

## Équation des géodésiques

Soit le lagrangiens suivant :

$$\mathcal{L}(x^\alpha, \dot{x}^\alpha) = [g_{\mu\nu}(x) \dot{x}^\mu \dot{x}^\nu]^{\frac{1}{2}}$$

avec  $\dot{x}^\alpha = \frac{dx^\alpha}{ds}$

Ainsi on n'a :

$$\frac{\partial \mathcal{L}}{\partial x^\alpha} = \frac{1}{2} [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} \frac{\partial}{\partial x^\alpha} (g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu)$$

car  $\dot{x}^\mu$  et  $\dot{x}^\nu$

sont constant

$$\begin{aligned} \text{face à } x^\alpha \rightarrow &= \frac{1}{2} [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} \frac{\partial g_{\mu\nu}}{\partial x^\alpha} \dot{x}^\mu \dot{x}^\nu \\ &= \frac{1}{2} [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} g_{\mu\nu, \alpha} \dot{x}^\mu \dot{x}^\nu \end{aligned}$$

où on note :  $g_{\mu\nu, \alpha} = \frac{\partial g_{\mu\nu}}{\partial x^\alpha}$

Et

$$\begin{aligned} \frac{\partial \mathcal{L}}{\partial \dot{x}^\alpha} &= \frac{1}{2} [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} \frac{\partial}{\partial \dot{x}^\alpha} (g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu) = \frac{1}{2} [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} g_{\mu\nu} \frac{\partial}{\partial \dot{x}^\alpha} (\dot{x}^\mu \dot{x}^\nu) \\ &= \frac{1}{2} [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} g_{\mu\nu} \left( \underbrace{\dot{x}^\nu \frac{\partial}{\partial \dot{x}^\alpha} (\dot{x}^\mu)}_{=\delta_\alpha^\mu} + \underbrace{\dot{x}^\mu \frac{\partial}{\partial \dot{x}^\alpha} (\dot{x}^\nu)}_{=\delta_\alpha^\nu} \right) \\ &= \frac{1}{2} [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} g_{\mu\nu} \left( \dot{x}^\nu \delta_\alpha^\mu + \underbrace{\dot{x}^\mu \delta_\alpha^\nu}_{=\dot{x}^\nu \delta_\alpha^\mu \text{ car } \mu \text{ et } \nu \text{ sont muet}} \right) \\ &= \cancel{\frac{1}{2}} [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} g_{\mu\nu} \cdot \cancel{2} \dot{x}^\nu \delta_\alpha^\mu \\ &= [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} g_{\alpha\nu} \dot{x}^\nu \end{aligned}$$

Ainsi par les équation d'Euler-La Grange, on obtient l'équation suivante :

$$\frac{d}{ds} \left( [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} g_{\alpha\nu} \dot{x}^\nu \right) = \frac{1}{2} [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} g_{\mu\nu, \alpha} \dot{x}^\mu \dot{x}^\nu$$

En posant :

$$d\lambda = [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{\frac{1}{2}} ds$$

$$\text{Donc } \frac{d}{d\lambda} = [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} \frac{d}{ds}$$

Ainsi, en réécrivant l'équation précédente :

$$\frac{d}{ds} \left( [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} g_{\alpha\nu} \frac{d}{ds} x^\nu \right) = \frac{1}{2} [g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu]^{-\frac{1}{2}} g_{\mu\nu, \alpha} \frac{d}{ds} x^\mu \frac{d}{ds} x^\nu$$

$$\begin{aligned}
\text{donc } [g_{\mu\nu}\dot{x}^\mu\dot{x}^\nu]^{\frac{1}{2}} \frac{d}{d\lambda} \left( [\cancel{g_{\mu\nu}\dot{x}^\mu\dot{x}^\nu}]^{-\frac{1}{2}} g_{\alpha\nu} [\cancel{g_{\mu\nu}\dot{x}^\mu\dot{x}^\nu}]^{\frac{1}{2}} \frac{d}{d\lambda} x^\nu \right) &= \frac{1}{2} [g_{\mu\nu}\dot{x}^\mu\dot{x}^\nu]^{-\frac{1}{2}} [g_{\mu\nu}\dot{x}^\mu\dot{x}^\nu] g_{\mu\nu,\alpha} \frac{d}{d\lambda} x^\mu \frac{d}{d\lambda} x^\nu \\
\text{donc } [\cancel{g_{\mu\nu}\dot{x}^\mu\dot{x}^\nu}]^{\frac{1}{2}} \frac{d}{d\lambda} \left( g_{\alpha\nu} \frac{d}{d\lambda} x^\nu \right) &= \frac{1}{2} [\cancel{g_{\mu\nu}\dot{x}^\mu\dot{x}^\nu}]^{\frac{1}{2}} g_{\mu\nu,\alpha} \frac{d}{d\lambda} x^\mu \frac{d}{d\lambda} x^\nu \\
\text{donc } \frac{d}{d\lambda} \left( g_{\alpha\nu} \frac{d}{d\lambda} x^\nu \right) &= \frac{1}{2} g_{\mu\nu,\alpha} \frac{d}{d\lambda} x^\mu \frac{d}{d\lambda} x^\nu
\end{aligned}$$

Si on écrit l'action par rapport à  $s$ , on a :

$$S_1 = \int \left[ g_{\mu\nu} \frac{d}{ds} x^\mu \frac{d}{ds} x^\nu \right]^{\frac{1}{2}} ds$$

Alors en opérant le changement de variable  $S_1$  devient :

$$\begin{aligned}
S_1 &= \int [\cancel{g_{\mu\nu}\dot{x}^\mu\dot{x}^\nu}]^{\frac{1}{2}} \left[ g_{\mu\nu} \frac{d}{d\lambda} x^\mu \frac{d}{d\lambda} x^\nu \right]^{\frac{1}{2}} [\cancel{g_{\mu\nu}\dot{x}^\mu\dot{x}^\nu}]^{-\frac{1}{2}} d\lambda \\
&= \int \left[ g_{\mu\nu} \frac{d}{d\lambda} x^\mu \frac{d}{d\lambda} x^\nu \right]^{\frac{1}{2}} d\lambda
\end{aligned}$$

Donc comme  $S_1$  est invariante par la transformation  $s \rightarrow \lambda$  et que celle-ci est un difféomorphisme, on a alors :

$$\frac{d}{ds}(g_{\alpha\nu}\dot{x}^\nu) = \frac{1}{2}g_{\mu\nu,\alpha}\dot{x}^\mu\dot{x}^\nu$$

Donc en calculant le terme de gauche :

$$\begin{aligned}
\frac{d}{ds}(g_{\alpha\nu}\dot{x}^\nu) &= \dot{x}^\nu \frac{d}{ds} g_{\alpha\nu} + g_{\alpha\nu}\ddot{x}^\nu \\
&= g_{\alpha\nu}\ddot{x}^\nu + \dot{x}^\nu \underbrace{\frac{\partial g_{\alpha\nu}}{\partial \mu} \frac{dx^\mu}{ds}}_{=g_{\alpha\nu,\mu}=\dot{x}^\mu} \\
&= g_{\alpha\nu}\ddot{x}^\nu + g_{\alpha\nu,\mu}\dot{x}^\nu\dot{x}^\mu
\end{aligned}$$

Ainsi :

$$\begin{aligned}
\underbrace{g^{\alpha\beta}g_{\alpha\nu}}_{=\delta_\nu^\beta} \ddot{x}^\nu &= \frac{1}{2}g_{\mu\nu,\alpha}\dot{x}^\mu\dot{x}^\nu - g_{\alpha\nu,\mu}\dot{x}^\nu\dot{x}^\mu \\
\text{Donc } \ddot{x}^\beta &= \frac{1}{2}g^{\alpha\beta}(g_{\mu\nu,\alpha} - 2g_{\alpha\nu,\mu})\dot{x}^\nu\dot{x}^\mu \\
\text{Donc } \ddot{x}^\beta &= \frac{1}{2}g^{\alpha\beta}(g_{\mu\nu,\alpha} - g_{\alpha\nu,\mu} - g_{\alpha\mu,\nu})\dot{x}^\nu\dot{x}^\mu
\end{aligned}$$

En introduisant les symboles de Cristoffel  $\Gamma_{\mu\nu}^\beta = \frac{1}{2}g^{\beta\alpha}(g_{\mu\alpha,\nu} + g_{\nu\alpha,\mu} - g_{\mu\nu,\alpha})$  on a :

$$\ddot{x}^\beta = -\Gamma_{\mu\nu}^\beta\dot{x}^\mu\dot{x}^\nu$$

$$\text{Soit } \ddot{x}^\beta + \Gamma_{\mu\nu}^\beta\dot{x}^\mu\dot{x}^\nu = 0$$

## Commuteur des Co-dérivé

On définis les dérivé covariante par :

$$\nabla_{\mu} v^{\nu} = \partial_{\mu} v^{\nu} + \Gamma_{\mu\rho}^{\nu} v^{\rho}$$

Ainsi, soit  $v^{\rho}$  un vecteur :

$$\begin{aligned} [\nabla_{\mu}; \nabla_{\nu}] v^{\rho} &= \nabla_{\mu} (\nabla_{\nu} v^{\rho}) - \nabla_{\nu} (\nabla_{\mu} v^{\rho}) \\ &= \nabla_{\mu} (\partial_{\nu} v^{\rho} + \Gamma_{\nu\alpha}^{\rho} v^{\alpha}) - \nabla_{\nu} (\partial_{\mu} v^{\rho} + \Gamma_{\mu\alpha}^{\rho} v^{\alpha}) \\ &= \partial_{\mu} (\partial_{\nu} v^{\rho} + \Gamma_{\nu\alpha}^{\rho} v^{\alpha}) + \Gamma_{\mu\beta}^{\rho} (\partial_{\nu} v^{\beta} + \Gamma_{\nu\alpha}^{\beta} v^{\alpha}) - \partial_{\nu} (\partial_{\mu} v^{\rho} + \Gamma_{\mu\alpha}^{\rho} v^{\alpha}) - \Gamma_{\nu\beta}^{\rho} (\partial_{\mu} v^{\beta} + \Gamma_{\mu\alpha}^{\beta} v^{\alpha}) \\ &= \cancel{\partial_{\nu\mu} v^{\rho}} + \partial_{\mu} \Gamma_{\nu\alpha}^{\rho} v^{\alpha} + \Gamma_{\mu\beta}^{\rho} \partial_{\nu} v^{\beta} + \Gamma_{\mu\beta}^{\rho} \Gamma_{\nu\alpha}^{\beta} v^{\alpha} - \cancel{\partial_{\mu\nu} v^{\rho}} - \partial_{\nu} \Gamma_{\mu\alpha}^{\rho} v^{\alpha} - \Gamma_{\nu\beta}^{\rho} \partial_{\mu} v^{\beta} - \Gamma_{\nu\beta}^{\rho} \Gamma_{\mu\alpha}^{\beta} v^{\alpha} \\ &= \left( \Gamma_{\mu\beta}^{\rho} \Gamma_{\nu\alpha}^{\beta} - \Gamma_{\nu\beta}^{\rho} \Gamma_{\mu\alpha}^{\beta} + \partial_{\mu} \Gamma_{\nu\alpha}^{\rho} - \partial_{\nu} \Gamma_{\mu\alpha}^{\rho} \right) v^{\alpha} + \underbrace{\Gamma_{\mu\beta}^{\rho} \partial_{\nu} v^{\beta} - \Gamma_{\nu\beta}^{\rho} \partial_{\mu} v^{\beta}}_{=0} \end{aligned}$$

Ainsi on reconnais le tenseur de Riemann définis par :  $R_{\mu\nu\kappa}^{\sigma} = \partial_{\nu} \Gamma_{\mu\kappa}^{\sigma} - \partial_{\kappa} \Gamma_{\mu\nu}^{\sigma} + \Gamma_{\nu\lambda}^{\sigma} \Gamma_{\mu\kappa}^{\lambda} - \Gamma_{\kappa\lambda}^{\sigma} \Gamma_{\mu\nu}^{\lambda}$

Donc on obtient la relation suivante :

$$[\nabla_{\mu}; \nabla_{\nu}] v^{\rho} = R_{\alpha\mu\nu}^{\rho} v^{\alpha}$$