

QUANTUM EXTREMAL TRANSITIONS AND SPECIAL L-VALUES

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ABSTRACT. A threefold extremal transition $Y \searrow X$ consists of a crepant extremal contraction $\phi: Y \rightarrow \bar{Y}$ with curve class $\ell \in \text{NE}(Y)$, followed by a smoothing $\bar{Y} \leadsto X$. We consider the Type II case that ϕ contracts a divisor E to a point and prove that the quantum cohomology $QH(X)$ is obtained from $QH(Y)$ via *analytic continuation, regularization, and specialization* in Q^ℓ . Besides roots of unity, special L-values appear in $\lim Q^\ell$ whenever \bar{Y} admits more than one smoothings.

Further techniques are employed and explored beyond known tools in Gromov–Witten theory including (i) the *canonical local B model* attached to $Y \searrow X$, (ii) existence of *semistable reduction of double point type* for the smoothing, (iii) the *modularity of the extremal function* $\mathbb{E} := E^3 / \langle E, E, E \rangle^Y$, and (iv) *periods integrals of Eisenstein series*. Our study provides a geometric framework linking classifications of del Pezzo surfaces, Ramanujan’s theta functions, and Zagier’s special ODE list via Type II transitions.

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Notation and convention.

- In this paper the coefficients for cohomology $H^k(M)$ are in \mathbb{C} if not stated otherwise.
- To shorten notation, we set $H(M) := H^{\text{ev}}(M)$, the even part of the cohomology groups.
- $\text{NE}(M)$ is the integral Mori cone, and $\text{NE}(M)_{\mathbb{R}} = \text{NE}(M) \otimes_{\mathbb{Z}} \mathbb{R}$.
- q^β is the Novikov variable for $\beta \in \text{NE}(M)$.
- $P^\beta = q^\beta e^{(t,\beta)}$, where $t \in H^2(M)$ is the small parameter of the I -function.
- $Q^\beta = q^\beta e^{(\tau,\beta)}$, where $\tau \in H^2(M)$ is the small parameter of the J -function.

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- \hat{D} is the quantization operator of D , defined by $\hat{D}P^\beta = z(D, \beta)P^\beta$.
- δ^D is the power operator of D , defined by $\delta^D Q^\beta = (D, \beta)Q^\beta$.
- $\theta = P^\ell \frac{\partial}{\partial P^\ell}$ where ℓ is the extremal curve class for the contraction $\phi: Y \rightarrow \bar{Y}$.
- $q = e^{2\pi i \tau_q}$ where $\tau_q \in \mathbb{H}$.
- $[z^{-k}]F$ denotes the coefficient of z^{-k} in F .
- $\mathbf{1}_N$ denotes the trivial Dirichlet character modulo N .
- $\omega_N = e^{2\pi i/N}$ is the N -th root of unity.
- \mathfrak{S}_n denotes the symmetric group of degree n .
- $R[\![x]\!]$ is the formal power series ring with coefficients in R .
- $R(\!(x)\!)$ is the formal Laurent series ring with coefficients in R (field of fraction of $R[\![x]\!]$).

0. INTRODUCTION

0.1. Statement of main results. We study the relation of (big) quantum cohomologies of smooth projective threefolds Y and X when they are related by an extremal transition

$$Y \searrow X \quad \text{via} \quad Y \xrightarrow{\phi} \bar{Y} \rightsquigarrow X,$$

where $\phi: Y \rightarrow \bar{Y}$ is a primitive (birational) crepant contraction (so $\rho(Y/\bar{Y}) = 1$) and $\bar{Y} \rightsquigarrow X$ is a smoothing $\mathfrak{X} \rightarrow \Delta$ of $\bar{Y} = \mathfrak{X}_0$ with the fiber $X = \mathfrak{X}_t$ for some $t \neq 0$.

Let $\ell \in \text{NE}(Y)$ be the extremal curve class for ϕ , the contraction of $\mathbb{R}_{\geq 0}\ell$. Following Wilson [50], ϕ is of Type I if it contracts finitely many curves to points $p_i \in \bar{Y}$, of Type II if it contracts a divisor to a point, and of Type III if it contracts a divisor to a curve $C \subseteq \bar{Y}$. We thus call the transition $Y \searrow X$ being of Type I, II or III accordingly.

It is well known that for Type I transitions a symplectic deformations of it leads to a conifold transition (all p_i 's are ordinary double points) where the question was fully understood by Li–Ruan [26] around 2000. In the Calabi–Yau case it states that

$$\langle - \rangle_\beta^X = \sum_{\gamma \mapsto \beta} \langle - \rangle_\gamma^Y. \tag{0.1}$$

The sum of Gromov–Witten (GW) invariants is in fact finite and the RHS can be understood by setting $q^\ell = 1$ in the Novikov variables. Moreover the notion of linked GW invariants was developed in [24] in order to read out the individual terms on Y from those on X .

For Type III transitions, Wilson had shown that C is a smooth curve [50, 51] and if $g(C) \neq 0$ then one may deform the primitive contraction ϕ to a Type I case with $2g(C) - 2$ singular fibers in an appropriate length counting (cf. [50, Proposition 4.2]). Since GW invariants are invariant under deformations, it is not hard to see that (0.1) is still valid in this case.

Thus the remaining cases are Type II transitions and Type III with $g(C) = 0$. It is our goal to solve both cases. Due to the difference of techniques involved, we will solve the Type II case in the current paper and defer the Type III genus zero case to a subsequent work.

For a Type II transition, ϕ contracts a (possibly non-normal) del Pezzo surface $E \subseteq Y$ to a rational Gorenstein singularity $p \in \bar{Y}$. Based on earlier works of Gross and Wilson [13, 51], we proved that there is also a symplectic deformation on $Y \searrow X$ so that E is deformed to a

smooth del Pezzo surface S_d of degree $d = E^3$. Based on Fujita–Iskovskikh’s classification theory of del Pezzo surfaces and threefolds (Appendix A), as well as Reid’s classifications of isolated Gorenstein singularities [45], we are able to construct projective local model of the transition, which we call a del Pezzo transition $Y_d \searrow X_d$ where X_d is a del Pezzo threefold of degree d . Since \bar{Y}_d is smoothable, S_d is not isomorphic to \mathbb{P}^2 or the Hirzebruch surface \mathbb{F}_1 . In particular, we have $1 \leq d \leq 8$. Up to deformations, the smoothing is unique for $d \neq 6$. For $d = 6$, there are two smoothings of \bar{Y}_6 named by $X_{6\text{I}}$ and $X_{6\text{II}}$.

We are able to construct semistable reductions on both the *Kähler degeneration* $\phi: Y \rightarrow \bar{Y}$ and the *complex degeneration* $\mathfrak{X} \rightarrow \Delta$ of $\bar{Y} = \mathfrak{X}_0$ so that we may try to prove a statement on GW invariants first in the local setting $Y_d \searrow X_d$ and then for any $Y \searrow X$ by the degeneration formula of GW theory. Furthermore, again by the classification theory of S_d and X_d , we find directly the degree of the finite base change needed for the semistable reduction of $\mathfrak{X} \rightarrow \Delta$, avoiding toroidal embeddings (see Proposition 1.10 and Theorem 4.1).

We may state our main results informally as the following: we provide analytic continuations of quantum cohomology $QH(Y)$ in Q^ℓ under Type II extremal transitions $Y \searrow X$ of threefolds and recover $QH(X)$ after regularization and specialization of Q^ℓ to certain value Q_r^ℓ (see (5.9)). Exotic behavior of Q_r^ℓ is discovered: under the contraction ϕ we have $\phi_*(\ell) = 0$, however Q_r^ℓ might not even be a root of unity. It could be a special L-value!

Now the formal description: a canonical local B model, i.e., a pair of I -functions (I^Y, I^X) with a distinguished variable P^ℓ mirror to the quantum variable Q^ℓ is constructed. It allows us to study analytic continuations of $QH(Y)$ via Picard–Fuchs equations and asymptotic analysis near $P^\ell = \infty$. Namely, $QH(X)$ is obtained from $QH(Y)$ by a regularized holomorphic limit (see Definition 4.5) as $P^\ell \rightarrow \infty$:

Theorem 0.1 (= Corollary 4.4 + Corollary 4.6). *Let $n \geq 1$. For any insertions $\vec{b} \in H(X)^{\oplus n}$ and $0 \neq \tilde{\beta} \in \text{NE}(X)$ with lifting $\tilde{\beta} \in \text{NE}(Y)$ such that $(E, \tilde{\beta}) = 0$, we have*

$$\lim_{P^\ell \rightarrow \infty} \left(\langle \phi^* \vec{b} \rangle_{\tilde{\beta}}^{Y \searrow X} Q^{-\tilde{\beta}} \right)^{\text{reg}} = \langle \vec{b} \rangle_{\tilde{\beta}}^X. \quad (0.2)$$

Moreover, the regularization is superfluous precisely when $d \in \{1, 2, 3, 4, 7\}$.

The limit of Q^ℓ is more subtle and requires tools from modular forms. We show that (see Theorem 5.9, and Table 6 for the congruence subgroups Γ_d):

Theorem 0.2. *For $d \neq 7$, there exists a natural identification of $\mathbb{P}_{P^\ell}^1$ as a modular curve $X(\Gamma_d)$ (a double cover of $X(\Gamma(1))$ if $d = 1$).*

The coordinate change is expressed via the extremal function

$$\mathbb{E}^{Y \searrow X} = \mathbb{E}_d := \frac{E^3}{\sum \langle E, E, E \rangle_{m\ell}^{Y \searrow X} Q^{m\ell}} = \frac{d \log Q}{d \log q}, \quad (0.3)$$

where $q = e^{2\pi i \tau_q}$ and $Q = -q + O(q^2)$. It is an Eisenstein series of weight 3.

Moreover $\langle \phi^* \vec{b} \rangle_{\tilde{\beta}}^{Y \searrow X}$ is a quasi-modular form of weight 0 with value $\langle \vec{b} \rangle_{\tilde{\beta}}^X$ at a special cusp or elliptic point, depending on the monodromy of $\mathbb{E}^{Y \searrow X}$ being logarithmic or finite.

We call the elliptic point/cusp of q corresponding to $P^\ell \rightarrow \infty$ the *transition point* (Definition 5.6). The corresponding limiting values of Q^ℓ are necessarily path dependent.¹

Theorem 0.3 (= Proposition 5.1 + Theorem 5.17 (i)). *For $d \in \{1, 2, 3, 4, 7\}$, $QH(Y)$ can be analytically continued along $P^\ell \in [0, \infty)$ and $Q^\ell \rightarrow 1$ as $P^\ell \rightarrow \infty$.*

In general, for $d \neq 7$, whenever \bar{Y}_d has unobstructed deformations, namely the smoothing is unique up to deformations, we have $Q^\ell \rightarrow 1$ or a root of unity.

This is of course expected since the curve class ℓ is contracted. However, when \bar{Y} has more than one smoothings, namely $d = 6$, we show that:

Theorem 0.4 (= Theorem 5.17 (ii)). *The limit of Q^ℓ in case 6I is still a root of unity while in case 6II is given by*

$$Q_r^\ell := \omega \cdot e^{-cL'(-1, x_{3,2})} \quad (0.4)$$

for some root of unity ω , where the cusp $r = \frac{a}{c} \in \mathbb{Q}$ with $\gcd(a, c) = 1$ and $c \equiv 2 \pmod{6}$, and L is a Dirichlet L-function. In particular, Q_r^ℓ is not a root of unity.

To our knowledge, this is the first appearance of special L-values as specialization points in the study of Gromov–Witten theory. Our proof relies on detailed analysis on period integrals of Eisenstein series.

Some of the results extend to certain higher dimensional transitions which contract a divisor to a singular point. Also the period integral of Eisenstein series can be carried out for q approaching any of the elliptic points or cusps. These extensions will be contained in the first author’s 2026 PhD thesis.

0.2. Outline of the paper. In this section we discuss the main ideas behind the proofs of our main results. The details are referred to the corresponding locations in this paper.

There are classifications of smooth del Pezzo surfaces (Theorem A.1) according to its degree $d := E^3 \in \mathbb{N}$ as well as of smoothable isolated rational Gorenstein threefold singularities (by Gross, Namikawa and Reid [13, 41, 45]). It turns out to admit a perfect match on both and we get for each $1 \leq d \leq 8$, $E = S_d$ (smooth del Pezzo surface of degree d), a local model of Type II extremal transition (Example 1.9):

$$Y_d = \mathbb{P}_{S_d}(K_{S_d} \oplus \mathcal{O}) \searrow X_d,$$

where X_d is the del Pezzo threefold of degree d arising from smoothing of \bar{Y}_d .

The smoothing family is unique if $d \neq 6$. For $d = 6$ there are two smoothings denoted by cases $d = 6\text{I}$ and $d = 6\text{II}$. So we let $d \in \{1, 2, 3, 4, 5, 6\text{I}, 6\text{II}, 7, 8\}$.

The classification of del Pezzo threefolds (Theorem A.5) allows us to find a smooth polarized variety $(T_d, \mathcal{O}_{T_d}(1))$ with known $QH(T_d)$ and a convex vector bundle \mathcal{V}_d on T_d such that $X_d = Z(s_d) \subseteq \mathcal{V}_d$ for $s_d \in \Gamma(T_d, \mathcal{V}_d)$, and $\mathcal{O}_{X_d}(1) = \mathcal{O}_{T_d}(1)|_{X_d}$, see Table 4.

¹The modularity for $d = 7$ has a different nature which will be discussed elsewhere.

All these geometric data fit into the following basic diagram

$$\begin{array}{ccccc} Y_d & & \mathcal{V}_d \\ j_d \downarrow \pi_d & & s_d \downarrow \\ S_d & \xhookrightarrow{i_d} & X_d & \hookrightarrow & T_d \end{array}$$

which allows us to setup a *canonical local B model* of $Y_d \setminus X_d$, i.e. a pair of closely related I -functions I^{Y_d} and I^{X_d} , to compare GW invariants of Y_d and X_d .

The annihilator of the I -function is the Picard–Fuchs ideal PF generated by equations \square_β over curve classes β . For $d \neq 7$, we show that PF^{X_d} is generated by $\square_{\bar{\gamma}} I^{X_d} = 0$ with

$$\square_{\bar{\gamma}} = \hat{h}^4 - P^{\bar{\gamma}}(\kappa\hat{h}^2 + \kappa z\hat{h} + \lambda z^2) - \mu P^{2\bar{\gamma}},$$

and PF^{Y_d} is generated by $\square_\gamma I^{Y_d} = \square_\ell I^{Y_d} = 0$ with

$$\begin{aligned} \square_\gamma &= \hat{E}\hat{H} - P^\gamma, \\ \square_\ell &= \hat{F}^3 - P^\ell(\kappa\hat{F}^2 + \kappa z\hat{F} + \lambda z^2)\hat{E} - \mu P^{2\ell}(\hat{F} + z)\hat{E}(\hat{E} - z). \end{aligned}$$

(See Proposition 3.9.) Here the constants $(\kappa, \lambda, \mu) = (\kappa_d, \lambda_d, \mu_d)$ are listed below:²

d	1	2	3	4	5	6I	6II	8
κ_d	432	64	27	16	11	7	10	0
λ_d	60	12	6	4	3	2	3	0
μ_d	0	0	0	0	1	8	-9	16

TABLE 1. Coefficients in \square_ℓ .

To simplify notations, the subscript d will often be omitted throughout this section, except X_d and Y_d . Since X_d is Fano, we have $J^{X_d} = I^{X_d}$. For J^{Y_d} , we showed that the mirror transform $\tau = t - g(P)E$ depends only on the canonical variables $P = P^\ell$ and $Q = Q^\ell$ via

$$Q = Pe^{g(P)}, \tag{0.5}$$

where $f := 1 + \theta g$ is uniquely determined by the equation

$$u\theta^2 f + \theta u\theta f + (\lambda P - \mu P^2)f = 0 \tag{0.6}$$

with $u = u(P) := 1 + \kappa P - \mu P^2$, $\theta = P \frac{d}{dP}$.

²Technically speaking, the contraction $Y_d \rightarrow \bar{Y}_d$ has $\rho(Y_d/\bar{Y}_d) > 1$ and we need to perform identifications to get the single equation \square_ℓ . For $d \neq 7$ we symmetrize the Picard–Fuchs equations and GW theory to achieve this in §3.2, with the group G_{S_d} being the Weyl group of the root system E_{9-d} in Table 3. However, for $d = 7$ the symmetrization is not working and we take an ad hoc approach instead (see §3.3).

In principle, knowing f near $P = 0$ determines all genus zero GW invariants of Y_d by mirror theorem and reconstructions. The key player to make the reconstruction manageable is the 3-point function $\langle E, E, E \rangle^{Y_d}$ which is determined by *the extremal function* in P via

$$\mathbb{E} := \frac{E^3}{\langle E, E, E \rangle^{Y_d}} = u f^3. \quad (0.7)$$

To actually compare GW invariants on Y_d with GW invariants of X_d we have to take advantage of the canonical local B model and the comparison turns out has to take place at $P = \infty$. A computation shows that for 1-point and 2-point insertions \vec{b} in $H(Y_d)$, the invariant $\langle \phi^* \vec{b} \rangle^{Y_d}$ lies in $\mathbb{C} + O(P^{-1})$ precisely when $\mu = 0$.

On the other hand, $\langle \phi^* \vec{b} \rangle^{Y_d}$ might contain terms like $v := \frac{\theta f}{f} = -1 + O(\frac{1}{\log P})$ near $P = \infty$ if $\mu \neq 0$ (i.e., $d \in \{5, 6I, 6II, 8\}$) and a holomorphic regularization needs to be introduced in order to perform specializations. More precisely, the local exponents of (0.6) at $y := P^{-1} = 0$ is $(1, 1)$, hence $f \in y \mathbb{C}[[y]] \oplus y \log y \mathbb{C}[[y]]$. Let $f^{\text{reg}} \in y \mathbb{C}[[y]]$ be the non-zero *monodromy invariant* solution (unique up to a scalar) of (0.6) then it is clear that

$$v^{\text{reg}} := \frac{\theta f^{\text{reg}}}{f^{\text{reg}}} = -1 + O(y)$$

and the regularization we have in mind is basically the replacement of v by v^{reg} . The point is that we need to make it compatible with reconstructions in GW theory. For this purpose, let $S = \mathbb{C}(y)[v]$ and $w = v - v^{\text{reg}}$. Consider the subring $R \subseteq S$ and the ideal I of R :

$$R = S \cap (\mathbb{C}[[y]] + w \cdot \mathbb{C}((y))[w]), \quad I = S \cap (y \cdot \mathbb{C}[[y]] + w \cdot \mathbb{C}((y))[w]).$$

We first prove that for certain initial 1-point and 2-point insertions \vec{b} in $H(Y_d)$, the invariant $\langle \phi^* \vec{b} \rangle^{Y_d} \in R[Q^{\tilde{\gamma}}]$ where $\tilde{\gamma}$ is the canonical lifting of γ in Y and

$$\langle \phi^* \vec{b} \rangle^{Y_d} - \phi^* \langle \vec{b} \rangle^{X_d} \in I[Q^{\tilde{\gamma}}].$$

Once this is achieved, by working out the delicate comparison of the two reconstruction procedure on both Y_d and X_d , we discovered a mysterious cancellation via \mathbb{E} to get rid off logarithmic singularities and the procedure can be carried over. This eventually leads to the proof of Theorem 0.1 in the local case $Y_d \searrow X_d$. (For $d = 7$ see §3.3.)

To globalize the result we construct semistable reductions of both the Kähler and complex degenerations in §4.1. They must be of double point type in order for the GW degeneration formula to be applicable. It is crucial to control the base change degree n_d for complex degeneration $\mathfrak{X} \rightarrow \Delta$ so that we may first perform the base change, then perform one weighted blowup, hence conclude that the semistable model is of double point type (see Proposition 1.10 and Theorem 4.1):

$$\mathcal{X}_0 = Y' \cup X_d$$

with $Y' \cap X_d = E'$ and E' is a del Pezzo surface of degree d . Moreover, Y and Y' are symplectic deformation equivalent of each other. We emphasize that the result seems too good to be true at first sight!

On the other hand, for Kähler degeneration the semistable reduction is simply the deformations to the normal cone which is naturally of double point type. Now by applying both symplectic and algebraic degeneration formula of GW invariants the global case $Y \setminus X$ is reduced to the local models. This completes the outline of proof of Theorem 0.1.³

It is also worth mentioning that, after the change of coordinates $y = P^{-1}$, the Picard–Fuchs ideal PF^Y has an extension $\widetilde{\text{PF}}^Y$ over $y = 0$. However, the plausible restriction identity $\widetilde{\text{PF}}^Y|_{y=0} = \text{PF}^X$ holds only when $\mu = 0$ (see Remark 3.11). This also explains the necessity of regularization.

While the regularization/specialization formula in Theorem 0.1 is accurate, it will be satisfactory only if we understand the limit of the intrinsic quantum variables Q^β on Y and determine the Riemann surface (*natural domain of definition*) of \mathbb{E} and of $QH(Y)$.

To study the problem, we proceed to show that $f = f_d$ and $\mathbb{E} = \mathbb{E}_d$ are Γ_d modular forms in $q = e^{2\pi i \tau_q}$ of weight one and three respectively, under the change of variable

$$\frac{d \log P}{d \log q} = u f^2 \quad \text{or equivalently} \quad \frac{d \log Q}{d \log q} = \mathbb{E} = u f^3 \quad (0.8)$$

with $q = -P + O(P^2) = -Q + O(Q^2)$. For f the result is essentially due to Maier [34] and Zagier [54] where $f(P(\tau_q))$ is known as Ramanujan's theta functions (cf. Cooper's book [8, §4–§6]).

For example for $d = 1$, near $q = 0$ we have $P(\tau_q) = (\mathcal{E}_6(\tau_q)/\mathcal{E}_4^{3/2}(\tau_q) - 1)/864$ and

$$f(P(\tau_q)) = {}_2F_1\left(\frac{1}{6}, \frac{5}{6}; 1; -432P(\tau_q)\right) = \mathcal{E}_4(\tau_q)^{1/4}.$$

Since $j(\tau_q)^{-1} = -P(\tau_q)(1 + 432P(\tau_q))$, the multi-valued function P parametrizes a double cover of $X(\Gamma(1))$. For $d \geq 2$ and $d \neq 7$, P indeed parametrize $X(\Gamma_d)$ and similar expressions are listed in Proposition 5.5. Since P and hence u is a weight 0 meromorphic modular function and f is weight 1, we see that $\mathbb{E} = u f^3$ is a weight 3 modular form.

To test modularity of GW n -point functions, the reconstructions makes use of differentiations which generally destroys the genuine modularity. This leads us to invoke the notion of *meromorphic quasi-modular forms* instead (see Definition 5.4).

Let $D_q = \frac{d}{d \log q} = q \frac{d}{dq}$. It follows from (0.8) that

$$v = \frac{\theta f}{f} = \frac{D_q f}{\mathbb{E}},$$

which is a quasi-modular of weight 0 and depth 1 (with root and monodromy if $d = 1$). This implies that every element in $S = \mathbb{C}(y)[v]$ is quasi-modular of weight 0.

For $d \in \{4, 5, 6I, 6II, 8\}$, the transition point is a cusp. Take a representative $r \in \mathbb{Q}$ and an element $A = \begin{pmatrix} a & * \\ c & * \end{pmatrix} \in \text{SL}(2, \mathbb{Z})$ such that $A \cdot i\infty = r$. A brief computation leads to

$$v^{\text{reg}} = \frac{\theta f^{\text{reg}}}{f^{\text{reg}}} = \mathbb{E}^{-1} \cdot \left(D_q f - \frac{1}{2\pi i} \frac{c}{(a - c\tau_q)} f \right)$$

³Nevertheless, logically we need to firstly figure out the correct regularization statement before we can run any induction on reconstruction or degeneration analysis.

as well as $v^{\text{reg}}(A\tau_q) = v_0^A$. By taking into account the regularization procedure, it implies that given $0 \neq \bar{\beta} \in \text{NE}(X)$ with lifting $\tilde{\beta} \in \text{NE}(Y)$ such that $(E, \tilde{\beta}) = 0$,

$$\langle \phi^* \vec{b} \rangle_{\bar{\beta}}^{Y \searrow X} Q^{-\tilde{\beta}} = \sum_m \langle \phi^* b \rangle_{\tilde{\beta} + m\ell} Q^{m\ell}$$

is a meromorphic quasi-modular form of weight 0 which coincides with $\langle b \rangle_{\bar{\beta}}^X$ at the transition point. This completes the outline of proof of Theorem 0.2 (cf. Theorem 5.9).

Consequently, we see that the GW invariants on Y in the extremal direction, which are polynomials in v , are quasi-modular forms of weight 0 on Γ_d . Since non-constant quasi-modular forms all have logarithmic singularities near the real line, it follows that the unit disk $\mathbb{D} = \{q \mid |q| < 1\}$ is exactly the Riemann surface we want (cf. Remark 5.10).

As a byproduct, we can compute the limit of Q by integrating \mathbb{E}_d along some path in \mathbb{D} ending at a representative of the transition point. It turns out that \mathbb{E}_d is an Eisenstein series, and the integration can be expressed explicitly as a combination of L-values.

To start with, f admits an analytic continuation to a multi-valued function on $\mathbb{P}^1 \setminus \text{Sing}_f$ (Sing_f is contained in the set of singular points of (0.6)). For any path $p: [0, 1] \rightarrow \mathbb{P}^1$ with $p(0) = 0$, $p(1) = \infty$, and $p(s) \in \mathbb{P}^1 \setminus \text{Sing}_{f_d}$ for $s \in (0, 1)$, the limit of Q_p^ℓ along the path p is

$$Q_p^\ell := \lim_{s \rightarrow 1^-} Q_p^\ell(s) = \lim_{\varepsilon \rightarrow 0^+} p(\varepsilon) \exp\left(\int_\varepsilon^1 f_d(p(s)) \frac{dp(s)}{p(s)}\right).$$

For $d \in \{1, 2, 3, 4\}$, we may take the path p to be the line $[0, \infty)$ and the integral can be evaluated by techniques in complex function theory. It is based on the fact that

$$f(P) = {}_2F_1\left(\tfrac{1}{n}, 1 - \tfrac{1}{n}; 1; -\kappa P\right),$$

where $n = 6, 4, 3, 2$ for $d = 1, 2, 3, 4$ ($n = n_d$ is precisely the degree of base change needed in Proposition 1.10), and Euler's integral representation of hypergeometric functions. The result is $Q_p^\ell = 1$. The case $d = 7$ is treated separately, and it is in fact easier. Thus the first half of Theorem 0.3 is essentially classical (see Proposition 5.1).

For the second half of Theorem 0.3, as well as for Theorem 0.4, we develop period integrals of Eisenstein series in §5.3 to evaluate the limit. The first ingredient is the following extension of a similar formula in [1] for the case $\chi = \mathbf{1}_1$, using the method in [52] (see Proposition 5.14). Let $k \geq 2$ be an integer. For any triple (χ, ψ, n) with $\chi\psi(-1) = (-1)^k$, define

$$I_k^{\chi, \psi, n} = n^{-1} \lim_{s \rightarrow 0} L(s+1, \chi) L(s+2-k, \psi).$$

Let $a_{i\infty, 0}$ and $a_{0, 0}$ be the cusp values of $E_k^{\chi, \psi, n}(\tau_q)$ at $\tau_q = i\infty$ and $\tau_q = 0$, respectively. Then

$$\int_{i\infty}^0 (E_k^{\chi, \psi, n}(\tau_q) - a_{i\infty, 0} - a_{0, 0}\tau_q^{-k}) d(2\pi i\tau_q) = I_k^{\chi, \psi, n}.$$

The second ingredient is the translation formula of Eisenstein series we proved in Proposition 5.12: Let $r = \frac{a}{c} \in \mathbb{Q}$ with $\gcd(a, c) = 1$. Let $E_r^\psi(\tau_q) = E_3^{\mathbf{1}_1, \psi, 1}(\tau_q + r)$. Then

$$E_r^\psi = \sum_{c=c_1 c_2 c_3} \sum_{\chi \in \widehat{(\mathbb{Z}/c_3 \mathbb{Z})^\times}} \frac{\chi(a) \mathfrak{g}(\bar{\chi}) \psi(c_1) c_1^2}{\varphi(c_3)} \cdot E_3^{\chi, \mathbf{1}_{c_2 c_3} \chi \psi, c_1 c_2},$$

where \mathfrak{g} is the Gauss sum and φ is Euler's totient function. It is then clear that the period integrals of Eisenstein series are reduced to sums over arithmetic functions.

We conclude this paper by working out some explicit examples and to conclude the proofs of both Theorem 0.3 and Theorem 0.4, We also include a complete table in Remark 5.18.

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1. GEOMETRY OF TYPE II EXTREMAL TRANSITIONS

1.1. Extremal rays and Type II contractions. We recall some basic terminology from Mori theory, see [22] for a complete treatment.

Let Y be a smooth projective threefold. A *birational contraction* on Y is a projective birational morphism $\phi: Y \rightarrow \bar{Y}$ onto a normal variety \bar{Y} such that $\phi_* \mathcal{O}_Y = \mathcal{O}_{\bar{Y}}$. We are interested in the case that ϕ is *divisorial*, i.e, the exceptional set $\text{Ex}(\phi)$ has codimension one.

Let $\text{NE}(Y)$ be the *integral Mori cone* of effective one cycles, let $\text{NE}(Y)_\mathbb{R} = \text{NE}(Y) \otimes_{\mathbb{Z}} \mathbb{R}$. A divisorial contraction ϕ is called *extremal* in the sense of Mori theory if for every irreducible curve C in Y , we have $\phi_*[C] = 0$ if and only if $[C]$ lies in some fixed ray $\mathbb{R}_{\geq 0} \ell$ in $\text{NE}(Y)_\mathbb{R}$. Such ray $\mathbb{R}_{\geq 0} \ell$ is an *extremal ray* of $\text{NE}(Y)_\mathbb{R}$. It is called *crepant* if the canonical divisor K_Y is ϕ -trivial.

We recall the following terminology from [50, Definition 2.1].

Definition 1.1. A divisorial extremal contraction ϕ is called primitive of Type II if it is crepant and $\dim \phi(\text{Ex}(\phi)) = 0$.

Remark 1.2. Notice that $E = \text{Ex}(\phi)$ is irreducible and therefore contracted to a point [22, Proposition 2.5].

In Wilson's terminology, a birational contraction ϕ is *primitive* if it cannot be factored in the algebraic category [50, Definition 1.2]. As a standard argument shows (cf. [22, Theorem 3.25]), ϕ is primitive if and only if it is a contraction of an extremal ray corresponding to $K_Y + E$. Wilson classified a primitive birational contraction according to the dimensions of its exceptional set and its image. *Types I and III contractions* are small contractions and divisorial contractions to curves, respectively [50, Definition 2.1].

Recall that a *generalized del Pezzo surface* is a projective Gorenstein surface with anti-ample dualizing sheaf (Remark A.3).

Proposition 1.3 ([13, 50, 51]). *Let $\phi: Y \rightarrow \bar{Y}$ be a primitive Type II contraction. Then \bar{Y} is \mathbb{Q} -factorial and Gorenstein, and the exceptional divisor E is a generalized del Pezzo surface. Moreover, if $d := E^3 > 3$ then E is either*

- (1) *a normal and rational del Pezzo surface of degree $d \leq 9$, or*
- (2) *a non-normal del Pezzo surface of degree $d = 7$ whose normalization is a smooth rational ruled surface.*

Proof. Since ϕ is a crepant contraction of an extremal ray, \bar{Y} is \mathbb{Q} -factorial [22, Corollary 3.18]. According to the fact that K_Y is trivial in a neighborhood of E , it descends to \bar{Y} , and thus \bar{Y} is Gorenstein.

Recall that E is irreducible. By the adjunction formula and ϕ being crepant, we get $K_E = E|_E$ and hence $-K_E$ is ample, because $-E$ is ϕ -ample. Also, E is Gorenstein since it is an effective divisor on a smooth threefold. Hence E is a generalized del Pezzo surface (see also [45, Proposition 2.13]). The second part follows from [51, Lemma 2.3] (or from [13, Theorem 5.2] when Y is Calabi–Yau). \square

Remark 1.4. With notation as in Proposition 1.3, let $\phi(E) = \{p\}$. If $d \leq 3$, then the singularity (\bar{Y}, p) is a hypersurface singularity with an equation of the form (cf. [45, Corollary 2.10]):

- (1) $d = 1$, and $x_1^2 + x_2^3 + x_2 f(x_3, x_4) + g(x_3, x_4) = 0$ where f (resp. g) is a sum of monomials of degree at least 4 (resp. 6).
- (2) $d = 2$, and $x_1^2 + f(x_2, x_3, x_4) = 0$ where f is a sum of monomials of degree at least 4.
- (3) $d = 3$, and $x_1^3 + x_1 f(x_2, x_3, x_4) + g(x_2, x_3, x_4) = 0$ where f (resp. g) is a sum of monomials of degree at least 2 (resp. 3).

In this case the generalized del Pezzo surface E is rational or is a cone over an elliptic curve, see [15, Theorem 2.2] and [13, Remark 5.3].

The following proposition will be used in the proof of Proposition 1.11 and the construction of (symplectic) semistable degenerations (§4.1).

Proposition 1.5 ([13, 51]). *Let $\phi: Y \rightarrow \bar{Y}$ be a primitive Type II contraction with exceptional divisor E and $d = E^3$. Then there exists a neighborhood U of E and a holomorphic deformation of the complex structure on U such that E deforms to a smooth del Pezzo surface of degree d .*

Proof. Let $\phi(E) = \{p\}$. If $d \leq 3$, then the singularity (\bar{Y}, p) is a hypersurface singularity with an equation of the form given in Remark 1.4.

Let α be the weight

$$\alpha(x_1, x_2, x_3, x_4) = \begin{cases} (3, 2, 1, 1) & \text{if } d = 1, \\ (2, 1, 1, 1) & \text{if } d = 2, \\ (1, 1, 1, 1) & \text{if } d = 3, \end{cases} \quad (1.1)$$

and $\mathbb{P}(\alpha)$ the weighted projective space. Then the morphism ϕ is the weighted blow-up of the singularity with weight α . If we denote by h_a the weighted homogeneous part of degree a in a power series h , then E is naturally embedded in $\mathbb{P}(\alpha)$ with the equation:

$$\begin{cases} x_1^2 + x_2^3 + x_2 f_4(x_3, x_4) + g_6(x_3, x_4) = 0 & \text{if } d = 1, \\ x_1^2 + f_4(x_2, x_3, x_4) = 0 & \text{if } d = 2, \\ x_1^3 + x_1 f_2(x_2, x_3, x_4) + g_3(x_2, x_3, x_4) = 0 & \text{if } d = 3. \end{cases}$$

It is easily seen that there exists a deformation of the hypersurface singularity (\bar{Y}, p) such that the exceptional locus E' of the deformed singularity is a smooth del Pezzo surface of degree d in $\mathbb{P}(\alpha)$. Note that the singularities can be resolved in the family by means of the weighted blow-up described above.

We may take a good representative $\bar{\mathcal{U}} \rightarrow \Delta$ of the deformation of (\bar{Y}, p) such that on the boundary the family is differentiably trivial (see [33, Theorem 2.8]). If we take $\mathcal{U} \rightarrow \Delta$ to be the corresponding simultaneous resolution⁴, Ehresmann fibration theorem (with boundary) then applies to show that $\mathcal{U} \rightarrow \Delta$ is itself differentiably trivial. Therefore $\mathcal{U} \rightarrow \Delta$ is regarded as a family of holomorphic deformations of the complex structure on the (fixed) neighborhood U of E in Y , and therefore we can deform E to a smooth del Pezzo surface E' .

From now on we will assume that $d > 3$. Let $\iota: E \rightarrow (Y, E)$ be the inclusion map of E into the germ (Y, E) and $\text{Def}((Y, E), E, \iota)$ the deformation space of the inclusion map. By Proposition 1.3 and [13, Lemma 5.6 (iii)], the generalized del Pezzo surface E is smoothable. Moreover, the argument from Theorem 5.8 [13, pp.217-218] shows that the maps of deformation spaces $\text{Def}((Y, E), E, \iota) \rightarrow \text{Def}(E)$ and $\text{Def}((Y, E), E, \iota) \rightarrow \text{Def}(Y, E)$ are surjective. Hence we deduce that a smoothing E' of E can be achieved by holomorphically deforming the complex structure on some neighborhood U of E . \square

We will also need the existence of the holomorphic tubular neighborhood.

Proposition 1.6 ([13, Proposition 5.4]). *Suppose that (\bar{Y}, p) is an isolated rational Gorenstein threefold and that $Y \rightarrow \bar{Y}$ is the blow-up of \bar{Y} at p with exceptional divisor E . If Y and E are smooth and $E^3 \geq 5$, then (\bar{Y}, p) is analytically isomorphic to a cone over E , that is, there exists a neighborhood $U \subseteq Y$ of E which is biholomorphic to a neighborhood of the zero section in the total space of the normal bundle $N_{E/Y}$.*

⁴It is a family of manifolds with boundary where the boundary is differentiably trivial over the disc Δ .

1.2. Type II extremal transitions. We first review the definition of geometric transitions.

Definition 1.7. Given smooth projective threefolds X and Y , we say that X and Y are related by a *geometric transition* if there exists a crepant contraction $\phi: Y \rightarrow \bar{Y}$ followed by a smoothing $\mathfrak{X} \rightarrow \Delta$ of $\bar{Y} = \mathfrak{X}_0$ with the fiber $X = \mathfrak{X}_t$ (for some $t \neq 0$). We write $Y \searrow X$ or $Y \xrightarrow{\phi} \bar{Y} \rightsquigarrow X$ for this process.

A geometric transition $Y \searrow X$ is called:

- (i) a *del Pezzo transition* if $\text{Ex}(\phi)$ is a smooth divisor and $\phi(\text{Ex}(\phi))$ is a point.
- (ii) a *Type II extremal transition* if ϕ is a primitive Type II contraction. In particular, the relative Picard number of ϕ is one.

If the self-intersection number of the exceptional divisor $\text{Ex}(\phi)$ is d , we say that $Y \searrow X$ is of degree d .

Note that the exceptional divisor $\text{Ex}(\phi)$ is a smooth del Pezzo surface of degree d if $Y \searrow X$ is a del Pezzo transition of degree d .

Remark 1.8. If the singularity \bar{Y} is isomorphic to the cone over \mathbb{P}^2 or the Hirzebruch surface \mathbb{F}_1 then it is not smoothable (see [13, (5.5)]). Therefore these two surfaces do not appear in the exceptional divisor of Type II extremal transitions.

The following are the standard examples of del Pezzo transitions, which we will study how the GW invariants are related under such transitions in §3.

Example 1.9. For $1 \leq d \leq 8$, we denote S_d by a smooth del Pezzo surface of degree d which is isomorphic to the blow-up of $\mathbb{P}^1 \times \mathbb{P}^1$ at $8 - d$ points, see Remark A.2. We shall construct a del Pezzo transition $Y_d \searrow X_d$ with exceptional divisor S_d for each d (cf. Remark 1.8).

Let α be the weight as in (1.1) for $1 \leq d \leq 3$ and

$$\alpha(x_1, \dots, x_{d+1}) = (1, \dots, 1) \quad (1.2)$$

for $4 \leq d \leq 8$. Then we have an embedding $S_d \hookrightarrow \mathbb{P}(\alpha)$ with

$$-K_{S_d} = \mathcal{O}_{\mathbb{P}(\alpha)}(1)|_{S_d}. \quad (1.3)$$

Let $\text{Proj}_p: \mathbb{P}(\alpha, 1) \dashrightarrow \mathbb{P}(\alpha)$ be the rational map of weighted projective spaces given by the projection from the point $p = [0 : \dots : 0 : 1] \in \mathbb{P}(\alpha, 1)$. It is easily seen that the projective bundle $\mathbb{P}_{\mathbb{P}(\alpha)}(\mathcal{O}(-1) \oplus \mathcal{O})$ over $\mathbb{P}(\alpha)$ is the graph of the rational map Proj_p . Note that $\mathbb{P}_{\mathbb{P}(\alpha)}(\mathcal{O}(-1) \oplus \mathcal{O})$ is also isomorphic to the weighted blow-up of $\mathbb{P}(\alpha, 1)$ at p with the weight $(\alpha, 1)$.

Now we define \bar{Y}_d the projective cone over S_d with vertex p in $\mathbb{P}(\alpha, 1)$ and

$$Y_d = \mathbb{P}_{S_d}(K_{S_d} \oplus \mathcal{O}).$$

According to (1.3), it follows that Y_d is the weighted blow-up $\mathbb{P}_{S_d}(\mathcal{O}(-1) \oplus \mathcal{O})$ of \bar{Y}_d at p with the weight α and thus S_d is the exceptional divisor of $Y_d \rightarrow \bar{Y}_d$. Note that $Y_d \rightarrow \bar{Y}_d$ is

the restriction of the weighted blow-up $\mathbb{P}_{\mathbb{P}(\alpha)}(\mathcal{O}(-1) \oplus \mathcal{O}) \rightarrow \mathbb{P}(\alpha, 1)$.

$$\begin{array}{ccccc}
Y_d & \xhookrightarrow{\quad} & \mathbb{P}_{\mathbb{P}(\alpha)}(\mathcal{O}(-1) \oplus \mathcal{O}) & \xleftarrow{\quad} & \\
\downarrow & & \swarrow & \searrow & \\
\bar{Y}_d & \xhookrightarrow{\quad} & \mathbb{P}(\alpha, 1) & \xrightarrow{\text{Proj}_p} & \mathbb{P}(\alpha) \xleftarrow{\quad} S_d
\end{array}$$

Let X_d be a smooth del Pezzo threefold of degree d in $\mathbb{P}(\alpha, 1)$. Then X_d is a smoothing of the singular del Pezzo threefold \bar{Y}_d and therefore we get the desired del Pezzo transition $Y_d \setminus X_d$.

We remark that, when $d = 6$ there are two cases (6I) and (6II) in Theorem A.5, we will denote $X_{6\text{I}}$ and $X_{6\text{II}}$, respectively. It will be distinguished by the topological Euler characteristics, see Table 2.

We shall construct the (local) semistable projective degenerations, which will be used in §4.1.

Proposition 1.10. *Let E be a smooth del Pezzo surface which is isomorphic to a blow-up of $\mathbb{P}^1 \times \mathbb{P}^1$ at $8 - d$ points. Assume that (\bar{Y}, p) is an isolated singularity which is analytically isomorphic to a cone over E . Let $\mathfrak{V} \rightarrow \Delta$ be a good representative of a smoothing of (\bar{Y}, p) with the special fiber $\bar{U} := \mathfrak{V}_0$. Given the weight α as in (1.1) and (1.2), we denote $\phi: U \rightarrow \bar{U}$ by the weighted blow-up at p with weight α . Set $n_d = 6, 4, 3$ for $d = 1, 2, 3$, respectively, and $n_d = 2$ for $d \geq 4$.*

If U is smooth, then there exists a semistable reduction $\mathcal{V} \rightarrow \Delta$ of $\mathfrak{V} \rightarrow \Delta$ which is obtained by a degree n_d base change $\mathcal{V}' \rightarrow \Delta$ allowed by the weighted blow-up at $p \in \mathcal{V}'$ with weight $(\alpha, 1)$. Moreover, the special fiber $\mathcal{V}_0 = U \cup X_d$ is a simple normal crossing divisor such that X_d is a smooth del Pezzo threefold of degree d and $U \cap X_d = E$.

Proof. Note that the exceptional divisor of ϕ is the smooth del Pezzo surface E of degree d by assumption. Let $C(E)$ be the affine cone with vertex o over E , that is,

$$C(E) = \text{Proj} \left(\bigoplus_{m \geq 0} H^0(E, \mathcal{O}(-mK)) \right). \quad (1.4)$$

By Proposition 1.6 and Theorem A.1, we have

$$\hat{\mathcal{O}}_{\bar{Y}, p} \cong \hat{\mathcal{O}}_{C(E), o} \cong \mathbb{C}[\underline{x}] / \mathcal{I}$$

and E is defined by the generator of \mathcal{I} . Here we write \underline{x} a short form for indeterminates $\{x_m\}$. Moreover, $\mathcal{I} = (F)$ if $1 \leq d \leq 3$, and \mathcal{I} is generated by $d(d-3)/2$ quadrics $\{Q_k\}$ if $d \geq 4$.

Let us denote by t the coordinate of the disk Δ and let t be weight one. We let $\mathcal{V}' \rightarrow \Delta$ be a degree n_d base change of $\mathfrak{V} \rightarrow \Delta$, and let \mathcal{V} be the weighted blow-up of \mathcal{V}' at p with weight $(\alpha, 1)$. Note that $\mathcal{V} \rightarrow \mathcal{V}'$ is an ordinary blow-up at p if $d \geq 3$. We claim that $\mathcal{V} \rightarrow \Delta$ is the

desired degeneration. Indeed, we observe that

$$\hat{\mathcal{O}}_{\mathfrak{V}, p} \cong \begin{cases} \mathbb{C}[[x_1, x_2, x_3, x_4, t]]/(F - t) & \text{if } 1 \leq d \leq 3, \\ \mathbb{C}[[x_1, \dots, x_{d+1}, t]]/(Q_1 - t, \dots, Q_{d(d-3)/2} - t) & \text{if } 4 \leq d \leq 8. \end{cases}$$

We will denote X_d by the exceptional divisor of the weighted blow-up \mathcal{V} .

For $1 \leq d \leq 3$, we find that the defining equation of (\mathcal{V}', p) is $F - t^{n_d} = 0$, which is a weighted homogeneous polynomial of degree n_d with weights $(\alpha, 1)$. Then, by construction, the exceptional divisor X_d of $\mathcal{V} \rightarrow \mathcal{V}'$ is the hypersurface in $\mathbb{P}(\alpha, 1)$ defined by $F - t^{n_d} = 0$ and therefore X_d is a smooth del Pezzo threefold of degree d . We remark that if t is regarded as a section $t \in H^0(X_d, \mathcal{O}(1))$ then it defines the del Pezzo surface $E = (t = 0)$ in X_d . Also, according to that U is the weighted blow-up of \bar{U} at p with weight α , it follows that the proper transform of $\mathcal{V}'_0 = \bar{U}$ in \mathcal{V} is the open neighborhood U of E . Hence we conclude that the special fiber \mathcal{V}_0 of $\mathcal{V} \rightarrow \Delta$ over $t = 0$ consists of two irreducible components U and X_d , which are smooth and $U \cap X_d = (F = 0)$ is the del Pezzo surface E .

Similarly, for $d \geq 4$, the exceptional divisor X_d of $\mathcal{V} \rightarrow \mathcal{V}'$ is defined by quadrics $Q_k - t^2$ in \mathbb{P}^{d+1} . It is easily seen that X_d is a smooth threefold, which contains the del Pezzo surface E of degree d defined by the section $t \in H^0(X_d, \mathcal{O}(1))$. By the Lefschetz theorem, the restriction $\text{Pic}(X_d) \rightarrow \text{Pic}(E)$ is injective. Hence, using the adjunction formula, we get that $\mathcal{O}_{X_d}(-K) \cong \mathcal{O}_{X_d}(2)$, and thus X_d is del Pezzo of degree d . The remaining part follows from the same argument as above. \square

1.3. Global topology. In this subsection we describe topological properties of Type II extremal transitions.

Proposition 1.11. *If $Y \searrow X$ is a Type II extremal transition of degree d , then*

$$\chi_{\text{top}}(Y) - \chi_{\text{top}}(X) = 2\chi_{\text{top}}(S_d) - \chi_{\text{top}}(X_d),$$

where S_d (resp. X_d) is a smooth del Pezzo surface (resp. threefold) of degree d and $\chi_{\text{top}}(-)$ is the topological Euler characteristic.

Proof. Let $\phi: Y \rightarrow \bar{Y}$ be the crepant contraction of $Y \searrow X$, E the exceptional divisor of ϕ and $\phi(E) = \{p\}$. Since the question is topological, we apply Proposition 1.5 and replace an open neighborhood of E by another open set that is diffeomorphic to it. So we may assume that E is a smooth del Pezzo surface of degree d and (\bar{Y}, p) is analytically isomorphic to the cone (1.4) over E .

Let B be the Milnor fiber of the smoothing of (\bar{Y}, p) corresponding to X . Note that B is homotopy equivalent to a finite cell complex of $\dim \leq 3$ by [33, (5.6)]. It is checked that B is homeomorphic to $X_d \setminus (X_d \cap A)$, where X_d is a smooth del Pezzo threefold of degree d and $A \in |-K_{X_d}|$ a general hyperplane. Write $S_d = X_d \cap A$, a smooth del Pezzo surface of degree d , and let $\iota: X_d \setminus S_d \rightarrow X_d$ be the inclusion. By using the Gysin sequence

$$\cdots \rightarrow H^k(X_d) \xrightarrow{\iota^*} H^k(X_d \setminus S_d) \xrightarrow{\text{Res}} H^{k-1}(S_d) \rightarrow H^{k+1}(X_d) \rightarrow \cdots,$$

we have

$$\chi_{\text{top}}(B) = \chi_{\text{top}}(X_d) - \chi_{\text{top}}(S_d). \quad (1.5)$$

Also, by Theorems A.1 and A.5, we get the Betti numbers $b_0(B) = 1$, $b_2(B) = h^{1,1}(X_d) - 1$,

$$b_3(B) = 2h^{2,1}(X_d) - h^{1,1}(X_d) + 10 - d$$

and $b_k(B) = 0$ for other k . Note that $b_1(B) = 0$ by [12, Theorem 2].

Since \bar{Y} is Y with E collapsed to the point p (\bar{Y} is homeomorphic to Y/E), the relative cohomology $H^i(Y, E)$ is isomorphic to the reduced cohomology $\widetilde{H}^k(Y/E) \cong \widetilde{H}^k(\bar{Y})$ for all k . By the long exact sequence of the pair (Y, E) , we get

$$\cdots \rightarrow \widetilde{H}^k(\bar{Y}) \xrightarrow{\phi^*} H^k(Y) \rightarrow H^k(E) \rightarrow \widetilde{H}^{k+1}(\bar{Y}) \rightarrow \cdots. \quad (1.6)$$

On the other hand, one can construct a homeomorphism between \bar{Y} and X/B_d . We denote the collapsing (continuous) map by

$$c: X \rightarrow \bar{Y}. \quad (1.7)$$

Hence there also exists the long exact sequence

$$\cdots \rightarrow \widetilde{H}^k(\bar{Y}) \xrightarrow{c^*} H^k(X) \rightarrow H^k(B_d) \rightarrow \widetilde{H}^{k+1}(\bar{Y}) \rightarrow \cdots. \quad (1.8)$$

By (1.6) and (1.8), we get

$$\chi_{\text{top}}(X) - \chi_{\text{top}}(B) = \chi_{\text{top}}(Y) - \chi_{\text{top}}(E).$$

Then the proposition follows from $\chi_{\text{top}}(E) = \chi_{\text{top}}(S_d)$ and (1.5). \square

Write $\Delta\chi_{\text{top}} := \chi_{\text{top}}(Y) - \chi_{\text{top}}(X)$, the topological difference of a Type II extremal transition $Y \searrow X$ of degree d . We give $\Delta\chi_{\text{top}}$ in Table 2 by using Theorem A.1 and Table 9.

d	1	2	3	4	5	6I	6II	7	8
$\Delta\chi_{\text{top}}$	60	36	24	16	10	6	4	4	4

TABLE 2. The topological difference $\Delta\chi_{\text{top}}$ of $Y \searrow X$.

Remark 1.12. The differences $\frac{1}{2}\Delta\chi_{\text{top}}$ are exactly the dual Coxeter numbers of E_{9-d} (cf. [39, pp.241–242]). The root system is summarized in Table 3. Here we set $E_1 = E_2 = A_1$, and root systems E_7 , E_6 , E_5 , E_4 and E_3 are obtained from E_8 by deleting one by one the vertices from the long end of its Dynkin diagram. For $d \leq 7$, these root systems are associated to a smooth del Pezzo surfaces of degree d , see Proposition 25.2 and Theorem 25.4 in [35].

For the case of $d = 6$, the transition of degree $d = 6I$ (resp. $d = 6II$) corresponds to the direct summand A_2 (resp. A_1) of the root system $A_2 \times A_1$. The dual Coxeter number of A_2 and A_1 are 3 and 2, respectively.

d	1	2	3	4	5	6	7	8
Type of E_{9-d}	E_8	E_7	E_6	D_5	A_4	$A_2 \times A_1$	A_1	A_1

TABLE 3. Root systems and del Pezzo surfaces.

The following proposition relates cohomology groups of X and Y .

Proposition 1.13. *Let $Y \xrightarrow{\phi} \bar{Y} \rightsquigarrow X$ be a Type II extremal transition with exceptional divisor E .*

(i) *Let $\phi(E) = \{p\}$, and let B be the Milnor fiber of the smoothing of (\bar{Y}, p) corresponding to X . Then there exists an exact sequence*

$$0 \rightarrow H^2(B) \rightarrow H^3(\bar{Y}) \rightarrow H^3(X) \rightarrow H^3(B) \rightarrow 0,$$

and $H^k(\bar{Y}) \xrightarrow{\sim} H^k(X)$ for all $k \neq 3$.

(ii) *There exist exact sequences*

$$0 \rightarrow H^2(\bar{Y}) \xrightarrow{\phi^*} H^2(Y) \rightarrow \mathbb{C}[K_E] \subseteq H^2(E), \quad (1.9)$$

$$0 \rightarrow H^2(E)/\mathbb{C}[K_E] \rightarrow H^3(\bar{Y}) \rightarrow H^3(Y) \rightarrow H^3(E) \rightarrow 0, \quad (1.10)$$

$$0 \rightarrow H^4(\bar{Y}) \rightarrow H^4(Y) \rightarrow H^4(E) = \mathbb{C}[\text{pt}] \rightarrow 0, \quad (1.11)$$

and $H^k(\bar{Y}) \xrightarrow{\sim} H^k(Y)$ if $k \neq 2, 3, 4$.

Proof. To prove (i), we use the long exact sequence (1.8) induced by the topological pair (X, B) . Recall that B is homotopy equivalent to a finite cell complex of $\dim \leq 3$, and $H^1(B) = 0$. Then we get $H^k(\bar{Y}) \xrightarrow{\sim} H^k(X)$ for $k = 0, 1, 5, 6$, and $H^k(\bar{Y}) \rightarrow H^k(X)$ is injective (resp. surjective) for $k = 2$ (resp. $k = 4$). By Poincaré duality, we get

$$b_2(\bar{Y}) \leq b_2(X) = b_4(X) \leq b_4(\bar{Y}). \quad (1.12)$$

Since \bar{Y} has rational singularities and is \mathbb{Q} -factorial, we have

$$b_2(\bar{Y}) = b_4(\bar{Y}) \quad (1.13)$$

by [43, Theorem A]. Then (i) follows from (1.12) and (1.13).

The proof of (ii) uses the other long exact sequence (1.6) induced by the topological pair (Y, E) . Since ϕ is an extremal contraction and \bar{Y} has rational singularities, we have $\phi^*: H^2(\bar{Y}) \rightarrow H^2(Y)$ is injective with $\dim \text{Coker}(\phi^*) = 1$. Therefore the image of the restriction map $H^2(Y) \rightarrow H^2(E)$ is generated by the class of $E|_E = K_E$ and (1.9) follows.

On the other hand, since \bar{Y} has only an isolated singularity, the transition $Y \searrow X$ does not change the fundamental group and thus the first Betti number, i.e., $b_1(X) = b_1(Y)$. Also, $b_1(\bar{Y}) = b_1(X)$ by (i). Therefore the injective map $H^1(\bar{Y}) \rightarrow H^1(Y)$ is an isomorphism. The assertion for $k = 0$ is trivial. Observe that $H^4(Y) \rightarrow H^4(E)$ is surjective, since it is a nonzero map and $H^4(E) = \mathbb{C}[\text{pt}]$. Then $H^k(\bar{Y}) \xrightarrow{\sim} H^k(Y)$ for $k = 5, 6$. To finish (i),

we claim that $b_4(\bar{Y}) = b_4(Y) - 1$, this will give (1.10) and (1.11). To see this, we already know $b_2(\bar{Y}) = b_2(Y) - 1$. The claim then follows from Poincaré duality, $b_2(Y) = b_4(Y)$, and (1.13). \square

Remark 1.14. For a Type II extremal transition $Y \searrow X$, we have

$$h^{1,1}(Y) - h^{1,1}(X) = 1, \quad (1.14)$$

$$h^{2,1}(X) - h^{2,1}(Y) = \frac{1}{2}\Delta\chi_{\text{top}} - 1. \quad (1.15)$$

by Propositions 1.11 and 1.13. Indeed, the transition $Y \searrow X$ preserves canonical bundles and thus the Hodge number $h^{3,0}$. Note that it also preserves the Hodge numbers $h^{p,0} = h^0(\Omega^p)$ for $p = 1, 2$.

When X and Y are smooth Calabi–Yau threefolds and the exceptional divisor E of the transition is smooth, (1.14) is Proposition 3.1 in [20] by applying Kollar–Mori’s results in [21, §12.1 and 12.2] (see also [14, Proposition 3.1] and [40, Proposition 6.1]). In addition, (1.15) was proved in [20, Theorem 3.3, Remark 3.6] (see also [41, p.761] for $E^3 = 6$). Moreover, if $E^3 \geq 5$, then the miniversal Kuranishi space of \bar{Y} is the product of the miniversal space of the singularity (\bar{Y}, p) , analytically isomorphic to the cone (1.4) over E , with an appropriate germ of a linear space [20, Remark 3.6].

Remark 1.15. It is interesting to compare (1.15) with the situation of conifold transitions, see Lemma 1.6 in [24].

Corollary 1.16. *Let $Y \xrightarrow{\phi} \bar{Y} \rightsquigarrow X$ be a Type II extremal transition. Then there exists an embedding*

$$H(X) \hookrightarrow H(Y) \quad (1.16)$$

of cohomology rings which we abuse notation slightly and denote by ϕ^* , i.e., it preserves the cup product.

Proof. By Proposition 1.13 (i), the collapsing (continuous) map $c: X \rightarrow \bar{Y}$ (1.7) induces a ring isomorphism $c^*: H(\bar{Y}) \rightarrow H(X)$. On the other hand, the ring homomorphism $\phi^*: H(\bar{Y}) \rightarrow H(Y)$ is injective by Proposition 1.13 (ii). Therefore the composition $\phi^* \circ (c^*)^{-1}$ gives the desired embedding (1.16). \square

Remark 1.17. Let $\mathfrak{X} \rightarrow \Delta$ be the smoothing of the transition $Y \searrow X$. By Proposition 1.13 (i) and the fact that \bar{Y} is a deformation retract of \mathfrak{X} , we get

$$H(\mathfrak{X}) \cong H(\bar{Y}) \xrightarrow{\sim} H(X). \quad (1.17)$$

Hence every (even) cohomology class of X has a lifting in $H(\mathfrak{X})$.

Remark 1.18. Using the argument of the proof of Proposition 1.13, one can infer that $b_4(\bar{Y}_d) > b_2(\bar{Y}_d)$ for the canonical local model $Y_d \rightarrow \bar{Y}_d \rightsquigarrow X_d$, see Example 1.9, for all $1 \leq d \leq 8$. In particular, \bar{Y}_d is non- \mathbb{Q} -factorial (cf. [42, Theorem 3.2] and [43, Theorem A]).

2. PRELIMINARIES ON GROMOV–WITTEN THEORY

2.1. Gromov–Witten invariants. We consider only *genus zero* GW invariants.

Let M be a smooth projective variety. For an effective curve class $\beta \in \mathrm{NE}(M)$, we denote by $\overline{\mathcal{M}}_{0,n}(M, \beta)$ the moduli space of n -pointed stable maps $f: (C, p_1, \dots, p_n) \rightarrow X$ with arithmetic genus $g(C) = 0$ and with degree $f_*[C] = \beta$. Let ψ_ν be the first Chern class of the ν -th universal cotangent line bundle on $\overline{\mathcal{M}}_{0,n}(M, \beta)$ and $\mathrm{ev}_\nu: \overline{\mathcal{M}}_{0,n}(M, \beta) \rightarrow M$ the evaluation map $[f] \mapsto f(p_\nu)$. For $\vec{a} = (a_1, \dots, a_n) \in H(M)^{\oplus n}$ and integers $k_1, \dots, k_n \geq 0$, we define the (genus zero, descendant) GW invariant as

$$\langle \psi^{k_1} a_1, \dots, \psi^{k_n} a_n \rangle_\beta^M = \int_{[\overline{\mathcal{M}}_{0,n}(M, \beta)]^\mathrm{vir}} \prod_{\nu=1}^n \mathrm{ev}_\nu^* a_\nu \cup \psi_\nu^{k_\nu}$$

where $[\overline{\mathcal{M}}_{0,n}(M, \beta)]^\mathrm{vir}$ denotes the *virtual fundamental class* on $\overline{\mathcal{M}}_{0,n}(M, \beta)$ [4, 29]. The virtual dimension is

$$\mathrm{vdim}_\beta := (c_1(M), \beta) + n + (\dim M - 3).$$

Let q^β be the (formal) Novikov variable corresponding to $\beta \in \mathrm{NE}(M)$. For a parameter $\tau \in H^2(M)$, set $Q^\beta = q^\beta e^{(\tau, \beta)}$. We define the *n -point function* (in τ) by

$$\langle a_1, \dots, a_n \rangle^M = \sum_{\beta} \langle a_1, \dots, a_n \rangle_\beta^M Q^\beta.$$

Given $D \in H^2(M)$, we define the *power operator* δ^D by setting

$$\delta^D Q^\beta = (D, \beta) Q^\beta \quad \text{for each } \beta \in \mathrm{NE}(M)$$

and extending linearly.

Applying [25, Corollary 1] and the WDVV equations, we have the following relations between multipoint functions, known as the reconstruction formula.

Theorem 2.1 ([25]). *Let $(a_1, a_2, \dots, a_n) \in H(M)^{\oplus n}$ and $D \in H^2(M)$. Pick a basis $\{T_\mu\}$ of $H^*(M)$ with dual basis $\{T^\mu\}$. By abuse of notation, we denote the insertions $\bigotimes_{\nu \in I} a_\nu$ briefly by \vec{a}_I for any subset I of $\{1, \dots, n\}$.*

(a) *For $n \geq 2$, we have*

$$\begin{aligned} \langle Da_1, a_2, \dots, a_n \rangle^M &= \langle a_1, Da_2, \dots, a_n \rangle^M + \delta^D \langle a_1, \psi a_2, \dots, a_n \rangle^M \\ &\quad - \sum_{I_L \sqcup I_R = \{3, \dots, n\}} \delta^D \langle a_1, \vec{a}_{I_L}, T_\mu \rangle^M \cdot \langle T^\mu, a_2, \vec{a}_{I_R} \rangle^M. \end{aligned}$$

(b) For $n \geq 3$, we have

$$\begin{aligned} \langle Da_1, a_2, \dots, a_n \rangle^M &= \langle a_1, Da_2, \dots, a_n \rangle^M \\ &\quad + \sum_{3 \in I_L, 2 \in I_R} \langle a_1, \vec{a}_{I_L}, T_\mu \rangle^M \cdot \delta^D \langle T^\mu, \vec{a}_{I_R} \rangle^M \\ &\quad - \sum_{2, 3 \in I_R} \delta^D \langle a_1, \vec{a}_{I_L}, T_\mu \rangle^M \cdot \langle T^\mu, \vec{a}_{I_R} \rangle^M, \end{aligned}$$

where $I_L \sqcup I_R = \{2, 3, \dots, n\}$ is a partition.

2.2. Symmetrized Gromov–Witten. Let M be a smooth projective variety. Suppose that there is a smooth projective algebraic family $\mathcal{M} = \{M_t\}$ such that $M_0 \cong M \cong M_1$. Then \mathcal{M} induces an automorphism on $H(M)$:

$$\sigma: H(M) \cong H(M_0) \xleftarrow{\sim} H(\mathcal{M}) \xrightarrow{\sim} H(M_1) \cong H(M).$$

Lemma 2.2. For each $\vec{a} \in H(M)^{\oplus n}$, we have

$$\langle \sigma(\vec{a}) \rangle^M = \sigma \langle \vec{a} \rangle^M,$$

where $\sigma(Q^\beta) = Q^{\sigma(\beta)}$.

Proof. By comparing the coefficient of $Q^{\sigma(\beta)}$ on both sides, we need to show that $\langle \sigma(\vec{a}) \rangle_{\sigma(\beta)}^M = \langle \vec{a} \rangle_\beta^M$ holds for each β . This follows from the fact that the GW invariant is deformation invariant. \square

Let $G \subseteq \text{Aut}(H(M))$ be a finite group that consists of some of the automorphisms σ that came from deformations as above. Let $H(M)^G$ be the G -invariant part of $H(M)$ and K the kernel of the symmetrization

$$\begin{aligned} H(M) &\xrightarrow{\Sigma} H(M)^G \\ a &\longmapsto \frac{1}{|G|} \sum_{\sigma \in G} \sigma(a). \end{aligned}$$

Then $H(M) = H(M)^G \oplus K$ is an orthogonal direct sum with respect to the Poincaré pairing.

If $\tau^{(2)} \in H^2(M)^G$, then $e^{(\tau^{(2)}, \beta)} = e^{(\tau^{(2)}, \sigma(\beta))}$. So we may identify the Novikov variables $Q^\beta = Q^{\sigma(\beta)}$. We also denote the identification by Σ :

$$\begin{aligned} \mathbb{C}[[Q^{\text{NE}(M)}]] &\xrightarrow{\Sigma} \mathbb{C}[[Q^{\text{NE}(M)/G}]] \\ Q^\beta &\longmapsto Q^{G\beta}. \end{aligned}$$

Define the G -symmetrized invariant as

$$\langle \vec{a} \rangle_G^M = \Sigma \langle \vec{a} \rangle^M, \quad \vec{a} \in H(M)^{\oplus n}. \quad (2.1)$$

By Lemma 2.2, $\langle \vec{a} \rangle_G^M$ is G -invariant, i.e.,

$$\langle \sigma(\vec{a}) \rangle_G^M = \langle \vec{a} \rangle_G^M, \quad \vec{a} \in H(M)^{\oplus n}, \quad \sigma \in G. \quad (2.2)$$

Proposition 2.3. *The reconstruction formula for symmetrized invariants holds, i.e., if $(a_1, a_2, \dots, a_n) \in (H(M)^G)^{\oplus n}$, $D \in H^2(M)^G$ and $\{T_\mu\}, \{T^\mu\}$ are dual bases of $H(M)^G$, then,*

(a) *for $n \geq 2$, we have*

$$\begin{aligned} \langle Da_1, a_2, \dots, a_n \rangle_G^M &= \langle a_1, Da_2, \dots, a_n \rangle_G^M + \delta^D \langle a_1, \psi a_2, \dots, a_n \rangle_G^M \\ &\quad - \sum_{I_L \sqcup I_R = \{3, \dots, n\}} \delta^D \langle a_1, \vec{a}_{I_L}, T_\mu \rangle_G^M \cdot \langle T^\mu, a_2, \vec{a}_{I_R} \rangle_G^M; \end{aligned}$$

(b) *for $n \geq 3$, we have*

$$\begin{aligned} \langle Da_1, a_2, \dots, a_n \rangle_G^M &= \langle a_1, Da_2, \dots, a_n \rangle_G^M \\ &\quad + \sum_{3 \in I_L, 2 \in I_R} \langle a_1, \vec{a}_{I_L}, T_\mu \rangle_G^M \cdot \delta^D \langle T^\mu, \vec{a}_{I_R} \rangle_G^M \\ &\quad - \sum_{2, 3 \in I_R} \delta^D \langle a_1, \vec{a}_{I_L}, T_\mu \rangle_G^M \cdot \langle T^\mu, \vec{a}_{I_R} \rangle_G^M, \end{aligned}$$

where $I_L \sqcup I_R = \{2, 3, \dots, n\}$ is a partition.

Proof. We will only prove (a), since (b) follows from (a). Let $\{\tilde{T}_\nu\}, \{\tilde{T}^\nu\}$ be dual bases of K . Since $H(M) = H(M)^G \oplus^\perp K$, the original reconstruction formula (cf. Theorem 2.1) gives

$$\begin{aligned} \langle Da_1, a_2, \dots, a_n \rangle_G^M &= \langle a_1, Da_2, \dots, a_n \rangle_G^M + \delta^D \langle a_1, \psi a_2, \dots, a_n \rangle_G^M \\ &\quad - \sum_{I_L \sqcup I_R = \{3, \dots, n\}} \delta^D \langle a_1, \vec{a}_{I_L}, T_\mu \rangle_G^M \cdot \langle T^\mu, a_2, \vec{a}_{I_R} \rangle_G^M \\ &\quad - \sum_{I_L \sqcup I_R = \{3, \dots, n\}} \delta^D \langle a_1, \vec{a}_{I_L}, \tilde{T}_\nu \rangle_G^M \cdot \langle \tilde{T}^\nu, a_2, \vec{a}_{I_R} \rangle_G^M. \end{aligned}$$

To prove the statement, it suffices to show that for $\tilde{T} \in K$, $\langle \tilde{T}, \vec{a}_I \rangle_G^M = 0$. Indeed, by (2.2),

$$\langle \tilde{T}, \vec{a}_I \rangle_G^M = \frac{1}{|G|} \sum \langle \sigma(\tilde{T}), \vec{a}_I \rangle_G^M = \langle \Sigma(\tilde{T}), \vec{a}_I \rangle_G^M = 0. \quad \square$$

Recall that S_d is a del Pezzo surface of degree d (Example 1.9). Let \mathcal{C} be the minimal set of generators of $\text{NE}(S_d)$, the Mori cone of S_d . When $d \leq 7$, \mathcal{C} is the set of (-1) -curves on S_d by cone theorem. When $d = 8$, $\mathcal{C} = \{\ell_1, \ell_2\}$, the fibers of the projections $p_1, p_2: S_8 = \mathbb{P}^1 \times \mathbb{P}^1 \rightarrow \mathbb{P}^1$.

The following proposition will be used in §3.2.

Proposition 2.4. *Suppose that $d \neq 7$. For any pairs of (-1) -curves $\ell, \ell' \in \mathcal{C}$, there is a deformation S_d such that $S_{d,0} \cong S_d \cong S_{d,1}$ and the induced automorphism σ maps ℓ to ℓ' .*

In particular, there exists a finite group $G_{S_d} \subseteq \text{Aut}(H(S_d))$ such that G_{S_d} acts on \mathcal{C} transitively.

Proof. If $d = 8$, then $(\mathbb{P}^1 \times \mathbb{P}^1, \mathbb{P}^1 \times *)$ is clearly isomorphic to $(\mathbb{P}^1 \times \mathbb{P}^1, * \times \mathbb{P}^1)$ by interchanging two components.

If $d \leq 6$, let $(S_d, \ell) \rightarrow (S_{d+1}, p)$, $(S_d, \ell') \rightarrow (S'_{d+1}, p')$ be the contractions of ℓ, ℓ' , respectively. There is a deformation that links (S_{d+1}, p) and (S'_{d+1}, p') . Indeed, there is a family $\mathcal{S} = \{S(t)\}_t$ on the moduli space of del Pezzo surfaces of degree $d+1$ such that $S(0) = S_{d+1}$ and $S(1) = S'_{d+1}$. Then take any path $p(t) \in S(t)$ on the family \mathcal{S} that connects $p = p(0) \in S(0)$ and $p' = p(1) \in S(1)$. Let $\ell(t)$ be the exceptional curve of the blow up $\text{Bl}_{p(t)} S(t) \rightarrow S(t)$. Then the family $(\text{Bl}_{p(t)} S(t), \ell(t))$ is the desired deformation. \square

Remark 2.5. For $1 \leq d \leq 6$, the group G_{S_d} is the Weyl group $W(E_{9-d})$ of the root system E_{9-d} in Table 3, see for example [35, Theorem 23.9].

2.3. I-function and mirror theorem. For $\tau = \tau^{(0)} + \tau^{(2)} \in H^0(M) \oplus H^2(M)$, we define the (small) *J-function* of M to be the generating function

$$\begin{aligned} J^M(\tau, z^{-1}) &:= e^{\tau/z} \sum_{\beta \geq 0} J_\beta^M(z^{-1}) Q^\beta \\ &:= e^{\tau/z} \left(1 + \sum_{\beta > 0} \text{ev}_{1*} \left(\frac{1}{z(z - \psi_1)} \cap [\overline{\mathcal{M}}_{0,1}(M, \beta)]^{\text{vir}} \right) Q^\beta \right) \\ &= e^{\tau/z} \left(1 + \sum_{\beta > 0, k \geq 0} \frac{1}{z^{k+2}} \langle \psi^k T_\mu \rangle_\beta T^\mu Q^\beta \right), \end{aligned}$$

where $Q^\beta = e^{(\tau^{(2)}, \beta)}$.

Let $\mathcal{V} = \bigoplus_m \mathcal{L}_m$ be a vector bundle over M and $i: Z \hookrightarrow M$ a smooth subvariety defined by a section of \mathcal{V} . We further assume that \mathcal{V} is *convex*, i.e., if $H^1(\mathbb{P}^1, \varphi^* \mathcal{V}) = 0$ for any holomorphic map $\varphi: \mathbb{P}^1 \rightarrow M$. For $t = t^{(0)} + t^{(2)} \in H^0(M) \oplus H^2(M)$, we define the *I-function* of $Z \hookrightarrow M$ to be the generating function

$$\begin{aligned} I^Z(t, z) &:= e^{t/z} \sum_{\beta \geq 0} I_\beta^Z(z) P^\beta \\ &:= e^{t/z} \sum_{\beta \geq 0} (J_\beta^M(z^{-1}) \cdot \mathcal{V}^{\bar{\beta}}) P^\beta \\ &:= e^{t/z} \sum_{\beta \geq 0} \left(J_\beta^M(z^{-1}) \cdot \prod_m (c_1(\mathcal{L}_m))^{\overline{(c_1(\mathcal{L}_m), \beta)}} \right) P^\beta, \end{aligned}$$

where $P^\beta = e^{(t^{(2)}, \beta)}$ and

$$(D)^{\bar{n}} := \prod_{\nu=1}^n (D + \nu z) := \frac{\prod_{\nu=-\infty}^n (D + \nu z)}{\prod_{\nu=-\infty}^0 (D + \nu z)}.$$

The following classical mirror theorem allows us to compute the GW invariants of $Z \subseteq M$.

Theorem 2.6 ([31, 6]). *Suppose that $c_1(M) \geq c_1(\mathcal{V}) = \sum_m c_1(\mathcal{L}_m)$. Then*

$$i_* J^Z(\tau, z^{-1}) = e(\mathcal{V}) \cdot I^Z(t, z)$$

under some mirror transform $\tau = \tau(t)$, where $i_* Q^\beta = Q^{i_* \beta}$ for $\beta \in \text{NE}(Z)$. More precisely, $\tau^{(0)} = t^{(0)} + zg^{(0)}(P)$, $\tau^{(2)} = t^{(2)} + g^{(2)}(P)$ for some power series $g^{(0)}$, $g^{(2)}$ that is fully determined by I^Z .

When M is a weighted projective space with singularities, we can still use

Theorem 2.7 ([7, 44]). *Suppose that M is a weighted projective space $\mathbb{P}(\alpha_1, \dots, \alpha_{k+1})$, then Theorem 2.6 holds if Z is a smooth complete intersection of hypersurfaces Z_1, \dots, Z_r which do not intersect the singular locus of M with*

$$J_\beta^M(z^{-1}) := \frac{1}{\prod_{n=1}^{k+1} (\alpha_n h)^{(\alpha_n h, \beta)}}, \quad \mathcal{V} = \bigoplus_{m=1}^r \mathcal{O}(Z_m),$$

where $h = c_1(\mathcal{O}(1))$ is the hyperplane class of M .

We will use the following theorem due to Brown, which give us the I -function of projective bundle $\mathbb{P}_M(\mathcal{V} \oplus \mathcal{O})$ of a split vector bundle $\mathcal{V} \oplus \mathcal{O} = \bigoplus_m \mathcal{L}_m \oplus \mathcal{O}$.

Theorem 2.8 ([5, Corollary 2]). *Let γ be the fiber curve class of $\pi: \mathbb{P}_M(\mathcal{V} \oplus \mathcal{O}) \rightarrow M$, and let $\xi = c_1(\mathcal{O}_{\mathbb{P}_M(\mathcal{V} \oplus \mathcal{O})}(1)) = [\mathbb{P}_M(\mathcal{V})]$ be the infinity hyperplane class. Assume that $-c_1(\mathcal{L}_m) \geq 0$ for all m and $c_1(M) + c_1(\mathcal{V}) \geq 0$. Then*

$$J^{\mathbb{P}_M(\mathcal{V} \oplus \mathcal{O})}(\tau, z^{-1}) = I^{\mathbb{P}_M(\mathcal{V} \oplus \mathcal{O})}(t, z)$$

under some mirror transform $\tau = \tau(t)$, where

$$I^{\mathbb{P}_M(\mathcal{V} \oplus \mathcal{O})} = e^{t/z} \sum_{\beta} \pi^* J_\beta^M(z^{-1}) \sum_{n \geq 0} \frac{1}{\prod_i (\xi + \pi^* c_1(\mathcal{L}_m))^{n+(c_1(\mathcal{L}_m), \beta)}} P^{j_* \beta + n\gamma},$$

and $j: M \rightarrow \mathbb{P}_M(\mathcal{O}) \subseteq \mathbb{P}_M(\mathcal{V} \oplus \mathcal{O})$ is the zero section.

3. CANONICAL LOCAL MODELS

In this section, we will omit the Novikov variables q^β since all the summation will be convergent.

3.1. Geometric realizations. We shall recall realizations of del Pezzo threefolds X_d (Appendix A), and the del Pezzo transition $Y_d \searrow X_d$ (Example 1.9).

For $1 \leq d \leq 8$, we denote X_d with the polarization $\mathcal{O}_{X_d}(1)$ by the del Pezzo threefold of degree d appearing in Theorem A.5. Recall that, when $d = 6$ there are two cases (6I) and (6II), we denote $X_{6\text{I}}$ and $X_{6\text{II}}$, respectively. Therefore the zero locus S_d defined by a geneal global section of $\mathcal{O}_{X_d}(1)$ is a smooth del Pezzo surface of degree d .

Using Theorem A.5, we are going to define an ambient space T_d containing X_d , and to associate to X_d a vector bundles \mathcal{V}_d .

Proposition 3.1. *There exist a smooth polarized variety $(T_d, \mathcal{O}_{T_d}(1))$ and a vector bundle \mathcal{V}_d on T_d such that X_d is the zero locus of a general global section of \mathcal{V}_d and $\mathcal{O}_{X_d}(1) = \mathcal{O}_{T_d}(1)|_{X_d}$.*

All of the T_d and \mathcal{V}_d are given in Table 4. Here we define $\mathcal{O}_{T_8}(\frac{1}{2}) := \mathcal{O}_{\mathbb{P}^4}(1)$ and follow the notation of Remark A.6.

d	$(T_d, \mathcal{O}_{T_d}(1))$	\mathcal{V}_d
1	$(\mathbb{P}(3, 2, 1, 1, 1), \mathcal{O}(1))$	$\mathcal{O}_{T_1}(6)$
2	$(\mathbb{P}(2, 1, 1, 1, 1), \mathcal{O}(1))$	$\mathcal{O}_{T_2}(4)$
3	$(\mathbb{P}^4, \mathcal{O}(1))$	$\mathcal{O}_{T_3}(3)$
4	$(\mathbb{P}^5, \mathcal{O}(1))$	$\mathcal{O}_{T_4}(2)^{\oplus 2}$
5	$(\text{Gr}(2, 5), \mathcal{O}(1))$	$\mathcal{O}_{T_5}(1)^{\oplus 3}$
6I	$(\mathbb{P}^2 \times \mathbb{P}^2, \mathcal{O}(1, 1))$	$\mathcal{O}_{T_{6I}}(1)$
6II	$(\mathbb{P}^1 \times \mathbb{P}^1 \times \mathbb{P}^1, \mathcal{O}(1, 1, 1))$	0
7	$(\text{Bl}_{\text{pt}} \mathbb{P}^3, \mathcal{O}(h_1 + h_2))$	0
8	$(\mathbb{P}^4, \mathcal{O}(2))$	$\mathcal{O}_{T_8}(\frac{1}{2})$

TABLE 4. Ambient spaces of del Pezzo threefolds.

Remark 3.2. Each del Pezzo surface S_d is defined by a general global section of $\mathcal{V}_d \oplus \mathcal{O}_{T_d}(1)$. We also have $T_d = X_d$ for $d = 6\text{II}$ and 7.

Recall that Y_d is the projective bundle $\mathbb{P}_{S_d}(K_{S_d} \oplus \mathcal{O})$ with projection $\pi_d: Y_d \rightarrow S_d$. Let $i_d: S_d \rightarrow X_d$ denote the inclusion and $j_d: S_d \rightarrow Y_d$ the infinity section. These maps fit into the following diagram:

$$\begin{array}{ccc} & Y_d & \\ j_d \swarrow & \downarrow \pi_d & \\ S_d & \xhookrightarrow{i_d} & X_d \end{array}$$

We denote by h the restriction of $c_1(\mathcal{O}_{T_d}(1))$ on X_d . We denote by E (resp. H) the zero section $\mathbb{P}_{S_d}(\mathcal{O})$ (resp. the infinity section $\mathbb{P}_{S_d}(K_{S_d}) = j_d(S_d)$) of Y_d . Then $F := H - E$ is the pullback $\pi_d^* i_d^* h$. It is easily seen that the first Chern classes of X_d and Y_d are

$$c_1(X_d) = 2h \quad \text{and} \quad c_1(Y_d) = 2H, \tag{3.1}$$

since X_d is del Pezzo and $Y_d = \mathbb{P}_{S_d}(K_{S_d} \oplus \mathcal{O})$.

3.2. Gromov–Witten invariants for $d \neq 7$. In this subsection, we aim to compute the symmetrized GW invariants (§2.2) of both X_d and Y_d , and investigate how they are related. To do this, we shall introduce the groups acting on their cohomology.

For each $d \in \{1, 2, 3, 4, 5, 8\}$, it is clear that $H(X_d) = \mathbb{C}[h]/(h^4)$, so we set the group G_{X_d} to be trivial. We define $G_{X_{6I}} := \mathfrak{S}_2$ and $G_{X_{6II}} := \mathfrak{S}_3$. The group action of $G_{X_{6I}}$ (resp. $G_{X_{6II}}$) on $T_{6I} = \mathbb{P}^2 \times \mathbb{P}^2$ (resp. $T_{6II} = \mathbb{P}^1 \times \mathbb{P}^1 \times \mathbb{P}^1$) induces automorphisms on $H(X_d)$ such that

$$H(X_d)^{G_{X_d}} = \mathbb{C}[h]/(h^4) = \mathbb{C} \cdot 1 \oplus \mathbb{C} \cdot h \oplus \mathbb{C} \cdot h^2 \oplus \mathbb{C} \cdot \text{pt}.$$

Since $d \neq 7$, Proposition 2.4 tells us that there is a group action of G_{S_d} on $H(S_d)$, and hence on $H(Y_d)$, such that

$$H(S_d)^{G_{S_d}} = \mathbb{C}[i_d^*h]/((i_d^*h)^3) = \mathbb{C} \cdot 1 \oplus \mathbb{C} \cdot i_d^*h \oplus \mathbb{C} \cdot \text{pt},$$

$$H(Y_d)^{G_{S_d}} = \mathbb{C}[E, H]/(EH, F^3) = \mathbb{C} \cdot 1 \oplus \mathbb{C} \cdot E \oplus \mathbb{C} \cdot H \oplus \mathbb{C} \cdot H^2 \oplus \mathbb{C} \cdot E^2 \oplus \mathbb{C} \cdot \text{pt}.$$

Remark 3.3. The groups $G_{X_{6I}}$ and $G_{X_{6II}}$ are the Weyl group of A_1 and A_2 , respectively—the two components of the reducible root system $E_3 = A_1 \times A_2$.

From now on, we will work within the invariant subspaces, and all GW invariants are symmetrized (cf. (2.1)).

Convention 3.4. In the remainder of this section, we will omit the groups in the notation for simplicity. We will write $H(X_d)$, $H(S_d)$, and $H(Y_d)$ for $H(X_d)^{G_{X_d}}$, $H(S_d)^{G_{S_d}}$, and $H(Y_d)^{G_{S_d}}$, respectively, and similarly for the notation of symmetrized GW invariants.

Choose $\bar{\gamma} \in H_2(X_d)$ such that $(h, \bar{\gamma})_{X_d} = 1$. When $d \leq 6$ (resp. $d = 8$), $\bar{\gamma}$ (resp. $2\bar{\gamma}$) generate the cone $\text{NE}(X_d)/G_{X_d}$.

Let $\gamma \in H_2(Y_d)$ denote the fiber class of the projection $Y_d \rightarrow S_d$. Choose $\ell \in H_2(Y_d)$ such that $(H, \ell)_{Y_d} = 0$, $(F, \ell)_{Y_d} = 1$. When $d \leq 6$ (resp. $d = 8$), γ and ℓ (resp. 2ℓ) generate the cone $\text{NE}(Y_d)/G_{Y_d}$.

Write $\mathcal{W}_d = \mathcal{V}_d \oplus \mathcal{O}_{T_d}(1)$. It follows from Theorem 2.6, Theorem 2.8 and Remark 3.2 that for $d \leq 6$,

$$\begin{aligned} J^{X_d}(\tau, z^{-1}) &= I^{X_d}(t, z) := e^{t/z} \sum_{m \geq 0} \left(J_{m\bar{\gamma}}^{T_d}(z^{-1}) \cdot \mathcal{V}_d^{\overline{m\bar{\gamma}}} \right) P^{m\bar{\gamma}}, \\ J^{Y_d}(\tau, z^{-1}) &= I^{Y_d}(t, z) := e^{t/z} \sum_{m, n \geq 0} \left(\pi_d^* i_d^* \left(J_{m\bar{\gamma}}^{T_d}(z^{-1}) \cdot \mathcal{W}_d^{\overline{m\bar{\gamma}}} \right) \frac{1}{(E)^{\overline{n-m}}(H)^{\overline{n}}} \right) P^{m\ell+n\gamma} \\ &= e^{t/z} \sum_{m, n \geq 0} \left(\pi_d^* i_d^* I_{m\bar{\gamma}}^{X_d}(z) \cdot \frac{(F)^{\overline{m}}}{(E)^{\overline{n-m}}(H)^{\overline{n}}} \right) P^{m\ell+n\gamma} \end{aligned}$$

under some mirror transform; for $d = 8$, we have

$$\begin{aligned} J^{X_8}(\tau, z^{-1}) &= I^{X_8}(t, z) := e^{t/z} \sum_{n \geq 0} \left(J_{m(2\bar{\gamma})}^{T_8}(z^{-1}) \cdot \mathcal{V}_8^{\overline{m(2\bar{\gamma})}} \right) P^{m(2\bar{\gamma})}, \\ J^{Y_8}(\tau, z^{-1}) &= I^{Y_8}(t, z) := e^{t/z} \sum_{m, n \geq 0} \left(\pi_8^* i_8^* \left(J_{m(2\bar{\gamma})}^{T_8}(z^{-1}) \cdot \mathcal{W}_8^{\overline{m(2\bar{\gamma})}} \right) \frac{1}{(E)^{\overline{n-2m}}(H)^{\overline{n}}} \right) P^{m(2\ell)+n\gamma} \\ &= e^{t/z} \sum_{m, n \geq 0} \left(\pi_8^* i_8^* I_{m(2\bar{\gamma})}^{X_8} \cdot \frac{(\frac{1}{2}F)^{\overline{m}}}{(E)^{\overline{n-2m}}(H)^{\overline{n}}} \right) P^{m(2\ell)+n\gamma} \end{aligned}$$

under some mirror transform.

For a divisor D , we define its quantization operator \hat{D} via

$$\hat{D}(e^{t/z} P^\beta) = (D + (D, \beta)z)e^{t/z} P^\beta,$$

i.e., the derivative $z \frac{\partial}{\partial t_D}$ along the D -direction.

Lemma 3.5. *For $d = 1, 2, 3, 4$, let $(\kappa_d, \lambda_d) = (432, 60), (64, 12), (27, 6), (16, 4)$, respectively. Then $I^{X_d}(t)$ satisfies the Picard–Fuchs equation $\square_{\bar{\gamma}} I^{X_d} = 0$, where*

$$\square_{\bar{\gamma}} = \hat{h}^4 - P^{\bar{\gamma}}(\kappa_d \hat{h}^2 + \kappa_d z \hat{h} + \lambda_d z^2).$$

Proof. For $d \in \{1, 2, 3\}$, X_d is a hypersurface of degree $n_d = \sum \alpha_i - 1$ in $\mathbb{P}(\alpha, 1)$. By [44],

$$I_{m\bar{\gamma}}^{X_d} = \frac{(n_d h)^{\overline{n_d m}}}{(h)^{\overline{m}} \prod_{i=0}^3 (\alpha_i h)^{\overline{\alpha_i m}}}.$$

It is straightforward to verify case by case that

$$\begin{aligned} \frac{I_{(m+1)\bar{\gamma}}^{X_d}}{I_{m\bar{\gamma}}^{X_d}} &= \frac{n_d^{n_d-2}}{\prod_{1 < e | n_d} (n_d/e)^e} \frac{(n_d h + (n_d m + 1)z)(n_d h + (n_d m + (n_d - 1))z)}{(h + (m + 1)z)^4} \\ &= \kappa_d \cdot \frac{(h + (m + \frac{1}{n_d})z)(h + (m + \frac{n_d-1}{n_d})z)}{(h + (m + 1)z)^4} = \frac{\kappa_d(h + mz)^2 + \kappa_d z(h + mz) + \lambda_d z^2}{(h + (m + 1)z)^4}. \end{aligned}$$

Consequently, $\square_{\bar{\gamma}} I^{X_d} = 0$.

In the case $d = 4$, we have

$$I_{m\bar{\gamma}}^{X_d} = \frac{((2h)^{\overline{2m}})^2}{((h)^{\overline{m}})^6}.$$

So

$$\frac{I_{(m+1)\bar{\gamma}}^{X_d}}{I_{m\bar{\gamma}}^{X_d}} = \frac{16(h + (m + \frac{1}{2})z)^2}{(h + (m + 1)z)^4} = \frac{\kappa_4(h + mz)^2 + \kappa_4 z(h + mz) + \lambda_4 z^2}{(h + (m + 1)z)^4},$$

and hence, $\square_{\bar{\gamma}} I^{X_d} = 0$. \square

Lemma 3.6. *For $d = 5, 6\text{I}, 6\text{II}$, let $(\kappa_d, \lambda_d, \mu_d) = (11, 3, 1), (7, 2, 8), (10, 3, -9)$, respectively. Then $I^{X_d}(t)$ satisfies the Picard–Fuchs equation $\square_{\bar{\gamma}} I^{X_d} = 0$, where*

$$\square_{\bar{\gamma}} = \hat{h}^4 - P^{\bar{\gamma}}(\kappa_d \hat{h}^2 + \kappa_d z \hat{h} + \lambda_d z^2) - \mu_d P^{2\bar{\gamma}}.$$

Proof. We prove the statement case by case.

For $d = 5$, it follows from [3, p.659] that $J^{\text{Gr}(2,5)}(t)$ satisfies the Picard–Fuchs equation $\square_{\bar{\gamma}}^{\text{Gr}(2,5)} J^{\text{Gr}(2,5)} = 0$, where

$$\square_{\bar{\gamma}}^{\text{Gr}(2,5)} = \hat{h}^7(\hat{h} - z)^3 - P^{\bar{\gamma}} \hat{h}^3(11\hat{h}^2 + 11z\hat{h} + 3z^2) - P^{2\bar{\gamma}},$$

i.e., for each $m \geq 0$,

$$\begin{aligned} (h + (m + 1)z)^7(h + mz)^3 J_{(m+1)\bar{\gamma}}^{\text{Gr}(2,5)} \\ = (h + mz)^3(11(h + mz)^2 + 11z(h + mz) + 3z^2)J_{m\bar{\gamma}}^{\text{Gr}(2,5)} + J_{(m-1)\bar{\gamma}}^{\text{Gr}(2,5)}. \end{aligned}$$

Together with $I_{m\bar{\gamma}}^{X_5} = J_{m\bar{\gamma}}^{\text{Gr}(2,5)}((h)^{\bar{m}})^3$, we get

$$(h + (m+1)z)^4 I_{(m+1)\bar{\gamma}}^{X_5} = (11(h+mz)^2 + 11z(h+mz) + 3z^2) I_{m\bar{\gamma}}^{X_5} + I_{(m-1)\bar{\gamma}}^{X_5}.$$

Hence, I^{X_5} satisfies the Picard–Fuchs equation $\square_{\bar{\gamma}} I^{X_5} = 0$.

For $d = 6I$, let $h_1 = c_1(\mathcal{O}(1,0))$, $h_2 = c_1(\mathcal{O}(0,1))$ be the hyperplane classes of $T_{6I} = \mathbb{P}^2 \times \mathbb{P}^2$, and let $\bar{\gamma}_1 = h_1 h^2$, $\bar{\gamma}_2 = h_2^2 h_2$ be the generators of $\text{NE}(\mathbb{P}^2 \times \mathbb{P}^2)$. Then I^{X_d} satisfies the Picard–Fuchs equation $\square_{\bar{\gamma}_1} I^{X_d} = \square_{\bar{\gamma}_2} I^{X_d} = 0$, where

$$\begin{aligned}\square_{\bar{\gamma}_1} &= \hat{h}_1^3 - (\hat{h}_1 + \hat{h}_2) P^{\bar{\gamma}_1}, \\ \square_{\bar{\gamma}_2} &= \hat{h}_2^3 - (\hat{h}_1 + \hat{h}_2) P^{\bar{\gamma}_2}.\end{aligned}$$

One can also check that I^{X_d} is also annihilated by

$$\square_{\text{top}} = \hat{h}_1^2 - \hat{h}_1 \hat{h}_2 + \hat{h}_2^2 - (P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2}).$$

Indeed,

$$\square_{\bar{\gamma}_1} + \square_{\bar{\gamma}_2} = (\hat{h}_1 + \hat{h}_2) \square_{\text{top}}.$$

We want to symmetrize the differential equations. Let $\hat{h} = \hat{h}_1 + \hat{h}_2$ and let $\hat{k} = \hat{h}_1 - \hat{h}_2$. Then

$$\begin{aligned}\square_{\bar{\gamma}_1} &= \frac{1}{8}(\hat{h} + \hat{k})^3 - \hat{h} P^{\bar{\gamma}_1} = \frac{1}{8}\hat{h}^3 + \frac{3}{8}\hat{h}^2\hat{k} + \frac{3}{8}\hat{h}\hat{k}^2 + \frac{1}{8}\hat{k}^3 - \hat{h} P^{\bar{\gamma}_1}, \\ \square_{\bar{\gamma}_2} &= \frac{1}{8}(\hat{h} - \hat{k})^3 - \hat{h} P^{\bar{\gamma}_2} = \frac{1}{8}\hat{h}^3 - \frac{3}{8}\hat{h}^2\hat{k} + \frac{3}{8}\hat{h}\hat{k}^2 - \frac{1}{8}\hat{k}^3 - \hat{h} P^{\bar{\gamma}_2}, \\ \square_{\text{top}} &= \frac{1}{4}((\hat{h} + \hat{k})^2 - (\hat{h} + \hat{k})(\hat{h} - \hat{k}) + (\hat{h} - \hat{k})^2) - (P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2}) \\ &= \frac{1}{4}(\hat{h}^2 + 3\hat{k}^2) - (P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2}).\end{aligned}$$

In order to cancel out the \hat{k}^3 term, we consider

$$3(\square_{\bar{\gamma}_1} - \square_{\bar{\gamma}_2}) - \hat{k} \square_{\text{top}} = 2\hat{h}^2\hat{k} - 3\hat{h}(P^{\bar{\gamma}_1} - P^{\bar{\gamma}_2}) + \hat{k}(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2}).$$

So, modulo the left ideal generated by $\square_{\bar{\gamma}_1}$, $\square_{\bar{\gamma}_2}$, and \square_{top} , we get

$$\begin{aligned}\hat{h}^4 &= \hat{h}^2(4(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2}) - 3\hat{k}^2) \\ &= 4(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2})(\hat{h} + z)^2 - \frac{3}{2}\hat{k}(3\hat{h}(P^{\bar{\gamma}_1} - P^{\bar{\gamma}_2}) - \hat{k}(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2})) \\ &= 4(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2})(\hat{h} + z)^2 - \frac{9}{2}(P^{\bar{\gamma}_1}(\hat{k} + z) - P^{\bar{\gamma}_2}(\hat{k} - z))(\hat{h} + z) \\ &\quad + \frac{3}{2}(P^{\bar{\gamma}_1}(\hat{k} + z)^2 + P^{\bar{\gamma}_2}(\hat{k} - z)^2) \\ &= 4(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2})(\hat{h} + z)^2 - \frac{9}{2}(P^{\bar{\gamma}_1} - P^{\bar{\gamma}_2})\hat{k}(\hat{h} + z) - \frac{9}{2}(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2})z(\hat{h} + z) \\ &\quad + \frac{1}{2}(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2})(4(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2}) - \hat{h}^2) + \frac{3}{2}(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2})z^2.\end{aligned}$$

Restrict this equation on the diagonal $P^{\bar{\gamma}} = P^{\bar{\gamma}_1} = P^{\bar{\gamma}_2}$, we get

$$\begin{aligned}\hat{h}^4 &= 8P^{\bar{\gamma}}(\hat{h} + z)^2 - 9P^{\bar{\gamma}}z(\hat{h} + z) + P^{\bar{\gamma}}(8P^{\bar{\gamma}} - \hat{h}^2) + 3P^{\bar{\gamma}}z^2 \\ &= P^{\bar{\gamma}}(7\hat{h}^2 + 7z\hat{h} + 2z^2) + 8P^{2\bar{\gamma}},\end{aligned}$$

which is what we want.

For $d = 6\mathbb{I}$, let $h_1 = c_1(\mathcal{O}(1, 0, 0))$, $h_2 = c_1(\mathcal{O}(0, 1, 0))$, $h_3 = c_1(\mathcal{O}(0, 0, 1))$ be the hyperplane classes of $T_{6\mathbb{I}} = \mathbb{P}^1 \times \mathbb{P}^1 \times \mathbb{P}^1$, and let $\bar{\gamma}_1 = h_2h_3$, $\bar{\gamma}_2 = h_3h_1$, $\bar{\gamma}_3 = h_1h_2$ be the generators of $\text{NE}(\mathbb{P}^1 \times \mathbb{P}^1 \times \mathbb{P}^1)$. Then I^{X_d} satisfies the Picard–Fuchs equation $\square_{\bar{\gamma}_1} I^{X_d} = \square_{\bar{\gamma}_2} I^{X_d} = \square_{\bar{\gamma}_3} I^{X_d} = 0$, where

$$\square_{\bar{\gamma}_i} = \hat{h}_i^2 - P^{\bar{\gamma}_i}.$$

Let $\hat{h} = \hat{h}_1 + \hat{h}_2 + \hat{h}_3$, $\hat{k}_+ = \hat{h}_1 + \omega\hat{h}_2 + \omega^2\hat{h}_3$, $\hat{k}_- = \hat{h}_1 + \omega^2\hat{h}_2 + \omega\hat{h}_3$, where $\omega = \omega_3 = e^{2\pi i/3}$ is a third root unity. Then

$$\begin{aligned}\square_{\bar{\gamma}_1} + \square_{\bar{\gamma}_2} + \square_{\bar{\gamma}_3} &= \frac{1}{3}\hat{h}^2 + \frac{2}{3}\hat{k}_+\hat{k}_- - (P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3}), \\ \square_{\bar{\gamma}_1} + \omega\square_{\bar{\gamma}_2} + \omega^2\square_{\bar{\gamma}_3} &= \frac{1}{3}\hat{k}_-^2 + \frac{2}{3}\hat{h}\hat{k}_+ - (P^{\bar{\gamma}_1} + \omega P^{\bar{\gamma}_2} + \omega^2 P^{\bar{\gamma}_3}), \\ \square_{\bar{\gamma}_1} + \omega^2\square_{\bar{\gamma}_2} + \omega\square_{\bar{\gamma}_3} &= \frac{1}{3}\hat{k}_+^2 + \frac{2}{3}\hat{h}\hat{k}_- - (P^{\bar{\gamma}_1} + \omega^2 P^{\bar{\gamma}_2} + \omega P^{\bar{\gamma}_3}).\end{aligned}$$

So, modulo the left ideal generated by $\square_{\bar{\gamma}_1}$, $\square_{\bar{\gamma}_2}$, and $\square_{\bar{\gamma}_3}$, we get

$$\begin{aligned}\hat{h}^4 &= \hat{h}^2(3(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3}) - 2\hat{k}_+\hat{k}_-) \\ &= 3(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3})(\hat{h} + z)^2 - \hat{h}\hat{k}_+(-\hat{k}_+^2 + 3(P^{\bar{\gamma}_1} + \omega^2 P^{\bar{\gamma}_2} + \omega P^{\bar{\gamma}_3})) \\ &= 3(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3})(\hat{h} + z)^2 + \frac{1}{2}\hat{k}_+^2(-\hat{k}_-^2 + 3(P^{\bar{\gamma}_1} + \omega P^{\bar{\gamma}_2} + \omega^2 P^{\bar{\gamma}_3})) \\ &\quad - 3(P^{\bar{\gamma}_1}(\hat{k}_+ + z) + \omega^2 P^{\bar{\gamma}_2}(\hat{k}_+ + \omega z) + \omega P^{\bar{\gamma}_3}(\hat{k}_+ + \omega^2 z))(\hat{h} + z) \\ &= 3(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3})(\hat{h} + z)^2 - \frac{1}{2}\hat{k}_+^2\hat{k}_-^2 \\ &\quad + \frac{3}{2}(P^{\bar{\gamma}_1}(\hat{k}_+ + z)^2 + \omega P^{\bar{\gamma}_2}(\hat{k}_+ + \omega z)^2 + \omega^2 P^{\bar{\gamma}_3}(\hat{k}_+ + \omega^2 z)^2) \\ &\quad - 3(P^{\bar{\gamma}_1} + \omega^2 P^{\bar{\gamma}_2} + \omega P^{\bar{\gamma}_3})\hat{k}_+(\hat{h} + z) - 3z(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3})(\hat{h} + z).\end{aligned}\tag{3.2}$$

Since

$$\begin{aligned}\hat{k}_+^2\hat{k}_-^2 &= \frac{1}{2}\hat{k}_+\hat{k}_-(-\hat{h}^2 + 3(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3})) \\ &= -\frac{1}{4}\hat{h}^2(-\hat{h}^2 + 3(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3})) \\ &\quad + \frac{3}{2}(P^{\bar{\gamma}_1}(\hat{k}_+ + z)(\hat{k}_- + z) + P^{\bar{\gamma}_2}(\hat{k}_+ + \omega z)(\hat{k}_- + \omega^2 z) + P^{\bar{\gamma}_3}(\hat{k}_+ + \omega^2 z)(\hat{k}_- + \omega z)) \\ &= \frac{1}{4}\hat{h}^4 - \frac{3}{4}(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3})(\hat{h} + z)^2 \\ &\quad + \frac{3}{4}(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3})(-\hat{h}^2 + 3(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3})) \\ &\quad + \frac{3}{2}(P^{\bar{\gamma}_1} + \omega^2 P^{\bar{\gamma}_2} + \omega P^{\bar{\gamma}_3})z\hat{k}_+ + \frac{3}{2}(P^{\bar{\gamma}_1} + \omega P^{\bar{\gamma}_2} + \omega^2 P^{\bar{\gamma}_3})z\hat{k}_- + \frac{3}{2}(P^{\bar{\gamma}_1} + P^{\bar{\gamma}_2} + P^{\bar{\gamma}_3})z^2,\end{aligned}$$

it follows that the restriction of (3.2) on the diagonal $P^{\bar{\gamma}} = P^{\bar{\gamma}_1} = P^{\bar{\gamma}_2} = P^{\bar{\gamma}_3}$ is

$$\begin{aligned}\hat{h}^4 &= 9P^{\bar{\gamma}}(\hat{h} + z)^2 - \frac{1}{2}(\frac{1}{4}\hat{h}^4 - \frac{9}{4}P^{\bar{\gamma}}(\hat{h} + z)^2 - \frac{9}{4}P^{\bar{\gamma}}\hat{h}^2 + \frac{81}{4}P^{2\bar{\gamma}} + \frac{9}{2}P^{\bar{\gamma}}z^2) \\ &\quad + \frac{9}{2}P^{\bar{\gamma}}z^2 - 9P^{\bar{\gamma}}z(\hat{h} + z),\end{aligned}$$

or, equivalently,

$$\hat{h}^4 = P^{\bar{\gamma}}(10\hat{h}^2 + 10z\hat{h} + 3z^2) - 9P^{2\bar{\gamma}}.\tag*{\square}$$

Lemma 3.7. *If $I^X(t)$ satisfies the Picard–Fuchs equation $\square_{\bar{\gamma}} I^X$, where*

$$\square_{\bar{\gamma}} = \hat{h}^4 - P^{\bar{\gamma}}(\kappa\hat{h}^2 + \kappa z\hat{h} + \lambda z^2) - \mu P^{2\bar{\gamma}},$$

and $I_{m\ell+n\gamma}^Y = \pi^* i^* I_{m\bar{\gamma}}^X \cdot \frac{(F)^{\bar{m}}}{(E)^{n-m}(H)^{\bar{n}}}$, then

$$I^Y(t) = e^{t/z} \sum_{m,n \geq 0} I_{m\ell+n\gamma}^Y P^{m\ell+n\gamma}$$

satisfies the Picard–Fuchs equation $\square_\gamma I^Y = \square_\ell I^Y = 0$, where

$$\begin{aligned} \square_\gamma &= \hat{E}\hat{H} - P^\gamma, \\ \square_\ell &= \hat{F}^3 - P^\ell(\kappa\hat{F}^2 + \kappa z\hat{F} + \lambda z^2)\hat{E} - \mu P^{2\ell}(\hat{F} + z)\hat{E}(\hat{E} - z). \end{aligned}$$

Proof. It follows from

$$\frac{I_{m\ell+(n-1)\gamma}^Y}{I_{m\ell+n\gamma}^Y} = \frac{(E)^{\bar{n}-m}(H)^{\bar{n}}}{(E)^{(n-1)-m}(H)^{n-1}} = (E + (n-m)z)(H + nz),$$

that $\square_\gamma I^Y = 0$. The equation $\square_{\bar{\gamma}} I^{X_d} = 0$ gives

$$(h + (n+1)z)^4 I_{(n+1)\bar{\gamma}}^X = (\kappa(h + nz)^2 + \kappa z(h + nz) + \lambda z^2) I_{n\bar{\gamma}}^X + \mu I_{(n-1)\bar{\gamma}}^X.$$

This implies

$$\begin{aligned} (F + (m+1)z)^3 I_{(m+1)\ell+n\gamma}^Y &= (\kappa(F + mz)^2 + \kappa z(F + mz) + \lambda z^2)(E + (n-m)z) I_{m\ell+n\gamma}^Y \\ &\quad + \mu(F + mz)(E + (n-m)z)(E + (n-m+1)z) I_{(m-1)\ell+n\gamma}^Y, \end{aligned}$$

or, equivalently, $\square_\ell I^Y = 0$. \square

Lemma 3.8. *The I -function $I^{X_8}(t)$ satisfies the Picard–Fuchs equation $\square_{\bar{\gamma}} I^{X_8} = 0$, where*

$$\square_{\bar{\gamma}} = \hat{h}^4 - 16P^{2\bar{\gamma}};$$

$I^{X_8}(t)$ satisfies the Picard–Fuchs equation $\square_\gamma I^{X_8} = \square_\ell I^{X_8} = 0$, where

$$\begin{aligned} \square_\gamma &= \hat{E}\hat{H} - P^\gamma, \\ \square_\ell &= \hat{F}^3 - 16P^{2\ell}(\hat{F} + z)\hat{E}(\hat{E} - z). \end{aligned}$$

Proof. Since $I_{n(2\bar{\gamma})}^{X_8} = \frac{1}{((\frac{1}{2}h)^{\bar{n}})^4}$,

$$\frac{I_{(n+1)(2\bar{\gamma})}^{X_8}}{I_{n(2\bar{\gamma})}^{X_8}} = \frac{1}{(\frac{1}{2}h + (n+1)z)^4} = \frac{16}{(h + 2(n+1)z)^4}.$$

This means that $\square_{\tilde{\gamma}} I^{X_8} = 0$. Since $I_{m(2\ell)+n\gamma}^{Y_8} = \frac{(F)^{\overline{2m}}}{((\frac{1}{2}F)^{\overline{m}})^4(E)^{\overline{n-2m}}(H)^{\overline{n}}}$, we have

$$\begin{aligned} \frac{I_{m(2\ell)+(n-1)\gamma}^{Y_8}}{I_{m(2\ell)+n\gamma}^{Y_8}} &= (E + (n - 2m)z)(H + nz) \\ \frac{I_{(m+1)(2\ell)+n\gamma}^{Y_8}}{I_{m(2\ell)+n\gamma}^{Y_8}} &= \frac{(F + (2m + 2)z)(F + (2m + 1)z)(E + (n - 2m)z)(E + (n - 2m - 1)z)}{(\frac{1}{2}F + (m + 1)z)^4} \\ &= \frac{16((F + 2mz) + z)(E + (n - 2m)z)(E + (n - 2m - 1)z)}{(\frac{1}{2}F + (m + 1)z)^3}, \end{aligned}$$

or, equivalently, $\square_{\gamma} I^{Y_8} = \square_{\ell} I^{Y_8} = 0$. \square

Summarizing Lemmas 3.5, 3.6, 3.7 and 3.8, we have:

Proposition 3.9. *For each $d \neq 7$, $\square_{\tilde{\gamma}} I^{X_d} = 0$ and $\square_{\gamma} I^{Y_d} = \square_{\ell} I^{Y_d} = 0$, where*

$$\begin{aligned} \square_{\tilde{\gamma}} &= \hat{h}^4 - P^{\tilde{\gamma}}(\kappa_d \hat{h}^2 + \kappa_d z \hat{h} + \lambda_d z^2) - \mu_d P^{2\tilde{\gamma}}, \\ \square_{\gamma} &= \hat{E} \hat{H} - P^{\gamma}, \\ \square_{\ell} &= \hat{F}^3 - P^{\ell}(\kappa_d \hat{F}^2 + \kappa_d z \hat{F} + \lambda_d z^2) \hat{E} - \mu_d P^{2\ell}(\hat{F} + z) \hat{E} (\hat{E} - z), \end{aligned}$$

where κ_d , λ_d and μ_d are listed in Table 1.

Remark 3.10. The constants $(\kappa_d, \lambda_d, \mu_d)$ appeared in Zagier's list [54], which was obtained by performing a computer search for integral sequences that satisfy a certain three-term recurrence relation. For example, when $d \in \{1, 2, 3, 4, 8\}$, the constants correspond to hypergeometric solutions (cf. [54, pp.352-353]). The case $d \in \{5, 6\text{I}, 6\text{II}\}$ correspond to the new label **D**, **A**, **C** in [54, Table 2], respectively. The solutions **A-F** in [54, Table 2] are now called Zagier's sporadic sequences (cf. [8, §5.10]).

Remark 3.11. In [38], it is shown that the quantum D -module of $Y_{\text{loc}} = Y_d$, after analytic continuation and restriction, recovers the quantum D -module of $X_{\text{loc}} = X_d$, in the case of toric ambient spaces. We now compute the Picard–Fuchs ideal in our setting. Let $\tilde{\gamma} = \gamma + \ell$ be the lifting of γ such that $(E, \tilde{\gamma}) = 0$. In the coordinates $(P^{\tilde{\gamma}}, P^{-\ell})$, we have

$$\hat{F} = z\theta_{P^{\tilde{\gamma}}} - z\theta_{P^{-\ell}}, \quad \hat{E} = z\theta_{P^{-\ell}}, \quad \hat{H} = z\theta_{P^{\tilde{\gamma}}},$$

where $\theta_{P^{\tilde{\gamma}}} = P^{\tilde{\gamma}} \frac{d}{dP^{\tilde{\gamma}}}$, $\theta_{P^{-\ell}} = P^{-\ell} \frac{d}{dP^{-\ell}}$. In particular, $\hat{F} P^{\gamma} = P^{\gamma} \hat{F}$.

Modulo the left ideal PF^{Y_d} generated by \square_{γ} and \square_{ℓ} , we have

$$\begin{aligned} \hat{F}^3 \hat{H} &= P^{\ell}(\kappa_d \hat{F}^2 + \kappa_d z \hat{F} + \lambda_d z^2) \hat{E} \hat{H} + \mu_d P^{2\ell}(\hat{F} + z)(\hat{E} - z) \hat{E} \hat{H} \\ &= P^{\tilde{\gamma}}(\kappa_d \hat{F}^2 + \kappa_d z \hat{F} + \lambda_d z^2) + \mu_d P^{\tilde{\gamma}+\ell}(\hat{F} + z) \hat{E}. \end{aligned} \tag{3.3}$$

When $\mu_d = 0$, i.e., $d \in \{1, 2, 3, 4\}$, the restriction of (3.3) to the locus $P^{-\ell} = 0$ is

$$\theta_{P^{\tilde{\gamma}}}^4 = z^2 P^{\tilde{\gamma}}(\kappa_d \theta_{P^{\tilde{\gamma}}}^2 + \kappa_d \theta_{P^{\tilde{\gamma}}} + \lambda_d)$$

which is precisely the pullback of $\square_{\bar{\gamma}}$. However, when $\mu_d \neq 0$, i.e., $d \in \{5, 6I, 6II, 8\}$, the term

$$\mu_d P^{\bar{\gamma}+\ell} (\hat{F} + z) \hat{E} = z^2 \mu_d P^{\bar{\gamma}} P^\ell (\theta_{P^{\bar{\gamma}}} - \theta_{P^{-\ell}} + 1) \theta_{P^{-\ell}}$$

fails to be specialized to $\square_{\bar{\gamma}}$. In fact it is not specializable on the locus $P^{-\ell} = 0$.

Nevertheless, the conclusion of [38] still holds: there exists a two parameter extension $I^{X_{\text{loc}}}(tH, y)$ of $I^{X_{\text{loc}}}(th)$ which is a monodromy invariant solution of $\widetilde{\text{PF}}^Y$ near $y = P^{-\ell} = 0$.

Lemma 3.12.

- (i) *The mirror transform $t \mapsto \tau$ for $J^{X_d}(\tau) = I^{X_d}(t)$ is trivial; that is, $\tau = t$.*
- (ii) *The mirror transform $t \mapsto \tau$ for $J^{Y_d}(\tau) = I^{Y_d}(t)$ is given by $\tau = t - g_d(P^\ell)E$, for a unique power series $g_d(P^\ell) \in P^\ell \mathbb{Q}[[P^\ell]]$. Moreover, letting $\theta = P^\ell \frac{d}{dP^\ell} = -\frac{1}{z} \hat{E}$, the function $f_d := 1 + \theta g_d$ satisfies the differential equation*

$$(1 + \kappa_d P^\ell - \mu_d P^{2\ell}) \theta^2 f_d + (\kappa_d P^\ell - 2\mu_d P^{2\ell}) \theta f_d + (\lambda_d P^\ell - \mu_d P^{2\ell}) f_d = 0. \quad (3.4)$$

Proof. It follows from (3.1) that

$$(c_1(X_d), n\bar{\gamma}) = (2h, \bar{\gamma}) = 2n, \quad (c_1(Y_d), m\ell + n\gamma) = (2H, m\ell + n\gamma) = 2n.$$

For X_d , this shows that $I_{n\bar{\gamma}}^{X_d} = O(z^{-2n})$, and hence $I^{X_d}(t) = e^{t/z}(1 + O(z^{-2})) = 1 + tz^{-1} + O(z^{-2})$. Therefore, $\tau = t$.

For Y_d , this shows that $I_{m\ell+n\gamma}^{Y_d} = O(z^{-2n})$, and hence

$$I^{Y_d}(t) = e^{t/z} \left(\sum_{m \geq 0} I_{m\ell}^{Y_d} P^{m\ell} + O(z^{-2}) \right).$$

By $\square_\ell I^{Y_d} = 0$, we have

$$\begin{aligned} (F + (m+1)z)^3 I_{(m+1)\ell}^{Y_d} &= (\kappa_d(F+mz)^2 + \kappa_d z(F+mz) + \lambda_d z^2)(E-mz) I_{m\ell}^{Y_d} \\ &\quad + \mu_d(F+mz)(E-mz)(E-(m-1)z) I_{(m-1)\ell}^{Y_d}. \end{aligned}$$

The initial condition $I_{-\ell} = 0$, $I_0 = 1$ gives

$$I_\ell = \lambda_d E z^{-1} + O(z^{-2}), \quad I_{2\ell} = -\frac{1}{2}(2\kappa_d + \lambda_d + \mu_d) E z^{-1} + O(z^{-2}). \quad (3.5)$$

By induction on m , one obtains $I_{m\ell} = -g_m E z^{-1} + O(z^{-2})$ and the sequence $\{g_m\}_{m \geq 1}$ satisfies the following recursion relation

$$(m+1)^3 g_{m+1} + (\kappa_d(m^2+m) + \lambda_d) m g_m - \mu_d m^2 (m-1) g_{m-1} = 0.$$

It follows that

$$I^{Y_d}(t) = 1 + \left(t - \sum_{m \geq 1} g_m P^{m\ell} \cdot E \right) z^{-1} + O(z^{-2}).$$

Set $g_d(P^\ell) = \sum_{m \geq 1} g_m P^{m\ell}$. Then $f_d(P^\ell) = \sum_{m \geq 0} f_m P^{m\ell} := 1 + \theta g_d$ satisfies

$$(m+1)^2 f_{m+1} + (\kappa_d(m^2+m) + \lambda_d) f_m - \mu_d m^2 f_{m-1} = 0.$$

Note that this also holds for $m = 1$, as follows from (3.5). There fore f_d satisfies the differential equation

$$\theta^2 f + P^\ell(\kappa_d(\theta + 1)\theta + \lambda_d)f - \mu_d P^{2\ell}(\theta + 1)^2 = 0,$$

which coincides with equation (3.4). \square

Since $Q^\beta = e^{(\tau, \beta)}$, $P^\beta = e^{(t, \beta)}$, we obtain the following:

Corollary 3.13. *We have the identities*

$$Q^{\bar{\gamma}} = P^{\bar{\gamma}}, \quad Q^\gamma = P^\gamma e^{-g(P^\ell)}, \quad Q^\ell = P^\ell e^{g(P^\ell)}.$$

In particular, $\delta^E = -f^{-1}\theta$.

Proof. The identities follow directly from the definitions of the variables involved. For the final statement, since

$$\delta^E Q^{\gamma+\ell} = 0 = -f^{-1}\theta P^{\gamma+\ell} = -f^{-1}\theta Q^{\gamma+\ell},$$

it suffices to prove that $\delta^E P^\ell = -f^{-1}\theta P^\ell$. Observe that

$$\delta^E P^\ell = -Q^\ell \frac{dP^\ell}{dQ^\ell} = -Q^\ell \left(\frac{dQ^\ell}{dP^\ell} \right)^{-1}.$$

Since $Q^\ell = P^\ell e^{g(P^\ell)}$, we compute

$$\frac{dQ^\ell}{dP^\ell} = e^{g(P^\ell)} \cdot (1 + \theta g(P^\ell)).$$

Hence,

$$\delta^E P^\ell = -Q^\ell (e^{g(P^\ell)} \cdot (1 + \theta g(P^\ell)))^{-1} = -f^{-1} P^\ell,$$

as claimed. \square

Proposition 3.14. *Let $d = 1, 2, 3, 4$, and let $n_d = 6, 4, 3, 2$, respectively. Then the function $f_d(P^\ell)$ is given by the hypergeometric function*

$${}_2F_1\left(\frac{1}{n_d}, 1 - \frac{1}{n_d}; 1; -\kappa_d P^\ell\right).$$

Additionally, when $d = 8$, we have

$$f_8(P^\ell) = {}_2F_1\left(\frac{1}{2}, \frac{1}{2}; 1; 16P^{2\ell}\right) = f_4(-P^{2\ell}).$$

Proof. This follows immediately from the differential equation (3.4) and the initial condition $f_d(P^\ell) = 1 + O(P^\ell)$. \square

We compute the 1-point and 2-point invariants of X_d explicitly:

Proposition 3.15. *Suppose $I^X(t)$ satisfies the Picard–Fuchs equation $\square_{\bar{\gamma}} I^X = 0$, where*

$$\square_{\bar{\gamma}} = \hat{h}^4 - P^\ell(\kappa \hat{h}^2 + \kappa z \hat{h} + \lambda z^2) - \mu P^{2\bar{\gamma}}.$$

Then

$$(i) \langle \text{pt} \rangle^X = \lambda Q^{\bar{\gamma}};$$

- (ii) $\langle h^2, h^2 \rangle^X = d(\kappa - 2\lambda)Q^{\bar{\gamma}}$;
- (iii) $\langle \text{pt}, \text{pt} \rangle^X = \frac{\lambda^2 + \mu}{2d}Q^{2\bar{\gamma}}$.

Proof. It follows from the Picard–Fuchs equations that for $\beta = n\bar{\gamma}$,

$$(h + nz)^4 I_\beta^X = (\kappa h^2 + \kappa(2n - 1)zh + (\kappa n^2 - \kappa n + \lambda)z^2)I_{\beta-\bar{\gamma}}^X + \mu I_{\beta-2\bar{\gamma}}.$$

Together with the initial condition $I_0 = 1$, we get

$$\begin{aligned} I_{\bar{\gamma}}^X &= \frac{\kappa h^2 + \kappa zh + \lambda z^2}{((h)^{\bar{1}})^4}, \\ I_{2\bar{\gamma}}^X &= \frac{(2\kappa + \lambda)\lambda + \mu}{16z^4} + O(z^{-5}). \end{aligned} \tag{3.6}$$

These coefficients can also be computed directly from the I -functions.

Since

$$e^{\tau/z} \left(1 + \sum_{k \geq 0} \frac{1}{z^{k+2}} \langle \psi^k T_\mu \rangle^X T^\mu \right) = J^X(\tau) = I^X(t) = e^{t/z} \sum_{\beta} I_\beta^X P^\beta$$

and $\tau = t$,

$$\langle \psi^k T \rangle^X = [z^{-(k+2)}] \int_T \sum_{\beta} I_\beta^X Q^\beta,$$

where $[z^{-k}]F$ denotes the coefficient of z^{-k} in F . So for (i), it follows from (3.6) that

$$\langle \text{pt} \rangle^X = [z^{-2}] \int_{\text{pt}} \sum_{\beta} I_\beta^X Q^\beta = [z^{-2}] \int_{\text{pt}} I_{\bar{\gamma}}^X Q^{\bar{\gamma}} = \lambda Q^{\bar{\gamma}}.$$

By Proposition 2.3 and divisor equation,

$$\begin{aligned} \langle h^2, h^2 \rangle^X &= \delta^h \langle h^3 \rangle^X + \delta^h \langle h, \psi h^2 \rangle^X - \delta^h \langle h, T_\mu \rangle^X \langle T^\mu, h^2 \rangle^X \\ &= \langle h^3 \rangle^X + \langle h^3 \rangle^X + \langle \psi h^2 \rangle^X \\ &= 2d \langle \text{pt} \rangle^X + [z^{-3}] \int_{h^2} I_{\bar{\gamma}}^X Q^{\bar{\gamma}} \\ &= 2d \cdot \lambda Q^{\bar{\gamma}} + d(\kappa - 4\lambda)Q^{\bar{\gamma}} = d(\kappa - 2\lambda)Q^{\bar{\gamma}}, \end{aligned}$$

which is (ii).

Finally, it follows from Proposition 2.3, (i), (ii), and (3.6) that

$$\begin{aligned} d \langle \text{pt}, \text{pt} \rangle^X &= \delta^H \langle H^2, \psi \text{pt} \rangle^X - \delta^H \langle H^2, T_\mu \rangle^X \langle T^\mu, \text{pt} \rangle^X \\ &= 2(\delta^H \langle H, \psi^2 \text{pt} \rangle^X - \delta^H \langle H, T_\mu \rangle^X \langle T^\mu, \psi \text{pt} \rangle^X) - \frac{1}{d} \langle H^2, H^2 \rangle^X \langle H, \text{pt} \rangle^X \\ &= 8 \langle \psi^2 \text{pt} \rangle^X - 2 \langle \text{pt} \rangle^X \langle 1, \psi \text{pt} \rangle^X - (\kappa - 2\lambda)Q^{\bar{\gamma}} \cdot \lambda Q^{\bar{\gamma}} \\ &= 8 \cdot \frac{(2\kappa + \lambda)\lambda + \mu}{16} Q^{2\bar{\gamma}} - 2(\lambda Q^{\bar{\gamma}})^2 - (\kappa - 2\lambda)\lambda Q^{2\bar{\gamma}} \\ &= \frac{\lambda^2 + \mu}{2} Q^{2\bar{\gamma}}. \end{aligned}$$

which is (iii). \square

We compute the 1-point and 2-point invariants of Y_d explicitly. As a consequence we are able to determine the 3-point function $\langle E, E, E \rangle^Y$ (and hence the extremal function \mathbb{E}).

Proposition 3.16. *Suppose $I^Y(t)$ satisfies the Picard–Fuchs equation $\square_\gamma I^Y = \square_\ell I^Y = 0$, where*

$$\begin{aligned}\square_\gamma &= \hat{E}\hat{H} - P^\gamma, \\ \square_\ell &= \hat{F}^3 - P^\ell(\kappa\hat{F}^2 + \kappa z\hat{F} + \lambda z^2)\hat{E} - \mu P^{2\ell}(\hat{F} + z)\hat{E}(\hat{E} - z).\end{aligned}$$

Let $\tilde{\gamma} = \gamma + \ell$ be the lifting of γ such that $(E, \tilde{\gamma}) = 0$. Let $f(P^\ell) \in 1 + \mathbb{C}[[P^\ell]]$ be the (unique) solution of the differential equation

$$(1 + \kappa P^\ell - \mu P^{2\ell})\theta^2 f + (\kappa P^\ell - 2\mu P^{2\ell})\theta f + (\lambda P^\ell - \mu P^{2\ell})f = 0, \quad \theta = P^\ell \frac{d}{dP^\ell},$$

and let $u = u(P^\ell) = 1 + \kappa P^\ell - \mu P^{2\ell}$. Then

- (i) $\langle \text{pt} \rangle^Y = P^\gamma + \lambda Q^{\tilde{\gamma}}$;
- (ii) $\langle H^2, H^2 \rangle^Y = d(P^\gamma + (\kappa - 2\lambda)Q^{\tilde{\gamma}})$;
- (iii) $\langle \text{pt}, \text{pt} \rangle^Y = \frac{1}{d}(\frac{\lambda^2 - \mu}{2}Q^{2\tilde{\gamma}} + \lambda P^{2\gamma + \ell} + u \frac{\theta f}{f} P^{2\gamma})$;
- (iv) $\langle H^2, E^2 \rangle^Y = -du(f + \theta f)P^\gamma$;
- (v) $\langle E, E, E \rangle^Y = \frac{d}{uf^3}$.

Proof. We will prove the statements in the order (i), (ii), (iv), (v) and finally (iii).

It follows from the Picard–Fuchs equations that for $\beta = m\ell + n\gamma$,

$$\begin{aligned}0 &= (E + (n-m)z)(H + nz)I_\beta^Y - I_{\beta-\gamma}^Y, \\ 0 &= (F + mz)^3 I_\beta^Y \\ &\quad - (\kappa F^2 + \kappa(2m-1)zF + (\kappa m^2 - \kappa m + \lambda)z^2)(E + (n-m+1)z)I_{\beta-\ell}^Y \\ &\quad - \mu(F + (m-1)z)(E + (n-m+2)z)(E + (n-m+1)z)I_{\beta-2\ell}^Y.\end{aligned}$$

Together with the initial condition $I_0 = 1$, we get

$$\begin{aligned}I_\gamma^Y &= \frac{1}{(E)^{\bar{1}}(H)^{\bar{1}}}, \\ I_{\ell+\gamma}^Y &= \frac{\kappa F^2 + \kappa zF + \lambda z^2}{((F)^{\bar{1}})^3(H)^{\bar{1}}}\end{aligned}\tag{3.7}$$

and $I_{m\ell+\gamma}^Y$ is a multiple of E for $m \geq 2$;

$$\begin{aligned}I_{2\gamma}^Y &= \frac{1}{(E)^{\bar{2}}(H)^{\bar{2}}}, \\ I_{\ell+2\gamma}^Y &= \frac{\kappa F^2 + \kappa zF + \lambda z^2}{((F)^{\bar{1}})^3(E)^{\bar{1}}(H)^{\bar{2}}}, \\ I_{2\ell+2\gamma}^Y &= \frac{(2\kappa + \lambda)\lambda + \mu}{16z^4} + O(z^{-5})\end{aligned}\tag{3.8}$$

and $I_{m\ell+2\gamma}^Y$ is a multiple of E for $m \geq 3$. These coefficients can also be computed directly from the I -functions.

Since

$$e^{\tau/z} \left(1 + \sum_{k \geq 0} \frac{1}{z^{k+2}} \langle \psi^k T_\mu \rangle^Y T^\mu \right) = J^Y(\tau) = I^Y(t) = e^{t/z} \sum_{\beta} I_\beta^Y P^\beta,$$

and $\tau = t - gE$,

$$\langle \psi^k T \rangle^Y = [z^{-(k+2)}] \int_T e^{gE} \sum_{\beta} I_\beta^Y P^\beta.$$

So for (i), it follows from (3.7) that

$$\langle \text{pt} \rangle^Y = [z^{-2}] \int_{\text{pt}} e^{gE} \sum_{\beta} I_\beta^Y P^\beta = [z^{-2}] \int_{\text{pt}} I_\gamma^Y P^\gamma + I_{\ell+\gamma}^Y P^{\ell+\gamma} = P^\gamma + \lambda Q^{\tilde{\gamma}}.$$

Let $a \in H^4(Y_d)$ and $D \in H^2(Y_d)$. Then, by Proposition 2.3,

$$\langle HD, a \rangle^Y = \delta^D \langle Ha \rangle^Y + \delta^H \langle D, \psi a \rangle^Y - \delta^H \langle D, T_\mu \rangle^Y \langle T^\mu, a \rangle^Y = \delta^D \langle Ha \rangle^Y + \langle Da \rangle^Y + \delta^D \langle \psi a \rangle^Y.$$

If we take $a = H^2$ and $D = H$, then

$$\begin{aligned} \langle H^2, H^2 \rangle^Y &= 2d \langle \text{pt} \rangle^Y + [z^{-3}] \int_{H^2} e^{gE} \sum_{\beta} I_\beta^Y P^\beta \\ &= 2d(P^\gamma + \lambda Q^{\tilde{\gamma}}) + [z^{-3}] \int_{H^2} I_\gamma^Y P^\gamma + I_{\ell+\gamma}^Y P^{\ell+\gamma} \\ &= 2d(P^\gamma + \lambda Q^{\tilde{\gamma}}) + (-d)P^\gamma + d(\kappa - 4\lambda)Q^{\tilde{\gamma}} = d(P^\gamma + (\kappa - 2\lambda)Q^{\tilde{\gamma}}), \end{aligned}$$

which is (ii).

If we take $a = E^2$ and $D = E$, then

$$0 = \langle HE, a \rangle^Y = \langle E^3 \rangle^Y + \delta^E \langle \psi E^2 \rangle^Y.$$

This implies

$$\delta^E \langle \psi E^2 \rangle^Y = f \cdot d(P^\gamma + \lambda Q^{\tilde{\gamma}}) = \delta^E (-du(f + \theta f)P^\gamma),$$

i.e., $\langle \psi E^2 \rangle^Y = -du(f + \theta f)P^\gamma + cQ^{\tilde{\gamma}}$ for some constant c . Comparing the $Q^{\tilde{\gamma}}$ -coefficients of both sides we get $c = 0$. Hence, we get (iv):

$$\langle H^2, E^2 \rangle^Y = \langle \psi E^2 \rangle^Y = -du(f + \theta f)P^\gamma.$$

On the other hand,

$$\langle E^2, H^2 \rangle^Y = \delta^E \langle E, \psi H^2 \rangle^Y - \delta^E \langle E, T_\mu \rangle^Y \langle T^\mu, H^2 \rangle^Y = (\delta^E)^2 \langle \psi H^2 \rangle^Y - \delta^E \langle E, T_\mu \rangle^Y \langle T^\mu, H^2 \rangle^Y$$

implies

$$\frac{1}{d} \langle E, E, E \rangle^Y \langle E^2, H^2 \rangle^Y = (\delta^E)^2 (dP^\gamma) = -d \frac{f+\theta f}{f^2} P^\gamma.$$

Together with (iv), this gives (v).

Finally, it follows from Proposition 2.3, (i), (ii), (iv) and (3.8) that

$$\begin{aligned}
d\langle \text{pt}, \text{pt} \rangle^Y &= \delta^H \langle H^2, \psi \text{pt} \rangle^Y - \delta^H \langle H^2, T_\mu \rangle^Y \langle T^\mu, \text{pt} \rangle^Y \\
&= 2(\delta^H \langle H, \psi^2 \text{pt} \rangle^Y - \delta^H \langle H, T_\mu \rangle^Y \langle T^\mu, \psi \text{pt} \rangle^Y) \\
&\quad - \frac{1}{d}(\langle H^2, H^2 \rangle^Y \langle H, \text{pt} \rangle^Y + \langle H^2, E^2 \rangle^Y \langle E, \text{pt} \rangle^Y) \\
&= 8\langle \psi^2 \text{pt} \rangle^Y - 2\langle \text{pt} \rangle^Y \langle 1, \psi \text{pt} \rangle^Y \\
&\quad - ((P^\gamma + (\kappa - 2\lambda)Q^{\tilde{\gamma}})(P^\gamma + \lambda Q^{\tilde{\gamma}}) - u(f + \theta f)P^\gamma(\frac{P^\gamma}{f})) \\
&= 8\left(\frac{1}{4}P^{2\gamma} + \frac{\lambda}{2}P^{\ell+2\gamma} + \frac{(2\kappa+\lambda)\lambda+\mu}{16}P^{2\ell+2\gamma}\right) - 2(P^\gamma + \lambda Q^{\tilde{\gamma}})^2 \\
&\quad - (P^\gamma + (\kappa - 2\lambda)Q^{\tilde{\gamma}})(P^\gamma + \lambda Q^{\tilde{\gamma}}) + (P^{2\gamma} + \kappa P^{\ell+2\gamma} - \mu P^{2\ell+2\gamma}) + u\frac{\theta f}{f}P^{2\gamma} \\
&= \frac{\lambda^2-\mu}{2}Q^{2\tilde{\gamma}} + \lambda P^{\ell+2\gamma} + u\frac{\theta f}{f}P^{2\gamma},
\end{aligned}$$

which is (iii). \square

Consider the lifting

$$\begin{aligned}
\mathbb{C}[[Q^{\text{NE}(X_d)}]] &\xrightarrow{\phi^*} \mathbb{C}[[Q^{\text{NE}(Y_d)}]] \\
Q^\gamma &\longmapsto Q^{\tilde{\gamma}} = P^\gamma P^\ell.
\end{aligned}$$

Comparing Proposition 3.15 and Proposition 3.16, we observe that

$$\begin{aligned}
\langle \text{pt} \rangle^Y - \phi^* \langle \text{pt} \rangle^X &= P^\gamma = P^{-\ell} Q^{\tilde{\gamma}}, \\
\langle H^2, H^2 \rangle^Y - \phi^* \langle h^2, h^2 \rangle^X &= dP^\gamma = dP^{-\ell} Q^{\tilde{\gamma}}, \\
\langle \text{pt}, \text{pt} \rangle^Y - \phi^* \langle \text{pt}, \text{pt} \rangle^X &= \frac{1}{d}(-\mu + \lambda P^{-\ell} + u\frac{\theta f}{f}P^{-2\ell})Q^{2\tilde{\gamma}} \\
&= \frac{1}{d}(-\mu(1 + \frac{\theta f}{f}) + (\lambda + \kappa\frac{\theta f}{f})P^{-\ell} + \frac{\theta f}{f}P^{-2\ell})Q^{2\tilde{\gamma}}.
\end{aligned}$$

For each d , denote by Sing_{f_d} the set of regular singular points of the differential equation (3.4). At each singularity, the corresponding pair of local exponents can be computed explicitly. These are summarized in Table 5.

d	Sing_{f_d}	local exponents
1	$\{0, \infty, -\frac{1}{432}\}$	$(0, 0), (\frac{1}{6}, \frac{5}{6}), (0, 0)$
2	$\{0, \infty, -\frac{1}{64}\}$	$(0, 0), (\frac{1}{4}, \frac{3}{4}), (0, 0)$
3	$\{0, \infty, -\frac{1}{27}\}$	$(0, 0), (\frac{1}{3}, \frac{2}{3}), (0, 0)$
4	$\{0, \infty, -\frac{1}{16}\}$	$(0, 0), (\frac{1}{2}, \frac{1}{2}), (0, 0)$
5	$\{0, \infty, \frac{11-5\sqrt{5}}{2}, \frac{11+5\sqrt{5}}{2}\}$	$(0, 0), (1, 1), (0, 0), (0, 0)$
6I	$\{0, \infty, -\frac{1}{8}, 1\}$	$(0, 0), (1, 1), (0, 0), (0, 0)$
6II	$\{0, \infty, -\frac{1}{9}, -1\}$	$(0, 0), (1, 1), (0, 0), (0, 0)$
8	$\{0, \infty, -\frac{1}{4}, \frac{1}{4}\}$	$(0, 0), (1, 1), (0, 0), (0, 0)$

TABLE 5. Regular singular points of f_d and their corresponding local exponents.

Let $y = P^{-\ell}$. Then, near $y = 0$, we can rewrite (3.4) as

$$(y^2 + \kappa_d y - \mu_d)(y \frac{d}{dy})^2 f_d - (\kappa_d y - 2\mu_d)(y \frac{d}{dy}) f_d + (\lambda_d y - \mu_d) f_d = 0. \quad (3.9)$$

When $\mu_d = 0$, i.e., $d \in \{1, 2, 3, 4\}$, the local exponent at $y = 0$ is $(\frac{1}{n_d}, \frac{n_d-1}{n_d})$, where $n_d = 6, 4, 3, 2$ for $d = 1, 2, 3, 4$, respectively, just as in the proof of Lemma 3.5. This means that $f_d \in y^{1/n_d} \mathbb{C}[[y]] \oplus y^{1-1/n_d} \mathbb{C}[[y]]$ for $d \leq 3$ and $f_d \in y^{1/2} \mathbb{C}[[y]] \oplus y^{1/2} \log y \mathbb{C}[[y]]$ for $d = 4$.

When $\mu_d \neq 0$, i.e., $d \in \{5, 6I, 6II, 8\}$, the local exponent at $y = 0$ is $(1, 1)$. This means that $f_d \in y \mathbb{C}[[y]] \oplus y \log y \mathbb{C}[[y]]$. Let $f_d^{\text{reg}} \in y \mathbb{C}[[y]]$ be a nonzero monodromy invariant solution (unique up to a scalar) of (3.9). Then

$$\frac{\theta f_d^{\text{reg}}}{f_d^{\text{reg}}} = \frac{-y \frac{\partial}{\partial y} f_d^{\text{reg}}}{f_d^{\text{reg}}} = -1 + O(y).$$

By this observation, we found a suitable ring R_d and a suitable ideal I_d such that for $\vec{b} = (\text{pt}, (h^2, h^2), (\text{pt}, \text{pt}))$, the invariant $\langle \phi^* \vec{b} \rangle^{Y_d} \in R_d[Q^{\tilde{\gamma}}]$ and

$$\langle \phi^* \vec{b} \rangle^{Y_d} - \phi^* \langle \vec{b} \rangle^{X_d} \in I_d[Q^{\tilde{\gamma}}]. \quad (3.10)$$

Definition 3.17. Let $y = P^{-\ell}$ and let $v_d = \frac{\theta f_d}{f_d}$. Consider the ring $S_d = \mathbb{C}(y)[v_d] \subseteq \mathbb{C}(\!(y)\!)[v_d]$. We define the subring $R_d \subseteq S_d$ and its ideal I_d as follows:

- for $d \in \{1, 2, 3, 4\}$,

$$R_d = S_d \cap (\mathbb{C} + y \cdot \mathbb{C}[[y]][v_d]), \quad I_d = S_d \cap (y \cdot \mathbb{C}[[y]][v_d]);$$

- for $d \in \{5, 6I, 6II, 8\}$, let $v_d^{\text{reg}} = \frac{\theta f_d^{\text{reg}}}{f_d^{\text{reg}}} \in \mathbb{C}[[y]]$, $w_d = v_d - v_d^{\text{reg}}$,

$$R_d = S_d \cap (\mathbb{C}[[y]] + w_d \cdot \mathbb{C}(\!(y)\!)[w_d]), \quad I_d = S_d \cap (y \cdot \mathbb{C}[[y]] + w_d \cdot \mathbb{C}(\!(y)\!)[w_d]).$$

Lemma 3.18. The operator $\theta = P^\ell \frac{d}{dP^\ell}$ sends elements in R_d to I_d .

Proof. For $p(v_d) = \sum_n p_n(y)v_d^n \in \mathbb{C}(y)[v_d] = S_d$,

$$\theta p(v_d) = \sum (\theta p_n(y))v_d^n + \sum p_n(y)nv_d^{n-1}(\theta v_d).$$

To prove that $\theta p(v_d) \in S_d$, it suffices to prove that $\theta v_d \in S_d$. By (3.4) (or (3.9)),

$$\begin{aligned} \theta v_d &= \frac{\theta^2 f_d}{f_d} - \frac{(\theta f_d)^2}{f_d^2} \\ &= -\frac{\kappa_d y - 2\mu_d}{y^2 + \kappa_d y - \mu_d} v_d - \frac{\lambda_d y - \mu_d}{y^2 + \kappa_d y - \mu_d} - v_d^2 \in S_d \cap \mathbb{C}[[y]][v_d]. \end{aligned} \quad (3.11)$$

When $d = 1, 2, 3, 4$, if $p = p_0 + y \sum_n p_n(y)v_d^n \in \mathbb{C} + y \cdot \mathbb{C}[[y]][v_d]$, then

$$\theta p = y \sum p_n(y)v_d^n + y \sum (\theta p_n(y))v_d^n + y \sum p_n(y)nv_d^{n-1}(\theta v_d) \in y \cdot \mathbb{C}[[y]][v_d]$$

by (3.11). Hence, $\theta R_d \subseteq I_d$.

When $d = 5, 6I, 6II, 8$, if $p = p_0(y) + \sum_{n>0} p_n(y)w_d^n \in \mathbb{C}[[y]] + w_d \cdot \mathbb{C}((y))[w_d]$, then

$$\theta p = \theta p_0(y) + \sum_{n>0} (\theta p_n(y))w_d^n + \sum_{n>0} (\theta p_n(y))nw_d^{n-1}(\theta w_d).$$

To show that $\theta p \in I_d$, it suffices to show that $\theta w_d \in y \cdot \mathbb{C}[[y]] + w_d \cdot \mathbb{C}((y))[w_d]$. Indeed, by (3.11), which also holds for v_d^{reg} ,

$$\begin{aligned} \theta w_d &= \theta v_d^{\text{reg}} - \theta v_d \\ &= -\frac{\kappa_d y - 2\mu_d}{y^2 + \kappa_d y - \mu_d} (v_d^{\text{reg}} - v_d) - ((v_d^{\text{reg}})^2 - v_d^2) \\ &= \frac{\kappa_d y - 2\mu_d}{y^2 + \kappa_d y - \mu_d} w_d + w_d(w_d + 2v_d^{\text{reg}}) \in w_d \cdot \mathbb{C}((y))[w_d]. \end{aligned}$$

Hence, $\theta R_d \subseteq I_d$. □

We now extend (3.10) to any n -point GW invariants:

Proposition 3.19. *Fix $d \neq 7$. Consider the ring $R_{\text{loc}} = R_d[Q^{\tilde{\gamma}}]$ and its ideal $I_{\text{loc}} = I_d[Q^{\tilde{\gamma}}]$. Then for $n \geq 1$, $\vec{a} \in H(S_d)^{\oplus n}$, the invariant $\langle j_{d*}\vec{a} \rangle^{Y_d}$ lies in R_{loc} and*

$$\langle j_{d*}\vec{a} \rangle^{Y_d} \equiv \phi^* \langle i_{d*}\vec{a} \rangle^{X_d} \pmod{I_{\text{loc}}}.$$

Proof. Since $H(S_d)$ is generated by $i_d^* h^e$, $e = 0, 1, 2$, $i_{d*}H(S_d)$ is generated by $h^{e+1} = i_{d*}i_d^* h^e$ and $j_{d*}H(S_d)$ is generated by $H^{e+1} = j_{d*}i_d^* h^e$, $e = 0, 1, 2$.

Let $\vec{a} = (a_1, \dots, a_n)$, $b_\nu = i_{d*}a_\nu$, $c_\nu = j_{d*}a_\nu$. By divisor equation, we may assume that each $a_\nu = i_d^* h$ or $i_d^* h^2$. We prove the result by induction on n and the curve degree $\frac{1}{2} \sum \deg a_\nu$. For simplicity, we will omit “ d ” in the subscripts.

For $n \leq 2$, the result follows from the fundamental class axiom, the divisor equation, Proposition 3.15 and Proposition 3.16.

If i^*h appears in $\{a_1, \dots, a_n\}$, say $a_1 = i^*h$, then by the reconstruction formula Theorem 2.1 (or Proposition 2.3),

$$\begin{aligned}\langle b_1, \dots, b_n \rangle^X &= \langle h^2, b_2, b_3, \dots, b_n \rangle^X \\ &= \langle h, hb_2, b_3, \dots, b_n \rangle^X \\ &\quad + \sum_{3 \in I_L, 2 \in I_R} \langle h, \vec{b}_{I_L}, \bar{T}_\mu \rangle^X \cdot \delta^h \langle \bar{T}^\mu, \vec{b}_{I_R} \rangle^X - \sum_{2, 3 \in I_R} \delta^h \langle h, \vec{b}_{I_L}, \bar{T}_\mu \rangle^X \cdot \langle \bar{T}^\mu, \vec{b}_{I_R} \rangle^X, \\ \langle c_1, \dots, c_n \rangle^Y &= \langle H^2, c_2, c_3, \dots, c_n \rangle^Y \\ &= \langle H, Hc_2, c_3, \dots, c_n \rangle^Y \\ &\quad + \sum_{3 \in I_L, 2 \in I_R} \langle H, \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \delta^H \langle T^\mu, \vec{c}_{I_R} \rangle^Y - \sum_{2, 3 \in I_R} \delta^H \langle H, \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \langle T^\mu, \vec{c}_{I_R} \rangle^Y.\end{aligned}$$

Hence, to show that

$$\langle c_1, \dots, c_n \rangle^Y \equiv \phi^* \langle b_1, \dots, b_n \rangle^X \pmod{\mathbf{l}_{\text{loc}}},$$

it suffices to show that:

- (i) $\langle H, Hc_2, c_3, \dots, c_n \rangle^Y \equiv \phi^* \langle h, hb_2, b_3, \dots, b_n \rangle^X \pmod{\mathbf{l}_{\text{loc}}};$
- (ii) $\langle \vec{c}_I, T \rangle^Y \equiv \phi^* \langle \vec{b}_I, \bar{T} \rangle^X \pmod{\mathbf{l}_{\text{loc}}}$ for $(\bar{T}, T) = (h^e, H^e)$ with lower curve degree;
- (iii) $\langle \vec{c}_{I_L}, E \rangle^Y \langle E^2, \vec{c}_{I_R} \rangle^Y \in \mathbf{l}_{\text{loc}}.$

The (i) follows from divisor equation and the induction hypothesis:

$$\langle Hc_2, c_3, \dots, c_n \rangle^Y \equiv \phi^* \langle hb_2, b_3, \dots, b_n \rangle^X \pmod{\mathbf{l}_{\text{loc}}}.$$

The (ii) follows from the induction hypothesis, the fundamental class axiom and the divisor equation (when $e = 0, 1$). For (iii), we show that $\langle \vec{c}, E \rangle^Y \in f^{-1}\mathbf{l}_{\text{loc}}$ and $\langle E^2, \vec{c} \rangle^Y \in f\mathbf{R}_{\text{loc}}[v]$ for $\vec{c} \in j_*H(S)^{\oplus m}$, $1 \leq m < n$, so that

$$\langle \vec{c}_{I_L}, E \rangle^Y \langle E^2, \vec{c}_{I_R} \rangle^Y \in f^{-1}\mathbf{l}_{\text{loc}} \cdot f\mathbf{R}_{\text{loc}}[v] \subseteq \mathbf{l}_{\text{loc}}.$$

Indeed, by divisor equation, $E \cup j_*H(S_d) = 0$, Lemma 3.18 and the induction hypothesis $\langle \vec{c} \rangle^Y \in \mathbf{R}_{\text{loc}}$, we have

$$\langle \vec{c}, E \rangle^Y = \delta^E \langle \vec{c} \rangle^Y \in f^{-1}\theta\mathbf{R}_{\text{loc}} \subseteq f^{-1}\mathbf{l}_{\text{loc}}.$$

If $m = 1$, then $\langle E^2, \vec{c} \rangle \in f\mathbf{R}_{\text{loc}}[v]$ by Proposition 3.16 (iv). If $m \geq 2$, apply the reconstruction formula, we get the equation

$$\langle E^2, \vec{c} \rangle^Y = \sum_{1 \in I_L, 2 \in I_R} \langle E, \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \delta^E \langle T^\mu, \vec{c}_{I_R} \rangle^Y - \sum_{1, 2 \in I_R} \delta^E \langle E, \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \langle T^\mu, \vec{c}_{I_R} \rangle^Y.$$

When $I_L = \emptyset$, it follows from Proposition 3.16 (v) that

$$\begin{aligned}\delta^E \langle E, \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \langle T^\mu, \vec{c}_{I_R} \rangle^Y &= (\delta^E)^2 \langle T_\mu \rangle^Y \cdot \langle T^\mu, \vec{c} \rangle^Y \\ &= \left(\frac{1}{d} \langle E, E, E \rangle - 1\right) \langle E^2, \vec{c} \rangle^Y = \left(\frac{1}{uf^3} - 1\right) \langle E^2, \vec{c} \rangle^Y.\end{aligned}$$

So we get

$$\begin{aligned} \frac{1}{uf^3} \langle E^2, \vec{c} \rangle^Y &= \sum_{1 \in I_L, 2 \in I_R} \delta^E \langle \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \delta^E \langle T^\mu, \vec{c}_{I_R} \rangle^Y \\ &\quad - \sum_{I_L \neq \emptyset, 1, 2 \in I_R} (\delta^E)^2 \langle \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \langle T^\mu, \vec{c}_{I_R} \rangle^Y. \end{aligned} \quad (3.12)$$

For $d = 1, 2, 3, 4$, each term in the right hand side lies in one of the followings (by induction on m):

$$\begin{aligned} \delta^E \mathsf{R}_{\text{loc}} \cdot \delta^E \mathsf{R}_{\text{loc}} &\subseteq f^{-1} \mathsf{l}_{\text{loc}} \cdot f^{-1} \mathsf{l}_{\text{loc}} = f^{-2} \mathsf{l}_{\text{loc}}^2, \\ (\delta^E)^2 \mathsf{R}_{\text{loc}} \cdot \mathsf{R}_{\text{loc}} &\subseteq f^{-2} \mathsf{l}_{\text{loc}} \cdot \mathsf{R}_{\text{loc}} = f^{-2} \mathsf{l}_{\text{loc}}, \\ \delta^E(f^{-1} \mathsf{l}_{\text{loc}}) \cdot \delta^E(f \mathsf{R}_{\text{loc}}[v]) &\subseteq f^{-2} \mathsf{l}_{\text{loc}} \cdot \mathsf{R}_{\text{loc}}[v] = f^{-2} \mathsf{l}_{\text{loc}}, \\ (\delta^E)^2(f^{-1} \mathsf{l}_{\text{loc}}) \cdot f \mathsf{R}_{\text{loc}}[v] &\subseteq f^{-3} \mathsf{l}_{\text{loc}} \cdot f \mathsf{R}_{\text{loc}}[v] = f^{-2} \mathsf{l}_{\text{loc}}, \\ f^{-1} \mathsf{l}_{\text{loc}} \cdot (\delta^E)^2(f \mathsf{R}_{\text{loc}}[v]) &\subseteq f^{-1} \mathsf{l}_{\text{loc}} \cdot f^{-1} \mathsf{R}_{\text{loc}}[v] = f^{-2} \mathsf{l}_{\text{loc}}. \end{aligned}$$

Therefore, the right hand side of (3.12) lies in $f^{-2} \mathsf{l}_{\text{loc}}$, and hence

$$\langle E^2, c \rangle^Y \in u f \mathsf{l}_{\text{loc}} \subseteq f \mathsf{R}_{\text{loc}}[v].$$

For $d = 5, 6\mathbb{I}, 6\mathbb{II}, 8$, each term in the right hand side lies in one of the followings (by induction on m):

$$\begin{aligned} \delta^E \mathsf{R}_{\text{loc}} \cdot \delta^E \mathsf{R}_{\text{loc}} &\subseteq f^{-1} \mathsf{l}_{\text{loc}} \cdot f^{-1} \mathsf{l}_{\text{loc}} = f^{-2} \mathsf{l}_{\text{loc}}^2, \\ (\delta^E)^2 \mathsf{R}_{\text{loc}} \cdot \mathsf{R}_{\text{loc}} &\subseteq f^{-2} \mathsf{l}_{\text{loc}}^2 \cdot \mathsf{R}_{\text{loc}} = f^{-2} \mathsf{l}_{\text{loc}}^2, \\ \delta^E(f^{-1} \mathsf{l}_{\text{loc}}) \cdot \delta^E(f \mathsf{R}_{\text{loc}}[v]) &\subseteq f^{-2} \mathsf{l}_{\text{loc}}^2 \cdot \mathsf{R}_{\text{loc}}[v] = f^{-2} \mathsf{l}_{\text{loc}}^2, \\ (\delta^E)^2(f^{-1} \mathsf{l}_{\text{loc}}) \cdot f \mathsf{R}_{\text{loc}}[v] &\subseteq f^{-3} \mathsf{l}_{\text{loc}}^3 \cdot f \mathsf{R}_{\text{loc}}[v] = f^{-2} \mathsf{l}_{\text{loc}}^3, \\ f^{-1} \mathsf{l}_{\text{loc}} \cdot (\delta^E)^2(f \mathsf{R}_{\text{loc}}[v]) &\subseteq f^{-1} \mathsf{l}_{\text{loc}} \cdot f^{-1} \mathsf{l}_{\text{loc}} = f^{-2} \mathsf{l}_{\text{loc}}^2. \end{aligned}$$

Therefore, the right hand side of (3.12) lies in $f^{-2} \mathsf{l}_{\text{loc}}^2$, and hence

$$\langle E^2, \vec{c} \rangle^Y \in u f \mathsf{l}_{\text{loc}}^2 \subseteq f \mathsf{R}_{\text{loc}}.$$

If $a_\nu = i_d^* h^2$ for all ν , then, by the reconstruction formula and (i), (ii), (iii) above,

$$\begin{aligned} \langle c_1, \dots, c_n \rangle^Y &= \sum \langle H^2, \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \delta^H \langle T^\mu, \vec{c}_{I_R} \rangle^Y - \sum \delta^H \langle H^2, \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \langle T^\mu, \vec{c}_{I_R} \rangle^Y \\ &\equiv \phi^* \left(\sum \langle h^2, \vec{b}_{I_L}, \bar{T}_\mu \rangle^X \cdot \delta^h \langle \bar{T}^\mu, \vec{b}_{I_R} \rangle^X - \sum \delta^h \langle h^2, \vec{b}_{I_L}, \bar{T}_\mu \rangle^X \cdot \langle \bar{T}^\mu, \vec{b}_{I_R} \rangle^X \right) \\ &= \phi^* \langle b_1, \dots, b_n \rangle^X \pmod{\mathsf{l}_{\text{loc}}}, \end{aligned}$$

as desired. \square

3.3. Gromov–Witten invariants for $d = 7$. The case of $d = 7$ is slightly different from other cases since we are unable to identify the curve classes of $X_7 = \text{Bl}_{\text{pt}} \mathbb{P}^3$ only by its degree. We will follow the same approach as in the case $d \neq 7$ and Convention 3.4.

Since X_7 is a toric Fano variety, we can write down its J -function completely. Recall that h_1 is the pullback, via the blow-up $X_7 \rightarrow \mathbb{P}^3$, of the hyperplane class on \mathbb{P}^3 , e is the exceptional divisor in X_7 , $h_2 = h_1 - e$, and $h = h_1 + h_2 = c_1(\mathcal{O}_{X_7}(1))$ (Remark A.6). Then $H(X_7) = \mathbb{C}[h_1, h_2]/(h_1(h_1 - h_2), h_2^3)$. Let $\bar{\gamma}_1 = h_2^2$ and $\bar{\gamma}_2 = h_2(h_1 - h_2)$. Then $(h_i, \bar{\gamma}_j)_{X_7} = \delta_{ij}$ and the Mori cone $\text{NE}(X_7)$ of X_7 is generated by $\bar{\gamma}_1$ and $\bar{\gamma}_2$. The J -function of X_7 is then

$$J^{X_7}(t, z^{-1}) = e^{t/z} \sum_{m_1, m_2 \geq 0} \frac{1}{(h_1)^{\overline{m_1}} (h_1 - h_2)^{\overline{m_1 - m_2}} ((h_2)^{\overline{m_2}})^3} Q^{m_1 \bar{\gamma}_1 + m_2 \bar{\gamma}_2}.$$

Since S_7 is a zero locus of a general global section of the bundle $\mathcal{O}_{T_7}(1) = \mathcal{O}(h_1 + h_2)$ over X_7 , we can write down the I -function of Y_7 :

$$\begin{aligned} I^{Y_7}(t) &= e^{t/z} \sum_{n, m_1, m_2 \geq 0} \pi_7^* i_7^* J_{m_1 \bar{\gamma}_1 + m_2 \bar{\gamma}_2}^{X_7} \cdot \frac{(F)^{\overline{m_1 + m_2}}}{(E)^{\overline{n - m_1 - m_2}} (H)^{\overline{n}}} P^{m_1 \ell_1 + m_2 \ell_2 + n \gamma} \\ &= e^{t/z} \sum_{n, m_1, m_2 \geq 0} \frac{(F_1 + F_2)^{\overline{m_1 + m_2}}}{(F_1)^{\overline{m_1}} (F_1 - F_2)^{\overline{m_1 - m_2}} ((F_2)^{\overline{m_2}})^3 (E)^{\overline{n - m_1 - m_2}} (H)^{\overline{n}}} P^{m_1 \ell_1 + m_2 \ell_2 + n \gamma} \end{aligned} \quad (3.13)$$

where, as before, $F_1 = \pi_7^* i_7^* h_1$, $F_2 = \pi_7^* i_7^* h_2$, $F = F_1 + F_2 = \pi_7^* i_7^* h$, γ is the fiber class of the projective bundle $Y_7 \rightarrow S_7$, and $\ell_1, \ell_2 \in H_2(Y_7)^{G_{S_7}}$ are chosen such that $(H, \ell_i) = 0$, $(F_i, \ell_j) = \delta_{ij}$. Then γ , ℓ_1 , and ℓ_2 generate the cone $\text{NE}(Y_7)/G_{S_7}$. We will, again, omit the group $G_{S_7} \cong \mathfrak{S}_2$ in the notation for simplicity.

We first compute the 1-point and 2-point invariants of X_7 .

Proposition 3.20. *We have*

- (i) $\langle \text{pt} \rangle^{X_7} = Q^{\bar{\gamma}_1}$;
- (ii) for divisor D_1 and $D_2 \in H^2(X_7)$,

$$\langle hD_1, hD_2 \rangle^{X_7} = (hD_1 D_2, 1) Q^{\bar{\gamma}_1} + 2(D_1, \bar{\gamma}_1)(D_2, \bar{\gamma}_1) Q^{\bar{\gamma}_1} + (D_1, \bar{\gamma}_2)(D_2, \bar{\gamma}_2) Q^{\bar{\gamma}_2};$$

$$(iii) \langle \text{pt}, \text{pt} \rangle^{X_7} = Q^{\bar{\gamma}_1 + \bar{\gamma}_2}.$$

Proof. As before,

$$\langle \psi^k T \rangle^{X_7} = [z^{-(k+2)}] \int_T \sum_{\beta} J_{\beta}^{X_7} Q^{\beta}.$$

So, for (i),

$$\langle \text{pt} \rangle^{X_7} = [z^{-2}] \int_{\text{pt}} \sum_{\beta} J_{\beta}^{X_7} Q^{\beta} = [z^{-2}] \int_{\text{pt}} J_{\bar{\gamma}_1}^{X_7} Q^{\bar{\gamma}_1} + J_{\bar{\gamma}_2}^{X_7} Q^{\bar{\gamma}_2} = Q^{\bar{\gamma}_1}.$$

For further use, we also compute that

$$\langle \psi^2 \text{pt} \rangle = [z^{-4}] \int_{\text{pt}} J_{2\bar{\gamma}_1}^{X_7} Q^{2\bar{\gamma}_1} + J_{\bar{\gamma}_1+\bar{\gamma}_2}^{X_7} Q^{\bar{\gamma}_1+\bar{\gamma}_2} + J_{2\bar{\gamma}_2}^{X_7} Q^{2\bar{\gamma}_2} = \frac{1}{4} Q^{2\bar{\gamma}_1} + Q^{\bar{\gamma}_1+\bar{\gamma}_2}, \quad (3.14)$$

and for $D \in H^2(X_7)$,

$$\langle \psi h D \rangle = [z^{-3}] \int_{HD} J_{\bar{\gamma}_1}^{X_7} Q^{\bar{\gamma}_1} + J_{\bar{\gamma}_2}^{X_7} Q^{\bar{\gamma}_2} = (hD, -2h_1 + h_2) Q^{\bar{\gamma}_1} + (hD, h_1 - h_2) Q^{\bar{\gamma}_2}. \quad (3.15)$$

For (ii), by reconstruction formula,

$$\langle hD_1, hD_2 \rangle = \langle h, hD_1 D_2 \rangle + \delta^{D_1} \langle h, \psi h D_2 \rangle - \sum \delta^{D_1} \langle h, T_\mu \rangle \langle T^\mu, h D_2 \rangle.$$

It is easy to check that the sum is 0 by degree argument. So, by divisor equation, (i) and (3.15),

$$\begin{aligned} \langle hD_1, hD_2 \rangle &= \langle hD_1 D_2 \rangle + \delta^{D_1} \langle h^2 D_2 \rangle + \delta^{D_1} \delta^h \langle \psi h D_2 \rangle \\ &= (hD_1 D_2, 1) \langle \text{pt} \rangle + (D_2, h^2) \delta^{D_1} \langle \text{pt} \rangle \\ &\quad + \delta^{D_1} ((hD_2, -2h_1 + h_2) Q^{\bar{\gamma}_1} + (hD_2, h_1 - h_2) Q^{\bar{\gamma}_2}) \\ &= (hD_1 D_2, 1) Q^{\bar{\gamma}_1} + (D_2, h^2 + h(-2h_1 + h_2)) (D_1, \bar{\gamma}_1) Q^{\bar{\gamma}_1} \\ &\quad + (D_2, h(h_1 - h_2)) (D_1, \bar{\gamma}_2) Q^{\bar{\gamma}_2} \\ &= (hD_1 D_2, 1) Q^{\bar{\gamma}_1} + 2(D_1, \bar{\gamma}_1) (D_2, \bar{\gamma}_1) Q^{\bar{\gamma}_1} + (D_1, \bar{\gamma}_2) (D_2, \bar{\gamma}_2) Q^{\bar{\gamma}_2}. \end{aligned}$$

Finally, by reconstruction formula, (i), (ii) and (3.14),

$$\begin{aligned} 7 \langle \text{pt}, \text{pt} \rangle &= \langle h^3, \text{pt} \rangle = \delta^h \langle h^2, \psi \text{pt} \rangle - \delta^h \langle h^2, T_\mu \rangle \langle T^\mu, \text{pt} \rangle \\ &= 2(\delta^h \langle h, \psi^2 \text{pt} \rangle - \delta^h \langle h, T_\mu \rangle \langle T^\mu, \psi \text{pt} \rangle) - \langle h^2, \bar{\gamma}_1 \rangle \langle h_1, \text{pt} \rangle - \langle h^2, \bar{\gamma}_2 \rangle \langle h_2, \text{pt} \rangle \\ &= 8 \langle \psi^2 \text{pt} \rangle - 2(\delta^h)^2 \langle \text{pt} \rangle \langle 1, \psi \text{pt} \rangle - \langle h^2, h(-\frac{1}{2}h_1 + h_2) \rangle \delta^{h_1} \langle \text{pt} \rangle \\ &= 2Q^{2\bar{\gamma}} + 8Q^{\bar{\gamma}_1+\bar{\gamma}_2} - 2Q^{2\bar{\gamma}_1} - (h^2(-\frac{1}{2}h_1 + h_2), 1) Q^{2\bar{\gamma}_1} \\ &\quad - 2(h, \bar{\gamma}_1)((-\frac{1}{2}h_1 + h_2), \bar{\gamma}_1) Q^{2\bar{\gamma}_1} - (h, \bar{\gamma}_2)((-\frac{1}{2}h_1 + h_2), \bar{\gamma}_2) Q^{\bar{\gamma}_1+\bar{\gamma}_2} \\ &= 2Q^{2\bar{\gamma}} + 8Q^{\bar{\gamma}_1+\bar{\gamma}_2} - 2Q^{2\bar{\gamma}} - Q^{2\bar{\gamma}_1} + Q^{\bar{\gamma}_1+\bar{\gamma}_2} = 7Q^{\bar{\gamma}_1+\bar{\gamma}_2}, \end{aligned}$$

which is (iii). \square

Before computing the invariants of Y_7 , we first analyze the mirror map and the function f_7 which will be defined below.

It follows from (3.13) that the mirror transform is $\tau = t - g_7 E$, where

$$g_7 = g_7(P^{\ell_1}, P^{\ell_2}) := \sum_{\substack{m_1 \geq m_2 \\ m_1 + m_2 > 0}} \frac{(-1)^{m_1+m_2}}{m_1 + m_2} \frac{(m_1 + m_2)!^2}{m_1! (m_1 - m_2)! m_2!^3} P^{m_1 \ell_1 + m_2 \ell_2}. \quad (3.16)$$

This gives the following identities:

$$Q^\gamma = P^\gamma e^{-g_7}, \quad Q^{\ell_1} = P^{\ell_1} e^{g_7}, \quad Q^{\ell_2} = P^{\ell_2} e^{g_7}.$$

As before, define $\theta = -\frac{1}{z}\hat{E} = P^{\ell_1}\frac{d}{dP^{\ell_1}} + P^{\ell_2}\frac{d}{dP^{\ell_2}}$ and

$$f_7(P^{\ell_1}, P^{\ell_2}) = 1 + \theta g_7(P^{\ell_1}, P^{\ell_2}) = \sum_{m_1 \geq m_2} \frac{(-1)^{m_1+m_2}(m_1+m_2)!^2}{m_1!(m_1-m_2)!m_2!^3} P^{m_1\ell_1+m_2\ell_2}.$$

Let $\theta_1 = P^{\ell_1}\frac{d}{dP^{\ell_1}}$, $\theta_2 = P^{\ell_2}\frac{d}{dP^{\ell_2}}$. Then f_7 satisfies the following differential equations:

$$\begin{aligned} (\theta_1(\theta_1 - \theta_2) + P^{\ell_1}(\theta_1 + \theta_2 + 1)^2)f &= 0, \\ (\theta_2^3 + P^{\ell_2}(\theta_1 + \theta_2 + 1)^2(\theta_1 - \theta_2))f &= 0. \end{aligned} \quad (3.17)$$

Let $y = P^{-\ell_1}$ and let $r = P^{\ell_2-\ell_1} = Q^{\ell_2-\ell_1}$. Then $\theta_1 = -\theta_y - \theta_r$, $\theta_2 = \theta_r$, where we denote $\theta_y := y\frac{d}{dy}$, $\theta_r := r\frac{d}{dr}$. With this change of variables, the differential equation (3.17) becomes

$$((\theta_y - 1)^2 + y(\theta_y + \theta_r)(\theta_y + 2\theta_r))f_7 = 0, \quad (3.18)$$

$$((3 + 3y - 4r)\theta_y + (7y - 2 - 8r)\theta_r - 3)f_7 = 0. \quad (3.19)$$

Plugging (3.19) into (3.18), we obtain a second-order differential equation of the form

$$(A_2(y, r)\theta_y^2 + A_1(y, r)\theta_y + A_0(y, r))f_7 = 0, \quad (3.20)$$

where the coefficients are given by

$$A_2(y, r) = -2(1 + 4r)^3 + y + 72yr + 272yr^2 + 15y^2 - 162y^2r + 8y^2r^2 + 19y^3 - 15y^3r + 7y^4,$$

$$A_1(y, r) = 4(1 + 4r)^3 - 15y - 184yr - 496yr^2 - 42y^2 + 280y^2r - 23y^3 - y^3r,$$

$$A_0(y, r) = -2(1 + 4r)^3 + 12y + 108yr + 240yr^2 + 21y^2 - 126y^2r + 7y^3.$$

This shows that, for $r \neq -\frac{1}{4}, \infty$, the function $f_7|_r$ has local exponents $(1, 1)$ at $y = 0$. Therefore, we may assume that

$$f_7 = \sum_{k \geq 1} (a_k(r) \log y + b_k(r))y^k, \quad r \neq -\frac{1}{4}, \infty.$$

Substituting this into (3.19), we see that

$$((-4r)a_1(r) + (-2 - 8r)\theta_r a_1(r))y \log y + O(y) = 0,$$

or equivalently,

$$a_1(r) = \exp\left(-\int \frac{2}{1+4r} dr\right) = \frac{c_0}{(1+4r)^{1/2}}$$

for some constant c_0 . To determine c_0 , observe that when $r = 0$,

$$f_7|_{r=0} = \sum_{m \geq 0} (-1)^m P^{m\ell_1} = \frac{1}{1+P^{\ell_1}} = \frac{y}{1+y}.$$

This yields $c_0 = a_1(0) = 0$, and hence $a_1(r) = 0$ for all r . Since the solutions to (3.20) have local exponents $(1, 1)$, it follows that $a_k(r) \equiv 0$ for all k . Moreover, (3.19) together with the boundary condition again tell us that

$$f_7 = \sum_{k \geq 1} b_k(r)y^k = \frac{1}{(1+4r)^{1/2}}y - \frac{1-2r}{(1+4r)^{5/2}}y^2 + O(y^3), \quad (3.21)$$

where $\{b_k(r)\}$ satisfies the recurrence relation

$$k^2 b_{k+1}(r) + (k^2 + 3k\theta_r b_k(r) + 2\theta_r^2 b_k) = 0, \quad k \geq 1. \quad (3.22)$$

Since $f_7 = 1 - \theta_y g_7$, we have

$$g_7 = \log y - \sum_{k \geq 1} \frac{b_k(r)}{k} y^k - c(r)$$

for some function $c(r)$. Substituting $f_7 = 1 - \theta_y g_7$ into (3.18), we get

$$\begin{aligned} 0 &= (\theta_y^2 y^{-1} + (\theta_y + \theta_r)(\theta_y + 2\theta_r))(1 - \theta_y g_7) \\ &= -\theta_y(y^{-1} + (\theta_y y^{-1} \theta_y + (\theta_y + \theta_r)(\theta_y + 2\theta_r))g_7). \end{aligned}$$

A direct computation using (3.16) then gives

$$y^{-1} + (\delta_y y^{-1} \delta_y + (\delta_y + \delta_r)(\delta_y + 2\delta_r))g_7 = 0.$$

This implies that $\delta_r^2 c(r) = 0$, i.e., $c(r) = c_1 \log r + c_2$ for some c_1, c_2 . To determine these constants, note that

$$g_7|_{r=0} = \sum_{m \geq 1} \frac{(-1)^m}{m} P^{m\ell_1} = -\log(1 + P^{\ell_1}) = \log\left(\frac{y}{1+y}\right),$$

which shows that $c(r) = 0$ for all r . Hence,

$$g_7 = \log y - \sum_{k \geq 1} \frac{b_k(r)}{k} y^k, \quad (3.23)$$

$$g_{7,1} := \theta_{P^{\ell_1}} g_7 = -1 + \sum_{k \geq 1} \left(b_k(r) + \frac{\theta_r b_k(r)}{k} \right) y^k, \quad (3.24)$$

$$g_{7,2} := \theta_{P^{\ell_2}} g_7 = - \sum_{k \geq 1} \frac{\theta_r b_k(r)}{k} y^k. \quad (3.25)$$

In particular, both $\frac{g_{7,1}+1}{f}$ and $\frac{g_{7,2}}{f}$ are bounded as $y \rightarrow 0$. This suggests us to consider the divisors $\tilde{F}_1 = F_1 + E$, $\tilde{F}_2 = F_2$, for which

$$\delta^{\tilde{F}_1} = \theta_{\tilde{\gamma}_1} - \frac{g_{7,2}}{f} \theta_y - \theta_r, \quad \delta^{\tilde{F}_2} = \frac{g_{7,2}}{f} \theta_y + \theta_r \quad (3.26)$$

has coefficients that remain bounded near $y = 0$, where $\theta_{\tilde{\gamma}_1} = P^{\tilde{\gamma}_1} \frac{d}{dP^{\tilde{\gamma}_1}} = Q^{\tilde{\gamma}_1} \frac{d}{dQ^{\tilde{\gamma}_1}}$.

Proposition 3.21. *Let $\tilde{\gamma}_1 = \ell_1 + \gamma$ be the lifting of $\bar{\gamma}_1$ such that $(E, \tilde{\gamma}_1) = 0$. Take $F_0 = \frac{1}{2}F_1 - F_2$ so that $(EF_0, \tilde{F}_1) = (EF_0, \tilde{F}_2) = 0$ and $(EF_0, E) = 1$.*

When $r \neq -\frac{1}{4}, \frac{1}{2}, \infty$, we have:

(i) $\langle \text{pt} \rangle^{Y_7} = (1+y)Q^{\tilde{\gamma}_1}$;

(ii) for divisor $D_1, D_2 \in H^2(Y_7)$,

$$\langle HD_1, HD_2 \rangle^{Y_7} = ((HD_1 D_2, 1)(1+y) + 2(D_1, \tilde{\gamma}_1)(D_2, \tilde{\gamma}_1) + (D_1, \tilde{\gamma}_2)(D_2, \tilde{\gamma}_2)r)Q^{\tilde{\gamma}_1};$$

(iii) $\langle \text{pt}, \text{pt} \rangle^{Y_7} = rQ^{2\tilde{\gamma}_1} + O(y)$;

(iv) for divisor $D \in H^2(Y_7)$, $\langle HD, EF_0 \rangle^{Y_7} \in O(y)$;

(v) near $y = 0$,

$$\begin{aligned}\langle E, E, F_0 \rangle^{Y_7} &= -(1+4r)^{1/2}y^{-1} + O(1) \\ \langle E, F_0, F_2 \rangle^{Y_7} &= -\frac{3}{2} \cdot \frac{1-(1+4r)^{1/2}}{1+4r} + O(y)\end{aligned}$$

and all these invariants lies in the ring $\mathbb{C}[Q^{\tilde{\gamma}_1}, r, (1+4r)^{-1/2}, (1-2r)^{-1}](y)$.

Proof. As before,

$$\langle \psi^k T \rangle^{Y_7} = [z^{-(k+2)}] \int_T e^{gE} \sum_{\beta} I_{\beta}^{Y_7} P^{\beta}.$$

So, for (i),

$$\langle \text{pt} \rangle^{Y_7} = [z^{-2}] \int_{\text{pt}} e^{gE} \sum_{\beta} I_{\beta}^{Y_7} Q^{\beta} = [z^{-2}] \int_{\text{pt}} I_{\gamma}^{Y_7} P^{\gamma} + I_{\ell_1+\gamma}^{Y_7} P^{\ell_1+\gamma} = P^{\gamma} + P^{\ell_1+\gamma} = (1+y)Q^{\tilde{\gamma}_1}.$$

For further use, we also compute that

$$\begin{aligned}\langle \psi^2 \text{pt} \rangle^{Y_7} &= [z^{-4}] \int_{\text{pt}} e^{gE} \sum_{\beta} I_{\beta}^{Y_7} Q^{\beta} = \frac{1}{4}P^{2\gamma} + \frac{1}{2}P^{2\gamma+\ell_1} + Q^{\ell_1+\ell_2+2\gamma} + \frac{1}{4}Q^{2\ell_1+2\gamma} \\ &= (\frac{1}{4} + \frac{1}{2}y + \frac{1}{4}y^2 + r)Q^{2\tilde{\gamma}_1},\end{aligned}\tag{3.27}$$

and for $D \in H^2(X_7)$,

$$\begin{aligned}\langle \psi HD \rangle^{Y_7} &= [z^{-3}] \int_{HD} e^{gE} \sum_{\beta} I_{\beta}^{Y_7} Q^{\beta} \\ &= (HD, -2F_1 + F_2)P^{\ell_1+\gamma} + (HD, F_1 - F_2)P^{\ell_2+\gamma} - (HD, H)P^{\gamma} \\ &= ((HD, -2F_1 + F_2) - (HD, H)y + (HD, F_1 - F_2)r)Q^{\tilde{\gamma}_1}.\end{aligned}\tag{3.28}$$

For (ii), by reconstruction formula,

$$\langle HD_1, HD_2 \rangle^{Y_7} = \langle H, HD_1 D_2 \rangle^{Y_7} + \delta^{D_1} \langle H, \psi h D_2 \rangle^{Y_7} - \sum \delta^{D_1} \langle H, T_{\mu} \rangle^{Y_7} \langle T^{\mu}, HD_2 \rangle^{Y_7}.$$

It is easy to check that the sum is 0 by degree argument. So, by divisor equation, (i) and (3.28),

$$\begin{aligned}\langle HD_1, HD_2 \rangle^{Y_7} &= \langle HD_1 D_2 \rangle^{Y_7} + \delta^{D_1} \langle H^2 D_2 \rangle^{Y_7} + \delta^{D_1} \delta^H \langle \psi HD_2 \rangle^{Y_7} \\ &= (HD_1 D_2, 1) \langle \text{pt} \rangle^{Y_7} + (D_2, H^2) \delta^{D_1} \langle \text{pt} \rangle^{Y_7} \\ &\quad + \delta^{D_1} (((HD_2, -2F_1 + F_2) - (HD_2, H)y + (HD_2, F_1 - F_2)r)Q^{\tilde{\gamma}_1}) \\ &= (HD_1 D_2, 1)(1+y)Q^{\tilde{\gamma}_1} + (D_1, H^2 + H(-2F_1 + F_2))Q^{\tilde{\gamma}_1} \\ &\quad + (D_2, H(F_1 - F_2))(D_1, \tilde{\gamma}_2)Q^{\tilde{\gamma}_2} \\ &= ((HD_1 D_2, 1)(1+y) + 2(D_1, \tilde{\gamma}_1)(D_2, \tilde{\gamma}_1) + (D_1, \tilde{\gamma}_2)(D_2, \tilde{\gamma}_2)r)Q^{\tilde{\gamma}_1}.\end{aligned}$$

Write $\langle \psi EF_0 \rangle^{Y_7} = G(y, r)Q^{\tilde{\gamma}_1}$. Then for any divisor D ,

$$\begin{aligned}\langle HD, EF_0 \rangle^{Y_7} &= \delta^D \langle \psi EF_0 \rangle^{Y_7} + \langle H, EF_0 D \rangle^{Y_7} - \sum \delta^D \langle H, T_\mu \rangle^{Y_7} \langle T^\mu EF_0 \rangle^{Y_7} \\ &= \delta^D (GQ^{\tilde{\gamma}_1}) + (D, EF_0)(1+y)Q^{\tilde{\gamma}_1}.\end{aligned}$$

In particular, take $D = E$, we have

$$0 = \delta^E (GQ^{\tilde{\gamma}_1}) + (E, EF_0)(1+y)Q^{\tilde{\gamma}_1},$$

or, equivalently, $\theta_y G = f_7(1+y)$. By the expansion (3.21) of f_7 , we get

$$G = \sum_{k \geq 0} G_k(r)y^k, \quad G_k(r) = -\frac{b_k(r) + b_{k-1}(r)}{k} \quad \text{for } k \geq 1,$$

where $b_0(r)$ is defined to be 0.

Claim. We have $G|_{r=0} = -y$, and for each $n \geq 1$,

$$\left. \frac{\partial^n G}{\partial r^n} \right|_{r=0} = \frac{yV_n(y)}{(1+y)^{3n-1}}$$

for some polynomial $V_n(y)$ of at most degree n .

Proof of Claim. It follows from

$$GQ^{\tilde{\gamma}_1} = \langle \psi EF_0 \rangle^{Y_7} = [z^{-3}] \int_{EF_0} \sum_{\beta} I_\beta^{Y_7} P^\beta = -P^\gamma - Q^{\ell_2+\gamma} + \dots$$

and the identity

$$\theta_y G = (1+y)f = \sum_{m_1 \geq m_2} \frac{(-1)^{m_1+m_2} (m_1+m_2)!^2}{m_1! (m_1-m_2)! m_2!^3} (P^{m_1\ell_1+m_2\ell_2} + P^{(m_1-1)\ell_1+m_2\ell_2})$$

that

$$\begin{aligned}G &= -P^{-\ell_1} - P^{-\ell_1+\ell_2} + \sum_{\substack{m_1 \geq m_2 \\ m_1+m_2 \geq 2}} \frac{(-1)^{m_1+m_2} (m_1+m_2)!^2}{m_1! (m_1-m_2)! m_2!^3} \left(\frac{P^{m_1\ell_1+m_2\ell_2}}{m_1+m_2} + \frac{P^{(m_1-1)\ell_1+m_2\ell_2}}{m_1+m_2-1} \right) \\ &= -y - r + \sum_{\substack{m_1 \geq m_2 \\ m_1+m_2 \geq 2}} \frac{(-1)^{m_1+m_2} (m_1+m_2)!^2}{m_1! (m_1-m_2)! m_2!^3} \left(\frac{P^{(m_1+m_2)\ell_1}}{m_1+m_2} + \frac{P^{(m_1+m_2-1)\ell_1}}{m_1+m_2-1} \right) r^{m_2} \\ &= -y - \left(\sum_{m \geq 0} (-1)^m (3m+1) P^{m\ell_1} \right) r \\ &\quad + \sum_{m_2 \geq 2} \frac{(-r)^{m_2}}{m_2!^3} \left(\sum_{m_1 \geq m_2} \frac{(-1)^{m_1} (m_1+m_2)!^2}{m_1! (m_1-m_2)! (m_1+m_2)} P^{(m_1+m_2)\ell_1} \right. \\ &\quad \left. - \sum_{m_1 \geq m_2-1} \frac{(-1)^{m_1} (m_1+m_2+1)!^2}{(m_1+1)! (m_1+1-m_2)! (m_1+m_2)} P^{(m_1+m_2)\ell_1} \right).\end{aligned}$$

In particular,

$$G|_{r=0} = -y,$$

$$\frac{\partial G}{\partial r}\Big|_{r=0} = - \sum_{m \geq 0} (-1)^m (3m+1) P^{m\ell_1} = - \frac{1-3P^{\ell_1}}{(1+P^{\ell_1})^2} = \frac{y(3-y)}{(1+y)^2}.$$

For $n \geq 2$, we compute

$$\begin{aligned} \frac{\partial^n G}{\partial r^n}\Big|_{r=0} &= \frac{(-1)^n}{n!^2} \sum_{m \geq n-1} \left(\frac{(-1)^m (m+n)!(m+n-1)!}{(m+1)!(m+1-n)!} (-n)(n+3m+3) \right) P^{(m+n)\ell_1} \\ &= \frac{n P^{(2n-1)\ell_1}}{n!^2} \sum_{m \geq 0} \frac{(m+2n-1)!(m+2n-2)!}{(m+n)!m!} (4n+3m)(-P^{\ell_1})^m \\ &= \frac{n P^{(2n-1)\ell_1}}{n!^2} (4n+3\theta_{P^{\ell_1}})_2F_1(2n-1, 2n; n+1; -P^{\ell_1}). \end{aligned}$$

By the linear transformation formula for hypergeometric function,

$$\begin{aligned} {}_2F_1(2n-1, 2n; n+1; -P^{\ell_1}) &= (1+P^{\ell_1})^{(n+1)-(2n-1)-(2n)} {}_2F_1(2-n, 1-n; n+1; -P^{\ell_1}) \\ &= \frac{U_n(P^{\ell_1})}{(1+P^{\ell_1})^{3n-2}}, \end{aligned}$$

where $U_n(P^{\ell_1}) = {}_2F_1(2-n, 1-n; n+1; -P^{\ell_1})$ is a polynomial of degree $n-2$. Substituting $y = P^{-\ell_1}$, we obtain

$$\frac{\partial^n G}{\partial r^n}\Big|_{r=0} = \frac{ny^{-(2n-1)}}{n!^2} (4n-3\theta_y) \left(\frac{y^{3n-2} U_n(y^{-1})}{(1+y)^{3n-2}} \right) = \frac{y V_n(y)}{(1+y)^{3n-1}}$$

for some polynomial V_n of degree at most n , completing the proof. ■

It follows from the claim that $G_0(r) = 0$ for all r , and hence $G = -b_1(r)y + O(y^2)$. Thus, near $y = 0$, by (3.26) and the boundedness of $\frac{g_{7.2}}{f}$,

$$\begin{aligned} \langle H\tilde{F}_1, EF_0 \rangle^{Y_7} &= \delta^{\tilde{F}_1}(GQ^{\tilde{\gamma}_1}) + (\tilde{F}_1, EF_0)(1+y)Q^{\tilde{\gamma}_1} = \delta^{\tilde{F}_1}(GQ^{\tilde{\gamma}_1}) = O(y), \\ \langle H\tilde{F}_2, EF_0 \rangle^{Y_7} &= \delta^{\tilde{F}_2}(GQ^{\tilde{\gamma}_1}) + (\tilde{F}_2, EF_0)(1+y)Q^{\tilde{\gamma}_1} = \delta^{\tilde{F}_2}(GQ^{\tilde{\gamma}_1}) = O(y), \\ \langle HE, EF_0 \rangle^{Y_7} &= 0. \end{aligned}$$

Since D is a linear combination of $\tilde{F}_1, \tilde{F}_2, E$, we have $\langle HD, EF_0 \rangle^{Y_7} = O(y)$.

On the other hand, by reconstruction formula,

$$\langle HD, EF_0 \rangle^{Y_7} = \delta^E(\delta^{F_0} \langle \psi HD \rangle^{Y_7} + \langle F_0 HD \rangle^{Y_7}) - \delta^E \langle F_0, T_\mu \rangle^{Y_7} \langle T^\mu, HD \rangle^{Y_7}.$$

This shows that

$$\begin{aligned} \langle E, E, F_0 \rangle^{Y_7} \langle HD, EF_0 \rangle^{Y_7} + \langle E, F_0, F_2 \rangle^{Y_7} \langle H(\frac{3}{2}F_1 - 2F_2), HD \rangle^{Y_7} \\ = -(H^2 D) \delta^{F_0} \frac{y Q^{\tilde{\gamma}_1}}{f} + (F_0 HD) \frac{y Q^{\tilde{\gamma}_1}}{f}. \end{aligned}$$

So we can write $\langle E, E, F_0 \rangle^{Y_7}$ and $\langle E, F_0, F_2 \rangle^{Y_7}$ in terms of G by taking $D = \tilde{F}_1, \tilde{F}_2$:

$$\begin{aligned}\langle E, E, F_0 \rangle^{Y_7} &= \frac{(1+y-2r)Q^{\tilde{\gamma}_1}(-4\delta^{F_0}\frac{yQ^{\tilde{\gamma}_1}}{f_7} + \frac{yQ^{\tilde{\gamma}_1}}{f_7}) - (2-y)Q^{\tilde{\gamma}_1}(-3\delta^{F_0}\frac{yQ^{\tilde{\gamma}_1}}{f_7})}{(1+y-2r)\delta^{\tilde{F}_1}(GQ^{\tilde{\gamma}_1}) - (2-y)Q^{\tilde{\gamma}_1}\delta^{F_2}(GQ^{\tilde{\gamma}_1})}, \\ \langle E, F_0, F_2 \rangle^{Y_7} &= \frac{\delta^{\tilde{F}_1}(GQ^{\tilde{\gamma}_1})(-3\delta^{F_0}\frac{yQ^{\tilde{\gamma}_1}}{f_7}) - \delta^{F_2}(GQ^{\tilde{\gamma}_1})(-4\delta^{F_0}\frac{yQ^{\tilde{\gamma}_1}}{f_7} + \frac{yQ^{\tilde{\gamma}_1}}{f_7})}{(1+y-2r)\delta^{\tilde{F}_1}(GQ^{\tilde{\gamma}_1}) - (2-y)Q^{\tilde{\gamma}_1}\delta^{F_2}(GQ^{\tilde{\gamma}_1})}.\end{aligned}$$

Since

$$\begin{aligned}\delta^{\tilde{F}_1}(GQ^{\tilde{\gamma}_1}) &= (\theta_{\tilde{\gamma}_1} - \frac{g_{7,2}}{f_7}\theta_y - \theta_r)((-\frac{y}{(1+4r)^{1/2}} + O(y^2))Q^{\tilde{\gamma}_1}) = (-\frac{y}{(1+4r)^{1/2}} + O(y^2))Q^{\tilde{\gamma}_1}, \\ \delta^{F_2}(GQ^{\tilde{\gamma}_1}) &= (\frac{g_{7,2}}{f_7}\theta_y + \theta_r)((-\frac{y}{(1+4r)^{1/2}} + O(y^2))Q^{\tilde{\gamma}_1}) = O(y^2)Q^{\tilde{\gamma}_1}, \\ \frac{yQ^{\tilde{\gamma}_1}}{f_7} &= ((1+4r)^{1/2} + O(y))Q^{\tilde{\gamma}_1}, \\ \delta^{F_0}\frac{yQ^{\tilde{\gamma}_1}}{f_7} &= (\frac{1-2r}{2}(\frac{1-(1+4r)^{1/2}}{1+4r}) + O(y))Q^{\tilde{\gamma}_1},\end{aligned}$$

we have the following asymptotic expansion when $r \neq \frac{1}{2}$:

$$\begin{aligned}\langle E, E, F_0 \rangle^{Y_7} &= -(1+4r)^{1/2}y^{-1} + O(1), \\ \langle E, F_0, F_2 \rangle^{Y_7} &= -\frac{3}{2} \cdot \frac{1-(1+4r)^{1/2}}{1+4r} + O(y).\end{aligned}$$

Also, it follows from (3.22) and the initial condition $b_1(r) = (1+4r)^{-1/2}$, that the coefficients of f_7, f_7^{-1} , and hence $g_{7,2}$, all lies in $\mathbb{C}[r, (1+4r)^{-1/2}]$. This means that the invariants

$$\langle H\tilde{F}_1, EF_0 \rangle^{Y_7}, \langle H\tilde{F}_2, EF_0 \rangle^{Y_7}, \langle E, E, F_0 \rangle^{Y_7}, \langle E, F_0, F_2 \rangle^{Y_7}$$

all lies in the ring $\mathbb{C}[Q^{\tilde{\gamma}_1}, r, (1+4r)^{-1/2}, (1-2r)^{-1}](\!(y)\!)$.

Denote $\tilde{F}^1 = -\frac{1}{2}F_1 + F_2, \tilde{F}^2 = F_1 - F_2$ so that $(H\tilde{F}^i, F_j) = \delta_j^i$.

Finally, by reconstruction formula, (i), (ii), (iv) and (3.27),

$$\begin{aligned}7\langle pt, pt \rangle^{Y_7} &= \langle H^3, pt \rangle^{Y_7} = \delta^H \langle H^2, \psi pt \rangle^{Y_7} - \delta^H \langle H^2, T_\mu \rangle^{Y_7} \langle T^\mu, pt \rangle^{Y_7} \\ &= 2(\delta^H \langle H, \psi^2 pt \rangle^{Y_7} - \delta^H \langle H, T_\mu \rangle^{Y_7} \langle T^\mu, \psi pt \rangle^{Y_7}) \\ &\quad - \langle H^2, H\tilde{F}^1 \rangle^{Y_7} \langle \tilde{F}_1, pt \rangle^{Y_7} - \langle H^2, H\tilde{F}^2 \rangle^{Y_7} \langle \tilde{F}_2, pt \rangle^{Y_7} - \langle H^2, EF_0 \rangle^{Y_7} \langle E, pt \rangle^{Y_7} \\ &= 8\langle \psi^2 pt \rangle^{Y_7} - 2\langle pt \rangle \langle 1, \psi pt \rangle^{Y_7} - \langle H^2, H\tilde{F}^1 \rangle^{Y_7} \delta^{\tilde{F}_1} \langle pt \rangle^{Y_7} - \langle H^2, EF_0 \rangle^{Y_7} \frac{y}{f_7} Q^{\tilde{\gamma}_1} \\ &= 8(\frac{1}{4} + r)Q^{2\tilde{\gamma}_1} - 2Q^{2\tilde{\gamma}_1} - (H^2\tilde{F}^1, 1)Q^{2\tilde{\gamma}_1} \\ &\quad - 2(H, \tilde{\gamma}_1)(\tilde{F}^1, \tilde{\gamma}_1)Q^{2\tilde{\gamma}_1} - (H, \tilde{\gamma}_2)(\tilde{F}^2, \tilde{\gamma}_2)rQ^{2\tilde{\gamma}_1} + O(y) \\ &= 7rQ^{2\tilde{\gamma}_1} + O(y),\end{aligned}$$

which is (iii), and lies in the ring $\mathbb{C}[Q^{\tilde{\gamma}_1}, r, (1+4r)^{-1/2}, (1-2r)^{-1}](\!(y)\!)$. \square

Remark 3.22. The generating function $\langle E, E, E \rangle$ is still of the form $\frac{7}{u_7 f_7^3}$. In fact, one can show that

$$u_7(P^{\ell_1}, P^{\ell_2}) = \frac{(1+P^{\ell_1})^3 - P^{\ell_2} - 20P^{\ell_1}P^{\ell_2} + 8P^{2\ell_1}P^{\ell_2} + 16P^{\ell_1}P^{2\ell_2}}{1 - \frac{2}{7}P^{\ell_1} - \frac{8}{7}P^{\ell_2}},$$

which is equal to

$$\frac{1 + 2P^\ell - 17P^{2\ell} + 25P^{3\ell}}{1 - \frac{10}{7}P^\ell}$$

when restricting on the diagonal $P^\ell = P^{\ell_1} = P^{\ell_2}$.

Definition 3.23. We define the ring

$$R_7 = \mathbb{C}[Q^{\tilde{\gamma}_1}, r, (1+4r)^{-1/2}, (1-2r)^{-1}][[y]]$$

and its ideal $I_7 = y \cdot R_7$.

It is easy to see that the operators $\theta_{\tilde{\gamma}_1}$ and θ_r sends elements in R_7 to R_7 , and θ_y send elements in R_7 to I_7 .

Proposition 3.24. Consider the ring $R_{\text{loc}} = R_d[Q^{\tilde{\gamma}_1}]$ and its ideal $I_{\text{loc}} = I_d[Q^{\tilde{\gamma}_1}]$. Then for $n \geq 1$, $\vec{a} \in (H(S_7)^{G_{Y_7}})^{\oplus n}$, the invariant $\langle j_{7*}\vec{a} \rangle^{Y_7}$ lies in R_{loc} and

$$\langle j_{7*}\vec{a} \rangle^{Y_7} \equiv \phi_7^* \langle i_{7*}\vec{a} \rangle^{X_7} \pmod{I_{\text{loc}}}.$$

Proof. Since $H(S_7)^{G_{Y_7}}$ is generated by 1, $i_7^*h_1$, $i_7^*h_2$, $i_7^*h^2$, $i_{7*}H(S_7)^{G_{Y_7}}$ is generated by h , hh_1 , hh_2 , h^3 and $j_{7*}H(S_7)^{G_{Y_7}}$ is generated by H , $H\tilde{F}_1$, $H\tilde{F}_2$, H^3 .

Let $\vec{a} = (a_1, \dots, a_n)$, $b_\nu = i_{7*}a_\nu$, $c_\nu = j_{7*}a_\nu$. By divisor equation, we may assume that each $a_\nu = i_7^*h_1$, $i_7^*h_2$ or h^2 . We prove the result by induction on n and the curve degree $\frac{1}{2} \sum \deg a_\nu$.

For $n \leq 2$, the result follows from the fundamental class axiom, the divisor equation, Proposition 3.20 and Proposition 3.21.

As in the proof of Proposition 3.19, by divisor equation and reconstruction formula, it suffices to show that:

- (i) $\langle \tilde{F}_i, Hc_2, c_3, \dots, c_n \rangle^Y \equiv \phi^* \langle h_i, hb_2, b_3, \dots, b_n \rangle^X \pmod{I_{\text{loc}}}$;
- (ii) $\langle \vec{c}_I, T \rangle^Y \equiv \phi^* \langle b_I, \bar{T} \rangle^X \pmod{I_{\text{loc}}}$ for $(\bar{T}, T) = (1, 1)$, (h_1, \tilde{F}_1) , (h_2, \tilde{F}_2) , $(hh_1, H\tilde{F}_1)$, $(hh_2, H\tilde{F}_2)$, (pt, pt) with lower curve degree;
- (iii) $\langle \vec{c}_{I_L}, E \rangle^Y \langle EF_0, \vec{c}_{I_R} \rangle^Y \in I_{\text{loc}}$.

We note that if $B \in R_{\text{loc}}$ and $A \in \mathbb{C}[Q^{\tilde{\gamma}_1}, Q^{\tilde{\gamma}_2}]$ satisfies $B \equiv \phi^*A \pmod{I_{\text{loc}}}$, then

$$\begin{aligned} \delta^{\tilde{F}_1} B &= \theta_{\tilde{\gamma}_1} B - \frac{g_{7,2}}{f_7} \theta_y B - \theta_r B \equiv \phi^*((\theta_{\tilde{\gamma}_1} + \theta_{\tilde{\gamma}_2})A - \theta_{\tilde{\gamma}_2}A) = \phi^*(\delta^{h_1} A), \\ \delta^{\tilde{F}_1} B &= \frac{g_{7,2}}{f_7} \theta_y B + \theta_r B \equiv \phi^*(\theta_{\tilde{\gamma}_2}A) = \phi^*(\delta^{h_2} A). \end{aligned}$$

This, together with induction hypothesis, the fundamental class axiom and the divisor equation, gives (i) and (ii).

For (iii), we see that

$$\langle \vec{c}, E \rangle^Y = \delta^E \langle \vec{c} \rangle^Y = \delta^E \langle \vec{c} \rangle^Y \in f^{-1} \theta_y R_{\text{loc}} \subseteq R_{\text{loc}}.$$

If $m = 1$, then $\langle EF_0, \vec{c} \rangle^Y \in I_{\text{loc}}$ by Proposition 3.21 (iv). If $m \geq 2$, apply the reconstruction formula, we get the equation

$$\langle EF_0, \vec{c} \rangle^Y = \sum_{1 \in I_L, 2 \in I_R} \langle F_0, \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \delta^E \langle T^\mu, \vec{c}_{I_R} \rangle^Y - \sum_{1, 2 \in I_R} \delta^E \langle F_0, \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \langle T^\mu, \vec{c}_{I_R} \rangle^Y.$$

When $I_L = \emptyset$, it follows from

$$\langle E, F_0, \tilde{F}_1 \rangle^Y + \langle E, F_0, \tilde{F}_2 \rangle^Y = \langle E, F_0, H \rangle^Y = 0$$

that

$$\begin{aligned} & \delta^E \langle F_0, \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \langle T^\mu, \vec{c}_{I_R} \rangle^Y \\ &= \delta^E \delta^{F_0} \langle E \rangle^Y \cdot \langle EF_0, \vec{c} \rangle^Y + \delta^E \delta^{F_0} \langle \tilde{F}_1 \rangle^Y \cdot \langle H\tilde{F}^1, \vec{c} \rangle^Y + \delta^E \delta^{F_0} \langle \tilde{F}_2 \rangle^Y \cdot \langle H\tilde{F}^2, \vec{c} \rangle^Y \\ &= (\langle E, E, F_0 \rangle^Y - 1) \langle EF_0, \vec{c} \rangle^Y + \langle E, F_0, F_2 \rangle (\langle H\tilde{F}^2, \vec{c} \rangle^Y - \langle H\tilde{F}^1, \vec{c} \rangle^Y). \end{aligned}$$

So we get

$$\begin{aligned} \langle E, E, F_0 \rangle^Y \cdot \langle E^2, \vec{c} \rangle^Y &= \sum_{1 \in I_L, 2 \in I_R} \delta^E \langle \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \delta^E \langle T^\mu, \vec{c}_{I_R} \rangle^Y \\ &\quad - \sum_{I_L \neq \emptyset, 1, 2 \in I_R} (\delta^E)^2 \langle \vec{c}_{I_L}, T_\mu \rangle^Y \cdot \langle T^\mu, \vec{c}_{I_R} \rangle^Y \\ &\quad + \langle E, F_0, F_2 \rangle (\langle H\tilde{F}^2, \vec{c} \rangle^Y - \langle H\tilde{F}^1, \vec{c} \rangle^Y). \end{aligned} \tag{3.29}$$

Since $\delta^E = f^{-1}\theta_y$ sends elements in R_{loc} to R_{loc} , this, together with Proposition 3.21 (v) implies that

$$\langle E^2, \vec{c} \rangle^Y \in \frac{1}{\langle E, E, F_0 \rangle^Y} R_{loc} + \frac{\langle E, F_0, F_2 \rangle^Y}{\langle E, E, F_0 \rangle^Y} R_{loc} \subseteq I_{loc}.$$

We conclude that $\langle \vec{c}_{I_L}, E \rangle^Y \cdot \langle EF_0, \vec{c}_{I_R} \rangle^Y \in R_{loc} \cdot I_{loc} \subseteq I_{loc}$. \square

4. PROOF OF MAIN RESULTS IN THE GLOBAL CASES

4.1. Two (symplectic) semistable degenerations. To apply the degeneration analysis in the next subsection, we shall factor a Type II transition $Y \searrow X$ as a composition of two semistable (symplectic) degenerations $\mathcal{X} \rightarrow \Delta$ and $\mathcal{Y} \rightarrow \Delta$. As before, let $\phi: Y \rightarrow \bar{Y}$ be the crepant contraction of $Y \searrow X$, $E = \text{Ex}(\phi)$ and $\phi(E) = \{p\}$.

As a simple demonstration, we first assume that $Y \searrow X$ is a del Pezzo transition of degree d , that is, E is a smooth del Pezzo surface of degree d . In this case \mathcal{X} and \mathcal{Y} are semistable *projective* degenerations over the disk Δ .

The *Kähler degeneration*

$$\mathcal{Y} \rightarrow \Delta$$

is the deformation to the normal cone $\mathcal{Y} = \text{Bl}_{E \times \{0\}}(Y \times \Delta)$. Since E has codimension one in Y and $E|_E = K_E$, the special fiber $\mathcal{Y}_0 = Y \cup Y_d$ is a simple normal crossing divisor with $Y_d = \mathbb{P}_E(K_E \oplus \mathcal{O})$. The intersection $E = Y \cap Y_d$ is understood as the infinity divisor (or relative hyperplane section) of $Y_d \rightarrow E$.

The *complex degeneration*

$$\mathcal{X} \rightarrow \Delta$$

is the semistable reduction of the smoothing $\mathfrak{X} \rightarrow \Delta$ constructed in Proposition 1.10. Recall that, in the notation of the proof of Proposition 1.10, $\mathcal{X} \rightarrow \Delta$ is obtained by a degree n_d base change $\mathcal{X}' \rightarrow \Delta$ allowed by the weighted blow-up at $p \in \mathcal{X}'$ with weight $(\alpha, 1)$. The special fiber $\mathcal{X}_0 = Y \cup X_d$ is a simple normal crossing divisor with X_d being a smooth del Pezzo threefold of degree d . The intersection $E = Y \cap X_d$ is now understood as a general member of the linear system $| -K_{X_d}|$. Notice that the local model $Y_d \searrow X_d$ (Example 1.9) appears in the special fibers \mathcal{X}_0 and \mathcal{Y}_0 .

In general, let $Y \searrow X$ be a Type II extremal transition of degree d . If we work in the category of symplectic manifolds, we can still construct similar semistable degenerations to relate the local model $Y_d \searrow X_d$ as above.

To begin with, the smoothing $\mathfrak{X} \rightarrow \Delta$ with $\mathfrak{X}_t = X$ ($t \neq 0$) induces a holomorphic map

$$\sigma: \Delta \rightarrow \text{Def}(\bar{Y}) \rightarrow \text{Def}(\bar{Y}, p). \quad (4.1)$$

We may take two good representatives $\bar{\mathcal{U}} \rightarrow \Delta$ and $\tilde{\mathcal{U}} \rightarrow \Delta$ of the deformation corresponding to σ such that $\bar{\mathcal{U}}_t \supseteq \tilde{\mathcal{U}}_t$ are open subsets of X and $\bar{\mathcal{U}}_0 \supseteq \tilde{\mathcal{U}}_0$ are open neighborhoods of p . Set $V := \bar{\mathcal{U}}_t$ and $\bar{U} := \bar{\mathcal{U}}_0$.

Next, by Proposition 1.5, there exists an open neighborhood U of E in Y and a holomorphic deformation of the complex structure on U such that E deforms to a smooth del Pezzo surface E' . So $p \in \bar{U}$ can be deformed to $p' \in \bar{U}'$ such that $\phi': U' \rightarrow \bar{U}'$ is also the weighted blow-up at p' with weight α :

$$\begin{array}{ccccc} E \subseteq U & \rightsquigarrow & U' \supseteq E' \\ \downarrow & \downarrow \phi & \downarrow \phi' & \downarrow \\ p \in \bar{U} & \rightsquigarrow & \bar{U}' & \ni p'. \end{array} \quad (4.2)$$

Let $\tau: \Delta \rightarrow \text{Def}(\bar{Y}, p)$ be the corresponding map of the deformation of \bar{U} . Then there are a holomorphic map

$$\sigma': \Delta \rightarrow \Delta \times \Delta \xrightarrow{\sigma \times \tau} \text{Def}(\bar{Y}, p) \quad (4.3)$$

and a good representative $\mathfrak{V} \rightarrow \Delta$ of the deformation corresponding to σ' such that $\mathfrak{V}_0 = \bar{U}'$ and $\mathfrak{V}_t = V$. Therefore we get a local del Pezzo transition $U' \searrow V$ such that the exceptional divisor E' of $U' \rightarrow \bar{U}'$ is smooth. As before, applying the deformation to the normal cone and Proposition 1.10 to $U' \searrow V$, we get two semistable (projective) degenerations $\mathcal{U}' \rightarrow \Delta$ and $\mathcal{V} \rightarrow \Delta$ with the special fibers $\mathcal{U}'_0 = U' \cup Y_d$ and $\mathcal{V}_0 = U' \cup X_d$. Moreover, $\mathcal{U}' \setminus \mathcal{U}'_0$ is the trivial family $U' \times (\Delta \setminus \{0\})$ and $\mathcal{V}_t = V$.

Now we will use the trick in [26, 51] to construct semistable (symplectic) degenerations. Let J_Y be the (integrable) complex structure of Y . We fix an ample class of (Y, J_Y) and denote ω_Y by a corresponding Kähler form. Note that U' can be viewed as a deformed complex structure on U (4.2). We can patch this deformed complex structure on U together (in a C^∞ way by a partition of unity argument) with the original complex structure J_Y to produce an almost complex structure J'_Y on Y . The almost complex structure J'_Y is integrable in

a neighborhood of E . Furthermore, it is ω_Y -tamed because the taming condition is open (cf. [37, p.153]). For abbreviation, we let Y' stand for (Y, J'_Y) . Notice that Y and Y' are symplectic deformation equivalent of each other.

Choose a suitable open neighborhood U'' of E in Y with $U'' \subseteq U$. Then we get a semistable (symplectic) degeneration, which we also denote by $\mathcal{Y} \rightarrow \Delta$, obtained by gluing the trivial family $(Y \setminus U'') \times \Delta$ and \mathcal{U}' . The family $\mathcal{Y}|_{\Delta \setminus \{0\}}$ is the trivial family with the fiber Y' . The special fiber \mathcal{Y}_0 consists of (Y, J'_Y) and Y_d which meet transversely and

$$Y' \cap Y_d = U' \cap Y_d = E'.$$

Here E' is the smooth del Pezzo surface of degree d . We also denote this degeneration by $Y' \rightsquigarrow Y' \cup_{E'} Y_d$.

Similarly, we can glue $\mathfrak{X} \setminus \tilde{\mathcal{U}}$ and \mathcal{V} to get a semistable (symplectic) degeneration, also denoted by $\mathcal{X} \rightarrow \Delta$, with the fiber $\mathcal{X}_t = X$. The special fiber \mathcal{X}_0 consists of Y' and X_d which meet transversely, and

$$Y' \cap X_d = U' \cap Y_d = E'.$$

The \mathcal{X} is also denoted by $X \rightsquigarrow Y' \cup_{E'} X_d$. We remark that, in contrast to the degeneration \mathcal{Y} with the general fiber $Y' = (Y, J'_Y)$, the general fiber of \mathcal{X} is the original smooth projective threefold X because $\mathcal{V}_t = \mathfrak{V}_t = V = \bar{\mathcal{U}}_t$ is the open subset of X by (4.1) and (4.3).

Now we summarize the above construction in the following proposition.

Theorem 4.1. *Let $Y \xrightarrow{\phi} \bar{Y} \rightsquigarrow X$ be a Type II extremal transition of degree d . Then there exist two semistable (symplectic) degeneration \mathcal{Y} and \mathcal{X} , i.e.,*

$$Y' \rightsquigarrow Y' \cup_{E'} Y_d \quad \text{and} \quad X \rightsquigarrow Y' \cup_{E'} X_d, \tag{4.4}$$

such that

- (1) Y and Y' are symplectic deformation equivalent of each other;
- (2) E' is a smooth del Pezzo surface of degree d and $Y_d = \mathbb{P}_{E'}(K_{E'} \oplus \mathcal{O})$;
- (3) X_d is a smooth del Pezzo threefold of degree d .

Moreover, up to a (local) deformation of the singularity $\text{Sing}(\bar{Y}) = \{p\}$, we have:

- (i) If $d = 1$, the \mathcal{X} is obtained by a degree six base change $\mathcal{X}' \rightarrow \Delta$ allowed by the weighted blow-up at $p \in \mathcal{X}'$ with weight $(3, 2, 1, 1, 1)$.
- (ii) If $d = 2$, the \mathcal{X} is obtained by a degree four base change $\mathcal{X}' \rightarrow \Delta$ allowed by the weighted blow-up at $p \in \mathcal{X}'$ with weight $(2, 1, 1, 1, 1)$.
- (iii) If $d = 3$ (resp. $d \geq 4$), the \mathcal{X} is obtained by a degree three (resp. two) base change $\mathcal{X}' \rightarrow \Delta$ allowed by the blow-up at $p \in \mathcal{X}'$.

Remark 4.2. For compact symplectic manifolds (M_i, ω_{M_i}) for $i = 0, \infty$, we have the symplectic sum construction. Indeed, assume that M_0 and M_∞ have a common compact symplectic divisor D and there exists an isomorphism

$$\Phi: N_{D/M_0} \otimes N_{D/M_\infty} \cong \mathcal{O}_D$$

of *complex* line bundles. The symplectic sum is a symplectic manifold (X, ω_Z) obtained from M_0 and M_∞ by gluing the complements of tubular neighborhoods of D in M_0 and M_∞ along their common boundary as directed by Φ . It produces a symplectic fibration $\mathcal{Z} \rightarrow \Delta$ with special fiber $\mathcal{Z}_0 = M_0 \cup_D M_\infty$, and it satisfies some conditions (see for example [48, §1.2] and references therein).

One can apply the symplectic sum construction to Y' and X_d with a common divisor E' , and similarly for $Y' \cup_{E'} Y_d$. However, since we assume that the global smoothing $\mathfrak{X} \rightarrow \Delta$ exists, we can construct the degeneration $X \rightsquigarrow Y' \cup_{E'} X_d$ directly as above.

4.2. Degeneration analysis revisited. To compare the (genus zero) GW invariants of $Y \searrow X$ for curve classes, we will use the degeneration formula to reduce the problem to local models $Y_d \searrow X_d$ as in [23].

Let (M, ω_M) be a compact symplectic manifold with a compact symplectic divisor D . For example, M is a smooth polarized variety and D is a smooth divisor. For a relative pair (M, D) , we let $\Gamma = (n, \beta, \rho, \mu)$ with $\mu = (\mu_1, \dots, \mu_\rho) \in \mathbb{N}^\rho$ a partition of the intersection number $(D, \beta) = |\mu| := \sum_{i=1}^\rho \mu_i$. For $A \in H(M)^{\otimes n}$ and $\varepsilon \in H(D)^{\otimes \rho}$, the relative invariant of stable maps with topological type Γ (i.e., with contact order μ_i in D at the i -th contact point) is

$$\langle A \mid \varepsilon, \mu \rangle_{\Gamma}^{(M, D)} := \int_{[\overline{\mathcal{M}}_{\Gamma}(M, D)]^{\text{vir}}} e_M^* A \cup e_D^* \varepsilon$$

where $e_M: \overline{\mathcal{M}}_{\Gamma}(M, D) \rightarrow M^n$ and $e_D: \overline{\mathcal{M}}_{\Gamma}(M, D) \rightarrow D^\rho$ are evaluation maps on marked points and contact points, respectively. The virtual dimension of $[\overline{\mathcal{M}}_{\Gamma}(M, D)]^{\text{vir}}$ is

$$\text{vdim}_{\Gamma} := (c_1(M), \beta) + \rho - |\mu| + n + (\dim M - 3).$$

Here the virtual class $[\overline{\mathcal{M}}_{\Gamma}(M, D)]^{\text{vir}}$ is constructed in [16, 26] in the symplectic case and in [27] in the algebraic case. In the symplectic setting, we refer the reader to [17, 49] for more details.

If $\Gamma = \bigsqcup_{\pi} \Gamma^{\pi}$, then the relative invariant with disconnected domain curve is defined by the product rule:

$$\langle A \mid \varepsilon, \mu \rangle_{\Gamma}^{\bullet(M, D)} := \prod \langle A \mid \varepsilon, \mu \rangle_{\Gamma^{\pi}}^{(M, D)}.$$

We apply the degeneration formula to the semistable (symplectic) degenerations (4.4) given in Theorem 4.1. For simplicity of notation, we denote these two degenerations by

$$\mathcal{W} \rightarrow \Delta.$$

It has a smooth fiber $W := \mathcal{W}_t$ ($t \neq 0$), a special fiber $\mathcal{W}_0 = M_0 \cup_D M_\infty$ and $D = M_0 \cap M_\infty$ a smooth divisor.

All cohomology classes $\vec{a} \in H(W)^{\oplus n}$ have global liftings and the restriction $\vec{a}(s)$ on \mathcal{W}_s is defined for all $s \in \Delta$, because $Y' \rightsquigarrow Y' \cup_{E'} Y_d$ is a degeneration of a trivial family and we apply Proposition 1.13 (i) to $X \rightsquigarrow Y' \cup_{E'} X_d$ (see (1.17)). Let $j_0: M_0 \hookrightarrow \mathcal{W}_0$ and $j_\infty: M_\infty \hookrightarrow \mathcal{W}_0$

be the inclusion maps. Let $\{e_i\}$ be a basis of $H(D)$, with $\{e^i\}$ its dual basis. Then $\{e_I\}$ forms a basis of $H(D^\rho)$ with dual basis $\{e^I\}$, where $|I| = \rho$ and $e_I = e_{i_1} \otimes \cdots \otimes e_{i_\rho}$. The *degeneration formula* expresses the absolute invariants of W in terms of the relative invariants of the two smooth pairs (M_0, D) and (M_∞, D) [26, 28, 17, 32, 48]:

$$\langle \vec{a} \rangle_\beta^W = \sum_I \sum_{\eta \in \Omega_\beta} C_\eta \langle j_0^* \vec{a}(0) \mid e_I, \mu \rangle_{\Gamma_0}^{\bullet(M_0, D)} \langle j_\infty^* \vec{a}(0) \mid e^I, \mu \rangle_{\Gamma_\infty}^{\bullet(M_\infty, D)}. \quad (4.5)$$

Here $\eta = (\Gamma_0, \Gamma_\infty, I_\rho)$ is an *admissible triple* which consists of (possibly disconnected) topological types

$$\Gamma_i = \bigsqcup_{\pi=1}^{|\Gamma_i|} \Gamma_i^\pi$$

with the same partition μ of the contact order under the identification I_ρ of the contact points. The gluing $\Gamma_0 +_{I_\rho} \Gamma_\infty$ have type (n, β) and is connected. In particular, $\rho = 0$ if and only if one of the Γ_i is empty. The total number of marked points n_i and the total degree $\beta_i \in \text{NE}(M_i)$ satisfy the splitting relations $n = n_0 + n_\infty$ and

$$\beta = \beta_0 + \beta_\infty,$$

where we omit the obvious pushforwards in the degree splitting in $\text{NE}(\mathcal{W})$, and $|\Gamma_0| + |\Gamma_\infty| = \rho + 1$. The constants $C_\eta = m(\mu)/|\text{Aut } \eta|$, where $m(\mu) = \prod \mu_i$ and $\text{Aut } \eta = \{\sigma \in \mathfrak{S}_\rho \mid \eta^\sigma = \eta\}$. We denote by Ω_β the set of equivalence classes of admissible triples with fixed degree β .

4.3. From local to global. Let $Y \xrightarrow{\phi} \bar{Y} \rightsquigarrow X$ be a Type II extremal transition of degree d . Let $\mathsf{R}_Y = \mathsf{R}_d[\![Q^{\phi^* \text{NE}(X)}]\!]$ and let $\mathsf{I}_Y = \mathsf{I}_d[\![Q^{\phi^* \text{NE}(X)}]\!]$ (see Definition 3.17 and Definition 3.23).

Theorem 4.3. *Let $n \geq 1$, $\vec{b} \in H(X)^{\oplus n}$ and $0 \neq \bar{\beta} \in \text{NE}(X)$. Consider the generating function*

$$\langle \phi^* \vec{b} \rangle_{\bar{\beta}}^{Y \searrow X} := \sum_{\beta \mapsto \bar{\beta}} \langle \phi^* \vec{b} \rangle_\beta^Y Q^\beta.$$

Then $\langle \phi^ \vec{b} \rangle_{\bar{\beta}}^{Y \searrow X}$ lies in R_Y and*

$$\langle \phi^* \vec{b} \rangle_{\bar{\beta}}^{Y \searrow X} \equiv \langle \vec{b} \rangle_{\bar{\beta}}^X Q^{\tilde{\beta}} \pmod{\mathsf{I}_Y}, \quad (4.6)$$

where $\tilde{\beta}$ is a lifting of $\bar{\beta}$ such that $(E, \tilde{\beta}) = 0$.

Proof. The proof is mainly based on the argument in [23, §4.4], using the method developed in [36].

We are going to apply the degeneration formula (4.5) to two semistable (symplectic) degenerations

$$Y' \rightsquigarrow Y' \cup_{E'} Y_d \quad \text{and} \quad X \rightsquigarrow Y' \cup_{E'} X_d$$

given in Theorem 4.1. Let $X_{\text{loc}} = X_d$ and $Y_{\text{loc}} = Y_d$ (see §3.1). For $\vec{b} = (b_1, \dots, b_n)$, without loss of generality we can assume $b_k \notin H^0(X)$ by the fundamental class axiom. We choose the support of b_k outside of $\phi'(E') = \{p'\}$, where ϕ' is the weighted blow-up in (4.2). Since we

construct semistable degenerations via weighted blow-ups at p' (see §4.1), the restriction of global lifting of b_k to X_{loc} is 0, and similarly for Y_{loc} . Denote $c_k = \phi^* b_k$ and $\vec{c} = (c_1, \dots, c_n)$. Since Y and Y' are symplectic deformation equivalent of each other, we know that $\langle \vec{c} \rangle_{\beta}^Y = \langle \vec{c} \rangle_{\beta}^{Y'}$ [30, 47]. By the degeneration formula (4.5), we obtain

$$\begin{aligned}\langle \vec{b} \rangle_{\bar{\beta}}^X &= \sum C_{\eta} \langle \vec{c} | e_I, \mu \rangle_{\bar{\Gamma}_0}^{\bullet(Y', E')} \cdot \langle \emptyset | e^I, \mu \rangle_{\bar{\Gamma}_{\infty}}^{\bullet(X_{\text{loc}}, h)}, \\ \langle \vec{c} \rangle_{\beta}^Y &= \sum C_{\eta} \langle \vec{c} | e_I, \mu \rangle_{\Gamma_0}^{\bullet(Y', E')} \cdot \langle \emptyset | e^I, \mu \rangle_{\Gamma_{\infty}}^{\bullet(Y_{\text{loc}}, H)},\end{aligned}\tag{4.7}$$

where $\bar{\Gamma}_0 = \Gamma_0 = (\beta_0, \rho, \mu)$, $\bar{\Gamma}_{\infty} = (\bar{\beta}_{\infty}, \rho, \mu)$, $\Gamma_{\infty} = (\beta_{\infty}, \rho, \mu)$ with

$$\bar{\beta} = \phi_* \beta_0, \quad \beta = \beta_0 + \pi_* \beta_{\infty}, \quad (E', \beta_0) = (h, \bar{\beta}_{\infty}) = (H, \beta_{\infty}) = |\mu|.$$

Here, $\pi_*: H_2(Y_{\text{loc}}) \rightarrow H_2(E') \rightarrow H_2(Y')$ is the natural pushforward. When $|\mu| = 0$, since $\bar{\beta} \neq 0$, we must have $\beta_{\infty} = 0$ and $(E', \beta) = (E', \beta_0) = 0$, and this implies $\beta = \bar{\beta}$.

The virtual dimensions satisfy

$$\begin{aligned}\text{vdim}_{\bar{\Gamma}_{\infty}} &= \text{vdim}_{\bar{\beta}_{\infty}} + \rho - |\mu| = \rho + |\mu| \geq 2\rho, \\ \text{vdim}_{\Gamma_{\infty}} &= \text{vdim}_{\beta_{\infty}} + \rho - |\mu| = \rho + |\mu| \geq 2\rho,\end{aligned}$$

which implies that the only non-vanishing contribution arises when $(e_I, \mu) = (1, 1)^{\rho}$, $(e^I, \mu) = (\text{pt}, 1)^{\rho}$. This also gives

$$\beta_0 = \tilde{\beta} - \rho\ell, \quad \bar{\beta}_{\infty} = \rho\bar{\gamma}, \quad \beta_{\infty} = \rho\tilde{\gamma} + m\ell,$$

where $m = -(E', \beta)$ satisfies $\beta = \tilde{\beta} + m\ell$.

We write $\Gamma_{\infty} \mapsto \bar{\Gamma}_{\infty}$ if $|\Gamma_{\infty}| \rightarrow |\bar{\Gamma}_{\infty}|$ is an isomorphism of graphs and $\beta_{\infty} \mapsto \bar{\beta}_{\infty}$ and define

$$\langle \emptyset | (\text{pt}, 1)^{\rho} \rangle_{\bar{\Gamma}_{\infty}}^{\bullet(Y_{\text{loc}}, H) \searrow (X_{\text{loc}}, h)} := \sum_{\Gamma_{\infty} \mapsto \bar{\Gamma}_{\infty}} \langle \emptyset | (\text{pt}, 1)^{\rho} \rangle_{\Gamma_{\infty}}^{\bullet(Y_{\text{loc}}, H)} Q^{\beta_{\infty} - \rho\tilde{\gamma}}.$$

Let $\beta_{\min} \in \text{NE}(Y)$ be the minimal lifting of $\bar{\beta}$, and let $e = (E', \beta_{\min})$. Then (4.7) induces

$$\begin{aligned}\langle \vec{b} \rangle_{\bar{\beta}}^X Q^{\tilde{\beta}} &= \sum_{\rho=1}^e C_{\eta} \langle \vec{c} | (1, 1)^{\rho} \rangle_{\bar{\Gamma}_0}^{\bullet(Y', E')} Q^{\tilde{\beta}} \langle \emptyset | (\text{pt}, 1)^{\rho} \rangle_{\bar{\Gamma}_{\infty}}^{\bullet(X_{\text{loc}}, h)} + \langle \vec{c} | \emptyset \rangle_{\tilde{\beta}}^{(Y', E')} Q^{\tilde{\beta}} \\ \sum_{\beta \mapsto \bar{\beta}} \langle \vec{c} \rangle_{\beta}^Y Q^{\beta} &= \sum_{m \geq -e} \langle \vec{c} \rangle_{\tilde{\beta} + m\ell}^{Y'} Q^{\tilde{\beta} + m\ell} \\ &= \sum_{\rho=1}^e C_{\eta} \langle \vec{c} | (1, 1)^{\rho} \rangle_{\Gamma_0}^{\bullet(Y', E')} Q^{\tilde{\beta}} \langle \emptyset | (\text{pt}, 1)^{\rho} \rangle_{\bar{\Gamma}_{\infty}}^{\bullet(Y_{\text{loc}}, H) \searrow (X_{\text{loc}}, h)} + \langle \vec{c} | \emptyset \rangle_{\tilde{\beta}}^{(Y', E')} Q^{\tilde{\beta}}\end{aligned}$$

Define

$$\begin{aligned}[\rho]^{Y_{\text{loc}}} &:= \sum_{m \geq -\rho} \langle \emptyset | (\text{pt}, 1)^{\rho} \rangle_{\rho\tilde{\gamma} + m\ell}^{(Y_{\text{loc}}, H)} Q^{m\ell} \\ [\rho]^{X_{\text{loc}}} &:= \langle \emptyset | (\text{pt}, 1)^{\rho} \rangle_{\rho\tilde{\gamma}}^{(X_{\text{loc}}, h)}\end{aligned}$$

By decomposing each $\bar{\Gamma}_\infty$ into its connected component, it suffices to prove that for each $\rho > 0$,

$$[\rho]^{Y_{\text{loc}}} \in \mathsf{R}_d \quad \text{and} \quad [\rho]^{Y_{\text{loc}}} \equiv [\rho]^{X_{\text{loc}}} \pmod{\mathsf{I}_d}. \quad (4.8)$$

Denote $[\rho]_\Gamma^{\bullet Y_{\text{loc}}} = \prod [\rho_i]^{Y_{\text{loc}}}$ and $[\rho]_\Gamma^{\bullet X_{\text{loc}}} = \prod [\rho_i]^{X_{\text{loc}}}$. Applying the degeneration formula again (4.5), we get

$$\begin{aligned} \langle \text{pt}^{\otimes \rho} \rangle_{\rho\tilde{\gamma}}^{X_{\text{loc}}} Q^{\rho\tilde{\gamma}} &= \sum_{\rho'=1}^{\rho} C_\eta \langle \text{pt}^{\otimes \rho} \mid (1, 1)^{\rho'} \rangle_{\bar{\Gamma}_0}^{\bullet(Y_{\text{loc}}, E')} Q^{\rho\tilde{\gamma}} \cdot [\rho']_{\bar{\Gamma}_\infty}^{\bullet X_{\text{loc}}} + \langle \text{pt}^{\otimes \rho} \mid \emptyset \rangle_{\rho\tilde{\gamma}}^{(Y_{\text{loc}}, E')} Q^{\rho\tilde{\gamma}}, \\ \sum_{m \geq -\rho} \langle \text{pt}^{\otimes \rho} \rangle_{\rho\tilde{\gamma}+m\ell}^{Y_{\text{loc}}} Q^{\rho\tilde{\gamma}+m\ell} &= \sum_{\rho'=1}^{\rho} C_\eta \langle \text{pt}^{\otimes \rho} \mid (1, 1)^{\rho'} \rangle_{\bar{\Gamma}_0}^{\bullet(Y_{\text{loc}}, E')} Q^{\rho\tilde{\gamma}} \cdot [\rho']_{\bar{\Gamma}_\infty}^{\bullet Y_{\text{loc}}} + \langle \text{pt}^{\otimes \rho} \mid \emptyset \rangle_{\rho\tilde{\gamma}}^{(Y_{\text{loc}}, E')} Q^{\rho\tilde{\gamma}}, \end{aligned}$$

where $\bar{\Gamma}_0 = \Gamma_0 = (\rho\tilde{\gamma} - \rho'\ell, \rho', 1^{\rho'})$ and $\bar{\Gamma}_\infty = (\rho'\tilde{\gamma}, \rho', 1^{\rho'})$.

Now, by Proposition 3.19 (if $d \neq 7$) or Proposition 3.24 (if $d = 7$), we know that

$$\langle \text{pt}^{\otimes \rho} \rangle_{\rho\tilde{\gamma}}^{X_{\text{loc}}} Q^{\rho\tilde{\gamma}} \equiv \sum_{m \geq -\rho} \langle \text{pt}^{\otimes \rho} \rangle_{\rho\tilde{\gamma}+m\ell}^{Y_{\text{loc}}} Q^{\rho\tilde{\gamma}+m\ell} \pmod{\mathsf{I}_{\text{loc}}}.$$

Moreover, among the relative invariants appearing in the degeneration formula, the only term contributing to the coefficient of $[\rho]^{X_{\text{loc}}}$ (and $[\rho]^{Y_{\text{loc}}}$) is

$$\langle \text{pt}^{\otimes \rho} \mid (1, 1)^\rho \rangle_{\rho\tilde{\gamma}-\rho\ell}^{\bullet(Y_{\text{loc}}, E')} = \langle \text{pt}^{\otimes \rho} \mid (1, 1)^\rho \rangle_{\rho\tilde{\gamma}}^{\bullet(Y_{\text{loc}}, E')} = \left(\langle \text{pt} \mid (1, 1) \rangle_{\gamma}^{(Y_{\text{loc}}, E')} \right)^\rho,$$

which is nonzero (cf. [36, §1.2], or using the identity $\langle \text{pt} \mid (1, 1) \rangle_{\gamma}^{(Y_{\text{loc}}, E')} [1]_{\gamma}^{(Y_{\text{loc}}, H)} = \langle \text{pt} \rangle_{\gamma} = 1$). All other terms $[\rho']_{\bar{\Gamma}_\infty}^{\bullet X_{\text{loc}}}$ and $[\rho']_{\bar{\Gamma}_\infty}^{\bullet Y_{\text{loc}}}$ satisfy either $\rho' < \rho$, or has at least two connected components. Therefore (4.8) follows by induction on ρ . \square

Corollary 4.4. Suppose $d \in \{1, 2, 3, 4, 7\}$. Let $n \geq 1$, $\vec{b} \in H(X)^{\oplus n}$ and $0 \neq \bar{\beta} \in \text{NE}(X)$. Then $\langle \phi^* \vec{b} \rangle_{\bar{\beta}}^{Y \searrow X} Q^{-\bar{\beta}}$ is a multi-valued analytic function of $P^\ell \in \mathbb{P}^1$, and

$$\lim_{P^\ell \rightarrow \infty} \langle \phi^* \vec{b} \rangle_{\bar{\beta}}^{Y \searrow X} Q^{-\bar{\beta}} = \langle \vec{b} \rangle_{\bar{\beta}}^X. \quad (4.9)$$

Proof. Since $v_d = \frac{\theta f_d}{f_d}$ is bounded near $P^\ell = \infty$, elements in $\mathsf{I}_d = \mathsf{S}_d \cap y \cdot \mathbb{C}[[y]]v_d$ tend to 0 as $P^\ell \rightarrow \infty$. Therefore (4.9) follows from (4.6). \square

Recall that f_d^{reg} is a nonzero holomorphic solution of (3.9) near $P^\ell = \infty$ and $v_d^{\text{reg}} = \frac{\theta f_d^{\text{reg}}}{f_d}$ (see Definition 3.17).

Definition 4.5. For $d \in \{5, 6\text{I}, 6\text{II}, 8\}$, we define the *regularization map*

$$\begin{aligned} \mathsf{S}_d = \mathbb{C}(y)[v_d] &\xrightarrow{(-)^{\text{reg}}} \mathbb{C}(y)[v_d^{\text{reg}}] =: \mathsf{S}_d^{\text{reg}} \subseteq \mathbb{C}(\!(y)\!) \\ v_d &\longmapsto v_d^{\text{reg}}, \end{aligned}$$

and extend it to $\mathsf{S}_d[[Q^{\phi^* \text{NE}(X)}]] \rightarrow \mathsf{S}_d^{\text{reg}}[[Q^{\phi^* \text{NE}(X)}]]$, which we also denote by $(-)^{\text{reg}}$.

The following corollary follows by a similar argument to Corollary 4.4.

Corollary 4.6. Suppose $d \in \{5, 6\mathbb{I}, 6\mathbb{II}, 8\}$. Let $n \geq 1$, $\vec{b} \in H(X)^{\oplus n}$ and $0 \neq \bar{\beta} \in \text{NE}(X)$. Then the regularization $(\langle \phi^* \vec{b} \rangle_{\bar{\beta}}^{Y \setminus X} Q^{-\bar{\beta}})^{\text{reg}}$ is holomorphic near $P^\ell = \infty$, and

$$\lim_{P^\ell \rightarrow \infty} (\langle \phi^* \vec{b} \rangle_{\bar{\beta}}^{Y \setminus X} Q^{-\bar{\beta}})^{\text{reg}} = \langle \vec{b} \rangle_{\bar{\beta}}^X.$$

5. THE LIMIT OF Q^ℓ

Recall that we have the mirror transform $\tau = t - g_d(P^\ell)E$ so that $Q^\ell = P^\ell e^{g_d(P^\ell)}$ (see Corollary 3.13) and a function $f_d = 1 + \theta g_d$ that satisfies the differential equation (3.4) with the set Sing_{f_d} of regular singular points (see Table 5).

As an analytic function, f_d admits an analytic continuation to a multi-valued function on $\mathbb{P}^1 \setminus \text{Sing}_{f_d}$. We can take a path $p: [0, 1] \rightarrow \mathbb{P}^1$ such that $p(0) = 0$, $p(1) = \infty$, and $p(s) \in \mathbb{P}^1 \setminus \text{Sing}_{f_d}$ for $s \in (0, 1)$. Then, using $f_d = 1 + \theta g_d$,

$$\begin{aligned} Q^\ell(s) &:= p(s) e^{g_d(p(s))} = p(s) \exp\left(\int_0^s (f_d(p(s')) - 1) \frac{dp(s')}{p(s')}\right) \\ &= \lim_{\varepsilon \rightarrow 0^+} p(\varepsilon) \exp\left(\int_\varepsilon^s f_d(p(s')) \frac{dp(s')}{p(s')}\right) \end{aligned}$$

is well-defined for $s \in [0, 1]$ and we can define the limit of Q^ℓ along the path p to be

$$Q_p^\ell := \lim_{s \rightarrow 1^-} Q^\ell(s) = \lim_{\varepsilon \rightarrow 0^+} p(\varepsilon) \exp\left(\int_\varepsilon^1 f_d(p(s)) \frac{dp(s)}{p(s)}\right).$$

The goal of this section is to calculate Q_p^ℓ along some path p .

5.1. The ordinary case. We first consider the simplest cases, namely $d \in \{1, 2, 3, 4\}$, for which f_d is a hypergeometric function.

Proposition 5.1. For $d \in \{1, 2, 3, 4\}$, the limit of Q^ℓ along the positive real line $(0, \infty)$ is 1.

Proof. We need to show that

$$\lim_{\varepsilon \rightarrow 0^+} \varepsilon \exp\left(\int_\varepsilon^\infty f_d(p) \frac{dp}{p}\right) = 1.$$

This is equivalent to

$$\lim_{\varepsilon \rightarrow 0^+} \left(\log \varepsilon + \int_\varepsilon^\infty f_d(p) \frac{dp}{p} \right) = 0.$$

Recall from Proposition 3.14 that

$$f_d(p) = {}_2F_1\left(\frac{1}{n}, 1 - \frac{1}{n}; 1; -\kappa_d p\right),$$

where $n = n_d = 6, 4, 3, 2$ for $d = 1, 2, 3, 4$, respectively. By the Euler integral representation of the hypergeometric function, we have

$$f_d(p) = \frac{1}{B(1 - \frac{1}{n}, \frac{1}{n})} \int_0^1 \frac{ds}{s^{1/n} (1-s)^{1-1/n} (1+\kappa_d p s)^{1/n}}, \quad (5.1)$$

where $B(-, -)$ is the Beta function. Substitute this into the integral, we obtain

$$\begin{aligned} \int_{\varepsilon}^{\infty} f_d(p) \frac{dp}{p} &= \frac{1}{B(1 - \frac{1}{n}, \frac{1}{n})} \int_{\varepsilon}^{\infty} \int_0^1 \frac{ds dp}{s^{1/n}(1-s)^{1-1/n}(1+\kappa_d p s)^{1/n} p} \\ &= \frac{1}{B(1 - \frac{1}{n}, \frac{1}{n})} \int_0^1 \int_{\varepsilon}^{\infty} \frac{dp}{(1+\kappa_d p s)^{1/n} p} \frac{ds}{s^{1/n}(1-s)^{1-1/n}}. \end{aligned} \quad (5.2)$$

Make the substitution $x = (1 + \kappa_d p s)^{1/n}$, so that $p = \frac{x^{n-1}}{\kappa_d s}$ and $dp = \frac{n x^{n-1}}{\kappa_d s} dx$. Then

$$\int_{\varepsilon}^{\infty} \frac{dp}{(1+\kappa_d p s)^{1/n} p} = \int_{(1+\kappa_d \varepsilon s)^{1/n}}^{\infty} \frac{n x^{n-2} dx}{x^n - 1}.$$

Denote by ω_n the n -th root of unity $e^{2\pi i/n}$. Then, for R large,

$$\begin{aligned} \int_{(1+\kappa_d \varepsilon s)^{1/n}}^R \frac{n x^{n-2} dx}{x^n - 1} &= \sum_{j=0}^{n-1} \int_{(1+\kappa_d \varepsilon s)^{1/n}}^R \frac{\omega_n^{-j} dx}{x - \omega_n^j} \\ &= \int_{(1+\kappa_d \varepsilon s)^{1/n}}^R \frac{dx}{x - 1} + \sum_{j=1}^{n-1} \int_1^R \frac{\omega_n^{-j} dx}{x - \omega_n^j} + O(\varepsilon) \\ &= \left(\log R + \sum_{j=1}^{n-1} \omega_n^{-j} \log(R - \omega_n^j) \right) \\ &\quad - \log\left(\frac{\kappa_d \varepsilon s}{n}\right) - \sum_{j=1}^{n-1} \omega_n^{-j} \log(1 - \omega_n^j) + O(\varepsilon). \end{aligned}$$

As $R \rightarrow \infty$, we have

$$\log R + \sum_{j=1}^{n-1} \omega_n^{-j} \log(R - \omega_n^j) = \left(1 + \sum_{j=1}^{n-1} \omega_n^{-j} \right) \log R + o(1) \rightarrow 0.$$

Thus,

$$\int_{(1+\kappa_d \varepsilon s)^{1/n}}^{\infty} \frac{n x^{n-2} dx}{x^n - 1} = -\log\left(\frac{\kappa_d \varepsilon s}{n}\right) - \sum_{j=1}^{n-1} \omega_n^{-j} \log(1 - \omega_n^j) + O(\varepsilon).$$

Putting this back into the double integral (5.2), we obtain

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0^+} \left(\log \varepsilon + \int_{\varepsilon}^{\infty} f_d(p) \frac{dp}{p} \right) &= -\log\left(\frac{\kappa_d}{n}\right) - \sum_{j=1}^{n-1} \omega_n^{-j} \log(1 - \omega_n^j) \\ &\quad - \frac{1}{B(1 - \frac{1}{n}, \frac{1}{n})} \int_0^1 \frac{\log s ds}{s^{1/n}(1-s)^{1-1/n}}. \end{aligned} \quad (5.3)$$

It remains to evaluate the final integral. We have

$$\begin{aligned} \int_0^1 \frac{\log s \, ds}{s^{1/n}(1-s)^{1-1/n}} &= \frac{d}{d\alpha} \Big|_{\alpha=0} \int_0^1 \frac{s^\alpha \, ds}{s^{1/n}(1-s)^{1-1/n}} \\ &= \frac{d}{d\alpha} \Big|_{\alpha=0} B(1 - \frac{1}{n} + \alpha, \frac{1}{n}). \end{aligned}$$

Using the derivative of the Beta function

$$\frac{d}{d\alpha} \Big|_{\alpha=0} B(a + \alpha, b) = \frac{d}{d\alpha} \Big|_{\alpha=0} \frac{\Gamma(a + \alpha)\Gamma(b)}{\Gamma(a + b + \alpha)} = B(a, b)(\psi^{(0)}(a) - \psi^{(0)}(a + b)),$$

where $\psi^{(0)} = \Gamma'/\Gamma$ is the digamma function, we get

$$\frac{1}{B(1 - \frac{1}{n}, \frac{1}{n})} \int_0^1 \frac{\log s \, ds}{s^{1/n}(1-s)^{1-1/n}} = \psi^{(0)}(1 - \frac{1}{n}) - \psi^{(0)}(1).$$

By Gauss' digamma formula, we get

$$\psi^{(0)}(1 - \frac{1}{n}) - \psi^{(0)}(1) = \sum_{m=1}^{\infty} \left(\frac{1}{m} - \frac{1}{m - \frac{1}{n}} \right) = - \sum_{j=1}^{n-1} (1 - \omega_n^j) \log(1 - \omega_n^j).$$

Substituting this back into (5.3), we find

$$\lim_{\varepsilon \rightarrow 0^+} \left(\log \varepsilon + \int_{\varepsilon}^{\infty} f_d(p) \frac{dp}{p} \right) = -\log\left(\frac{\kappa_d}{n}\right) - \sum_{j=1}^{n-1} (\omega_n^j + \omega_n^{-j} - 1) \log(1 - \omega_n^j).$$

Finally, one verifies case by case that the right hand side is exactly 0 for $d \in \{1, 2, 3, 4\}$. \square

We can now rephrase Theorem 4.3 as follows:

Corollary 5.2. *Let $Y \searrow X$ be a Type II transition of degree $d \in \{1, 2, 3, 4, 7\}$ and let $\bar{\beta} \neq 0$. Then the analytic function*

$$\langle \phi^* \vec{b} \rangle_{\bar{\beta}}^{Y \searrow X} = \sum_{\beta \rightarrow \bar{\beta}} \langle \phi^* b \rangle_{\beta}^Y Q^{\beta}$$

has an analytic continuation along $Q^\ell \in [0, 1)$ and

$$\lim_{Q^\ell \rightarrow 1} \langle \phi^* \vec{b} \rangle_{\bar{\beta}}^{Y \searrow X} = \langle b \rangle_{\bar{\beta}}^X Q^{\tilde{\beta}},$$

where $\tilde{\beta}$ is any lifting of $\bar{\beta}$ under $\text{NE}(Y) \rightarrow \text{NE}(X)$.

Proof. For $d \in \{1, 2, 3, 4\}$,

$$Q^\ell = \lim_{\varepsilon \rightarrow 0^+} \varepsilon \exp\left(\int_{\varepsilon}^{P^\ell} f_d(p) \frac{dp}{p}\right)$$

is real-valued on the interval $(0, \infty)$, and its limit as $P^\ell \rightarrow \infty$ is equal to 1, as shown in Proposition 5.1. Since $v_d = \frac{\theta f_d}{f_d}$ is bounded near $P^\ell = \infty$, elements of \mathbf{l}_Y tend to 0 as $P^\ell \rightarrow \infty$. Hence,

$$\lim_{Q^\ell \rightarrow 1} \langle \phi^* \vec{b} \rangle_{\bar{\beta}}^{Y \searrow X} = \langle b \rangle_{\bar{\beta}}^X Q^{\tilde{\beta}}.$$

For $d = 7$, it follows from (3.23) that, as $y \rightarrow 0$,

$$Q^\ell = P^\ell e^{g(P^\ell)} = \exp\left(-\sum \frac{b_k(1)}{k} y^k\right) \rightarrow 1.$$

Since elements of I_Y tend to 0 as $y \rightarrow 0$, we also obtain

$$\lim_{Q^\ell \rightarrow 1} \langle \phi^* \vec{b} \rangle_{\tilde{\beta}}^{Y \searrow X} = \langle b \rangle_{\tilde{\beta}}^X Q^{\tilde{\beta}}.$$

Finally, note that we may choose any lifting $\tilde{\beta}$ of $\bar{\beta}$, since any two liftings differ by a multiple of ℓ , and Q^ℓ tends to 1. \square

The following is the first nontrivial case where the limit differs from 1, though, still a root of unity.

Corollary 5.3. *For $d = 8$, the limit of Q^ℓ along the positive imaginary line $(0, i\infty)$ is i .*

Proof. It follows from Proposition 3.14 that $f_8(P^\ell) = f_4(-P^{2\ell})$. Therefore, $g_8(P^\ell) = \frac{1}{2}g_4(-P^{2\ell})$, and hence $Q^\ell = P^\ell e^{g_8(P^\ell)} = P^\ell e^{\frac{1}{2}g_4(-P^{2\ell})}$. The result then follows easily from Proposition 5.1. \square

For $d \in \{5, 6I, 6II\}$, we will compute the corresponding integral using modular forms in the next subsection.

5.2. Relation to modular curves. Motivated by [2, 11], we establish a relation between meromorphic quasi-modular forms and the GW invariants of Y_{loc} , and consequently those of Y .

We briefly recall the definitions of modular and quasi-modular forms. Let $\mathbb{H} = \{\tau_q \in \mathbb{C} \mid \text{Im } \tau_q > 0\}$ denote the upper half plane, $\mathbb{H}^* = \mathbb{H} \cup \mathbb{Q} \cup \{i\infty\}$ denote the extended upper half plane, and set $q = e^{2\pi i \tau_q}$. The group $\text{SL}(2, \mathbb{Z})$ acts on \mathbb{H} via

$$A \cdot \tau_q = \begin{pmatrix} a_A & b_A \\ c_A & d_A \end{pmatrix} \cdot \tau_q = \frac{a_A \tau_q + b_A}{c_A \tau_q + d_A}.$$

The automorphy factor $j: \text{SL}(2, \mathbb{Z}) \times \mathbb{H} \rightarrow \mathbb{C}$ is defined by

$$j(A, \tau_q) = c_A \tau_q + d_A.$$

A subgroup $\Gamma \subseteq \text{SL}(2, \mathbb{Z})$ is called a *congruence subgroup* if $\Gamma \supseteq \Gamma(N)$ for some $N \geq 1$, where

$$\Gamma(N) = \left\{ A \in \text{SL}(2, \mathbb{Z}) \mid A \equiv \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \pmod{N} \right\}.$$

A cusp of Γ is a Γ -equivalence class in $\mathbb{Q} \cup \{i\infty\}$. A point $\tau \in \mathbb{H}$ is an elliptic point of Γ if the isotropy group Γ_τ is nontrivial in the sense that $\{\pm I\}\Gamma_\tau \neq \{\pm I\}$. We will use the

following congruence subgroups:

$$\begin{aligned}\Gamma_0(N) &= \left\{ A \in \mathrm{SL}(2, \mathbb{Z}) \mid A \equiv \begin{pmatrix} * & * \\ 0 & * \end{pmatrix} \pmod{N} \right\}, \\ \Gamma_1(N) &= \left\{ A \in \mathrm{SL}(2, \mathbb{Z}) \mid A \equiv \begin{pmatrix} 1 & * \\ 0 & 1 \end{pmatrix} \pmod{N} \right\}.\end{aligned}$$

A meromorphic function $F: \mathbb{H} \rightarrow \mathbb{P}^1$ is *weakly modular function* of weight k on Γ if for each $A \in \Gamma$,

$$F|_k^A(\tau_q) := \frac{F(A \cdot \tau_q)}{\mathbf{j}(A, \tau_q)^k} = F(\tau_q).$$

A holomorphic function $F: \mathbb{H} \rightarrow \mathbb{C}$ is a *modular form* of weight k on Γ if F is weakly modular and $F|_k^A(\tau_q)$ is holomorphic at $i\infty$ for all $A \in \mathrm{SL}(2, \mathbb{Z})$. The value of F at $A \cdot i\infty$ is defined to be $F|_k^A(i\infty)$.

The notion of quasi-modular forms was introduced by Kaneko and Zagier in [19] (see also [10]), where they are defined as the “constant term” of almost-holomorphic modular forms. In this paper, we adopt the following direct definition (see e.g., [53, p.58]), allowing meromorphicity:

Definition 5.4. For $A = \begin{pmatrix} a_A & b_A \\ c_A & d_A \end{pmatrix} \in \mathrm{SL}(2, \mathbb{Z})$, define

$$\mathbf{k}(A, \tau_q) = \frac{c_A}{\mathbf{j}(A, \tau_q)} = \frac{c_A}{c_A \tau_q + d_A}.$$

A meromorphic function $F: \mathbb{H} \rightarrow \mathbb{P}^1$ is a *meromorphic quasi-modular form* of weight k and depth $\leq s$ on Γ if

- (1) there exists meromorphic functions F_0, F_1, \dots, F_s on \mathbb{H} , meromorphic at $i\infty$, such that for each $A \in \Gamma$,

$$F|_k^A(\tau_q) = F_0(\tau_q) + \mathbf{k}(A, \tau_q)F_1(\tau_q) + \cdots + \mathbf{k}(A, \tau_q)^s F_s(\tau_q); \quad (5.4)$$

- (2) for each $A \in \mathrm{SL}(2, \mathbb{Z})$, there exists meromorphic functions $F_0^A, F_1^A, \dots, F_s^A$ on \mathbb{H} , meromorphic at $i\infty$, such that

$$F|_k^A(\tau_q) = F_0^A(\tau_q) + \mathbf{k}(A, \tau_q)F_1^A(\tau_q) + \cdots + \mathbf{k}(A, \tau_q)^s F_s^A(\tau_q).$$

The value of F at $A \cdot i\infty$ is defined to be $F_0^A(i\infty)$.

Applying $D_q = q \frac{d}{dq} = \frac{1}{2\pi i} \frac{d}{d\tau_q}$ to the identity (5.4), we get

$$(D_q F)|_{k+2}^A(\tau_q) = \sum_{n=0}^s \mathbf{k}(A, \tau_q)^n D_q F_n + \mathbf{k}(A, \tau_q)^{n+1} \frac{k-n}{2\pi i} F_n.$$

It follows that D_q maps a meromorphic quasi-modular form of weight k and depth $\leq s$ to a meromorphic quasi-modular form of weight $k+2$ and depth $\leq s+1$.

Recall that for $d \neq 7$, the function $f_d(P^\ell)$ satisfies the differential equation (3.4):

$$\theta(u_d \cdot \theta f_d) + (\lambda_d P^\ell - \mu_d P^{2\ell}) \cdot f_d = u_d \cdot \theta^2 f_d + \theta u_d \cdot \theta f_d + (\lambda_d P^\ell - \mu_d P^{2\ell}) \cdot f_d = 0$$

where $u_d = 1 + \kappa_d P^\ell - \mu_d P^{2\ell}$, $\theta = P^\ell \frac{d}{dP^\ell}$, and the triple $(\kappa_d, \lambda_d, \mu_d)$ is listed in Table 1.

For $d = 1$, consider the modular function (with root and monodromy)

$$P_1(\tau_q) = \frac{1}{864} \left(\frac{E_6(\tau_q)}{E_4(\tau_q)^{3/2}} - 1 \right),$$

defined near $q = 0$, where

$$E_4(\tau_q) = 1 + 240 \sum_{m \geq 1} \frac{m^3 q^m}{1 - q^m}, \quad E_6(\tau_q) = 1 - 504 \sum_{m \geq 1} \frac{m^5 q^m}{1 - q^m}$$

are the Eisenstein series. Then, near $q = 0$, we have (see [34, p.52], [2, §3], [8, Theorem 4.2])

$$f_1(P_1(\tau_q)) = {}_2F_1\left(\frac{1}{6}, \frac{5}{6}; 1; -432P_1(\tau_q)\right) = E_4(\tau_q)^{1/4}. \quad (5.5)$$

Since the j -invariant $j(\tau_q)$ parametrizes the modular curve $X(\Gamma(1)) = \Gamma(1) \backslash \mathbb{H}^*$ and

$$j(\tau_q)^{-1} = -P_1(1 + 432P_1),$$

the multi-valued function P_1 parametrizes a double cover $\tilde{X}(\Gamma(1))$ of $X(\Gamma(1))$, ramified at $j(\tau_q)^{-1} = 1728$ and ∞ , which correspond to the elliptic points $[\omega_4]$ and $[\omega_6]$, respectively. In Figure 1, a fundamental domain \mathcal{D}_1 of $\tilde{X}(\Gamma(1))$ (view as a double cover of a fundamental domain of $X(\Gamma(1))$) contains two cusps $[i\infty]$, $[0]$, which correspond to $P_1 = 0, -\frac{1}{432}$, respectively, and an elliptic point $[\omega_6]$, which corresponds to $P_1 = \infty$ (cf. [2, Figure 1]).

For $d \geq 2$, we have:

Proposition 5.5 ([34, 54]). *For each $d \in \{2, 3, 4, 5, 6, 8\}$, there is a congruence subgroup Γ_d and a meromorphic modular function $P_d(\tau_q) = -q + O(q^2)$ of weight 0 on Γ_d , such that*

- (i) if $d = 2$, then $f_d(P_d(\tau_q))^2$ is the weight 2 modular form $\Theta_{D_4}(\tau_q)$ on the group $\Gamma(2)$;
- (ii) if $3 \leq d \leq 6$, then $f_d(P_d(\tau_q))$ is a weight 1 modular form on Γ_d ;
- (iii) if $d = 8$, then $f_d(P_d(\tau_q))$ is a weight 1 modular form on $\Gamma_1(8)$.

Moreover, the induced map $P_d: X(\Gamma_d) = \Gamma_d \backslash \mathbb{H}^* \rightarrow \mathbb{P}^1$ is an isomorphism between Riemann surfaces. The points which are mapped to Sing_{f_d} by P_d correspond to cusps or elliptic points in $X(\Gamma_d)$, depending on whether the monodromy is infinite or finite, as listed in Table 6.

d	Γ_d	Sing_{f_d}	corresponding cusps or elliptic points
1	-	$\{0, \infty, -\frac{1}{432}\}$	$[\mathbf{i}\infty], [\omega_6], [0]$
2	$\Gamma(2)$	$\{0, \infty, -\frac{1}{64}\}$	$[\mathbf{i}\infty], [\frac{1+i}{2}], [0]$
3	$\Gamma_1(3)$	$\{0, \infty, -\frac{1}{27}\}$	$[\mathbf{i}\infty], [\frac{1+\omega_6}{3}], [0]$
4	$\Gamma_1(4)$	$\{0, \infty, -\frac{1}{16}\}$	$[\mathbf{i}\infty], [\frac{1}{2}], [0]$
5	$\Gamma_1(5)$	$\{0, \infty, \frac{11-5\sqrt{5}}{2}, \frac{11+5\sqrt{5}}{2}\}$	$[\mathbf{i}\infty], [\frac{2}{5}], [0], [\frac{1}{2}]$
6I	$\Gamma_1(6)$	$\{0, \infty, -\frac{1}{8}, 1\}$	$[\mathbf{i}\infty], [\frac{1}{3}], [0], [\frac{1}{2}]$
6II	$\Gamma_1(6)$	$\{0, \infty, -\frac{1}{9}, -1\}$	$[\mathbf{i}\infty], [\frac{1}{2}], [0], [\frac{1}{3}]$
8	$\Gamma_0(8)$	$\{0, \infty, -\frac{1}{4}, \frac{1}{4}\}$	$[\mathbf{i}\infty], [\frac{1}{4}], [0], [\frac{1}{2}]$

TABLE 6. Congruence subgroups Γ_d , singular points Sing_{f_d} , and corresponding points in $X(\Gamma_d)$.

Proof. This is a classical result, the functions $f_d(P_d(\tau_q))$ are known as Ramanujan's theta functions, see [8, §4-§6], [34], [54], and [55, Table 1.2] for $d = 2, 3, 4$. We summarize the expressions below for the reader's convenience.

As listed in [54] (with $P_d = -t$) and (5.5) for $d = 1$, the modular functions f_d and their corresponding P_d can be expressed as follows:

d	index	$f_d(P_d(\tau_q))$	$P_d(\tau_q)$
1	-	$\Theta_{E_8}^{1/4} = E_4(\tau_q)^{1/4}$	$-\frac{E_4(\tau_q)^{3/2} - E_6(\tau_q)}{864 E_4(\tau_q)^{3/2}}$
2	#20	$\Theta_{D_4}^{1/2} = (2 E_2(2\tau_q) - E_2(\tau_q))^{1/2}$	$-\frac{\eta(2\tau_q)^{24}}{\eta(\tau_q)^{24} + 64\eta(2\tau_q)^{24}}$
3	#14	$\Theta_{A_2} = \vartheta_3(\tau_q)$	$-\frac{\eta(3\tau_q)^{12}}{\eta(\tau_q)^{12} + 27\eta(3\tau_q)^{12}}$
4	#11	$\Theta_{\mathbb{Z}_2} = \vartheta_4(\tau_q)$	$-\frac{\eta(\tau_q)^8 \eta(4\tau_q)^{16}}{\eta(2\tau_q)^{24}}$
5	#9 = D	$f_5(\tau_q)$	$P_5(\tau_q)$
6I	#5 = A	$\frac{1}{3}\vartheta_3(\tau_q) + \frac{2}{3}\vartheta_3(2\tau_q)$	$-\frac{\eta(\tau_q)^3 \eta(6\tau_q)^9}{\eta(2\tau_q)^3 \eta(3\tau_q)^9}$
6II	#8 = C	$\frac{1}{2}\vartheta_3(\tau_q) + \frac{1}{2}\vartheta_3(2\tau_q)$	$-\frac{\eta(\tau_q)^4 \eta(6\tau_q)^8}{\eta(2\tau_q)^8 \eta(3\tau_q)^4}$
8	#2 = G	$\vartheta_4(2\tau_q)$	$-\frac{\eta(2\tau_q)^4 \eta(8\tau_q)^8}{\eta(4\tau_q)^{12}}$

TABLE 7. Ramanujan's theta functions f_d and their associated Hauptmoduln P_d .

Here,

$$E_2(\tau_q) = 1 - 24 \sum_{m \geq 1} \frac{mq^m}{1-q^m}$$

is the Eisenstein series,

$$\eta(\tau_q) = q^{1/24} \prod_{m=1}^{\infty} (1 - q^m)$$

is the Dedekind eta function, and

$$\vartheta_3(\tau_q) = 1 + 6 \sum_{m \geq 1} \chi_{3,2}(m) \frac{q^m}{1-q^m}, \quad \vartheta_4(\tau_q) = 1 + 4 \sum_{m \geq 1} \chi_{4,2}(m) \frac{q^m}{1-q^m}$$

are the theta series, with $\chi_{3,2} = \left(\frac{-3}{\cdot}\right)$ and $\chi_{4,2} = \left(\frac{-4}{\cdot}\right)$ denoting the unique nontrivial Dirichlet characters modulo 3 and 4, respectively.

For $d = 5$,

$$f_5(\tau_q) = 1 + \sum_{m \geq 1} \left(\frac{3-i}{2} \chi_{5,2}(m) + \frac{3+i}{2} \bar{\chi}_{5,2}(m) \right) \frac{q^m}{1-q^m}, \quad P_5(\tau_q) = -q \prod_{m=1}^{\infty} (1 - q^m)^{5\chi_{5,2}^2(m)}$$

where $\chi_{5,2}$ is the Dirichlet character modulo 5 with $\chi_{5,2}(2) = i$.

Using the explicit construction of P_d , we obtain in Table 8 the relation between $P_d(\tau_q)$ and the canonical Hauptmodul $t_d(\tau_q)$ that parametrizes the modular curve $X(\Gamma_0(d))$ (cf. [34, Table 2]):

d	2	3	4	5	6	8
$t_d(\tau_q)$	$-\frac{64^2 P_2}{1+64P_2}$	$-\frac{27^2 P_3}{1+27P_3}$	$-\frac{16^2 P_4}{1+16P_4}$	$-\frac{125 P_5}{1+11P_5-P_5^2}$	$-\frac{72 P_{6I}}{1+8P_{6I}}$	$= -\frac{72 P_{6II}}{1+9P_{6II}}$

TABLE 8. Relations between $P_d(\tau_q)$ and the canonical Hauptmodul $t_d(\tau_q)$.

These identities shows that P_d parametrize the modular curve $X(\Gamma_d)$, and moreover, by the values of t_d at cusps and elliptic points listed in [34, Table 2], each regular singular point corresponds to the desired cusp or elliptic point as listed in Table 6. \square

Definition 5.6. The point corresponding to $P^\ell = \infty$ in Table 6 (and shown in Figure 1) is called the *transition point* of degree d , denoted by $\tau_{q,d}^\infty$ ($d \neq 7$).

Since the coefficients of the differential equation (3.4) and the operator θ is invariant under the action of Γ_d , we see that $f_d(A\tau_q) = j(A, \tau_q)f_d(\tau_q)$ is also a solution for $A \in \Gamma_d$. In particular, $\tau_q f_d$ is a solution as well. This means that for any solution f , the following Wronskian determinant vanishes:

$$0 = \begin{vmatrix} f & f_d & \tau_q f_d \\ \theta f & \theta f_d & \theta(\tau_q f_d) \\ \theta^2 f & \theta^2 f_d & \theta^2(\tau_q f_d) \end{vmatrix} = W_2 \cdot \theta^2 f + W_1 \cdot \theta f + W_0 \cdot f.$$

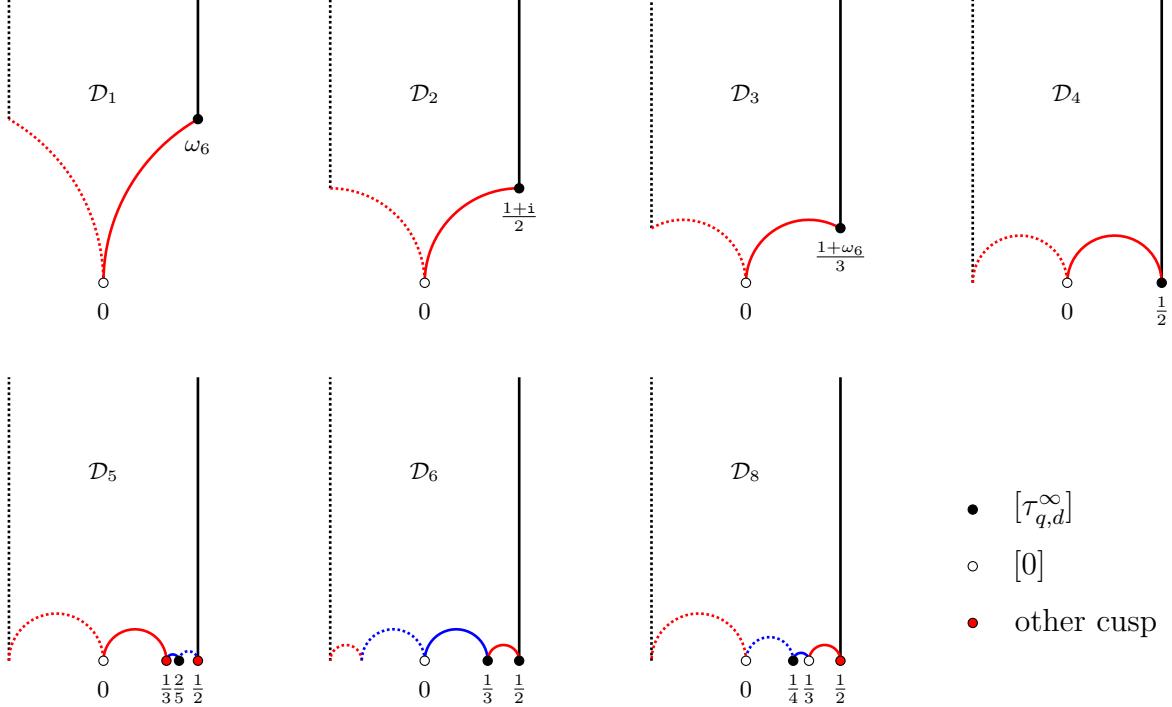


FIGURE 1. Fundamental domain \mathcal{D}_d of $X(\Gamma_d)$ with cusps and elliptic points.

Together with the differential equation (3.4), this shows that

$$\theta \log u_d = \frac{\theta u_d}{u_d} = \frac{W_1}{W_2} = \frac{-f_d \cdot (\theta^2 \tau_q \cdot f_d + 2\theta \tau_q \cdot \theta f_d)}{\theta \tau_q \cdot f_d^2} = -\frac{\theta^2 \tau_q}{\theta \tau_q} - 2 \frac{\theta f_d}{f_d} = -\theta \log(\theta \tau_q \cdot f_d^2).$$

Hence, the quantity

$$u_d f_d^2 \cdot \frac{d \log q}{d \log P^\ell} = 2\pi i \cdot u_d f_d^2 \cdot \theta \tau_q$$

is constant. Since $P^\ell = -q + O(q^2)$ near $q = 0$ and $u_d f_d^2 = 1 + O(P^\ell)$, it follows that

$$\frac{d \log P^\ell}{d \log q} = u_d f_d^2. \quad (5.6)$$

This implies

$$\frac{d \log Q^\ell}{d \log q} = \frac{d \log P^\ell}{d \log q} \cdot \frac{d \log Q^\ell}{d \log P^\ell} = u_d f_d^3.$$

Also, we have $Q^\ell = -q + O(q^2)$.

Note that, by Proposition 3.16 (v), $u_d f_d^3$ is equal to the ratio

$$\frac{d}{\langle E, E, E \rangle^{Y_d}} = \frac{E^3}{\sum_{m \geq 0} \langle E, E, E \rangle_{ml}^{Y_d} Q^{ml}} =: \mathbb{E}_d.$$

We are therefore led to the following definition:

Definition 5.7. For a Type II extremal transition $Y \searrow X$ with exceptional divisor E and exceptional curve ℓ , we define its *extremal function* to be

$$\mathbb{E}^{Y \searrow X} = \frac{E^3}{\sum_{m \geq 0} \langle E, E, E \rangle_{m\ell}^Y Q^{m\ell}}.$$

Then the *canonical modular coordinate* q of the transition is defined by

$$\frac{d \log q}{d \log Q^\ell} = \frac{1}{\mathbb{E}^{Y \searrow X}}, \quad q = -Q^\ell + O(Q^{2\ell}). \quad (5.7)$$

Proposition 5.8. If $Y \searrow X$ is a Type II extremal transition of degree d , then

$$\mathbb{E}^{Y \searrow X} = \mathbb{E}_d.$$

Proof. Let E be the exceptional divisor appearing in $Y \searrow X$. By the degeneration formula (4.5), we have

$$\sum_{m \geq 0} \langle E, E, E \rangle_{m\ell}^Y Q^{m\ell} = \sum_{m \geq m' \geq 0} \sum_{\eta} C_\eta \langle \emptyset \mid e_I, \mu \rangle_{\Gamma_0}^{\bullet(Y, E)} Q^{m'\ell} \langle E, E, E \mid e^I, \mu \rangle_{\Gamma_\infty}^{\bullet(Y_{\text{loc}}, H)} Q^{(m-m')\ell}.$$

Since $|\mu| = (E, m'\ell) = -m' \leq 0$, we must have $m' = 0$, which reduces the sum to

$$\sum_{m \geq 0} \langle E, E, E \mid \emptyset \rangle_{m\ell}^{(Y_{\text{loc}}, H)} Q^{m\ell}.$$

This quantity is independent of the choice of Y , and therefore equals $\sum_{m \geq 0} \langle E, E, E \rangle_{m\ell}^{Y_{\text{loc}}} Q^{m\ell}$. We thus conclude that $\mathbb{E}^{Y \searrow X} = \mathbb{E}^{Y_{\text{loc}} \searrow X_{\text{loc}}} = \mathbb{E}_d$. \square

Since u_d is a weight 0 meromorphic modular function and f_d is a weight 1 modular form (with root and monodromy if $d \leq 2$), we see that $\mathbb{E}_d = u_d f_d^3$ is a weight 3 modular form.

It follows from (5.6) that

$$v_d = \frac{\theta f_d}{f_d} = \frac{1}{f_d} \cdot \frac{d \log q}{d \log P^\ell} \cdot \frac{df_d}{d \log q} = \frac{D_q f_d}{u_d f_d^3}$$

where $D_q = \frac{d}{d \log q} = q \frac{d}{dq}$. Since f_d is a modular form of weight 1, $D_q f_d$ is a quasi-modular form of weight 3 and depth 1. Therefore, v_d is a quasi-modular form of weight 0 and depth 1 (with root and monodromy if $d = 1$). This implies that every element in $S_d = \mathbb{C}(y)[v_d]$ is a quasi-modular form of weight 0.

For the cases $d \in \{4, 5, 6\text{I}, 6\text{II}, 8\}$, the transition point $[\tau_{q,d}^\infty]$ is a cusp. Take a representative $r \in [\tau_{q,d}^\infty]$ and an element $A = \begin{pmatrix} a_A & b_A \\ c_A & d_A \end{pmatrix} \in \text{SL}(2, \mathbb{Z})$ such that $A \cdot i\infty = r$. Then, since

$$f_d|_1^A (A^{-1} \tau_q) = (a_A - c_A \tau_q) f_d(\tau_q) \quad \text{and} \quad f_d|_1^A (\tau_q) \in \mathbb{C}[[q]],$$

the holomorphic solution $f_d^{\text{reg}}(\tau_q)$ of (3.4) near $y = 0$ is proportional to $(a_A - c_A \tau_q) f_d(\tau_q)$. Hence,

$$v_d^{\text{reg}} = \frac{\theta f_d^{\text{reg}}}{f_d^{\text{reg}}} = \frac{1}{u_d f_d^2} \cdot \frac{D_q((a_A - c_A \tau_q) f_d(\tau_q))}{(a_A - c_A \tau_q) f_d(\tau_q)} = \frac{D_q f_d}{u_d f_d^3} - \frac{1}{2\pi i} \frac{c_A}{(a_A - c_A \tau_q) u_d f_d^2}$$

and therefore,

$$v_d^{\text{reg}}(A\tau_q) = \frac{D_q f_d|_3^A(\tau_q) - \frac{1}{2\pi i} \mathbf{k}(A, \tau_q) f_d}{u_d|_0^A(\tau_q) f_d|_1^A(\tau_q)^3} = \left(\frac{D_q f_d}{u_d f_d^3} \right)_0^A(\tau_q).$$

From this, we obtain the modular version of the main result:

Theorem 5.9. *Let $Y \searrow X$ be a Type II transition of degree $d \neq 7$ and let $\bar{\beta} \neq 0$. Let $\tilde{\beta}$ be a lifting of $\bar{\beta}$ such that $(E, \tilde{\beta}) = 0$. Then, in the canonical modular coordinate q of $Y \searrow X$,*

$$\langle \phi^* \vec{b} \rangle_{\tilde{\beta}}^{Y \searrow X} Q^{-\tilde{\beta}} = \sum_{\beta \rightarrow \bar{\beta}} \langle \phi^* b \rangle_{\beta}^Y Q^{\beta - \tilde{\beta}} = \sum_m \langle \phi^* b \rangle_{\tilde{\beta} + m\ell} Q^{m\ell}$$

is a meromorphic quasi-modular form of weight 0. Moreover, its value at the transition point $[\tau_{q,d}^\infty]$ coincides with the invariant $\langle b \rangle_{\tilde{\beta}}^X$.

Remark 5.10. Recall that for all $\vec{b} \in H(X)^{\oplus n}$, the coefficient of $\langle \phi^* \vec{b} \rangle_{\tilde{\beta}}^{Y \searrow X}$ lies in $S_d = \mathbb{C}(P^\ell)[v_d]$. Let \mathcal{D} be the maximal domain of the invariants. Then, for $d \in \{2, 3, 4, 5, 6, 8\}$, since v_d has to be defined on \mathcal{D} and v_d has logarithmic singularities at each rational point $\tau_q \in \mathbb{Q}$, we see that $\mathbb{H} \cup \{\text{i}\infty\}$ is a cover of \mathcal{D} . Since P^ℓ is well-defined on \mathcal{D} , we have the following diagram:

$$\begin{array}{ccc} \mathbb{H} \cup \{\text{i}\infty\} & \xrightarrow{\varphi} & \mathcal{D} \xrightarrow{\psi} X(\Gamma_d) \\ & & \searrow P^\ell \quad \downarrow l \\ & & \mathbb{P}^1. \end{array}$$

If $\varphi(\tau'_q) = \varphi(\tau''_q)$ for some $\tau'_q, \tau''_q \in \mathbb{H}$, then $\psi(\varphi(\tau'_q)) = \psi(\varphi(\tau''_q))$ implies that $\tau''_q = A\tau'_q$ for some $A \in \Gamma_d$. Since v_d is defined on \mathcal{D} , we get $v_d(A\tau_q) = v_d(\tau_q)$ near $\tau_q = \tau'_q$, and hence $v_d(A\tau_q) = v_d(\tau_q)$ for all $\tau_q \in \mathbb{H}$, as they are both holomorphic on \mathbb{H} . Using

$$v_d(A\tau_q) = v_d(\tau_q) + \frac{\mathbf{k}(A, \tau_q)}{2\pi i} \frac{1}{u_d f_d^2},$$

we must have $\mathbf{k}(A, \tau_q) = 0$, i.e., $A = \begin{pmatrix} 1 & b_A \\ 0 & 1 \end{pmatrix}$ for some b_A . This means that the maximal domain

$$\mathcal{D} = \left\{ \begin{pmatrix} 1 & b_A \\ 0 & 1 \end{pmatrix} \right\} \setminus (\mathbb{H} \cup \{\text{i}\infty\})$$

is exactly the unit disk $\mathbb{D} = \{q \mid |q| < 1\}$.

5.3. The wild case: periods of Eisenstein series. Let χ and ψ be Dirichlet characters modulo N_χ and N_ψ , respectively, so that $\chi\psi(-1) = -1$. Given a positive integer n , we denote by $E_k^{\chi, \psi, n}(\tau_q)$ the normalized Eisenstein series of weight $k \geq 3$ with parameter (χ, ψ, n) , which has the following Fourier expansion:

$$E_k^{\chi, \psi, n}(\tau_q) = \frac{1}{2} \delta_{\mathbf{1}_1, \chi} L(1-k, \psi) + \sum_{m \geq 1} \psi(m) m^{k-1} \sum_{\nu=1}^{N_\chi} \frac{\chi(\nu) q^{\nu mn}}{1 - q^{N_\chi mn}}, \quad (5.8)$$

where $\delta_{\mathbf{1}_1, \chi}$ is 1 if $\chi = \mathbf{1}_1$ and is 0 otherwise.

Proposition 5.11. *For $d \in \{3, 4, 5, 6\}$, the weight 3 modular form \mathbb{E}_d , as a function of τ_q , is the Eisenstein series*

$$E_{i\infty,3}^{\Gamma_1(d)}(\tau_q) := \sum_{(a,c)\equiv(1,0) \text{ mod } d} \frac{1}{(-c\tau_q + a)^3}$$

of weight 3 associate with the cusp $i\infty$ of the group $\Gamma_1(d)$. In particular, for $d \in \{3, 4, 5, 6, 8\}$, \mathbb{E}_d can be written in terms of the normalized Eisenstein series:

- (i) $\mathbb{E}_3 = -9 E_3^{1_1, \chi_{3,2}, 1};$
- (ii) $\mathbb{E}_4 = -4 E_3^{1_1, \chi_{4,2}, 1};$
- (iii) $\mathbb{E}_5 = (-1 + \frac{i}{2}) E_3^{1_1, \chi_{5,2}, 1} + (-1 - \frac{i}{2}) E_3^{1_1, \chi_{5,3}, 1};$
- (iv) $\mathbb{E}_6 = -E_3^{1_1, \chi_{3,2}, 1} - 8 E_3^{1_1, \chi_{3,2}, 2};$
- (v) $\mathbb{E}_8 = -4 E_3^{1_1, \chi_{4,2}, 2},$

where $\chi_{3,2} = \begin{pmatrix} -3 \\ \cdot \end{pmatrix}$ and $\chi_{4,2}$ are the unique nontrivial Dirichlet characters modulo 3 and 4, respectively, $\chi_{5,2}$ is the Dirichlet character modulo 5 defined by $\chi_{5,2}(2) = i$, and $\chi_{5,3} = \bar{\chi}_{5,2} = \chi_{5,2}^3$.

Proof. It follows from Table 5 and Table 6 that $\mathbb{E}_d = u_d f_d^3$ vanishes at all cusps except at $[i\infty]$. Since $\mathbb{E}_d = 1 + O(q)$, this means that the difference $\mathbb{E}_d - E_{i\infty,3}^{\Gamma_1(d)}$ is a weight 3 cusp form on $\Gamma_1(d)$.

By the dimension formula for spaces of cusp forms (see, e.g., [9, Theorem 3.6.1]), we have $\dim \mathcal{S}(\Gamma_1(d)) = 0$ for $d \in \{3, 4, 5, 6\}$. Hence, the difference must be zero, i.e., $\mathbb{E}_d = E_{i\infty,3}^{\Gamma_1(d)}$.

Using the explicit formula for Eisenstein series given in [9, Equation (4.6), Theorem 4.2.3], we obtain the expressions of $\mathbb{E}_d = E_{i\infty,3}^{\Gamma_1(d)}$ in (i) - (iv). For (v), we compute

$$\mathbb{E}_8(\tau_q) = u_8(\tau_q) f_8(\tau_q)^3 = u_4(2\tau_q) f_4(2\tau_q)^3 = \mathbb{E}_4(2\tau_q) = -4 E_3^{1_1, \chi_{4,2}, 2}. \quad \square$$

Any path $p: [0, 1] \rightarrow \mathbb{P}^1$ can be lifted to a path $\tilde{p}: [0, 1] \rightarrow \mathbb{H}^*$ via the projection $\mathbb{H}^* \rightarrow X(\Gamma_d) \cong \mathbb{P}^1$. Suppose $p(0) = 0, p(1) = \infty$; then $\tilde{p}(0) = i\infty$ and $\tilde{p}(1)$ will be a rational number $r \in \mathbb{Q}$ that represents the transition point $[\tau_{q,d}^\infty]$. It follows from the definition of q in (5.7) that, along the path \tilde{p} , we have

$$Q^\ell(s) = -e^{2\pi i \tilde{p}(s)} \exp \left(\int_{i\infty}^{\tilde{p}(s)} (\mathbb{E}_d(\tau_q) - 1) d(2\pi i \tau_q) \right).$$

Therefore, the limit of $Q^\ell(s)$ along the path p is

$$Q_p^\ell = -e^{2\pi i r} \exp \left(\int_{i\infty}^r (\mathbb{E}_d(\tau_q) - 1) d(2\pi i \tau_q) \right). \quad (5.9)$$

We denote (5.9) simply by Q_r^ℓ since the limit is independent of the choice of the path p .

Consider the translation $\tau_q \mapsto \tau_q + r$, we have

$$\int_{i\infty}^r (\mathbb{E}_d(\tau_q) - 1) d(2\pi i \tau_q) = \int_{i\infty}^0 (\mathbb{E}_d(\tau_q + r) - 1) d(2\pi i \tau_q).$$

By Proposition 5.11, when $d = 5, 6\mathbb{I}, 6\mathbb{II}, 8$, \mathbb{E}_d can be written as a linear combination of Eisenstein series $\sum_{\psi,n} \alpha_{\psi,n} E_3^{1_1,\psi,n}$. So

$$\int_{i\infty}^0 (\mathbb{E}_d(\tau_q + r) - 1) d(2\pi i \tau_q) = \int_{i\infty}^0 \left(\sum \alpha_{\psi,n} E_3^{1,\psi,n}(\tau_q + r) - 1 \right) d(2\pi i \tau_q)$$

The following proposition expresses the translated Eisenstein series $E_3^{1,\psi,n}(\tau_q + r)$ explicitly as a linear combination of Eisenstein series.

Proposition 5.12 (Translation formula). *Let $r = \frac{a}{c}$ be a rational number with $\gcd(a, c) = 1$. Define*

$$E_r^\psi(\tau_q) = E_3^{1_1,\psi,1}(\tau_q + r).$$

Then

$$E_r^\psi = \sum_{c=c_1c_2c_3} \sum_{\chi \in \widehat{(\mathbb{Z}/c_3\mathbb{Z})^\times}} \frac{\chi(a)\mathfrak{g}(\bar{\chi})\psi(c_1)c_1^2}{\varphi(c_3)} \cdot E_3^{\chi,1_{c_2c_3}\chi\psi,c_1c_2},$$

where $\mathfrak{g}(-)$ denotes the Gauss sum of a Dirichlet character and $\varphi(-)$ denotes the Euler's totient function.

Remark 5.13. The case $n > 1$ reduces to the case $n = 1$ by scaling: for any integer $n > 1$, we have

$$E_3^{1,\psi,n}(\tau_q + r) = E_3^{1,\psi}(n\tau_q + nr) = E_{nr}^\psi(n\tau_q).$$

Proof. Let $\omega = e^{2\pi ir}$. Define the non-constant part of the translated Eisenstein series by

$$E_{r,+}^\psi := E_r^\psi - a_0(E_r^\psi) = \sum \psi(m)m^2 \frac{\omega^m q^m}{1 - \omega^m q^m}.$$

Expanding the denominator, we write

$$\frac{\omega^m q^m}{1 - \omega^m q^m} = \sum_{n=1}^c \frac{\omega^{nm} q^{nm}}{1 - q^{cm}}.$$

Let $c_1 = \gcd(c, m)$ and write $c = c_1c'$, $m = c_1m'$. Then

$$\sum_{n=1}^c \omega^{nm} q^{nm} = \frac{1 - q^{cm}}{1 - q^{c'm}} \sum_{n=1}^{c'} e^{2\pi i anm'/c'} q^{nm}.$$

Next, for each n , let $c_2 = \gcd(n, c')$ and write $c' = c_2c_3$, $n = c_2n'$. Applying orthogonality of characters on $(\mathbb{Z}/c_3\mathbb{Z})^\times$, we get

$$\begin{aligned} \sum_{n=1}^{c'} e^{2\pi i anm'/c'} q^{nm} &= \sum_{c'=c_2c_3} \sum_{n' \perp c_3} e^{2\pi i an'm'/c_3} q^{c_1c_2n'm'} \\ &= \sum_{c'=de} \sum_{n' \perp c_3} \left(\frac{1}{\varphi(c_3)} \sum_{\chi \in \widehat{(\mathbb{Z}/c_3\mathbb{Z})^\times}} \chi(a)\chi(m')\mathfrak{g}(\bar{\chi}) \cdot \chi(n') \right) q^{c_1c_2n'm'}. \end{aligned}$$

Here we use the notation $n_1 \perp n_2$ to mean that n_1 and n_2 are coprime.

Putting everything together:

$$\begin{aligned}
E_{r,+}^\psi &= \sum \psi(m) m^2 \sum_{n=1}^c \frac{\omega^{nm} q^{nm}}{1 - q^{cm}} \\
&= \sum_{c=c_1 c_2 c_3} \sum_{m' \perp c_2 c_3} \sum_{n' \perp c_3} \sum_{\chi \in \widehat{(\mathbb{Z}/c_3 \mathbb{Z})^\times}} \frac{\chi(a) \mathfrak{g}(\bar{\chi}) \psi(c_1) c_1^2}{\varphi(c_3)} \cdot \psi(m') \chi(m') (m')^2 \cdot \frac{\chi(n') q^{c_1 c_2 n' m'}}{1 - q^{c_1 c_2 c_3 m'}} \\
&= \sum_{c=c_1 c_2 c_3} \sum_{\chi \in \widehat{(\mathbb{Z}/c_3 \mathbb{Z})^\times}} \frac{\chi(a) \mathfrak{g}(\bar{\chi}) \psi(c_1) c_1^2}{\varphi(c_3)} \cdot \sum_{m'=1}^{\infty} (\mathbf{1}_{c_2 c_3} \chi \psi)(m') (m')^2 \sum_{n'=1}^{c_3} \frac{\chi(n') (q^{c_1 c_2})^{n' m'}}{1 - (q^{c_1 c_2})^{c_3 m'}} \\
&= \sum_{c=c_1 c_2 c_3} \sum_{\chi \in \widehat{(\mathbb{Z}/c_3 \mathbb{Z})^\times}} \frac{\chi(a) \mathfrak{g}(\bar{\chi}) \psi(c_1) c_1^2}{\varphi(c_3)} \cdot E_{3,+}^{\chi, \mathbf{1}_{c_2 c_3} \chi \psi, c_1 c_2},
\end{aligned}$$

where $E_{3,+}^{\chi, \mathbf{1}_{c_2 c_3} \chi \psi, c_1 c_2}$ denotes the non-constant part of $E_{3,+}^{\chi, \mathbf{1}_{c_2 c_3} \chi \psi, c_1 c_2}$. Since both sides are non-constant parts of weight 3 modular forms, they must have the same constant term. This completes the proof. \square

Using the translation formula, we can write

$$E_d(\tau_q + r) = \sum \alpha_{\psi, n} E_3^{\mathbf{1}_1, \psi, n}(\tau_q + r) = \sum \alpha'_{\chi', \psi', n'} E_3^{\chi', \psi', n'}(\tau_q)$$

for some coefficients $\alpha'_{\chi', \psi', n'} \in \mathbb{C}$. It remains to compute the contribution of each $E_3^{\chi', \psi', n'}$ to the integral.

Proposition 5.14 ([1], extended by [52]). *Let $k \geq 2$ be an integer. For any triple (χ, ψ, n) with $\chi \psi(-1) = (-1)^k$, define*

$$I_k^{\chi, \psi, n} = n^{-1} \lim_{s \rightarrow 0} L(s+1, \chi) L(s+2-k, \psi).$$

Let $a_{i\infty, 0}$ and $a_{0, 0}$ denote the cusp values of $E_k^{\chi, \psi, n}(\tau_q)$ at $\tau_q = i\infty$ and $\tau_q = 0$, respectively. Then the integral

$$\int_{i\infty}^0 (E_k^{\chi, \psi, n}(\tau_q) - a_{i\infty, 0} - a_{0, 0} \tau_q^{-k}) d(2\pi i \tau_q) = I_k^{\chi, \psi, n}.$$

In particular, when χ is even and nontrivial, i.e., $\chi(-1) = 1$ but $\chi \neq \mathbf{1}_N$ for all N , the integral is equal to 0.

Proof. This identity was originally established by [1] in the special case $\chi = \mathbf{1}_1$, i.e., $a_{0, 0} = 0$. To handle the general case, we follow the method in [52].

Let $x = -\log q = -2\pi i \tau_q$. Define $E_{k,+}^{\chi,\psi,1}(\tau_q) = E_k^{\chi,\psi,1}(\tau_q) - a_{i\infty,0}$. Then, by the definition of Eisenstein series (5.8) and the Mellin transform on exponential function, we have

$$\begin{aligned} E_{k,+}^{\chi,\psi,1}(\tau_q) &= \sum_{m \geq 1} \psi(m) m^{k-1} \sum_{\nu \geq 1} \chi(\nu) e^{-m\nu x} \\ &= \sum_{m \geq 1} \psi(m) m^{k-1} \sum_{\nu \geq 1} \chi(\nu) \cdot \frac{1}{2\pi i} \int_{\Lambda-i\infty}^{\Lambda+i\infty} \Gamma(s) (m\nu x)^{-s} ds \\ &= \frac{1}{2\pi i} \int_{\Lambda-i\infty}^{\Lambda+i\infty} \sum_{m \geq 1} \frac{\psi(m) m^{k-1}}{m^s} \cdot \sum_{\nu \geq 1} \frac{\chi(\nu)}{\nu^s} \cdot \Gamma(s) x^{-s} ds \\ &= \frac{1}{2\pi i} \int_{\Lambda-i\infty}^{\Lambda+i\infty} L(s+1-k, \psi) L(s, \chi) \Gamma(s) x^{-s} ds, \end{aligned}$$

where $\Lambda > 0$ is a sufficiently large number. Thus, for $0 < \varepsilon < R < \infty$,

$$\begin{aligned} &n \int_{nRi/2\pi}^{n\varepsilon i/2\pi} E_{k,+}^{\chi,\psi,n}(\tau_q) d(2\pi i \tau_q) \\ &= \int_{Ri/2\pi}^{\varepsilon i/2\pi} E_{k,+}^{\chi,\psi,1}(\tau_q) d(2\pi i \tau_q) \\ &= \int_{\varepsilon}^R E_{k,+}^{\chi,\psi,1} dx \\ &= \frac{1}{2\pi i} \int_{\Lambda-i\infty}^{\Lambda+i\infty} L(s+1-k, \psi) L(s, \chi) \Gamma(s) \frac{\varepsilon^{1-s} - R^{1-s}}{s-1} ds \\ &= \frac{1}{2\pi i} \int_{\Lambda+1-i\infty}^{\Lambda+1+i\infty} L(s+2-k, \psi) L(s+1, \chi) \Gamma(s) (\varepsilon^{-s} - R^{-s}) ds. \end{aligned} \tag{5.10}$$

As $R \rightarrow \infty$, the R^{-s} term vanishes, and we are left with

$$\frac{1}{2\pi i} \int_{\Lambda+1-i\infty}^{\Lambda+1+i\infty} L(s+2-k, \psi) L(s+1, \chi) \Gamma(s) \varepsilon^{-s} ds.$$

To evaluate this integral, we shift the contour of integration left ward to $\operatorname{Re} s = -\delta$ for a small $\delta > 0$. In the area $-\delta < \operatorname{Re} s < \Lambda + 1$, there are two possible poles: if $\psi = \mathbf{1}_N$, then at $s = k - 1$,

$$\operatorname{Res}_{s=k-1} L(s+2-k, \psi) L(s+1, \chi) \Gamma(s) \varepsilon^{-s} = a \varepsilon^{-(k-1)}$$

for some constant a only depends on χ, ψ and k ; at $s = 0$,

$$\operatorname{Res}_{s=0} L(s+2-k, \psi) L(s+1, \chi) \Gamma(s) \varepsilon^{-s} = \lim_{s \rightarrow 0} L(s+2-k, \psi) L(s+1, \chi) = I_k^{\chi,\psi,1}.$$

By the residue theorem,

$$\begin{aligned} & \frac{1}{2\pi i} \int_{\Lambda+1-i\infty}^{\Lambda+1+i\infty} L(s+2-k, \psi) L(s+1, \chi) \Gamma(s) \varepsilon^{-s} ds \\ &= \frac{1}{2\pi i} \int_{-\delta-i\infty}^{-\delta+i\infty} L(s+2-k, \psi) L(s+1, \chi) \Gamma(s) \varepsilon^{-s} ds + a\varepsilon^{-(k-1)} + I_k^{\chi, \psi, 1}. \end{aligned}$$

As $\varepsilon \rightarrow 0$, the integral over the shifted path tends to 0. Hence, by (5.10),

$$\lim_{\varepsilon \rightarrow 0} \int_{i\infty}^{n\varepsilon i/2\pi} E_{k,+}^{\chi, \psi, n}(\tau_q) d(2\pi i \tau_q) - \frac{a}{n} \varepsilon^{-(k-1)} - I_k^{\chi, \psi, n} = 0.$$

Now observe that

$$\int_{i\infty}^{n\varepsilon i/2\pi} \tau_q^{-k} d(2\pi i \tau_q) = c_0 \varepsilon^{-(k-1)}$$

for some constant c_0 , and since the limit

$$\lim_{\varepsilon \rightarrow 0} \int_{i\infty}^{n\varepsilon i/2\pi} (E_{k,+}^{\chi, \psi, n}(\tau_q) - a_{0,0} \tau^{-k}) d(2\pi i \tau_q)$$

exists and is finite, we must have $\frac{a}{n} = a_{0,0} c_0$. Therefore,

$$\int_{i\infty}^{n\varepsilon i/2\pi} (E_k^{\chi, \psi, n}(\tau_q) - a_{i\infty, 0} - a_{0,0} \tau^{-k}) d(2\pi i \tau_q) = I_k^{\chi, \psi, n},$$

as claimed.

For the last statement, simply note that, in this case, $L(s, \chi)$ is bounded near $s = 1$ and $\psi(-1) = (-1)^k$, so

$$\lim_{s \rightarrow 0} L(s+1, \chi) L(s+2-k, \psi) = L(1, \chi) L(2-k, \psi) = 0. \quad \square$$

Let $a_{i\infty, 0}(E_3^{\chi', \psi', n'})$ and $a_{0,0}(E_3^{\chi', \psi', n'})$ denote the cusp values of $E_3^{\chi', \psi', n'}(\tau_q)$ at $i\infty$ and 0, respectively. Since $\mathbb{E}_d(\tau_q) = 1 + O(q)$ and $\mathbb{E}_d(\tau_q)$ vanishes at $r \in [\tau_{q,d}^\infty]$, we have

$$\sum \alpha'_{\chi', \psi', n'} a_{i\infty, 0}(E_3^{\chi', \psi', n'}) = 1, \quad \sum \alpha'_{\chi', \psi', n'} a_{0,0}(E_3^{\chi', \psi', n'}) = 0.$$

Therefore,

$$\begin{aligned} & \int_{i\infty}^0 \left(\sum \alpha'_{\chi', \psi', n'} E_3^{\chi', \psi', n'} - 1 \right) d(2\pi i \tau_q) \\ &= \int_{i\infty}^0 \sum \alpha'_{\chi', \psi', n'} \left(E_3^{\chi', \psi', n'} - a_{i\infty, 0}(E_3^{\chi', \psi', n'}) - a_{0,0}(E_3^{\chi', \psi', n'}) \right) d(2\pi i \tau_q) \\ &= \sum \alpha'_{\chi', \psi', n'} I_k^{\chi', \psi', n'}. \end{aligned}$$

This means that the limit can be computed explicitly in terms of special values of Dirichlet L-functions.

Example 5.15. When $d = 6$, we have

$$\mathbb{E}_6 = -E_3^{\mathbf{1}_1, \chi_{3,2}, 1} - 8 E_3^{\mathbf{1}_1, \chi_{3,2}, 2}.$$

To simplify notation, we will write $\chi = \chi_{3,2}$ throughout this example. Take $r_I = \frac{1}{3}$ that represents the transition point $[\tau_{q,6I}^\infty]$. Then, by Proposition 5.12, we have

$$\begin{aligned} E_3^{\mathbf{1}_1, \chi, 1}(\tau_q + r_I) &= E_{r_I}^\chi(\tau_q) = E^{\mathbf{1}_1, \chi, 3} - \frac{1}{2} E^{\mathbf{1}_3, \chi, 1} + \frac{\sqrt{3}i}{2} E^{\chi, \mathbf{1}_3, 1}. \\ E_3^{\mathbf{1}_1, \chi, 2}(\tau_q + r_I) &= E_{2r_I}^\chi(2\tau_q) = E^{\mathbf{1}_1, \chi, 6} - \frac{1}{2} E^{\mathbf{1}_3, \chi, 2} - \frac{\sqrt{3}i}{2} E^{\chi, \mathbf{1}_3, 2}. \end{aligned}$$

So by Proposition 5.14,

$$\begin{aligned} \int_{i\infty}^0 (\mathbb{E}_d(\tau_q + r_I) - 1) d(2\pi i \tau_q) &= -\left(I_3^{\mathbf{1}_1, \chi, 3} - \frac{1}{2} I_3^{\mathbf{1}_3, \chi, 1} + \frac{\sqrt{3}i}{2} I_3^{\chi, \mathbf{1}_3, 1}\right) \\ &\quad - 8\left(I_3^{\mathbf{1}_1, \chi, 6} - \frac{1}{2} I_3^{\mathbf{1}_3, \chi, 2} - \frac{\sqrt{3}i}{2} I_3^{\chi, \mathbf{1}_3, 2}\right). \end{aligned}$$

Since

$$\begin{aligned} I_3^{\mathbf{1}_1, \chi, 3} &= \frac{1}{3} I_3^{\mathbf{1}_1, \chi, 1}, \\ I_3^{\mathbf{1}_3, \chi, 1} &= \lim_{s \rightarrow 0} L(s+1, \mathbf{1}_3) L(s-1, \chi) \\ &= \lim_{s \rightarrow 0} (1 - 3^{-(s+1)}) L(s+1, \mathbf{1}_1) L(s-1, \chi) = \frac{2}{3} I_3^{\mathbf{1}_1, \chi, 1}, \end{aligned}$$

it follows that $I_3^{\mathbf{1}_1, \chi, 3} = \frac{1}{2} I_3^{\mathbf{1}_3, \chi, 1}$. Similarly, $I_3^{\mathbf{1}_1, \chi, 6} = \frac{1}{2} I_3^{\mathbf{1}_3, \chi, 2}$. Thus, it remains to compute

$$-\frac{\sqrt{3}i}{2} I_3^{\chi, \mathbf{1}_3, 1} + 8 \frac{\sqrt{3}i}{2} I_3^{\chi, \mathbf{1}_3, 2} = \frac{3\sqrt{3}i}{2} L(1, \chi) L(-1, \mathbf{1}_3),$$

which evaluates to $\frac{3\sqrt{3}i}{2} \cdot \frac{\pi}{3\sqrt{3}} \cdot \frac{1}{6} = \frac{\pi i}{12}$. Hence, the resulting limit is

$$Q_{r_I}^\ell = -e^{2\pi i \cdot r_I} \cdot e^{\frac{\pi i}{12}} = e^{-\frac{1}{8} \cdot 2\pi i},$$

which is an 8-th root of unity.

Now take $r_{\mathbb{II}} = \frac{1}{2}$ which represents another transition point $[\tau_{q,6\mathbb{II}}^\infty]$. Then, by Proposition 5.12,

$$\begin{aligned} E_3^{\mathbf{1}_1, \chi, 1}(\tau_q + r_{\mathbb{II}}) &= E_{r_{\mathbb{II}}}^\chi(\tau_q) = -4 E_3^{\mathbf{1}_1, \chi, 2} + E_3^{\mathbf{1}_1, \mathbf{1}_2 \chi, 2} - E_3^{\mathbf{1}_2, \mathbf{1}_2 \chi, 1}. \\ E_3^{\mathbf{1}_1, \chi, 2}(\tau_q + r_{\mathbb{II}}) &= E_{2r_{\mathbb{II}}}^\chi(2\tau_q) = E_3^{\mathbf{1}_1, \chi, 2}. \end{aligned}$$

Thus, by Proposition 5.14,

$$\begin{aligned} \int_{i\infty}^0 (\mathbb{E}_d(\tau_q + r_{\mathbb{II}}) - 1) d(2\pi i \tau_q) &= -(-4 I_3^{\mathbf{1}_1, \chi, 2} + I_3^{\mathbf{1}_1, \mathbf{1}_2 \chi, 2} - I_3^{\mathbf{1}_2, \mathbf{1}_2 \chi, 1}) - 8 I_3^{\mathbf{1}_1, \chi, 2} \\ &= -4 I_3^{\mathbf{1}_1, \chi, 2} - I_3^{\mathbf{1}_1, \mathbf{1}_2 \chi, 2} + I_3^{\mathbf{1}_2, \mathbf{1}_2 \chi, 1}. \end{aligned}$$

Similar to the case 6I, we have

$$\begin{aligned}
I_3^{1_1, \chi, 2} &= \frac{1}{2} I_3^{1_1, \chi, 1}, \\
I_3^{1_1, 1_2 \chi, 2} &= \frac{1}{2} \lim_{s \rightarrow 0} L(s+1, \mathbf{1}_1) L(s-1, \mathbf{1}_2 \chi) \\
&= \frac{1}{2} \lim_{s \rightarrow 0} (1 - \chi(2) 2^{-(s-1)}) L(s+1, \mathbf{1}_1) L(s-1, \chi) = \frac{3}{2} I_3^{1_1, \chi, 1}, \\
I_3^{1_2, 1_2 \chi, 1} &= \lim_{s \rightarrow 0} L(s+1, \mathbf{1}_2) L(s-1, \mathbf{1}_2 \chi) \\
&= \lim_{s \rightarrow 0} (1 - 2^{-(s+1)}) (1 - \chi(2) 2^{-(s-1)}) L(s+1, \mathbf{1}_1) L(s-1, \chi) = \frac{3}{2} I_3^{1_1, \chi, 1}.
\end{aligned}$$

Finally, since

$$I_3^{1_1, \chi, 1} = \lim_{s \rightarrow 0} L(s+1, \mathbf{1}_1) L(s-1, \chi) = L'(-1, \chi),$$

we conclude that the limit is

$$Q_{\mathbb{II}}^\ell = -e^{2\pi i \cdot r_{\mathbb{II}}} \cdot e^{-2L'(-1, \chi)} = e^{-2L'(-1, \chi)}.$$

Here, using the functional equation for Dirichlet L-series, we have

$$L'(-1, \chi_{3,2}) = \frac{3 \cdot \mathfrak{g}(\chi)}{2 \cdot 2\pi i} L(2, \bar{\chi}) = \frac{3\sqrt{3}}{4\pi} L(2, \chi) \approx 0.3230659\dots \quad (5.11)$$

Example 5.16. When $d = 5$, we have

$$\mathbb{E}_5 = (-1 + \frac{i}{2}) E_3^{1_1, \chi_{5,2}, 1} + (-1 - \frac{i}{2}) E_3^{1_1, \chi_{5,3}, 1}.$$

To simplify notation, let us write $\chi = \chi_{5,2}$ throughout this example. Then, by Proposition 5.12,

$$\begin{aligned}
E_3^{1_1, \chi, 1}(\tau_q + r) &= E_r^\chi(\tau_q) = E_3^{1_1, \chi, 5} - \frac{1}{4} E_3^{1_5, \chi, 1} + \frac{i}{4} \mathfrak{g}(\chi^3) E_3^{\chi, \chi^2, 1} - \frac{1}{4} \mathfrak{g}(\chi^2) E_3^{\chi^2, \chi^3, 1} - \frac{i}{4} \mathfrak{g}(\chi) E_3^{\chi^3, 1_5, 1}, \\
E_3^{1_1, \chi^3, 1}(\tau_q + r) &= E_r^{\chi^3}(\tau_q) = E_3^{1_1, \chi^3, 5} - \frac{1}{4} E_3^{1_5, \chi^3, 1} + \frac{i}{4} \mathfrak{g}(\chi^3) E_3^{\chi, 1_5, 1} - \frac{1}{4} \mathfrak{g}(\chi^2) E_3^{\chi^2, \chi, 1} - \frac{i}{4} \mathfrak{g}(\chi) E_3^{\chi^3, \chi^2, 1}.
\end{aligned}$$

As in the case $d = 6$, $I_3^{1_1, \chi, 5} = \frac{1}{4} I_3^{1_5, \chi, 1}$ and $I_3^{1_1, \chi^3, 5} = \frac{1}{4} I_3^{1_5, \chi^3, 1}$. Since χ^2 is even and nontrivial, we find that $I_3^{\chi^2, \chi, 1} = I_3^{\chi^3, \chi^2, 1} = 0$. We now compute the remaining constants

$$\begin{aligned}
\mathfrak{g}(\chi^3) L(1, \chi) &= -\pi i B_{1, \chi^3} = \frac{\pi(1 + 3i)}{5}, \\
\mathfrak{g}(\chi) L(1, \chi^3) &= -\pi i B_{1, \chi} = \frac{\pi(-1 + 3i)}{5}, \\
L(-1, \mathbf{1}_5) &= -\frac{B_{2, \mathbf{1}_5}}{2} = \frac{1}{3}, \\
L(-1, \chi^2) &= -\frac{B_{2, \chi^2}}{2} = -\frac{2}{5}.
\end{aligned}$$

Therefore,

$$\begin{aligned} \int_{i\infty}^0 (\mathbb{E}_5(\tau_q + r) - 1) d(2\pi i \tau_q) &= \frac{i}{4} \mathfrak{g}(\chi^3) L(1, \chi) \left((-1 + \frac{i}{2}) L(-1, \chi^2) + (-1 - \frac{i}{2}) L(-1, \mathbf{1}_5) \right) \\ &\quad - \frac{i}{4} \mathfrak{g}(\chi) L(1, \chi^3) \left((-1 + \frac{i}{2}) L(-1, \mathbf{1}_5) + (-1 - \frac{i}{2}) L(-1, \chi^2) \right) \\ &= \frac{7\pi i}{60}. \end{aligned}$$

Thus, we obtain

$$Q_r^\ell = -e^{2\pi ir} \cdot e^{7\pi i/60} = e^{-\frac{1}{24} \cdot 2\pi i},$$

which is a 24-th root of unity.

Theorem 5.17. *For $d = 4, 5, 6\mathbb{I}, 6\mathbb{II}, 8$ and $r \in [\tau_{q,d}^\infty]$, we have*

- (i) Q_r^ℓ is a root of unity if $d \neq 6\mathbb{II}$;
- (ii) if $d = 6\mathbb{II}$ and $r = \frac{a}{c}$ with $\gcd(a, c) = 1$ and $c \equiv 2 \pmod{6}$, then

$$Q_r^\ell = \omega \cdot e^{-c L'(-1, \chi_{3,2})}$$

for some root of unity ω . In particular, Q_r^ℓ is not a root of unity (cf. (5.11)).

Proof. We write $\mathbb{E}_d = \sum \alpha_{\psi,n} \mathbb{E}^{\mathbf{1}_1, \psi, n}$ as in Proposition 5.11, and write $r = \frac{a}{c}$ with $\gcd(a, c) = 1$.

For fixed n , write $c = \gcd(c, n) \cdot c_0$, $n = \gcd(c, n) \cdot n_0$. Then $nr = \frac{an_0}{c_0}$ with $\gcd(an_0, c_0) = 1$. By Proposition 5.14, the contribution from $\mathbb{E}_3^{\mathbf{1}_1, \psi, n}(\tau_q + r)$ is equal to

$$\begin{aligned} &\sum_{c_0=c_1c_2c_3} \frac{\mathbf{1}_{c_3}(n_0a)\mathfrak{g}(\mathbf{1}_{c_3})\psi(c_1)c_1^2}{\varphi(c_3)} \cdot (c_1c_2n)^{-1} \lim_{s \rightarrow 0} L(s+1, \mathbf{1}_{c_3}) L(s-1, \mathbf{1}_{c_2c_3}\psi) \\ &+ \sum_{c_0=c_1c_2c_3} \sum_{\substack{\chi \in \widehat{(\mathbb{Z}/c_3\mathbb{Z})^\times} \\ \chi(-1)=-1}} \frac{\chi(n_0a)\mathfrak{g}(\bar{\chi})\psi(c_1)c_1^2}{\varphi(c_3)} \cdot (c_1c_2n)^{-1} L(1, \chi) L(-1, \mathbf{1}_{c_2c_3}\chi\psi). \end{aligned}$$

We denote these two summations by $I_{r,\text{triv}}^{\psi,n}$ and $I_{r,\text{nontriv}}^{\psi,n}$, respectively. Then the limit

$$\begin{aligned} Q_r^\ell &= -\exp \left(\sum \alpha_{\psi,n} (I_{r,\text{triv}}^{\psi,n} + I_{r,\text{nontriv}}^{\psi,n}) \right) \\ &= -\exp \left(\sum \alpha_{\psi,n} I_{r,\text{triv}}^{\psi,n} \right) \exp \left(\sum \alpha_{\psi,n} I_{r,\text{nontriv}}^{\psi,n} \right). \end{aligned} \tag{5.12}$$

For the trivial part,

$$\begin{aligned} I_{r,\text{triv}}^{\psi,n} &= \sum_{c_0=c_1c_2c_3} \frac{\mu(c_3)\psi(c_1)c_1}{\varphi(c_3)c_2n} \cdot \prod_{p|c_3} (1 - p^{-1}) \cdot L'(-1, \mathbf{1}_{c_2c_3}\psi) \\ &= \sum_{c_0=c_1c_2c_3} \frac{\mu(c_3)\psi(c_1)c_1}{c_2c_3n} \cdot L'(-1, \mathbf{1}_{c_2c_3}\psi) \\ &= \sum_{c_1|c_0} \frac{\psi(c_1)c_1^2}{c_0n} L'(-1, \mathbf{1}_{c_0/c_1}\psi) \sum_{c_2c_3=c_0/c_1} \mu(c_3) = \frac{\psi(c_0)c_0}{n} L'(-1, \psi). \end{aligned}$$

Here μ is the Möbius function and we use the fact that $\sum_{m|n} \mu(m) = \delta_{1,n}$.

Thus, we find

$$\sum \alpha_{\psi,n} I_{r,\text{triv}}^{\psi,n} = \sum \alpha_{\psi,n} \frac{\psi(\frac{c}{\gcd(c,n)}) \frac{c}{\gcd(c,n)}}{n} L'(-1, \psi).$$

In particular, one can verify case by case that this sum is 0 when $d = 4, 5, 6I, 8$, and is equal to

$$-\frac{\chi_{3,2}(2)c}{1} L'(-1, \chi_{3,2}) - 8 \frac{\chi_{3,2}(1)\frac{c}{2}}{2} L'(-1, \chi_{3,2}) = -c L'(-1, \chi_{3,2})$$

when $d = 6II$.

By (5.12), it remains to show that $\sum \alpha_{\psi,n} I_{r,\text{nontriv}}^{\psi,n}$ lies in $2\pi i \mathbb{Q}$. We show that $\frac{1}{2\pi i} \sum \alpha_{\psi,n} I_{r,\text{nontriv}}^{\psi,n}$ is σ -invariant for any Galois action $\sigma: \widehat{\mathbb{Q}} \rightarrow \widehat{\mathbb{Q}}$.

Suppose $\chi \in \widehat{(\mathbb{Z}/c_3\mathbb{Z})^\times}$ with $\chi(-1) = -1$, and let its conductor be $N \mid c_3$, so that $\chi = \mathbf{1}_{c_3}\chi_N$ for some primitive character χ_N modulo N . Then the term

$$I_{c_1,c_2,c_3,n_0a}^{\chi,\psi,n} := \frac{1}{2\pi i} \frac{\chi(n_0a)\mathfrak{g}(\bar{\chi})\psi(c_1)c_1^2}{\varphi(c_3)} \cdot (c_1c_2n)^{-1} L(1, \chi) L(-1, \mathbf{1}_{c_2c_3}\chi\psi)$$

can be rewritten using properties of Gauss sums and special L-values as

$$\begin{aligned} & \frac{1}{2\pi i} \frac{\chi(n_0a)(\mathfrak{g}(\bar{\chi}_N)\mu(\frac{c_3}{N})\chi_N(\frac{c_3}{N}))\psi(c_1)c_1}{\varphi(c_3)c_2n} \cdot \prod_{p|c_3} (1 - p^{-1}\chi_N(p)) L(1, \chi_N) \cdot L(-1, \mathbf{1}_{c_2c_3}\chi\psi) \\ &= \frac{1}{2\pi i} \frac{\chi(n_0a)\mathfrak{g}(\bar{\chi}_N)\mu(\frac{c_3}{N})\chi_N(\frac{c_3}{N})\psi(c_1)c_1}{\varphi(c_3)c_2n} \prod_{p|c_3} (1 - p^{-1}\chi_N(p)) \left(\frac{\pi i \mathfrak{g}(\chi_N)B_{1,\bar{\chi}_N}}{N} \right) \left(-\frac{B_{2,\mathbf{1}_{c_2c_3}\chi\psi}}{2} \right) \\ &= -\frac{\chi(n_0a)\mu(\frac{c_3}{N})\chi_N(\frac{c_3}{N})\psi(c_1)c_1}{4\varphi(c_3)c_2n} \prod_{p|c_3} (1 - p^{-1}\chi_N(p)) B_{1,\bar{\chi}_N} B_{2,\mathbf{1}_{c_2c_3}\chi\psi} \end{aligned}$$

where, for a Dirichlet character $\chi' \in \widehat{(\mathbb{Z}/N'\mathbb{Z})^\times}$, $B_{m,\chi'}$ denotes the m -th generalized Bernoulli number associated to χ' , defined by the series expansion

$$\sum_{\nu=1}^{N'} \chi'(\nu) \frac{x e^{\nu x}}{e^{N'x} - 1} = \sum_{m=0}^{\infty} \frac{B_{m,\chi'}}{m!} x^m.$$

Because $\chi^\sigma = \sigma \circ \chi = \mathbf{1}_{c_3}\chi_N^\sigma$, $\sigma(B_{m,\chi'}) = B_{m,(\chi')^\sigma}$, and $\bar{\chi}_N^\sigma = \overline{\chi_N^\sigma}$, it follows that

$$\sigma(\alpha_{\psi,n} I_{c_1,c_2,c_3,n_0a}^{\chi,\psi,n}) = \sigma(\alpha_{\psi,n}) I_{c_1,c_2,c_3,n_0a}^{\chi^\sigma,\psi^\sigma,n} = \alpha_{\psi^\sigma,n} I_{c_1,c_2,c_3,n_0a}^{\chi^\sigma,\psi^\sigma,n}.$$

Hence, the sum $\frac{1}{2\pi i} \sum \alpha_{\psi,n} I_{r,\text{nontriv}}^{\psi,n}$ is σ -invariant, as desired. \square

Remark 5.18. Using the above method, along with induction on the denominator of r and change of coordinates $\tau_q \mapsto A\tau_q$ (as suggested by Y. Yang), we can compute the limit of Q^ℓ at any $r = \frac{a}{c} \notin [i\infty]$. Below are the explicit result in the cases $d = 3, 4, 5, 6$.

(i) For $d = 3$,

$$Q_r^\ell = (-1)^c e^{-9c L'(-1, \chi_{3,2})}, \quad r \in [0], \quad c \equiv 1 \pmod{3}.$$

(ii) For $d = 4$,

$$Q_r^\ell = \begin{cases} -e^{-4cL'(-1,\chi_{4,2})} & \text{if } r \in [0], c \equiv 1 \pmod{4}, \\ 1 & \text{if } r \in [\frac{1}{2}]. \end{cases}$$

(iii) For $d = 5$,

$$Q_r^\ell = \begin{cases} (-1)^c e^{c\operatorname{Re}((-2+1)L'(-1,\chi_{5,2}))} & \text{if } r \in [0], c \equiv 1 \pmod{5}, \\ e^{c\operatorname{Re}((-1-2i)L'(-1,\chi_{5,2}))} & \text{if } r \in [\frac{1}{2}], c \equiv 2 \pmod{5}, \\ \omega_{120}^{-c} & \text{if } r \in [\tau_{q,5}^\infty] = [\frac{2}{5}], c \equiv 0 \pmod{5}, a \equiv 2 \pmod{5}. \end{cases}$$

In particular, for some $r \in [\tau_{q,5}^\infty]$, for example, $r = \frac{7}{120}$, one has $Q_r^\ell = 1$. In other words, there exists a path on \mathbb{P}^1 connecting $P^\ell = 0$ and $P^\ell = \infty$ along which Q^ℓ tends to 1.

(iv) For $d = 6$,

$$Q_r^\ell = \begin{cases} -e^{-5cL'(-1,\chi_{3,2})} & \text{if } r = \frac{a}{c} \in [0], c \equiv 1 \pmod{6}, \\ \omega_{24}^{-c} & \text{if } r = \frac{a}{c} \in [\tau_{q,6I}^\infty] = [\frac{1}{3}], c \equiv 3 \pmod{6}, a \equiv 1 \pmod{3}, \\ e^{-cL'(-1,\chi_{3,2})} & \text{if } r = \frac{a}{c} \in [\tau_{q,6II}^\infty] = [\frac{1}{2}], c \equiv 2 \pmod{6}. \end{cases}$$

This shows that, in Theorem 5.17 (ii), the root of unity ω is actually superfluous.

Here, we plot the images of several straight line segments $[0, \omega_n)$ under the maps P_d and Q^ℓ , for $d = 6I$ and $6II$ (corresponding to Figure 2 and Figure 3, respectively), to illustrate how paths lift under the projection $\mathbb{D}^\times \rightarrow \mathbb{P}_{P^\ell}^1 \setminus \operatorname{Sing} f_d$, and how Q^ℓ tends to Q_r^ℓ , $r = 1/n$. Note that the function $Q^\ell(q)$ is the same in both figures.

Comparing the image curves $Q^\ell([0, \omega_N])$ with $N = 8, 10$ in Figure 3 and $N = 9$ in Figure 2 reveals the complexity of the Kähler moduli associated to the function $Q^\ell(q)$. We plan to study the geometry of the Kähler moduli in details in a subsequent work.

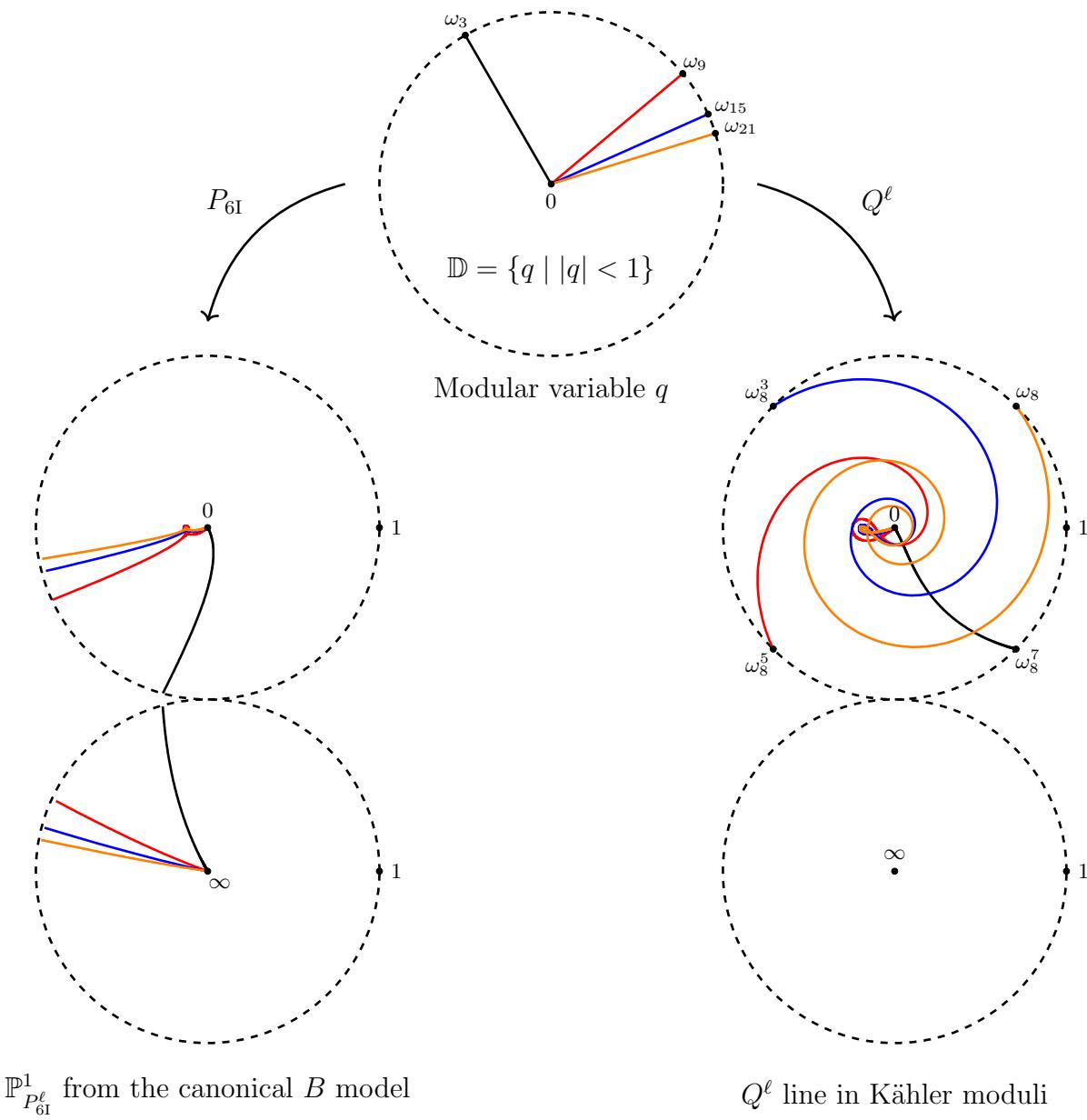


FIGURE 2. $d = 6I$

APPENDIX A. ON DEL PEZZO SURFACES AND THREEFOLDS

We briefly review the classification of del Pezzo manifolds of dimensions two and three. The standard reference is [18].

In classical geometry, a (smooth) *del Pezzo surface* is a smooth projective surface S such that the anti-canonical divisor $-K_S$ is ample. The self-intersection number $d = K_S^2$ is called the *degree* of the del Pezzo surface S . It satisfies $1 \leq d \leq 9$. The description of S is well-known as below, see [35, §24] for example.

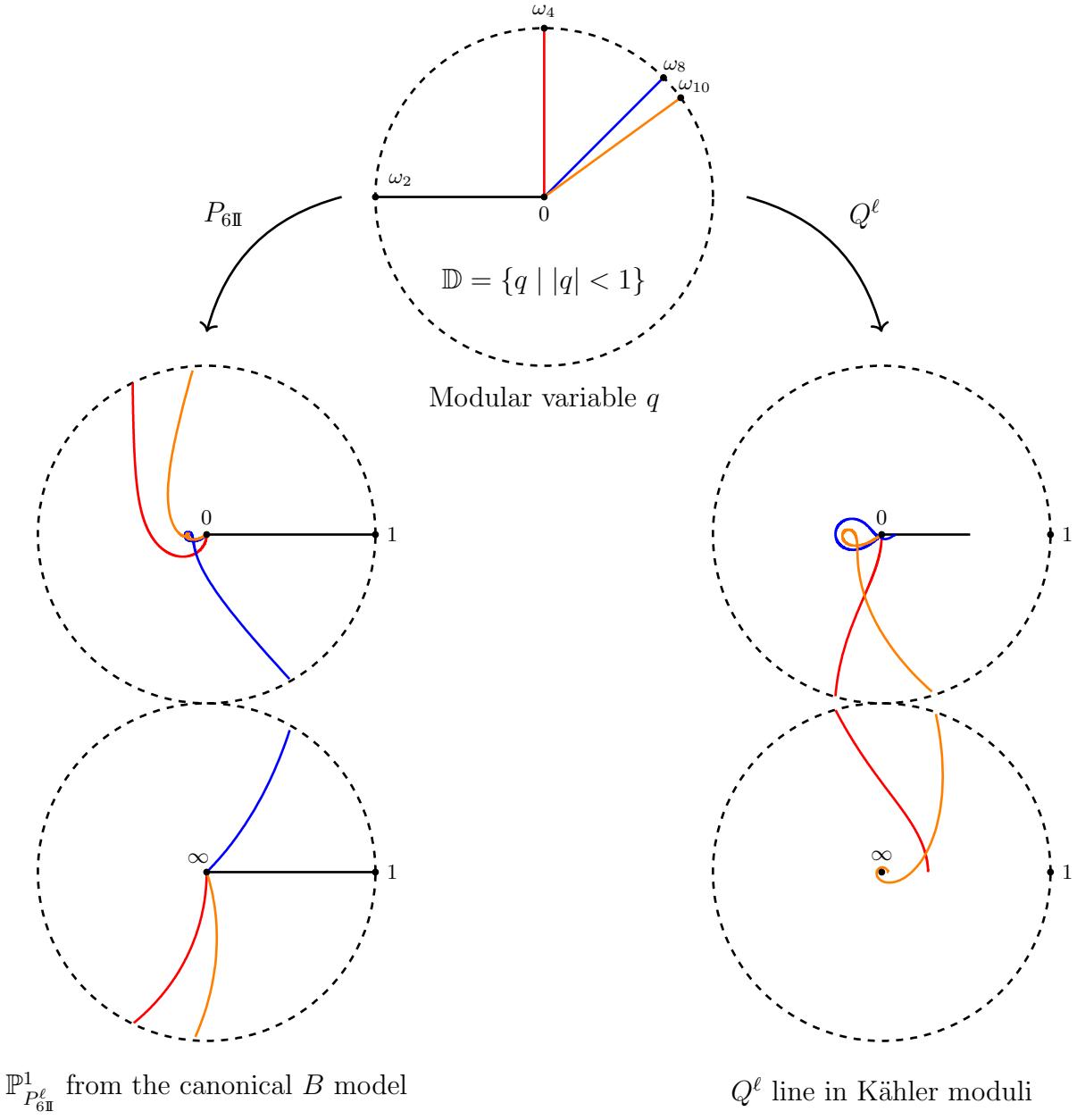


FIGURE 3. $d = 6\mathbb{II}$

Theorem A.1. Let S be a smooth del Pezzo surface of degree d . Set the anti-canonical algebra $R = \bigoplus_{m \geq 0} R_m$ with grading $R_m = H^0(S, \mathcal{O}(-mK))$. Then $S = \text{Proj } R$ is isomorphic to either a blow-up of \mathbb{P}^2 at $9 - d$ general points or $\mathbb{P}^1 \times \mathbb{P}^1$ with $d = 8$, and has Betti numbers

$$b_2(S) = h^{1,1}(S) = 10 - d,$$

$b_0(S) = b_4(S) = 1$ and $b_1(S) = b_3(S) = 0$. Furthermore, we have:

- (i) If $d = 1$, then R is generated by R_1 , R_2 and R_3 and S is a sextic $(6) \subseteq \mathbb{P}(3, 2, 1, 1)$. To be precise, $R \cong \mathbb{C}[x_1, x_2, x_3, x_4]/(F)$ with $x_3, x_4 \in R_1$, $x_2 \in R_2$ and $x_1 \in R_3$ and a weighted homogeneous polynomial $F = F(x_1, x_2, x_3, x_4)$ of weighted degree six.
- (ii) If $d = 2$, then R is generated by R_1 and R_2 and S is a quartic $(4) \subseteq \mathbb{P}(2, 1, 1, 1)$, that is, $R \cong \mathbb{C}[x_1, x_2, x_3, x_4]/(F)$ with $x_2, x_3, x_4 \in R_1$ and $x_1 \in R_2$ and F of weighted degree four.
- (iii) If $d = 3$, then $R \cong \mathbb{C}[x_1, x_2, x_3, x_4]/(F)$ with $x_i \in R_1$ ($1 \leq i \leq 4$) and a homogeneous polynomial F of degree three. That is, S is a cubic surface in \mathbb{P}^3 .
- (iv) If $d \geq 4$, then $R \cong \mathbb{C}[x_1, \dots, x_{d+1}]/\mathcal{I}$ with $x_i \in R_1$ ($1 \leq i \leq d+1$) and \mathcal{I} generated by $d(d-3)/2$ quadrics. In particular, if $d = 4$ then S is a complete intersection of two quadrics in \mathbb{P}^4 .

Remark A.2. If a del Pezzo surface S of degree d is a blow-up of \mathbb{P}^2 at more than one point, then it can also be described as a blow-up of $\mathbb{P}^1 \times \mathbb{P}^1$ at $8 - d$ points.

Remark A.3. The (i)-(iv) in Theorem A.1 still holds when S is normal and Gorenstein as shown by Hidaka and Watanabe [15]. This holds even for a (not necessarily normal) generalized del Pezzo surface as derived from the works of Reid [46].

Now, we recall the definition of del Pezzo manifolds, which is a natural higher dimensional polarized version of the notion of del Pezzo surfaces.

A smooth polarized variety $(X, \mathcal{O}_X(1))$ of dimension $n \geq 3$ is *del Pezzo* if $\mathcal{O}_X(-K) = \mathcal{O}_X(n-1)$. An ample divisor H is called a *fundamental divisor* of X if $\mathcal{O}_X(H) \cong \mathcal{O}_X(1)$. The corresponding linear system $|H|$ is called the *fundamental system* of X . The self-intersection number $d = H^n$ is also called the degree of the del Pezzo manifold X . Note that $H^i(X, \mathcal{O}_X) = 0$ for all $i > 0$ by Kodaira vanishing Theorem.

The classification of del Pezzo manifolds was obtained by Fujita and Iskovskikh. We are mainly interested in smooth del Pezzo threefolds in this paper.

For the following proposition, see [18, Propositions 3.2.3, 3.2.4] and the references therein.

Proposition A.4. *Let $(X, \mathcal{O}_X(1))$ be a smooth del Pezzo threefold of degree d . Then:*

- (1) *If S is the zero locus defined by a general global section of $\mathcal{O}_X(1)$, then it is a smooth del Pezzo surface and $\mathcal{O}_S(-K) \cong \mathcal{O}_X(1)|_S$.*
- (2) *If $d \geq 3$, then the fundamental system of X determines an embedding $X \hookrightarrow \mathbb{P}^{d+1}$ which is projectively normal and $\mathcal{O}_X(1) \cong \mathcal{O}_{\mathbb{P}^{d+1}}(1)|_X$.*
- (3) *If $d \geq 4$, then the image $X \hookrightarrow \mathbb{P}^{d+1}$ is an intersection of quadrics.*

We have the following classification of smooth del Pezzo threefolds, see [18, Theorem 3.3.1] and the references therein.

Theorem A.5. *Let $(X, \mathcal{O}_X(1))$ be a smooth del Pezzo threefold of degree d . Then X is one of the following:*

- (1) $d = 1$ and X is a sextic hypersurface in $\mathbb{P}(3, 2, 1, 1, 1)$.

- (2) $d = 2$ and X is a quartic hypersurface in $\mathbb{P}(2, 1, 1, 1, 1)$.
- (3) $d = 3$ and $X \hookrightarrow \mathbb{P}^4$ is a cubic hypersurface.
- (4) $d = 4$ and $X \hookrightarrow \mathbb{P}^5$ is a complete intersection of two quadrics.
- (5) $d = 5$ and $X \hookrightarrow \mathbb{P}^6$ is a linear section of Plücker embedding of $Gr(2, 5) \subseteq \mathbb{P}^9$ by codimension 3 subspace.
- (6I) $d = 6$ and $X \hookrightarrow \mathbb{P}^7$ is a divisor on $\mathbb{P}^2 \times \mathbb{P}^2$ of bidegree $(1, 1)$, and $\mathbb{P}^2 \times \mathbb{P}^2 \hookrightarrow \mathbb{P}^8$ by Segre embedding.
- (6II) $d = 6$ and $X = \mathbb{P}^1 \times \mathbb{P}^1 \times \mathbb{P}^1 \hookrightarrow \mathbb{P}^7$ by Segre embedding.
- (7) $d = 7$ and $X \hookrightarrow \mathbb{P}^8$ is a blow-up of \mathbb{P}^3 at a point.
- (8) $d = 8$ and $X = \mathbb{P}^3 \hookrightarrow \mathbb{P}^9$ with $\mathcal{O}_X(1) = \mathcal{O}_{\mathbb{P}^3}(2)$.

The Hodge numbers of del Pezzo threefolds are given in Table 9.

d	1	2	3	4	5	6I	6II	7	8
$h^{1,1}(X)$	1	1	1	1	1	2	3	2	1
$h^{2,1}(X)$	21	10	5	2	0	0	0	0	0

TABLE 9. Hodge numbers of del Pezzo threefolds.

Remark A.6. Let h_1 denote the pullback, via the blow-up $X_7 \rightarrow \mathbb{P}^3$, of the hyperplane class on \mathbb{P}^3 . Let e be the exceptional divisor in X_7 . Write $h_2 = h_1 - e$. Then $\mathcal{O}_{X_7}(1) = \mathcal{O}(h_1 + h_2)$.

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