

Solving under-constrained hyperelasticity without the null space

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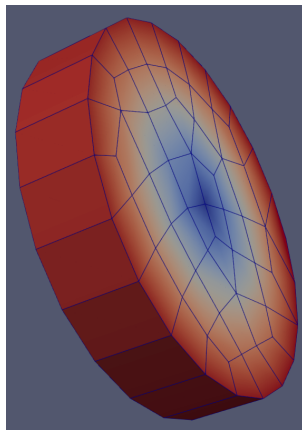
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Motivation & Goal

- Under-constrained hyperelasticity: free energy $f(\mathbf{x})$ has invariant motions
- Some motions (translations) are in the null space of the Hessian, but some (rotations) aren't
- How does this affect minimizing $f(\mathbf{u})$ using Newton-Krylov methods?



Our goal:

Robust Newton-Krylov for under-constrained hyperelasticity

Free energy: $f(\mathbf{x})$

- $\mathbf{g}(\mathbf{x}_k) = \nabla f(\mathbf{x}_k) \longleftarrow$ gradient
- $H(\mathbf{x}_k) = \nabla^2 f(\mathbf{x}_k) \longleftarrow$ Hessian

Key Step: Solve $H_k \delta \mathbf{x}_k = -\mathbf{g}_k$ using Krylov method

Objective: Find \mathbf{x}_* such that $f(\mathbf{x}_*)$ is local minimum

Important fact:

$$\text{if } \mathbf{g}_k \perp \ker(H_k) \longrightarrow \delta \mathbf{x}_k = -H_k^+ \mathbf{g}_k$$

Small Strain Invariants vs. Finite Strain Invariants

Small strain:

$$\boldsymbol{\varepsilon}(\mathbf{u}) = \frac{1}{2} (\nabla \mathbf{u} + \nabla \mathbf{u}^T)$$

The rigid body modes are in the null space of the symmetric gradient.

Finite strain:

$$\mathbf{e}(\mathbf{u}) = \frac{1}{2} \left(\nabla \mathbf{u} + \nabla \mathbf{u}^T + \underbrace{\nabla \mathbf{u} \nabla \mathbf{u}^T}_{\text{nonlinear}} \right)$$

The rigid body modes are **NOT** in the null space of the Green strain tensor.

Free Energy Restrictions: Objectivity & Invariance!

Coordinates of N particles of a free body without Dirichlet BCs:

$$\mathbf{x} = (x_1, y_1, z_1, x_2, y_2, z_2, \dots, x_N, y_N, z_N)^T$$

Translational Invariants:

$$\mathbf{t}_x = (1, 0, 0, 1, 0, 0, \dots, 1, 0, 0)^T$$

$$\mathbf{t}_y = (0, 1, 0, 0, 1, 0, \dots, 0, 1, 0)^T$$

$$\mathbf{t}_z = (0, 0, 1, 0, 0, 1, \dots, 0, 0, 1)^T$$

$$f(\mathbf{x} + \mathbf{t}) = f(\mathbf{x})$$

- is *subspace*:
 $\mathbf{t} \in \ker(H(\mathbf{x})), \mathbf{g} \perp \mathbf{t}$
- *doesn't* affect Krylov iteration

Rotational Invariants:

\mathbf{q} : an element in $SO(3)$

$\mathbf{R}(\mathbf{q}) = (I_N \otimes \mathbf{R}(\mathbf{q}))$: rotation matrix

$\mathbf{o}(\mathbf{q}, \mathbf{x}) = \mathbf{R}(\mathbf{q})\mathbf{x}$: orbit function

$$\mathbf{r}_x = (0, -z_1, y_1, 0, -z_2, y_2, \dots, 0, -z_n, y_n)^T$$

$$\mathbf{r}_y = (z_1, 0, -x_1, z_2, 0, -x_2, \dots, z_n, 0, -x_n)^T$$

$$\mathbf{r}_z = (-y_1, x_1, 0, -y_2, x_2, 0, \dots, -y_n, x_n, 0)^T$$

* \mathbf{r}_x , \mathbf{r}_y , and \mathbf{r}_z are the tangents of the orbit at \mathbf{x} .

$$f(\mathbf{o}(\mathbf{q}, \mathbf{x})) = f(\mathbf{x})$$

- is *curved*: $\mathbf{r} \notin \ker(H(\mathbf{x}))$
- *does* affect Krylov iteration

Plain Newton-Krylov & the Rotations!

An Implementation in PETSc:

Rigid body modes:

$$Z_{m \times 6} = [\mathbf{t}_x | \mathbf{t}_y | \mathbf{t}_z | \mathbf{r}_x | \mathbf{r}_y | \mathbf{r}_z]$$

$$Y_{m \times 6} = A_{m \times m} Z_{m \times 6}$$

$$D_{6 \times 6} = Y^T Y$$

$$D = U \sigma^2 V^T$$

$\rightarrow \hat{V}$: corresponding right singular vectors of the zero singular values

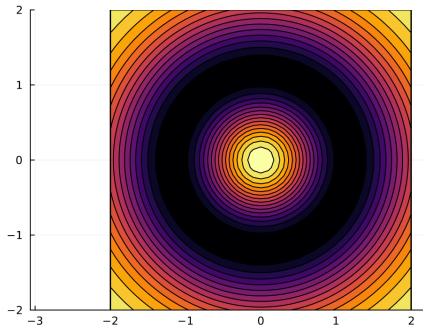
The desired (?) null space:

$$\hat{Z} = Z \hat{V}$$

```
newton_ad(x₀, strain_energy); # reproduced in Julia
```

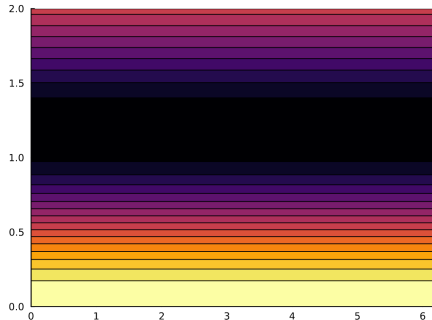
```
Iteration 1: energy 0.5402974677121607
             gradient norm 1.4105813620865468
             hessian rank 21
             effective condition number 61.58224740081209
             step length 1.0
Iteration 2: energy 0.005247987808287236
             gradient norm 0.11625451629789378
             hessian rank 21
             effective condition number 777.3217296711898
             step length 1.0
Iteration 3: energy 5.415541871298145e-7
             gradient norm 0.0015862456717649644
             hessian rank 21
             effective condition number 20594.65732797073
             step length 1.0
Iteration 4: energy 6.893112972106431e-15
             gradient norm 1.284767074447696e-7
             hessian rank 21
             effective condition number 1.8820447105892365e9
             step length 1.0
Iteration 5: energy 2.220446049250321e-16
             gradient norm 2.748665926407727e-15
             hessian rank 18
             effective condition number 15.000000000000001
```

Re-parameterize Rotations



```
[36]:  
H = ForwardDiff.hessian(f, randn(2))  
rank(H)
```

```
[36]:  
2
```



```
[38]:  
Ĥ = ForwardDiff.hessian(f̂, randn(2))  
rank(Ĥ)
```

```
[38]:  
1
```

Project out the Rotational Modes

U : Orthonormal basis spanning
 $\partial_{\mathbf{q}} o(\mathbf{q}, \mathbf{x}) : \mathbf{q} \in SO(3)$

Modified Hessian:
 $\tilde{H} = (I - UU^T)H(I - UU^T)$

Rotation-corrected Newton Krylov

It's easy to implement, but does it affect convergence?

If there is a neighborhood \mathcal{N} where:

- H is Lipschitz
- \tilde{H} is Lipschitz up to unitary transformation:
 $\min \tilde{H}(\mathbf{x})Q - Q\tilde{H}(\mathbf{y}) : Q^T Q = I \lesssim |\mathbf{x} - \mathbf{y}|$
- The curvature of orbits is bounded: $|D^2 o(\mathbf{q}, \mathbf{x})| \leq C$

Then the iteration $\mathbf{x}_{k+1} = \mathbf{x}_k - \tilde{H}(\mathbf{x}_k)^+ \mathbf{g}_k$ converges
 R -quadratically near a solution

Plain vs Corrected Newton-Krylov

```
newton_ad(x₀, strain_energy);
```

```
Iteration 1: energy 0.5402974677121604  
            gradient norm 1.4105813620865468  
            hessian rank 21  
            effective condition number 61.582247400812136  
            step length 1.0  
Iteration 2: energy 0.005247987808286681  
            gradient norm 0.11625451629789381  
            hessian rank 21  
            effective condition number 777.3217296712144  
            step length 1.0  
Iteration 3: energy 5.415541876854953e-7  
            gradient norm 0.0015862456717675895  
            hessian rank 21  
            effective condition number 20594.65732787524  
            step length 1.0  
Iteration 4: energy 6.44902375026225e-15  
            gradient norm 1.2847656622972564e-7  
            hessian rank 21  
            effective condition number 1.8820448808005939e9  
            step length 1.0  
Iteration 5: energy -1.221245327087671e-15  
            gradient norm 3.2984620808806433e-15  
            hessian rank 18  
            effective condition number 15.00000000000001
```

```
newton(x₀, strain_energy, x -> ForwardDiff.gradient(strain_energy, x),  
        x -> rigid_body_corrected_hessian(strain_energy, x));
```

```
Iteration 1: energy 0.5402974677121604  
            gradient norm 1.4105813620865468  
            hessian rank 18  
            effective condition number 15.55394648873718  
            step length 1.0  
Iteration 2: energy 0.004959838476031193  
            gradient norm 0.11654294393222203  
            hessian rank 18  
            effective condition number 15.20556531877761  
            step length 1.0  
Iteration 3: energy 6.684727159775428e-7  
            gradient norm 0.001280378513351747  
            hessian rank 18  
            effective condition number 15.002477767222615  
            step length 1.0  
Iteration 4: energy 9.788756571327347e-15  
            gradient norm 2.0455627721105363e-7  
            hessian rank 18  
            effective condition number 14.999999934516481  
            step length 1.0  
Iteration 5: energy 1.1102230246251573e-15  
            gradient norm 3.3453785797687322e-15  
            hessian rank 18  
            effective condition number 15.000000000000023
```

x —coordinate system

$$x_k = x_{k-1} + \delta x_{k-1}$$

$$\begin{aligned} f_{x_k}(x_k) \\ g_k = \nabla f_{x_k} \\ H_k = \nabla^2 f_{x_k} \end{aligned}$$

$$\delta x_k = -\tilde{H}_k^+ g_k$$

$$x_{k+1} = x_k + \delta x_k$$

$$\tilde{H} = (I - UU^T)H(I - UU^T)$$

① Changing coordinate system

\hat{x} —coordinate system

$$\varphi_x(\hat{x})$$

$$J = \partial x / \partial \hat{x}$$

$$\begin{aligned} \hat{f}_{x_k}(\hat{x}_k) &= f(\varphi_{x_k}(\hat{x}_k)) \\ \hat{g}_k &= \nabla_{\hat{x}} \hat{f}_{x_k}(0) = J^T g_k \\ \hat{H}_k &= \nabla_{\hat{x}}^2 \hat{f}_{x_k}(0) \end{aligned}$$

$$\tilde{H}_k = J_k^{-T} \hat{H}_k J_k^{-1}$$

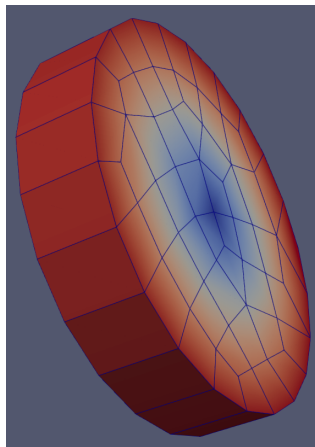
$$\delta x_k = J_k \delta \hat{x}_k$$

$$\delta \hat{x}_k = -\hat{H}_k^+ \hat{g}_k$$

② Projecting out rotational modes

What if the body is partially constrained?

- Rotation-corrected Newton should only project out tangents of invariant rotations
- Can the invariant rotations be determined automatically?



What if the set of invariant rotations is unknown?

A rotation \mathbf{q} is invariant if

$$f(\mathbf{o}(\epsilon \mathbf{q}, \mathbf{x})) = f(\mathbf{x}) \quad \text{for all } \epsilon$$

First order condition:

$$\partial_{\mathbf{q}} f(\mathbf{o}(\epsilon \mathbf{q}, \mathbf{x})) = 0 \Rightarrow \mathbf{g}(\mathbf{x})^T V_{\mathbf{o}}(\mathbf{x}) \mathbf{q} = 0 \text{ where, } \partial_{\mathbf{q}} \mathbf{o}(0, \mathbf{x}) = V_{\mathbf{o}}(\mathbf{x}) \mathbf{q}$$

Second order condition:

$$\partial_{\mathbf{q}}^2 f(\mathbf{o}(\epsilon \mathbf{q}, \mathbf{x})) = 0 \Rightarrow \mathbf{u}^T [\nabla_{\mathbf{x}} V_{\mathbf{o}}(\mathbf{x}) : \mathbf{g}(\mathbf{x})] \mathbf{q} + \mathbf{u}^T H(\mathbf{x}) V_{\mathbf{o}}(\mathbf{x}) \mathbf{q} = 0$$

Therefore \mathbf{q} is an invariant rotation if:

$$\begin{bmatrix} H V_{\mathbf{o}} + [\nabla_{\mathbf{x}} V_{\mathbf{o}}(\mathbf{x}) : \mathbf{g}(\mathbf{x})] \\ \mathbf{g}^T V_{\mathbf{o}} \end{bmatrix} \mathbf{q} = 0$$

Determining Invariant Rotations

If the orbit function $o(\mathbf{q}, \mathbf{x})$ of all rotations is known, the subset of invariant rotations can be determined at the cost (per Newton iteration) of

- 3 Hessian-vector products
- One tall-skinny $\mathbb{R}^{n \times 3}$ QR factorization

- Projecting out the whole rigid body mode from the Hessian
- Implementation in PETSc and our solid mechanics software (**Ratel**: <https://gitlab.com/micromorph/ratel>)