# On extremal eigenvalues that lays outside of the support of the spectral distribution

# **Table of Contents**

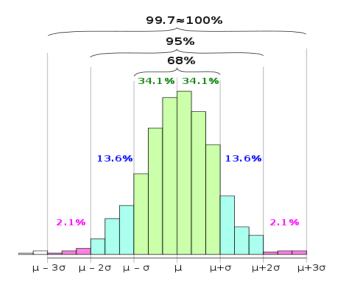
Opening Words	2
1 Basics and Notations	3
2 Introdunction to Random Matrices	4
2.1 Spectral Distribution	4
2.2 Classification	5
3 Extremal Eigenvalues' Fluctuation in Classical Models	6
3.1 Classical Wigner matrices	6
3.1.1 Asymptotics	7
3.1.2 All eigenvalues fluctuate in their definitional domains	7
3.1.3 Edge universality and Tracy-Widom law	8
3.2 Classical sample covariance matrices	10
3.2.1 Similarity and complexity	10
3.2.2 Asymptotics	11
3.2.3 Fluctuation in definitional domain	12
3.2.4 Edge universality and Tracy-Widom law	13
3.3 Summary:Universality at soft edges	14
4 Outliers in Deformed Models	14
4.1 Common deformations	14
4.2 Asymptotics	15
4.3 Non-universal fluctuations of outliers	17
5 References	1Ω

## **Opening Words**

This document is a little heavy in mathematics. Doing research usually follows the three steps: **concept clarification**, **criterion setting** and **system building**. I'll try to do the work according to this guidance and try to use the least mathematics for the topic.

There are four keywords in the title: **extremal eigenvalue**, **lay outside**, **support**, **spectral distribution**. It is not sufficient to only tell you that it belongs to the area of random matrix theory(RMT) and the document is to study extremal eigenvalues' fluctuation. Better to pay attention to the words "**random matrix**" and "**extremal eigenvalue**", since they are the origin and destination of the document, respectively.

For a preliminary understanding, it is better to see the graph illustrated. Just consider the common way in statistics to describe a data distribution, e.g. in univariate Gaussian distribution, let  $\mu$  and  $\sigma$  be the mean value and standard deviation. It is known that the data within one standard deviation of the mean account for about 68% of the set, while within two standard deviations account for about 95%. And within three standard deviations namely, in the domain  $[\mu \mp 3\sigma]$  which includes around 99.7% of the data, suppose we can see the domain as the main part of the data distribution(since it is nearly 100%), then the so-called extremal data outside it can be seen as outliers.



\*The figure is from

https://en.wikipedia.org/wiki/68%E2%80%9395%E2%80%9399.7\_rule

Similarly, but to be more complex, the study object in this document is also some extremal data but with more randomness.

#### 1 Basics and Notations

Before going into formal content, the necessary mathematics are listed here.

There are lots of beauty and profound mathematical properties, but only the needed results are put here, without so much in detailed explanation.

Not all the terms are explicitly used or mentioned, but they are foundational to the results in next parts and listed here

According to convention, the N-dimensional real number field is denoted by  $\mathbb{R}^n$ ; N-dimensional complex number field,  $\mathbb{C}^n$ 

**Deterministic Matrices**: usually in which the elements are deterministic functions or certain numbers. The elements of matrices are also called entries. Here the identity matrix is denoted by I

## **Eigenvalues and Eigenvectors**

For a square matrix A, a pair of scalar and non-zero vector  $(\lambda, \mathbf{v})$  satisfying

$$Av = \lambda v$$

is called an eigenvalue-eigenvector pair.

The eigenvalue of a square matrix is called the spectrum of a matrix.

**Transpose**: a matrix operator, which flips a matrix over its diagonal; that is, it switches the row and column indices of the matrix A by producing another matrix, often denoted by  $A^T$ 

**Trace**:, denoted by Tr, which is the sum of diagonal entries or eigenvalues, (trace invariance by similarity transformation)

**Symmetric matrix**: the square matrix that is equal to its transpose **Hermitian matrix**: the complex square matrix that is equal to its own conjugate transpose, for the reason, its diagonal entries are real. The two are both diagonalizable with real eigenvalues.

#### **Singular Values**

For a matrix *A*, if it is not diagonalizable, there is an alternative decomposition (SVD, singular value decomposition) as

$$A = VSU^T$$

where S is a non-negative diagonal matrix, whose elements are called the singular values of A, and U, V are two real, orthogonal matrices.

*A* is said to be **positive semi-definite** or **non-negative-definite**  $\Leftrightarrow \mathbf{x}^T A \mathbf{x} \geq 0, \forall \mathbf{x} \in \mathbb{R}^n$ 

**Dirac delta function**: also known as the unit impulse symbol, is a generalized function or distribution over the real numbers x, whose value is zero everywhere except at a base point  $\lambda$ , and whose integral over the entire real line is equal to one; denoted by  $\delta_{\lambda}(x)$ 

**almost surely**, denoted by a.s., for convergence, it means the measure of divergent points is zero.

**convergent weakly**: A sequence  $X_1, X_2, ...$ , of real-valued random variables is said to converge in distribution, or converge weakly, or converge in law to a random variable X if  $\lim_{n\to\infty} F_n(x) = F(x), \forall x\in\mathbb{R}$  at which F is continuous. Here  $F_n$  and F are the cumulative distribution functions of random variables  $X_n$  and X, respectively.

Support: the points of non-zero probability for a probability distribution, denoted by  $supp(\mathbb{P})$ 

#### 2 Introdunction to Random Matrices

**Random matrices**: some or all the entries are random variables, also called random ensembles

**Dyson index**: denoted by  $\beta = 1,2,4$ , which is referred to the random matrices with real entries, complex entries and quaternionic entries, respectively.

## 2.1 Spectral Distribution

For a Hermitian matrix H of size  $N \times N$ , we denote by  $\{\lambda_1 \ge ... \ge \lambda_N\}$  the set of eigenvalues of H, ranked in decreasing order,.

We define and denote by  $\rho_N = \frac{1}{N} \sum_{i=1}^{N} \delta_{\lambda_i(H)}$  the **empirical spectral distribution(ESD)**, which is also called the sample eigenvalue density.

The **Stieltjes transform** of  $\rho_N$  is denoted for  $z \in \mathbb{C} \setminus \mathbb{R}$  by

$$g_{\rho_N}(z) = \int_{R} \frac{1}{z - x} d\rho_N(x) = \frac{1}{N} Tr(zI - H)^{-1}$$

$$z \notin \{\lambda_i: 1 \leq i \leq N\}$$

If  $\rho_N(x)$  converges weakly to some probability measure  $\rho(x)$  as  $N \to \infty$ ,  $\rho(x)$  is called the **limiting spectral distribution (LSD)** of H, accordingly, the Stieltjes transform of  $\rho(x)$  is

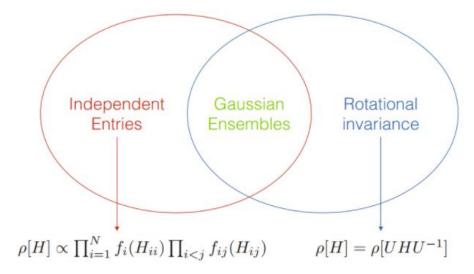
$$g_{\rho}(z) = \int_{supp\{\rho\}} \frac{1}{z - x} d\rho(x)$$

Both ESD and LSD are mean value, which are used to represent the average effect of eigenvalues.

\*Stieltjes transform is not explicitly used in document, but it is an important foundation for the results in next parts, so it is still put here.

#### 2.2 Classification

Here we resort to a simple way to classify random matrices, which is from section 3.2 in [1]



\*The figure is from ref[1],  $\prod_{i=1}^{N} f_i(H_{ii}) \prod_{i < j} f_i(H_{ij})$  is the joint distribution of a Wigner matrix

For the mathematically oriented readers, looking for more formal classifications of random matrix models, it is recommended to read the mini-review [47] and references therein.

#### a) Wigner matrices

Usually, the matrices have independent entries and some symmetric properties, including real symmetric matrices, complex Hermitian matrices..., e.g.

- adjacency matrices of random graphs
- Lévy matrices(independent power-law entries)
- power-law banded matrices
- ...

#### b) Rotational invariance matrices

The matrices, through a similarity transformation, the spectral distribution is invariant. As for similarity transformation, it includes:

- unitary transformation,
- biunitary transformation,

• ...

For simplicity, this property of rotational invariance means that any two matrices that are related via a similarity transformation occur in the ensemble with the same probability, and it indicates that the eigenvectors are not that important, as we can rotate our matrices as freely as we wish, and still leave their statistical weight unchanged. e.g.:

- Wishart-Laguerre ensembles (i.e., Laguerre unitary ensemble, LUE), which is the famous 'Wishart matrices'
- Jacobiclassical ensembles, the so-called "weakly-confined" ensembles
- ...

#### c)\*Gaussian Ensembles

The intersection of a) and b) is only Gaussian Ensembles, which was proved in [46] Due to the fine properties, Gaussian ensembles are often used as a good starting point for study. There are just the three:

- GOE: Gaussian Orthogonal Ensemble, real entries
- GUE: Gaussian Unitary Ensemble, complex entries
- GSE: Gaussian Symplectic Ensemble, quaternionic entries

#### d) Others

Others are exisitent, for instance, biorthogonal ensembles, which are non-invariant and with non-independent entries...

## 3 Extremal Eigenvalues' Fluctuation in Classical Models

In this part, we start from the two widely-studied and classical random matrix models, namely, wigner matrices and sample covariance matrices. The two corresponds to a) and b) in the classification part.

## 3.1 Classical Wigner matrices

The Wigner matrices here are real symmetric or complex Hermitian matrices whose entries are independent up to the symmetry condition, with the form:

$$W_N = \frac{1}{\sqrt{N}} X_N$$

where N is sample size,  $(X_N)_{ii}, \sqrt{2}Re\big((X_N)_{ij}\big)_{i < j}, \sqrt{2}Im\big((X_N)_{ij}\big)_{i < j}$  are i.i.d with distribution  $\tau$  with variance  $\sigma^2$  and  $mean\ 0$ ; if  $\tau = \mathcal{N}(0, \sigma^2)$ ,  $W_N^G$  is G.U.E.-matrix.

#### 3.1.1 Asymptotics

Here are the results on the asymptotic behaviors of the spectral measure[5,6] and extremal eigenvalues[7]

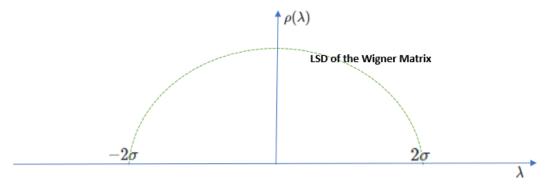
#### **Limiting Spectral Distribution:**

$$\mu_{W_N} := \frac{1}{N} \sum_{i=1}^{N} \delta_{\lambda_i} (W_N) \to \mu_{sc}(\lambda), a. s. N \to +\infty$$

where the spectral density:

$$\rho(\lambda) = \frac{d\mu_{sc}(\lambda)}{d\lambda} = \frac{1}{2\pi\sigma^2} \sqrt{4\sigma^2 - \lambda^2} 1[\lambda_-, \lambda_+]$$

here the left edge  $\lambda_{-}=-2\sigma$ , and the right dege  $\lambda_{+}=+2\sigma$ 



It can be seen that the LSD of the Wigner matrix is a semicircle with the radius of  $2\sigma$  It is worth mentioning that the above result can be called **bulk universality**[65],which concerns the local statistics of eigenvalues in the interior of the spectrum,we can refer to[18] for further discussion

**Extremal eigenvalues**: If  $\int x^4 d\tau(x) < \infty$ , then

$$\lambda_{max}(W_N) \to \lambda_+ \text{ and } \lambda_{min}(W_N) \to \lambda_-, a.s. N \to +\infty$$

We can see when  $N \to \infty$ , the spectral density of the Wigner matrix weakly converges to a semi-circle shape distribution with the bulk of  $[-2\sigma, 2\sigma]$ , and meanwhile, the largest eigenvalue( $\lambda_{min}$ ) and the smallest eigenvalue( $\lambda_{max}$ ) weakly converge to the right edge( $2\sigma$ ) and left edge( $-2\sigma$ ), respectively.

#### 3.1.2 All eigenvalues fluctuate in their definitional domains

\*The discussion here is actually not only limited to the case of Wigner matrices.

We can just think about it easily. In the process of  $N \to \infty$ , the edges of the bulk of the spectral density of the Wigner matrix is moving on the way to the limits, the

smallest and largest eigenvalues are also moving on the way to their limits, as said previously, the spectral distribution is merely spectral expectation, so there must exist deviations to it from other eigenvalues' moving behaviors, in other words, the distances between the moving edges and other eigenvalues are a dynamic process, which can be described through probability distribution.

Let's relax the process to a simpler case and fix the edges as the edges of the limiting spectral density (we can also call the kind of edge as the classical edge), namely in the Wigner case, they are  $-2\sigma$  and  $2\sigma$ . Theoretically, any eigenvalue is probable to occur in any place of its domain of definition. In fact, It has been observed that some eigenvalues are probable to walk beyond the one edge or two edges. That is to say, outside the bulk, the zone where the eigenvalues should not be existent in terms of spectral distribution, the probability of eigenvalues is non-zero, although it can be small. Thus, there exist the transition regions between the edges and some eigenvalues.

With regard to the largest and smallest eigenvalues, after randomly moving in transition regions, they finally converge to some points, which in the Wigner case are just the right and left edges, respectively. Intuitively, it is quite natural to consider the distance between an extremal eigenvalue and a near edge. In a general sense, the results on the distance between an edge and some kind of extremal eigenvalue have been investigated quite sufficiently.

Concretely, the fluctuating local density of eigenvalues can be computed from the number of eigenvalues within a distance, for concrete derivations, we can refer to page-75 in[2].

#### 3.1.3 Edge universality and Tracy-Widom law

#### Starting from the largest eigenvalue

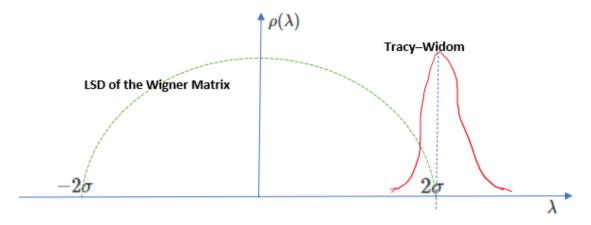
The first beautiful result was achieved by Tracy and Widom on the distribution of the largest eigenvalue of gaussian wigner matrices.

This result can be formally stated: the rescaled distribution of  $\lambda_{max} - \lambda_+$  converges, as  $N \to \infty$ , follows the Tracy–Widom distribution, usually denoted by  $TW_\beta$ :

$$\mathbb{P}\big(N^{2/3}(\lambda_{max}-\lambda_+)\leq s\big)=TW_{\beta\in(1,2)}(s)$$

here for the case,  $\lambda_+ = 2\sigma$ .

For intuitive understanding, see the graph below.



\*The curve of Tracy-Widom distribution in this and later graphs is only **schematic plot**, which is not in its actual size.

Simply speaking for the Wigner case, the largest eigenvalue  $\lambda_{max}$  does not fluctuate very far away from the right classical edge  $\lambda_+$ . Take for example N = 1000,and  $\lambda_{max}$  is within  $1000^{-2/3}=0.01$  away from  $\lambda_+=2\sigma$ . This indicates that the width of the region around  $\lambda_+$  within which one expects to observe the largest eigenvalue of a Wigner matrix goes down as  $N^{-2/3}$ .

For the **density of Tracy–Widom density**, we have known that  $tw_{\beta \in (1,2)}(s) = TW'_{\beta \in (1,2)}(s)$ , w.r.t the left and right far tails:

$$tw_{\beta\in(1,2)}(s) \propto -s^{3/2}$$
, as  $N\to\infty$ ;  $tw_{\beta\in(1,2)}(s) \propto -|s|^3$ , as  $N\to-\infty$ 

It is noted that the left tail is much thinner than the right tail, it shows pushing the largest eigenvalue inside the bulk is more difficult than pulling one eigenvalue away from  $\lambda_+$ . Using this analogy and the formalism in section 5.4.2 in [2], the large deviation regime of the Tracy–Widom problem (i.e.  $\lambda_{max} - \lambda_+ = O(1)$ ) can be obtained. Note that the result is exponentially small in N as the  $s^{3/2}$  behavior for  $s \to \infty$  combines with  $N^{2/3}$  to give a linear dependence in N

As discussed, the distribution of the largest eigenvalue follows the law of Tracy–Widom, which exhibits a type of universality, it's called the **edge universality**.

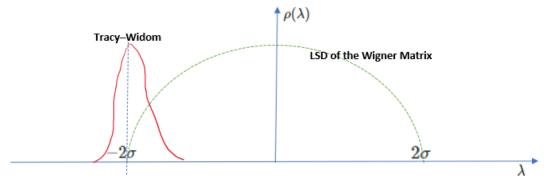
In fact, The Tracy–Widom law holds to more general Wigner matrices [8-13], **even in the case** that the symmetric assumption is partially removed [14,15]. The results can be found:

- Edge universality for Wigner matrices whose entries have vanishing third moments was proved in [42];
- Edge universality without moment matching was proved in [43] for Wigner matrices and in [44, 45] for generalized Wigner matrices;

- A necessary and sufficient condition on the entries' distributions for the edge universality to hold was given in [16].
- The fluctuations of the largest eigenvalue of a real symmetric or complex Hermitian Wigner matrix of size N converge to the Tracy–Widom laws has been proved at the rate of  $O(N^{-1/3+\omega})$ ,  $\forall \omega > 0$ , as  $N \to \infty$  in [48].

## Then what about the smallest eigenvalue?

The distribution of the smallest eigenvalue  $\lambda_{min}$  around the lower edge  $\lambda_{-}$  also follows the Tracy–Widom law. So is a similar graph:



As to the tails, still the lift is thinner than the right, but it means putting the smallest eigenvalue inside the bulk is easier than dragging outside.

#### For any ordered eigenvalue

For more **general cases**, in [20], the authors proved that after proper rescaling, the 1st, 2nd, 3rd, etc. eigenvalues of Wigner matrices weakly converge to Tracy–Widom distribution. This is to say, the distance between a classical edge( the left edge or the right edge) and any ordered eigenvalue follows the Tracy–Widom distribution as  $N \rightarrow \infty$ . Obviously, the case of extremal eigenvalues are included in.

**To recap**, in the Wigner matrices, the extremal eigenvalues weakly converge to the bulk edges of the limiting spectral distribution, and the corresponding fluctuations follow the Tracy–Widom distribution.

#### 3.2 Classical sample covariance matrices

After introducing Wigner matrices, with a similar introducing structure, let's go into a kind of invariance matrices--- sample covariance matrices

#### 3.2.1 Similarity and complexity

Although **MOST** of the laws on the fluctuation of eigenvalue hold in the part for Wigner matrices still hold for the sample covariance matrices here, there are some more complex conditions(mainly since dependent entries) and important exceptions of interest.

#### 3.2.2 Asymptotics

When the sample covariance matrix is assumed as

$$S_N = \frac{1}{p(N)} X_N X_N^*$$

where  $X_N$  is a complex random matrix with i.i.d. entries. It is assumed that  $X_N$  is a  $N\times p(N)$  and  $p(N)\leq N$  complex random matrix,  $Re\big((X_N)_{ij}\big), Im\big((X_N)_{ij}\big), i=1,\ldots,N, j=1,\ldots,p(N)$  are i.i.d,following distribution  $\tau$  with variance  $\frac{1}{2}$  and mean 0. Note that the spectra of  $\frac{1}{p(N)}X_NX_N^*$  and  $\frac{1}{p(N)}X_N^*X_N$  differ by |p(N)-N| zero eigenvalues. If  $\tau$  is Gaussian,  $S_N=:S_N^G$  is an L.U.E matrix,i.e.the Wishart matrix

From the definition, we can see, the sample covariance matrices are positive semidefinite.

The behavior of the spectral measure as  $N \to \infty$  follows Marcenko–Pastur distribution

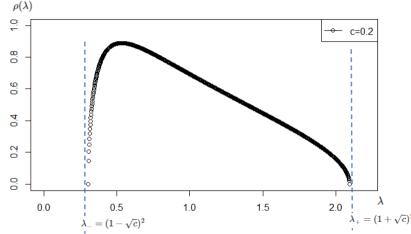
**Limiting Spectral Distribution** ([21]): If  $C_N := \frac{N}{p(N)} \to c \in (0,1]$  as  $N \to \infty$ , then

$$\mu_{SC} \rightarrow \mu_{MP}$$
, as  $N \rightarrow \infty$ 

where the spectral density

$$\rho(\lambda) = \frac{d\mu_{SC}}{d\lambda} = \frac{1}{2\pi c\lambda} \sqrt{(\lambda_{+} - \lambda)(\lambda - \lambda_{-})} 1_{[\lambda_{-}, \lambda_{+}]}$$

where  $\lambda_- = \left(1 - \sqrt{c}\right)^2$ ,  $\lambda_+ = \left(1 + \sqrt{c}\right)^2$ , c is called aspect ratio. The graph is an



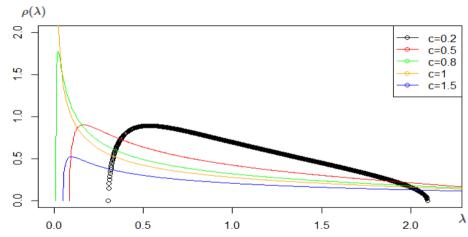
example with c = 0.2.

In[41],for some well-designed sample covariance matrices, the authors find a certain threshold  $c_+$  such that when  $c>c_+$ ,the limiting spectral distribution also exhibits convex decay at the right edge of the spectrum

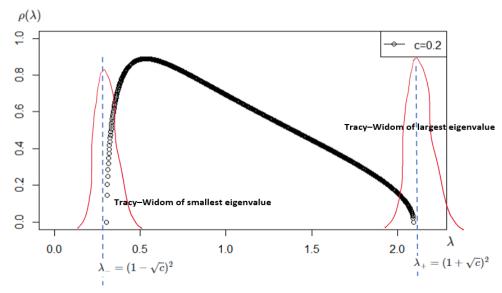
**Extremal eigenvalues**([2,22-24]): If  $\int x^4 d\tau(x) < \infty$ , then

$$\lambda_{max}(S_N) \rightarrow \lambda_+ \text{ and } \lambda_{min}(S_N) \rightarrow \lambda_-, a.s. N \rightarrow +\infty$$

For the two results, if we extend to more general conditions on entries, or let c be  $(0, \infty]$ , or choose other assumptions, some formalized results are also achieved, for further readings on the asymptotics of empirical spectral measure, we refer to [25]; more details, on the asymptotics of extremal eigenvalues of sample covariance matrices, are well-studied are in [26-29]. The graph below are the examples with more c values.



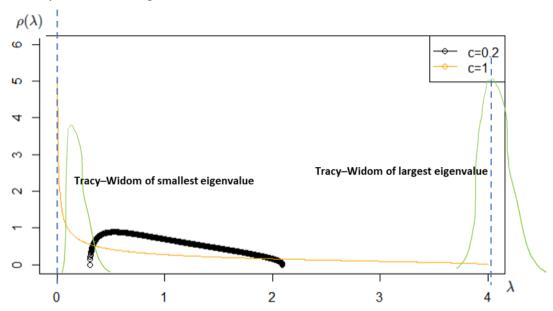
#### 3.2.3 Fluctuation in definitional domain



Like that of Wigner matrices, all the eigenvalues of sample covariance matrices

fluctuate and follows the Tracy–Widom law(details in next part), but under some condition, some eigenvalues only fluctuate inside the bulk, the corresponding edge is called **hard edge**,e.g. the fluctuation from the smallest eigenvalue to the left edge. Take the Wishart matrices as the example, when c=1, the lower edge  $\lambda_-=0$ , this is a "hard" edge since all eigenvalues of the empirical matrix must be non-negative. In contrast, the Wigner semicircle edges are "soft".





#### 3.2.4 Edge universality and Tracy-Widom law

Like wigner case, for the largest eigenvalue of Wishart matrices, as  $N \to \infty$ , the rescaled distribution of  $\lambda_{max} - \lambda_+$  converges to the Tracy–Widom distribution, with the similar form

$$\mathbb{P}(N^{2/3}(\lambda_{max} - \lambda_+) \le \gamma s) = TW_{\beta \in (1,2)}(s)$$

here for the case,  $\lambda_{+} = (1 + \sqrt{c})^{2}$ ,  $\gamma = \sqrt{c}\lambda_{+}^{2/3}$ .

In fact, this law on the largest eigenvalue is not restricted to the Wishart ensemble and it is universal for general sample covariance matrices, even removing the condition on c (see [30]). As to the Tracy–Widom law on the smallest eigenvalue, it still holds when  $c \neq 1$ ([31]).

#### The results can be found:

- Edge universality for sample covariance matrices when  $\sqrt{N}X_N$ )<sub>ij</sub> have vanishing third moments was proved in [32] under the condition  $c \neq 1$ ;
- Edge universality without moment matching for sample covariance matrices was proved in [33] when  $c \neq 1$ ;

- A necessary and sufficient condition on the entries' distributions for the edge universality to hold was given in [34].
- When c = 1 (i.e. N = p(N)), the smallest eigenvalue exhibits a different asymptotic behavior, which is referred to as the **hard edge** case, and the corresponding universality in general sample covariance matrices was studied in [35,36].
- When c > 0, the convergent rate to its Tracy–Widom limit has been improved to  $N^{-1/3}$  in [37].

To recap, compared to Wigner matrices, since the increased complexity in elements, relation, structure and interaction in sample covariance matrices, the asymptotics of extremal values presents more difficulties for study. At least, it is known that in a general sense, the convergence and fluctuation of extremal values have achieved the comparable results to Wigner matrices

## 3.3 Summary: Universality at soft edges

With regard to the universality of the Tracy-Widom fluctuations at soft edges, it can extend to quite general deformed Wigner matrices or sample covariance matrices without outliers. The involved methods pursue a Green function comparison strategy ([40, 49, 50]) or make use of anisotropic local laws [51].

Two keywords can be found: 'deformed' and 'outlier', so we go to the next section.

#### 4 Outliers in Deformed Models

Here we draw a line on outliers.

For a square random matrix with size  $N \times N$ , as  $N \to \infty$ , if an eigenvalue converge to outside the bulk and all edges, the eigenvalue can be seen as an outlier, although it can fluctuate inside the bulk.

Later in the deformed models, the clear cause and mechanism for outlier's occurrence can be found.

#### 4.1 Common deformations

Deformed models here means the perturbation or distortion of some classical ensembles (such as Wigner matrices, sample covariance matrices...), here we introduce three:

- i) Additive perturbation,  $M_N = W_N + A_N$ ,
  - $W_N$  is a Wigner matrix,
  - $A_N$  is an  $N \times N$  Hermitian matrix
- ii) Multiplicative perturbation,  $M_N = A_N^{\frac{1}{2}} S_N A_N^{\frac{1}{2}}$ 
  - $S_N$  is a sample covariance matrix,
  - $A_N$  is a  $N \times N$  non negative Hermitian matrix

- iii) Information-Plus-Noise,  $M_N = \left(\sigma \frac{X_N}{\sqrt{p}} + A_N\right) \left(\sigma \frac{X_N}{\sqrt{p}} + A_N\right)^*$ 
  - $A_N$  is a rectangular  $N \times p(N)$  matrix,
  - $X_N$  is a rectangular  $N \times p(N)$  matrix, with i.i.d. entries
  - $\sigma$  is some positive real numbers
- Other random perturbations

Usually, $M_N$  is called **perturbation** matrix; $A_N$ , **deformation** matrix.

Additionally, these three kinds of deformations have been also considered for isotropic models, if replacing in i) and ii)  $W_N$ ,  $S_N$  by  $UBU^*$  with U Haar distributed and B deterministic, and in iii)  $X_N$  by a random matrix whose distribution is biunitarily invariant.

Obviously, the complexity to study deformed models is increased quite a lot in the deformed models. On the basis of original methods and by the empowerment of **free convolution and analytic subordination** from free probability theory(see more in[52,1,2]), lots of important results are attained elegantly.

### 4.2 Asymptotics

Setting different conditions on  $A_N$  can render different results, the usual settings are on rank, the fixed number of eigenvalues, the fixed multiplicity of some independent eigenvalues; these fixed eigenvalues are called **spikes**.

When  $A_N$  is finite rank, LSD of  $M_N$  is still the same as that of  $W_N$  or  $S_N$ . But to an extremal eigenvalue of  $M_N$ , it **does not** hold always. There is a critical threshold, if a corresponding eigenvalue in  $A_N$  goes beyond the threshold, the extremal eigenvalue does not stick to the bulk of  $M_N$  and become an outlier, which is said to be generated by a spiked eigenvalue of  $A_N$ . In fact, this law can be extended to any eigenvalue of  $M_N$ .

In a general sense, the relevant criterion for a spiked eigenvalue of  $A_N$  to generate an outlier in the spectrum of the deformed model is to belong to some set related to the subordination functions, which can provide the definite supports of bulk and outliers, respectively.

For the asymptotic behavior of the deformed model, the location of outliers, the property of the corresponding eigenvectors, in [0], the authors give a unified presentation on the results:

$M_{N} = A_{N} + Y_{N}$ $\mu_{A_{N}} \to_{N \to +\infty} \nu$ $\mu_{Y_{N}} \to_{N \to +\infty} \mu$ $\theta \in \operatorname{Spect}(A_{N})$ $\theta \text{ multiplicity } k_{i}$ $\theta \notin \operatorname{supp}(\nu)$	$M_{N} = A_{N}^{1/2} Y_{N} A_{N}^{1/2}$ $\mu_{A_{N} A_{N}^{*}} \to_{N \to +\infty} \nu$ $\mu_{Y_{N}} \to_{N \to +\infty} \mu$ $\theta \in \operatorname{Spect}(A_{N})$ $\theta \text{ multiplicity } k_{i}$ $\theta > 0, \theta \notin \operatorname{supp}(\nu)$	$M_{N} = (A_{N} + Y_{N})(A_{N} + Y_{N})^{*}$ $\mu_{A_{N}A_{N}^{*}} \to_{N \to +\infty} \nu$ $\mu_{Y_{N}Y_{N}^{*}} \to_{N \to +\infty} \mu$ $\sqrt{\mu} \text{ or } \sqrt{\nu}  \boxplus_{c} \text{ infinitely divisible}$ $\theta \in \text{ Spect}(A_{N}A_{N}^{*})$ $\theta \text{ multiplicity } k_{i}$ $\theta > 0, \theta \notin \text{ supp}(\nu)$
$\mu_{M_N} \to_{N \to +\infty} \mu \boxplus \nu$	$\mu_{M_N} \to_{N \to +\infty} \mu \boxtimes \nu$	$\mu_{M_N} \to_{N \to +\infty} (\sqrt{\mu} \boxplus_c \sqrt{\nu})^2$
$g_{ au}(z) = \int_{\mathbb{R}} rac{d au(x)}{z-x}$	$\Psi_{\tau}(z) = \frac{1}{z}g_{\tau}(\frac{1}{z}) - 1$	$H_{\sqrt{\tau}}^{(c)} = \frac{c}{z} g_{\tau}(\frac{1}{z})^2 + (1 - c)g_{\tau}(\frac{1}{z})$
$g_{\mu\boxplus u}(z) = g_{ u}(\omega_{\mu, u}(z))$	$\Psi_{\mu\boxtimes\nu}(z) = \Psi_{\nu}(F_{\mu,\nu}(z))$	$H_{\sqrt{\mu}\boxplus_{c}\sqrt{\nu}}^{(c)}(z) = H_{\sqrt{\nu}}^{(c)}(\Omega_{\mu,\nu}(z))$
$k_i$ outliers of $M_N$ in the neighborhood of each $\rho$ s.t $\omega_{\mu,\nu}(\rho) = \theta$	$k_i$ outliers of $M_N$ in the neighborhood of each $\rho$ s.t $\frac{1}{F_{\mu,\nu}(1/\rho)} = \theta$	$k_i$ outliers of $M_N$ in the neighborhood of each $\rho$ s.t $\frac{1}{\Omega_{\mu,\nu}(1/\rho)} = \theta$
$\xi \text{ eigenvector of } M_N$ associated to an outlier in the neighborhood of $\rho$ s.t $\omega_{\mu,\nu}(\rho) = \theta$ $\ P_{\mathrm{Ker}(\theta I - A)}\xi\ ^2 \to_{N \to +\infty} \frac{1}{\omega'_{\mu,\nu}(\rho)}$	$\begin{array}{c} \xi \text{ eigenvector of } M_N \\ \text{associated to an outlier} \\ \text{in the neighborhood} \\ \text{of } \rho \text{ s.t } \frac{1}{F_{\mu,\nu}(1/\rho)} = \theta \\ \ P_{\text{Ker}(\theta I - A)} \xi\ ^2 \to_{N \to +\infty} \frac{\rho F_{\mu,\nu}(1/\rho)}{F_{\mu,\nu}^{\prime}(1/\rho)} \end{array}$	$\begin{array}{l} \xi \ \text{eigenvector of} \ M_N \\ \text{associated to an outlier} \\ \text{in the neighborhood} \\ \text{of} \ \rho \ \text{s.t.} \frac{1}{\Omega_{\mu,\nu}(1/\rho)} = \theta \\ \ P_{\text{Ker}(\theta I - A)}\xi\ ^2 \to_{N \to +\infty} \frac{\rho^2 g_{(\sqrt{\mu} \boxplus_{\mathbb{C}} \sqrt{\nu})^2}(\rho)}{\theta^2 g_{\nu}(\theta) \Omega'_{\mu,\nu}(1/\rho)} \end{array}$

However, it is impossible to give a comprehensive interpretation on the table in the document, since there are lots of theories, especially free probability theory, and technical assumptions to introduce, and the notations in the table are quite different. If possible later, I'll write a new document on the part.

Still, an intuitive explanation is here, and please allow me skip some conceptual explanations:

- The 1st row shows the basic set of each deformed model
- The 2nd row gives the LSD of each deformed model
- The 3rd and 4th are two different integral transformation(based on Stieltjes transform, free convolution and analytic subordination) for each deformed model
- The 5th gives the sets of outliers' occurrence of each deformed model
- The last row indicates the limiting projection of the eigenvectors associated to outliers' kernel space.
- Concerning the limiting projection of the eigenvectors associated to outliers of Information-Plus-Noise type models, has been proved only in the iid case for diagonal perturbation  $A_N$  and in the isotropic case for nite rank perturbation  $A_N[0]$

#### 4.3 Non-universal fluctuations of outliers

The phenomenon is existent for the fluctuations of the outliers: the limiting distribution can depend on the distribution of the **entries** (thus, the dependence is a kind of non- universality), according to the localization/delocalization of the eigenvectors of  $A_N$ . There are some discussions under different model assumptions:

- Gaussian case, see[58,59], on random perturbations, see [60,61]
- Wigner case, see[54,55,56,57]
- Sample covariance case, see[53,62,63]

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