## Problem Set # 1

Math262a: Quantum theory from a geometric viewpoint I

These problems are supplementary to the lectures and the lecture notes. I hope you tackle many of them, even if you don't complete them. Many are open-ended. They provide a gateway to engage with the material in the class. I suggest you form small groups to discuss the class material, including these problems. Those getting a grade in the course should regularly hand in some solutions. Please come to office hours and the discussion section to discuss these problems, and also take advantage of Discord as a platform for discussion.

## **Problems**

- In this problem you will derive the symplectic form on the space of classical trajectories of a particle, for simplicity a particle of mass m moving on the standard Euclidean line E¹ with coordinate x.
   Let V: E¹ → R be a smooth function, which is the potential energy of the particle. Let F = Map(R, E¹) be the space of smooth possible particle trajectories.
  - (a) What is the space  $N \subset \mathcal{F}$  of classical trajectories, i.e., trajectories that satisfy Newton's law? Consider the special cases V = 0 and  $V(x) = kx^2/2$  for some k > 0.
  - (b) Review the derivation of Newton's law from the calculus of variations (Euler-Lagrange equations) applied to the *action* function

$$S(\gamma) = \int \left[ \frac{1}{2} m \dot{x}(t)^2 - V(x(t)) \right] dt$$

where  $x: \mathbb{R} \to \mathbb{E}^1$  is a smooth motion.

(c) Rephrase this computation using calculus on  $\mathcal{F} \times \mathbb{R}$ . Observe that the Euler-Lagrange equations only depend integrand in the action—the lagrangian—not its integral. Write  $\delta$  for the de Rham differential on  $\mathcal{F}$  and d for the de Rham differential  $d = dt \cdot \partial/\partial t$  on  $\mathbb{R}$ . Then the total de Rham differential on  $\mathcal{F} \times \mathbb{R}$  is  $\delta + d$ . In particular,  $(\delta + d)^2 = 0$ . A differential form on  $\mathcal{F} \times \mathbb{R}$  has  $type\ (p,q),\ p \in \mathbb{Z}^{\geqslant 0},\ q \in \{0,1\}$ , if it is nonzero only when evaluated on p vectors tangent to  $\mathcal{F}$  and q vectors tangent to  $\mathbb{R}$ . Calculus on the Cartesian product  $\mathcal{F} \times \mathbb{R}$  is depicted in the diagram

(If one writes the lagrangian and action as a density rather than a differential form—use |dt| in place of dt—which we should do since we should not use an orientation on  $\mathbb{R}$  to define this

mechanical system, then the formatting of the square makes more sense: starting from the top, the vertical degrees are then written as '|0|' and '|-1|'.) Let

$$e: \mathfrak{F} \times \mathbb{R} \longrightarrow \mathbb{E}^1$$
 $x, t \longmapsto x(t)$ 

be the evaluation function. Show that the integrand above can be written as

$$L = \left[\frac{1}{2}m\dot{e}^2 - V \circ e\right]dt$$

and that it is an element of  $\Omega^{0,1}(\mathcal{F}\times\mathbb{R})$ . (Be sure to define  $\dot{e}$ , which is a directional derivative of e along a certain vector field on  $\mathcal{F} \times \mathbb{R}$ .)

- (d) Compute  $\delta L \in \Omega^{1,1}(\mathfrak{F} \times \mathbb{R})$ . Carry out the integration by parts in the following form: find  $\theta \in \Omega^{1,0}(\mathcal{F} \times \mathbb{R})$  so that  $\delta L + d\theta$  is a function times  $\delta e \wedge dt$ , i.e., it has no  $\delta \dot{e} \wedge dt$  term.
- (e) Show that  $\delta L + d\theta$  vanishes on N, and in fact that vanishing defines N.
- (f) Show that the restriction of  $\delta\theta \in \Omega^{2,0}(\mathcal{F} \times \mathbb{R})$  to  $N \times \{t_0\}$  is independent of  $t_0$ . Show that this restriction is a symplectic form. Identify it with the symplectic form you may have run into before in this context.
- 2. In the previous problem specialize to  $V(x) = kx^2/2$  for fixed k > 0.
  - (a) Fix  $t \in \mathbb{R}$  and consider the observables (functions on N) given by  $\mathcal{O}_1(x) = x(0)$  and  $\mathcal{O}_2(x) = x(t)$ . Choose a pure state  $\sigma$ , which is a point of N. Compute the expectation values  $\langle \mathcal{O}_1 \rangle_{\sigma}$  and  $\langle \mathcal{O}_2 \rangle_{\sigma}$ . Compute the correlation function  $\langle \mathcal{O}_2 \mathcal{O}_1 \rangle_{\sigma}$ . Investigate the dependence on t and on  $\sigma$ .
  - (b) Now consider the mixed Gibbs state, which is a probability measure on the space N of classical solutions. Its probability density is a constant times  $e^{-\beta E}$ , where  $\beta > 0$  is a constant and  $E: N \to \mathbb{R}$  is the energy function  $E = m\dot{x}^2/2 + kx^2/2$ . Compute the Gibbs state (a function times the standard measure on the  $(x, \dot{x})$ -plane) and evaluate the expectation values in part (a) in the Gibbs state.
  - 3. Verify the axioms in §1.2 of the lecture notes for the case of a classical mechanical system.
  - 4. In the context of a general mechanical system, let  $\sigma_1, \sigma_2$  be states and let A, A' be observables. Fix  $t \in [0,1]$ . Prove that for the convex combination  $\sigma = t\sigma_1 + (1-t)\sigma_2$  the standard deviations  $\Delta_{\rho}B$  of the probability distributions  $\rho_{B}$  obtained from the various states  $\rho$  and observables B satisfy

$$\Delta_{\sigma} A \, \Delta_{\sigma} A' \geqslant t(\Delta_{\sigma_1} A) \, (\Delta_{\sigma_1} A') + (1 - t)(\Delta_{\sigma_2} A) \, (\Delta_{\sigma_2} A')$$

In particular, conclude that

$$\Delta_{\sigma}A \geqslant \min_{2}(\Delta_{\sigma_{1}}A, \Delta_{\sigma_{2}}A).$$

(Recall that the standard deviation of a probability measure on  $\mathbb{R}$  is the square root of the integral of  $(\lambda - \mu)^2$  times the probability measure on the real line  $\mathbb{R}_{\lambda}$  with coordinate  $\lambda$ , where  $\mu$  is the *mean* or *expected value*, i.e.,  $\mu$  is the integral of  $\lambda$  times the probability measure.)