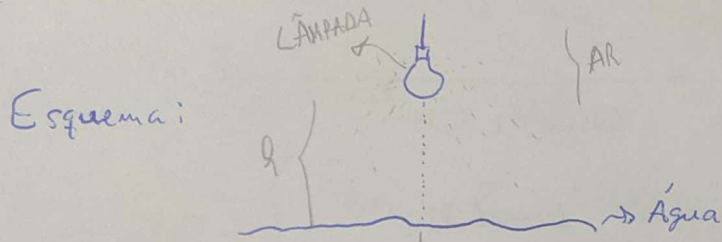


LISTA 2 | Eletromag 2 | LEONARDO C. ROSSATO

01 $n_{\text{água}} = 1,33$; $n_{\text{ar}} = 1$; $h = 1$; $P = 60 \text{ W}$



$h = 1 \left\{ \begin{array}{l} \text{LÂMPADA} \\ \text{Água} \end{array} \right\} I = \frac{P}{A}$
 $A_{\text{inc}} = \pi r^2$, $r \approx \text{lâmpada}$

a)

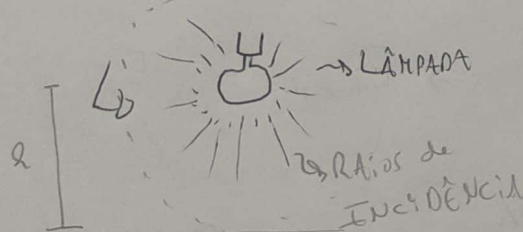
$\rightarrow I = \frac{E_{\text{rms}}^2}{Z} \therefore E_{\text{rms}}^2 = IZ = \frac{PZ}{A} \Rightarrow E_{\text{rms}} = \sqrt{\frac{PZ}{A}}$

$\rightarrow \text{Se } Z_{\text{AR}} \approx Z_0 = 377 \text{ e } A \approx S_{\text{esfera PONTUAL}} = 4\pi r^2 \text{ com } r=1 \quad (*)$

$\therefore E_{\text{rms}} = \sqrt{\frac{60 \cdot 377}{4\pi}} \approx 42,427 \frac{\text{V}}{\text{m}}$

$\rightarrow Z = \frac{E}{H} \Rightarrow H = \frac{E}{Z} \Rightarrow H_{\text{rms}} = \frac{E_{\text{rms}}}{Z} = \frac{42,427}{377} \approx 0,112 \text{ A/m}$

$(*)$ Superfície Esfera PONTUAL } ÁREA de INCIDÊNCIA da LÂMPADA



$\rightarrow A \approx S_{\text{esfera}} = \int dA = \int R^2 \sin\theta d\theta d\phi$
 $= 4\pi R^2$, para ESFERA PONTUAL

$h \approx \text{raio de iluminação}$

$P / \text{Área de Incidência}$

$\therefore A \approx 4\pi$

01) b) Equações Fresnel: $\theta_i = 0 \therefore \theta_r = 0$

$$\left\{ \begin{aligned} \left(\frac{E_{T,m}}{E_{I,m}} \right)_{\perp} &= \frac{2n_1 \cos \theta_i}{n_1 \cos \theta_i + n_2 \cos \theta_r} = \frac{2 \cdot 1 \cdot \cos(0)}{1 \cdot \cos(0) + 1,33 \cos(0)} = \frac{2}{2,33} \\ \left(\frac{E_{T,m}}{E_{I,m}} \right)_{\parallel} &= \frac{2n_1 \cos \theta_i}{n_1 \cos \theta_r + n_2 \cos \theta_i} = \frac{2}{\cos(0) + 1,33} = \frac{2}{2,33} \\ \left(\frac{E_{R,m}}{E_{I,m}} \right)_{\perp} &= \frac{n_1 \cos \theta_i - n_2 \cos \theta_r}{n_1 \cos \theta_i + n_2 \cos \theta_r} = \frac{1 - 1,33}{1 + 1,33} = -\frac{0,33}{2,33} \\ \left(\frac{E_{R,m}}{E_{I,m}} \right)_{\parallel} &= \frac{n_1 \cos \theta_r - n_2 \cos \theta_i}{n_1 \cos \theta_r + n_2 \cos \theta_i} = \frac{1 - 1,33}{1 + 1,33} = -\frac{0,33}{2,33} \end{aligned} \right.$$

→ PARA ENCONTRAR OS VALORES RMS da Onda Refletida e transmitida:

$$\therefore \left\{ \begin{aligned} \frac{E_{T,m}}{E_{I,m}} &= \frac{E_T e^{i(\omega t - kz)}}{E_I e^{i(\omega t - kz)}} = \frac{E_T}{E_I} = \frac{\sqrt{2} E_{rms,T}}{\sqrt{2} E_{rms,I}} = \frac{E_{rms,T}}{E_{rms,I}} \\ \frac{E_{R,m}}{E_{I,m}} &= \frac{E_R e^{i(\omega t - kz + \pi)}}{E_I e^{i(\omega t - kz)}} = \frac{E_R}{E_I} e^{i\pi} = \frac{E_R}{E_I} [\cos \pi + i \sin \pi] = -\frac{E_R}{E_I} = -\frac{E_{rms,R}}{E_{rms,I}} \end{aligned} \right.$$

$$\therefore \left\{ \begin{aligned} \left(\frac{E_{T,m}}{E_{I,m}} \right) &= \frac{2}{2,33} \Leftrightarrow \frac{E_{rms,T}}{E_{rms,I}} = \frac{2}{2,33} \Leftrightarrow E_{rms,T} = \frac{2}{2,33} E_{rms,I} \\ \left(\frac{E_{R,m}}{E_{I,m}} \right) &= -\frac{0,33}{2,33} \Leftrightarrow \frac{E_{rms,R}}{E_{rms,I}} = -\frac{0,33}{2,33} \Leftrightarrow E_{rms,R} = \frac{0,33}{2,33} E_{rms,I} \end{aligned} \right.$$

PORTANTO

$$\hookrightarrow \left\{ \begin{aligned} E_{rms,T} &= \frac{2}{2,33} \cdot 42,427 = 36,418 \text{ V/m} \\ E_{rms,R} &= \frac{0,33}{2,33} \cdot 42,427 = 6,009 \text{ V/m} \end{aligned} \right.$$

$$\left\{ \begin{aligned} H_{rms,T} &= \frac{E_{rms,T}}{Z} = \frac{36,418}{377} = 0,0966 \text{ A/m} \\ H_{rms,R} &= \frac{E_{rms,R}}{Z_0} = \frac{6,009}{377} = 0,0159 \text{ A/m} \end{aligned} \right.$$

RESPOSTAS
LETRA A

02 | $n_{\text{ar}} = 1$; $n_{\text{diamante}} = 2,42$

→ Construa Diagrama das Amplitudes dos $\{E_m, H_m\}$ em relação aos ângulos de incidência, tanto p/ Onda Refletida, quanto Transmitida.

→ PARA Onda TRANSMITIDA:

$$\left(\frac{E_{T,m}}{E_{I,m}} \right)_{\perp} = \frac{2 n_1 \cos \theta_i}{n_1 \cos \theta_i + n_2 \cos \theta_T} = \frac{2 \cos \theta_i}{\cos \theta_i + 2,42 \cos \theta_T}$$

$$\begin{aligned} \rightarrow \cos \theta_T &= \sqrt{1 - \sin^2 \theta_T} \quad \text{e pela Lei de Snell: } \left\{ \begin{aligned} \sin \theta_T &= \frac{n_1}{n_2} \sin \theta_i \\ &= \frac{1}{2,42} \sin \theta_i \end{aligned} \right. \\ &= \sqrt{1 - \left(\frac{1}{2,42} \right)^2 \sin^2 \theta_i} \\ &= \sqrt{1 - 0,17 \sin^2 \theta_i} \end{aligned}$$

$$\therefore \left(\frac{E_{T,m}}{E_{I,m}} \right)_{\perp} = \frac{2 \cos \theta_i}{\cos \theta_i + 2,42 \sqrt{1 - 0,17 \sin^2 \theta_i}} = \boxed{\Gamma}_{I,T}$$

$$\rightarrow \text{Como } \frac{E_{T,m}}{E_{I,m}} = \frac{E_T e^{-i(\omega t - K_T z)}}{E_I e^{-i(\omega t - K_I z)}} = \frac{E_T e^{-i\omega t} e^{+iK_T z}}{E_I e^{-i\omega t} e^{+iK_I z}} = \frac{E_T [\cos K_T z + i \sin K_T z]}{E_I [\cos K_I z + i \sin K_I z]}$$

$$\text{Como importa apenas: } \left\{ \text{Re} \left(\frac{E_{T,m}}{E_{I,m}} \right) = \frac{E_T \cos K_T z}{E_I \cos K_I z} \right\}, \text{ onde } \left\{ K_T = \frac{\omega_T}{c} = \frac{\omega_I}{c} = K_I \right.$$

$$\therefore \frac{E_T}{E_I} = \left(\frac{E_{T,m}}{E_{I,m}} \right) \quad \text{e como } E_{I,m} = E_I \left[\cos \frac{\omega z}{c} + i \sin \frac{\omega z}{c} \right], \text{ p/ } t=0$$

$$\rightarrow \text{Re}(E_{I,m}) = E_I \cos\left(\frac{\omega z}{c}\right)$$

$$\text{e } \text{Re}(E_{T,m}) = \boxed{\Gamma}_{I,T} \text{Re}(E_{I,m}) = \boxed{\Gamma}_{I,T} E_I \cos\left(\frac{\omega z}{c}\right)$$

→ PARA ONDA REFLETIDA:

$$\therefore \left(\frac{E_{R,m}}{E_{I,m}} \right)_{\perp} = \frac{n_1 \cos \theta_I - n_2 \cos \theta_T}{n_1 \cos \theta_I + n_2 \cos \theta_T} = \frac{\cos \theta_I - 2,42 \sqrt{1 - 0,17 \sin^2 \theta_I}}{\cos \theta_I + 2,42 \sqrt{1 - 0,17 \sin^2 \theta_I}} = \boxed{\Gamma_{I,R}}$$

$$\therefore E_{R,m} = \Gamma_{I,R} E_{I,m} = \Gamma_{I,R} E_I e^{-i(\omega t - kz)}, \quad k = \frac{\omega}{c}$$

→ Se $t=0$ e $\omega \approx c$:

*Obs: $E_{R,m}$ e $E_{I,m}$ ESTÃO DEFASADOS
↳ por " π "

$$\therefore \begin{cases} E_{I,m} = E_I e^{+i(\frac{\omega z}{c})} = E_I e^{+iz} = E_I (\cos(z) + i \sin(z)) \\ E_{R,m} = \Gamma_{I,R} E_I e^{i(z+\pi)} = \Gamma_{I,R} E_I (\cos(z+\pi) + i \sin(z+\pi)) \end{cases}$$

→ se usarmos só a parte real p/ fazer os Gráficos:

$$\therefore \begin{cases} \text{Re}(E_{I,m}) = E_I \cos z \\ \text{Re}(E_{R,m}) = \Gamma_{I,R} E_I \cos(z+\pi) \end{cases}$$

$$\therefore \left(\frac{E_{R,m}}{E_{I,m}} \right)_{\parallel} = \frac{n_2 \cos \theta_T - n_1 \cos \theta_I}{n_1 \cos \theta_T + n_2 \cos \theta_I} = \frac{\cos \theta_T - n_2 \cos \theta_I}{\cos \theta_T + n_2 \cos \theta_I} = \frac{\sqrt{1 - 0,17 \sin^2 \theta_I} - n_2 \cos \theta_I}{\sqrt{1 - 0,17 \sin^2 \theta_I} + n_2 \cos \theta_I} = \boxed{\Gamma_{I,R,\parallel}}$$

→ se $t=0$ e $\omega \approx c$:

$$\therefore \begin{cases} E_{I,m} = E_I e^{+i(\frac{\omega z}{c})} = E_I e^{iz} = E_I (\cos(z) + i \sin(z)) \\ E_{R,m} = E_R e^{i(\frac{\omega z}{c} + \pi)} = E_R [\cos(z+\pi) + i \sin(z+\pi)] \cdot \Gamma_{I,R,\parallel} \end{cases} = \boxed{\Gamma_{I,R,\parallel}}$$

→ se usarmos só a parte real:

$$\therefore \begin{cases} \text{Re}(E_{I,m}) = E_I \cos(z) \\ \text{Re}(E_{R,m}) = E_R \cos(z+\pi) \cdot \Gamma_{I,R,\parallel} \end{cases}$$

QUESTÃO 04) Eq. Fresnel: $\left\{ \left(\frac{E_{R,m}}{E_{I,m}} \right)_{||} = \frac{n_1 \cos \theta_r - n_2 \cos \theta_i}{n_1 \cos \theta_r + n_2 \cos \theta_i} \right.$

\rightarrow Se $\cos \theta_r = -i \sqrt{\sin^2 \theta_i - 1}$

$\therefore \left(\frac{E_{R,m}}{E_{I,m}} \right)_{||} = \frac{-(n_2/n_1) \cos \theta_i - i \sqrt{\sin^2 \theta_i - 1}}{(n_2/n_1) \cos \theta_i + i \sqrt{\sin^2 \theta_i - 1}} \quad (1)$

\rightarrow Usando a lógica do Capítulo 8 onde: $\left\{ a = -\frac{n_2}{n_1} \cos \theta_i ; b = \sqrt{\sin^2 \theta_i - 1} \right\}$

a Eq (1) muda pl essa forma:

$\therefore \left(\frac{E_{R,m}}{E_{I,m}} \right)_{||} = \frac{|a+ib| e^{iatg(b/a)}}{|a-ib| e^{-iatg(b/a)}} = e^{2iatg(b/a)}$

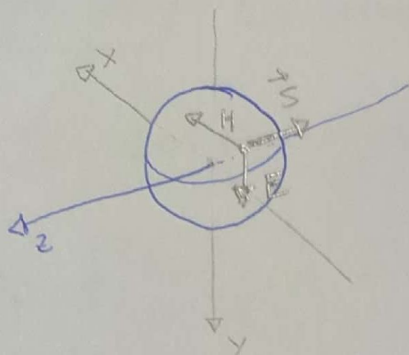
com isso $\left\{ \begin{aligned} |a+ib| &= \sqrt{\left(\frac{n_2}{n_1}\right)^2 \cos^2 \theta_i + \sin^2 \theta_i - 1} = M \\ |a-ib| &= \sqrt{\left(\frac{n_2}{n_1}\right)^2 \cos^2 \theta_i + \sin^2 \theta_i - 1} = M \end{aligned} \right.$

$\therefore \left(\frac{E_{R,m}}{E_{I,m}} \right)_{||} = \frac{M e^{itg\left(\frac{-\sqrt{\sin^2 \theta_i - 1}}{(n_2/n_1) \cos \theta_i}\right)}}{M e^{itg\left(\frac{-\sqrt{\sin^2 \theta_i - 1}}{-(n_2/n_1) \cos \theta_i}\right)}} = e^{-2itg\left(\frac{\sqrt{\sin^2 \theta_i - 1}}{(n_2/n_1) \cos \theta_i}\right)}$

$\rightarrow \phi_{||} = -2tg\left(\frac{\sqrt{\sin^2 \theta_i - 1}}{(n_2/n_1) \cos \theta_i}\right)$

$\therefore \left(\frac{E_{R,m}}{E_{I,m}} \right)_{||} = e^{i\phi_{||}}$

05)



$$dQ = e \rho dV = e \rho n^2 \sin \phi d\phi d\theta dn$$

FORÇA de LORENTZ:

$$dF = dQ (E_T + v \times B)$$

$$= dQ (E_T + \frac{\sigma}{\rho e} E_T \times \mu_0 H_T)$$

$$= dQ (E_T + \frac{\mu_0 \sigma}{\rho e} E_T \times H_T)$$

$$= dQ (E_T \hat{y} + \frac{\mu_0 \sigma}{\rho e} E_T H_T \hat{z})$$

$$= dQ \frac{\mu_0 \sigma}{\rho e} E_T H_T \hat{z}$$

$$= e \rho dV \frac{\mu_0 \sigma}{\rho e} E_T H_T \hat{z}$$

$$= \mu_0 \sigma E_T H_T \hat{z} dV \quad \Rightarrow \quad \underline{\mu_0 \sigma \vec{S} dV}$$

onde, $\vec{v} = \frac{\sigma E_T}{\rho e} \hat{y}$

$$\int dV = \int n^2 \sin \phi d\phi d\theta dn$$

$$= \int n^2 \sin \phi d\phi d\theta dn$$

$$= 4\pi \int_0^R n^2 dn$$

$$= \frac{4\pi R^3}{3}$$

$$\therefore F = \int \mu_0 \sigma E_T H_T \hat{z} dV = \mu_0 \sigma E_T H_T \hat{z} \int dV$$

$$= \mu_0 \sigma E_T H_T \hat{z} \left[\frac{4\pi R^3}{3} \right]$$

$$(E_T \hat{y}) \times (H_T \hat{x}) = E_T H_T \hat{z}$$

→ MAS
VETOR
POYNTING: $\vec{S} = \vec{E} \times \vec{H}$

$$\boxed{F_{\text{Radiação}} = \frac{4\pi R^3}{3} \mu_0 \sigma E_T H_T \hat{z}} \quad \text{ou} \quad \boxed{F = \frac{4\pi R^3}{3} \mu_0 \sigma \vec{S}}$$

DENSIDADE DE
FLUXO
DE MOMENTO

$\gamma = \frac{F}{A}$, Relacionando
com a Força
de Incidência $\Rightarrow \gamma = \frac{\bar{F}}{A}$, onde $\bar{F} \rightarrow$ Força
MÉDIA

$$\therefore \bar{F} = \frac{1}{2} \text{Re}(F) = \frac{1}{2} \text{Re} \left(\frac{4\pi R^3}{3} \mu_0 \sigma \vec{S} \right) = \frac{2\pi R^3}{3} \mu_0 \sigma \text{Re}(S) \quad (1)$$

$$\rightarrow \text{Re}(S) = \text{Re}(E \times H) = \text{Re}(E H^* \hat{z}) = \text{Re}(E_m H_m e^{-2\alpha z} \cos \pi/4) = \frac{\sqrt{2}}{2} E_m H_m e^{-2\alpha z}$$

$$(1) \Rightarrow \bar{F} = \frac{2\pi R^3}{3} \mu_0 \sigma \cdot \frac{\sqrt{2}}{2} E_m H_m e^{-2\alpha z} = \frac{2\sqrt{2}}{3} \pi R^3 \mu_0 \sigma |E_{rms}|^2 e^{-2\alpha z}$$

PARA BOM
CONDUCTOR

$$\therefore \gamma = \frac{\bar{F}}{A} = \frac{\left(\frac{2\sqrt{2}}{3} \right) \pi R^3 \mu_0 \sigma |E_{rms}|^2 e^{-2\alpha z}}{(4\pi R^2/4)} \Rightarrow \left| \frac{\sqrt{2}}{3} R \mu_0 \sigma |E_{rms}|^2 e^{-2\alpha z} = \gamma \right|$$

QUESTÃO 06 | Dado: $\vec{\sigma}_{ij} = \epsilon_0 \left(E_i E_j - \frac{1}{2} \delta_{ij} E^2 \right) + \frac{1}{\mu_0} \left(B_i B_j - \frac{1}{2} \delta_{ij} B^2 \right)$

$$\sigma_{xx} = \epsilon_0 \left(E_x E_x - \frac{1}{2} E^2 \right) + \frac{1}{\mu_0} \left(B_x B_x - \frac{1}{2} B^2 \right) = \epsilon_0 \left(E_x^2 - \frac{E^2}{2} \right) + \frac{1}{\mu_0} \left(B_x^2 - \frac{1}{2} B^2 \right)$$

\rightarrow So que $\vec{E} = (E_x, 0, 0)$; $\vec{B} = (0, B_y, 0)$;

$$\therefore \sigma_{xx} = \epsilon_0 \left(E^2 - \frac{1}{2} E^2 \right) + \frac{1}{\mu_0} \left(0 - \frac{1}{2} B^2 \right) = \left[\frac{\epsilon_0}{2} E^2 - \frac{B^2}{2\mu_0} = \sigma_{xx} \right]$$

$$\sigma_{xy} = \epsilon_0 \left(E_x E_y - \frac{1}{2} \delta_{xy} E^2 \right) + \frac{1}{\mu_0} \left(B_x B_y - \frac{1}{2} \delta_{xy} B^2 \right) = 0$$

\hookrightarrow Obs: Como os campos \vec{E} e \vec{B} dependem de apenas 1 coordenada, OS TERMOS CRUZADOS: $\{ E_i E_j, B_i B_j \text{ e } \delta_{ij} \}$ não todos iguais a ZERO.

$$\therefore i \neq j \Rightarrow E_i E_j = 0 = B_i B_j = \delta_{ij} = 0$$

\hookrightarrow Por isso: $\sigma_{xy} = \sigma_{xz} = \sigma_{yx} = \sigma_{yz} = \sigma_{zx} = \sigma_{zy} = 0$

$$\begin{aligned} \therefore \sigma_{yy} &= \epsilon_0 \left(E_y E_y - \frac{1}{2} E^2 \right) + \frac{1}{\mu_0} \left(B_y B_y - \frac{1}{2} B^2 \right) \\ &= \epsilon_0 \left(0 - \frac{1}{2} E^2 \right) + \frac{1}{\mu_0} \left(B^2 - \frac{1}{2} B^2 \right) \Rightarrow \boxed{\sigma_{yy} = -\frac{\epsilon_0}{2} E^2 + \frac{B^2}{2\mu_0}} \end{aligned}$$

$$\begin{aligned} \therefore \sigma_{zz} &= \epsilon_0 \left(E_z E_z - \frac{1}{2} E^2 \right) + \frac{1}{\mu_0} \left(B_z B_z - \frac{1}{2} B^2 \right) \\ &= \epsilon_0 \left(0 - \frac{1}{2} E^2 \right) + \frac{1}{\mu_0} \left(0 - \frac{1}{2} B^2 \right) \end{aligned}$$

$$\Rightarrow \boxed{\sigma_{zz} = -\frac{\epsilon_0}{2} E^2 - \frac{B^2}{2\mu_0}}$$