i>	double	unchanged	compulsory	$\Delta t$
<i>nrelax</i>	int	unchanged	optional/4	<i>max of iteration</i>
<i>res</i>	scalar*	<b>output</b>	optional/NULL	<i>RES</i>

### 2.4.2 Worth Mentioning Details

The function to assemble all the tools built in this headfile or in 'poisson.h' and solve the governing equation equation 1. Details about implicit iteration solver construction used here are explored in 'poisson.h Documentation' which shall not be repeated again.

### 2.4.3 Program Workflow

#### Solver Construction

##### input:

$$\mathbf{u} = \mathbf{u}' \quad \mathbf{mu} = \mu^{n+\frac{1}{2}} \quad \mathbf{rho} = \rho^{n+\frac{1}{2}}$$

$$\mathbf{dt} = \Delta t \quad \mathbf{res} = \text{empty}$$

$$\mathbf{nrelax} = \text{iteration number}$$

##### data assignment

*r* is created to store  $\mathbf{u}'$  and *u* now serves as  $\mathbf{u}^*$  which will be updated.

Then *mu* and *rho* is restricted to all level and *p* is defined as self made struct and to store  $\mu, \rho, \Delta t$ .

##### iteration solver

Solver is built using **mg\_solve** from 'poisson.h'

```

1 trace
2 mgstats viscosity (vector u, face vector mu, scalar rho, double dt, int
  ↳ nrelax = 4, scalar * res = NULL)
3 {
4     vector r[];
5     foreach()
6         foreach_dimension()
7             r.x[] = u.x[];
8 
```

```

9  restriction ({mu,rho});
10 struct Viscosity p = { mu, rho, dt };
11 return mg_solve ((scalar *){u}, (scalar *){r}, residual_viscosity,
    ↪ relax_viscosity, &p, nrelax, res);
12 }

```

## 2.5 viscosity\_explicit

### 2.5.1 Parameters

Name	Data type	Status	Option/Default	Representation (before/after)
<i>u</i>	vector	update	compulsory	$\mathbf{u}'/\mathbf{u}^*$
<i>mu</i>	face vector	unchanged	compulsory	$\mu^{n+\frac{1}{2}}$
<i>rho</i>	scalar	unchanged	compulsory	$\rho^{n+\frac{1}{2}}$
<i>dt</i>	double	unchanged	compulsory	$\Delta t$

### 2.5.2 Worth Mentioning Details

Different from implicit form shown in equation 1, the governing equation of this function is explicit

$$\rho_{n+\frac{1}{2}} \left[ \frac{\mathbf{u}^* - \mathbf{u}'}{\Delta t} \right] = \nabla \cdot [2\mu_{n+\frac{1}{2}} \mathbf{D}'] \quad (17)$$

Then

$$\mathbf{u}^* = \mathbf{u}' + \frac{\Delta t \nabla \cdot [2\mu_{n+\frac{1}{2}} \mathbf{D}']}{\rho_{n+\frac{1}{2}}} \quad (18)$$

According to section 2.3, if *a*, *b* of **viscosity\_residual** are set to be the same and equal to  $\mathbf{u}'$  then the output yields

$$RES = \frac{\Delta t \nabla \cdot [2\mu_{n+\frac{1}{2}} \mathbf{D}']}{\rho_{n+\frac{1}{2}}} \quad (19)$$

Hence

$$\mathbf{u}^* = \mathbf{u}' + \frac{\Delta t \nabla \cdot [2\mu_{n+\frac{1}{2}} \mathbf{D}']}{\rho_{n+\frac{1}{2}}} = \mathbf{u}' + RES \quad (20)$$

### 2.5.3 Program Workflow

**Solver Construction****input:**

$$\mathbf{u} = \mathbf{u}' \quad \mu = \mu^{n+\frac{1}{2}} \quad \rho = \rho^{n+\frac{1}{2}}$$

$$\Delta t = \Delta t$$

**explicit solver**

$$\mathbf{r} = \frac{\Delta t \nabla \cdot [2\mu_{n+\frac{1}{2}} \mathbf{D}']}{\rho_{n+\frac{1}{2}}}$$

 $\mathbf{u}$  then updates to  $\mathbf{u}^*$ 

```

1 trace
2 mgstats viscosity_explicit (vector u, face vector mu, scalar rho, double
  ↪ dt)
3 {
4   vector r[];
5   mgstats mg = {0};
6   struct Viscosity p = { mu, rho, dt };
7   mg.resb = residual_viscosity ((scalar *){u}, (scalar *){u}, (scalar
  ↪ *){r}, &p);
8   foreach()
9     foreach_dimension()
10      u.x[] += r.x[];
11   return mg;
12 }

```

## References

- [1] J. A. Van Hooft et al. “Towards adaptive grids for atmospheric boundary-layer simulations”. In: *Bound.-Layer Meteorol.* 167 (2018), pp. 421–443.