

数值代数中 Gummel 迭代方法的补充研究

报告人: 许笑颜

中国科学技术大学,数据科学(数学)

2023 年 12 月 28 日



- 广义 Gummel 线性化
- 2 MSP-Gummel 与近似牛顿迭代法的比较

Semiconductor Equation(With Density Gradient)



考虑引入量子修正效应的半导体器件方程:

$$\begin{cases} f_{1}: & \varepsilon \Delta \psi + q(p - n + N_{D}^{+} - N_{A}^{-}) + \rho_{s} = 0, \\ f_{2}: & \frac{1}{q} \nabla \cdot J_{n} - (U - G) = 0, \\ f_{3}: & -\frac{1}{q} \nabla \cdot J_{p} - (U - G) = 0, \\ f_{4}: & b_{n}(\Delta \ln n + \frac{1}{2} \varepsilon (\ln n)^{2}) - \Lambda_{n} = 0 \\ f_{5}: & b_{p}(\Delta \ln p + \frac{1}{2} \varepsilon (\ln p)^{2}) - \Lambda_{p} = 0. \end{cases}$$
(1)

现在已知第 k 次迭代时自变量 $(\psi,\phi_n,\phi_p,\Lambda_n,\Lambda_p)$ 的值 $(\psi_k,\phi_{n,k},\phi_{p,k},\Lambda_{n,k},\Lambda_{p,k})$,要 迭代求出第 k+1 次迭代时自变量 $(\psi,\phi_n,\phi_p,\Lambda_n,\Lambda_p)$ 的值 $(\psi_{k+1},\phi_{n,k+1},\Lambda_{p,k+1},\Lambda_{n,k+1},\Lambda_{p,k+1})$ 使以下方程组的残差接近于 0。

$$f_i(\psi_{k+1}, \phi_{n,k+1}, \phi_{p,k+1}, \Lambda_{n,k+1}, \Lambda_{p,k+1}) = 0, i = 1, 2, 3, 4, 5.$$



假设准费米势 ϕ_n, ϕ_p ,量子修正 Λ_n, Λ_p 是静电势 ψ 的函数:

$$n = n(\psi, \phi_n, \Lambda_n) \tag{2}$$

$$dn = \frac{\partial n}{\partial \psi} d\psi + \frac{\partial n}{\partial \phi_n} d\phi_n + \frac{\partial n}{\partial \Lambda_n} d\Lambda_n = \frac{\partial n}{\partial \psi} d\psi + \frac{\partial n}{\partial \phi_n} \frac{d\phi_n}{d\psi} d\psi + \frac{\partial n}{\partial \Lambda_n} \frac{d\Lambda_n}{d\psi} d\psi \tag{3}$$

$$= \frac{\partial n}{\partial \psi} d\psi + \frac{\partial n}{\partial \phi_n} \alpha_n^k d\psi + \frac{\partial n}{\partial \Lambda_n} \beta_n^k d\psi = \left(\frac{\partial n}{\partial \psi} + \frac{\partial n}{\partial \phi_n} \alpha_n^k + \frac{\partial n}{\partial \Lambda_n} \beta_n^k \right) d\psi \tag{4}$$

原 Possion 方程: $-\frac{\varepsilon \Delta \psi^{k+1}}{q} - p^{k+1} + n^{k+1} - N_D^+ + N_A^- + \frac{\rho_s}{q} = 0$. 普通 Gummel 迭代: 令 p^{k+1}, n^{k+1} 分别在 p^k, n^k 处对 ψ 一阶 Taylor 展开:

$$-\frac{\varepsilon \Delta \psi^{k+1}}{q} - \left(p^k + \frac{\partial p}{\partial \psi} \left(\psi^{k+1} - \psi^k\right)\right) + \left(n^k + \frac{\partial n}{\partial \psi} \left(\psi^{k+1} - \psi^k\right)\right) - N_D^+ + N_A^- + \frac{\rho_s}{q} = 0.$$
(5)



令 p^{k+1} , n^{k+1} 分别在 p^k , n^k 处对 ψ , ϕ_n , ϕ_p , Λ_n , Λ_p 一阶 Taylor 展开:

$$n^{k+1} = n^{k} \left(\psi^{k}, \phi_{n}^{k}, \Lambda_{n}^{k} \right) + \frac{\partial n}{\partial \psi} \left(\psi^{k+1} - \psi^{k} \right) + \frac{\partial n}{\partial \phi_{n}} \left(\phi_{n}^{k+1} - \phi_{n}^{k} \right) + \frac{\partial n}{\partial \Lambda_{n}} \left(\Lambda_{n}^{k+1} - \Lambda_{n}^{k} \right)$$

将
$$dn = \left(\frac{\partial n}{\partial \psi} + \frac{\partial n}{\partial \phi_n} \alpha_n^k + \frac{\partial n}{\partial \Lambda_n} \beta_n^k\right) d\psi$$
 代入:

$$\begin{split} & n^{k+1} = n^k \left(\psi^k, \phi_n^k, \Lambda_n^k \right) + \left(\frac{\partial n}{\partial \psi} + \frac{\partial n}{\partial \phi_n} \alpha_n^k + \frac{\partial n}{\partial \Lambda_n} \beta_n^k \right) d\psi, \\ & - \frac{\varepsilon \Delta \psi^{k+1}}{q} + n^k - p^k + \left(\frac{\partial n}{\partial \psi} - \frac{\partial p}{\partial \psi} \right) (\psi^{k+1} - \psi^k) + C + \\ & \left(\alpha_n^k \frac{\partial n}{\partial \phi_n} - \alpha_p^k \frac{\partial p}{\partial \phi_p} + \beta_n^k \frac{\partial n}{\partial \Lambda_n} - \beta_p^k \frac{\partial p}{\partial \Lambda_p} \right) (\psi^{k+1} - \psi^k) = 0 \end{split}$$



更新策略:

$$\begin{split} \phi_n^{k+1} &= \phi_n^k + \alpha_n^k (\psi^{k+1} - \psi^k), \quad \phi_p^{k+1} = \phi_p^k + \alpha_p^k (\psi^{k+1} - \psi^k), \\ \Lambda_n^{k+1} &= \Lambda_n^k + \beta_n^k (\psi^{k+1} - \psi^k), \quad \Lambda_p^{k+1} = \Lambda_p^k + \beta_p^k (\psi^{k+1} - \psi^k). \\ n_{k+1} &= n(\psi^{k+1}, \phi_n^{k+1}, \Lambda_n^{k+1}), \quad p_{k+1} = p(\psi^{k+1}, \phi_p^{k+1}, \Lambda_p^{k+1}), \\ R_{k+1} &= R(\psi^{k+1}, \phi_n^{k+1}, \phi_p^{k+1}, \Lambda_n^{k+1}, \Lambda_p^{k+1}). \end{split}$$

 $\alpha_p^{k+1}, \alpha_n^{k+1}, \beta_n^{k+1}, \beta_p^{k+1}$ 的更新策略:

$$dn = \left(\frac{\partial n}{\partial \psi} + \frac{\partial n}{\partial \phi_n} \alpha_n^k + \frac{\partial n}{\partial \Lambda_n} \beta_n^k\right) d\psi \Rightarrow \alpha_n^k = \frac{dn/d\psi - \partial n/\partial \psi - \partial n/\partial \Lambda_n \cdot \beta_n^k}{\partial n/\partial \phi_n}$$
$$\beta_n^k = \frac{dn/d\psi - \partial n/\partial \psi - \partial n/\partial \phi_n \cdot \alpha_n^k}{\partial n/\partial \Lambda_n}$$

广义 Gummel-算法流程



1. 对于 ψ 问题: 求解 ψ^{k+1} 使得

$$\begin{split} &-\frac{\varepsilon\Delta\psi^{k+1}}{q}+\mathbf{n}^k-\mathbf{p}^k+\left(\frac{\partial\mathbf{n}}{\partial\psi}-\frac{\partial\mathbf{p}}{\partial\psi}\right)(\psi^{k+1}-\psi^k)+\mathit{C}+\\ &\left(\alpha_{\mathbf{n}}^k\frac{\partial\mathbf{n}}{\partial\phi_{\mathbf{n}}}-\alpha_{\mathbf{p}}^k\frac{\partial\mathbf{p}}{\partial\phi_{\mathbf{p}}}+\beta_{\mathbf{n}}^k\frac{\partial\mathbf{n}}{\partial\Lambda_{\mathbf{n}}}-\beta_{\mathbf{p}}^k\frac{\partial\mathbf{p}}{\partial\Lambda_{\mathbf{p}}}\right)(\psi^{k+1}-\psi^k)=0. \end{split}$$

2. 对于 ϕ, Λ 问题,更新准费米势以及量子修正:

$$\phi_{n}^{k+1} = \phi_{n}^{k} + \alpha_{n}^{k}(\psi^{k+1} - \psi^{k}), \quad \phi_{p}^{k+1} = \phi_{p}^{k} + \alpha_{p}^{k}(\psi^{k+1} - \psi^{k}),$$

$$\Lambda_{n}^{k+1} = \Lambda_{n}^{k} + \beta_{n}^{k}(\psi^{k+1} - \psi^{k}), \quad \Lambda_{p}^{k+1} = \Lambda_{p}^{k} + \beta_{p}^{k}(\psi^{k+1} - \psi^{k}).$$

3. 更新载流子浓度以及 R:

$$n_{k+1} = n(\psi^{k+1}, \phi_n^{k+1}, \Lambda_n^{k+1}), \quad p_{k+1} = p(\psi^{k+1}, \phi_p^{k+1}, \Lambda_p^{k+1}),$$

$$R_{k+1} = R(\psi^{k+1}, \phi_n^{k+1}, \phi_p^{k+1}, \Lambda_n^{k+1}, \Lambda_p^{k+1}).$$



- 4. for $i \in \mathcal{I}$
 - 4.1 设 $\nabla \cdot \mathbf{J}_n qR = 0$, $\nabla \cdot \mathbf{J}_p + qR = 0$ 的离散格式,分别记为 f_n^I 以及 f_p^I ,及量子修正方程的离散格式,分别记为 f_n^Q 以及 f_n^Q 。
 - 4.2 更新节点 i 处的 n,p:

$$\begin{split} n_i^* &= f_n^J(N(i), n^{k+1}, R^{k+1}), \ p_i^* &= f_p^J(N(i), p^{k+1}, R^{k+1}), \\ \Lambda_{n,i}^* &= f_n^Q(N(i), n^{k+1}), \ \Lambda_{p,i}^* &= f_p^Q(N(i), p^{k+1}). \end{split}$$

5. 更新离散节点的 $\alpha_n, \alpha_p, \beta_n, \beta_p$:

$$\beta_{n}^{k+1} = \frac{\Lambda_{n}^{*} - \Lambda_{n}^{k}}{\psi^{k+1} - \psi^{k}} \simeq \frac{d\Lambda_{n}}{d\psi}, \alpha_{n}^{k+1} = \frac{\frac{n^{*} - n^{k}}{\psi^{k+1} - \psi^{k}} - \partial n/\partial \psi - \partial n/\partial \Lambda_{n} \cdot \beta_{n}^{k+1}}{\partial n/\partial \phi_{n}}$$
$$\beta_{p}^{k+1} = \frac{\Lambda_{p}^{*} - \Lambda_{p}^{k}}{\psi^{k+1} - \psi^{k}} \simeq \frac{d\Lambda_{p}}{d\psi}, \alpha_{p}^{k+1} = \frac{\frac{p^{*} - p^{k}}{\psi^{k+1} - \psi^{k}} - \partial p/\partial \psi - \partial p/\partial \Lambda_{p} \cdot \beta_{n}^{k+1}}{\partial p/\partial \phi_{p}}$$

MSP-Gummel 与近似牛顿迭代法的比较



观察方程形式:

$$\begin{cases}
-\frac{\varepsilon}{q}\Delta\psi + n - p + C = 0 \\
\nabla \cdot \mathbf{J}_n - qR = 0, \mathbf{J}_n = -q\mu_n n\nabla\phi_n \\
\nabla \cdot \mathbf{J}_p + qR = 0, \mathbf{J}_p = -q\mu_p p\nabla\phi_p
\end{cases} \tag{6}$$

- (1) 考虑对 ϕ_p^{k+1} , $d\phi_p$ 的求解:
 - **近似牛顿**: $\nabla \cdot [\mu_p p(\nabla d\phi_p + \nabla \phi_p)] = R + \frac{\partial R}{\partial \phi_p} d\phi_p.$
 - ► MSP-Gummel: $\nabla \cdot [\mu_p p \nabla \phi_p^{k+1}] = -qR$.
- (2) 考虑对 ψ_1 , $d\psi_1$ 的求解:
 - **上版似牛顿**: $-\frac{\varepsilon}{q} \left(\Delta d\psi_1 + \Delta \psi \right) + n p + C \frac{\partial p}{\partial \psi} d\psi_1 \frac{\partial p}{\partial \phi_p} = 0$
 - ► MSP-Gummel: $-\frac{\varepsilon}{q}\Delta\psi^{k+1} + n p + C = 0$.

MSP-Gummel 与近似牛顿迭代法的比较



- (3) 考虑对 $d\phi_n, \phi_n^{k+1}$ 的求解:
 - ▶ 近似牛顿:

$$\nabla \cdot [\mu_n n(\nabla d\phi_n + \nabla \phi_n)] = -R - \frac{\partial R}{\partial \phi_n} d\phi_n - \frac{\partial R}{\partial \phi_p} d\phi_p - \mu_n \nabla \cdot [\frac{\partial n}{\partial \psi} \nabla \phi_n d\psi_1].$$

- (4) 考虑对 $d\psi, \psi^{k+1}$ 的求解:
 - **上版似牛顿**: $-\frac{\varepsilon}{q}(\Delta d\psi + \Delta \psi) + n p + C = \frac{\partial p}{\partial \phi_p} d\phi_p \frac{\partial n}{\partial \phi_n} d\phi_n \frac{\partial n}{\partial \psi} d\psi + \frac{\partial p}{\partial \psi} d\psi$.
 - ► MSP-Gummel: $-\frac{\varepsilon}{q}\Delta\psi^{k+1} + n p + C = 0$.



Thanks for the help of X.H. Liu!

- C. Ringhofer and C. Schmeiser, "An approximate newton method for the solution of the basic semiconductor device equations," SIAM journal on numerical analysis, vol. 26, no. 3, p. 507–516, 1989.
- C. D. Falco, J. W. Jerome, and R. Sacco, "Quantum-corrected driftdiffusion models: Solution fixed point map and finite element approximation," Journal of Computational Physics, vol. 228, no. 5, pp. 17701789, 2009.
- M.S. Obrecht, M.I. Elmasry, and E.L. Heasell. Trasim: compact and efficient two-dimensional transient simulator for arbitrary planar semiconductor devices. IEEE Transactions on Computer-Aided Design of Integrated Circuits and Systems, 14(4):447–458, 1995.