

Basics of Linear Algebra in context of Quantum Computing

A quantum state is a 'superposition' of classical states. Mathematically, it is expressed as a vector of amplitudes. Nature has ensured that these amplitudes are complex numbers.

1. Vector Spaces

} in our case the
field is complex,
i.e., C

def: A vector space V over a field F is a set of objects called vectors that satisfy the following

- closed under scalar multiplication
 - i.e., if $w \in V$ and $c \in F$
then $cw \in V$. This holds for all $w \in V$ and all $c \in F$.
- closed under addition
 - i.e., if $v, w \in V$ then $v + w \in V$.
This holds for all $v, w \in V$.

Notation: If V is a vector space of d dimension then $V = F^d$ which is the set of all column vectors of d elements from F .

For e.g. let $w \in V$ and $F = C$ — ^{set of} complex numbers
and $V = F^d$

then we can write $w = \begin{pmatrix} c_1 \\ c_2 \\ \vdots \\ c_d \end{pmatrix}$ where each $c_i \in C$.

Linearly Independent Vectors

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A set of vectors $v_1, \dots, v_m \in V$ is 'linearly independent' if $\sum_{i=1}^m x_i v_i = 0$ only when $x_1 = x_2 = \dots = x_m = 0$.

Span one field \mathbb{F} of set of vectors

$S = \{v_1, v_2, \dots, v_m\} \subseteq V$ is the set $\text{Span}(S)$

that can be written as a linear combination

$$\sum_{i=1}^m d_i v_i \quad \text{with coefficients } d_i \in \mathbb{F} \quad i \in [m].$$

*very important
to our context*

Basis for V is a linearly independent set S of vectors such that $\text{Span}(S) = V$.

Exercise 1: Can you give an example of a basis for $V = \mathbb{C}^d$, which is a vector space of d -dimension?

There is a general way to define inner product space for a vector space V . But we will consider complex field \mathbb{C} .

2. Inner Product Space

We could add an extra constraint to our vector space to have 'inner product' defined between every vector in the space.

Vector Space + Inner product = Hilbert Space

Inner product between two vectors $v, w \in V = \mathbb{C}^d$

$$\text{Suppose } v = \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_d \end{pmatrix} \quad w = \begin{pmatrix} w_1 \\ w_2 \\ \vdots \\ w_d \end{pmatrix}$$

Inner product between v & w =

$$\sum_{i=1}^d v_i^* w_i$$

v_i^ is complex conjugate of v_i*

But, inner product between w & v

$$= \sum_{i=1}^d w_i^* v_i$$

Note that as for our context $\mathbb{F} = \mathbb{C}$

It is not necessary that $\sum_{i=1}^d v_i^* w_i = \sum_{i=1}^d w_i^* v_i$

In fact $\langle v | w \rangle = \overline{\langle w | v \rangle}$ conjugate denoted by $\langle \cdot, \cdot \rangle$.

Home Work: Prove it.

so some notations that help us.

To denote a vector $v \in V$ we use $|v\rangle$

Ket Notation

to we use $\langle v |$ to denote the row vector corresponding to v with complex conjugate of each element coordinate in the vector $|v\rangle$

For example. $v \in V = \mathbb{C}^d$

$$v = \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_d \end{pmatrix}$$

$$\text{we use } |v\rangle = \begin{pmatrix} v_1 \\ \vdots \\ v_d \end{pmatrix}$$

$$\text{and use } \langle v | = (v_1^*, v_2^*, \dots, v_d^*)$$

$$\text{Therefore } \langle v | v \rangle = \sum_{i=1}^d v_i^* v_i = \|v\|^2$$

norm of v .

Hence, $\langle v | v \rangle = \text{square of norm of } v$.

Let's define a Qubit.

It is a d -dimensional vector in complex field.

with unit norm.

$$|\phi\rangle = \begin{pmatrix} \alpha_1 \\ \alpha_2 \end{pmatrix} \text{ where } \alpha_1, \alpha_2 \in \mathbb{C}$$

$$\langle \phi | \phi \rangle = 1 \Rightarrow \alpha_1^* \alpha_1 + \alpha_2^* \alpha_2 = 1 \Rightarrow |\alpha_1|^2 + |\alpha_2|^2 = 1$$

2 qubits

Consider the Basis for 1 qubit space
 $\{|0\rangle, |1\rangle\}$

Now two qubits can be written as.

linear combination of

$$\alpha_1|100\rangle + \alpha_2|01\rangle + \alpha_3|110\rangle + \alpha_4|111\rangle \text{ with } \sum_{i=1}^4 |\alpha_i|^2 = 1$$

where $|100\rangle = |10\rangle \otimes |0\rangle$, similarly for other elements as well.

Exercise 2:

Show that $\{|100\rangle, |101\rangle, |110\rangle, |111\rangle\}$

form a basis for ^{vector} space expressing 2 qubits.

Also show the basis as Basis vectors.

Exercise 3: Show that a state

$$|\Phi\rangle = \frac{1}{\sqrt{2}}|100\rangle + \frac{1}{\sqrt{2}}|111\rangle \text{ is entangled.}$$

i.e., $|\Phi\rangle \neq |\Psi_A\rangle \otimes |\Psi_B\rangle$ for any

$|\Psi_A\rangle$ and $|\Psi_B\rangle$ in \mathbb{C}^2 .

Exercise 4: show that there exists no unitary ^{Later exercise!}
 U that can clone a state $|\Phi\rangle$ in \mathbb{C}^2 .

i.e., show that it is not possible to have
 $a \ U \ s.t \quad U|\Phi\rangle|0\rangle \rightarrow |\Phi\rangle|\Phi\rangle$

Hint: Use linearity of U and entanglement ~~as~~ definition.

Homework: Show that there is no unitary U that can delete information.

i.e. No unitary U exists that does $U|\Phi\rangle|\Phi\rangle \rightarrow |\Phi\rangle|0\rangle$

The Unitary Matrices

To evolve a vector $v \in V \subset \mathbb{C}^d$ we use 'unitary' matrices of dimension $d \times d$.

A matrix U is unitary if $U^{-1} = U^T$

Inverse
of a
matrix

Complex Conjugate
+
Transpose

Some properties

1. $U^T U = I$ — easy to see from the defn.

2. $\langle Uv | Uw \rangle = \langle v | w \rangle$ for all $v, w \in V$
Preserves inner product

3. $\langle Uv | Uv \rangle = \langle v | v \rangle$ for all $v \in V$
Preserves norm

4. Linearity $U|\alpha u + \beta v\rangle = \alpha U|u\rangle + \beta U|v\rangle$ for all $\alpha, \beta \in \mathbb{C}$
 and $u, v \in V$.

Tensor Product

If we have two vector spaces of dimension d & d' respectively then $V = V_A \otimes V_B$ is a $d \cdot d'$ dimensional vector space.

Additionally if $v \in V_A$ and $w \in V_B$

$$|v\rangle \otimes |w\rangle = \begin{pmatrix} v_1 \\ \vdots \\ v_d \end{pmatrix} \otimes \begin{pmatrix} w_1 \\ \vdots \\ w_{d'} \end{pmatrix} = \begin{pmatrix} v_1 w_1 \\ v_1 w_2 \\ \vdots \\ v_1 w_{d'} \\ v_2 w_1 \\ v_2 w_2 \\ \vdots \\ v_d w_1 \\ v_d w_{d'} \end{pmatrix}$$

(d · d'
dimensional
vector.)

Examples of some cute single qubit unitaries

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Example of a 2 qubit unitary

$$\text{CNOT} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

$$\text{CNOT } |0b\rangle \rightarrow |0\bar{b}\rangle \quad \text{Here } b \in \{0, 1\}$$

$$\text{CNOT } |1b\rangle \rightarrow |1\bar{b}\rangle \quad \bar{b} = b \oplus 1 \\ b \in \{0, 1\}.$$

Exercise: How to create an EPR pair from $|00\rangle$ state.

$$\text{EPR pair: } \frac{1}{\sqrt{2}}|00\rangle + \frac{1}{\sqrt{2}}|11\rangle$$

$$\text{Exercise: } H \otimes H : |0\rangle|0\rangle \rightarrow ?$$

$$\text{Exercise: } H^{\otimes n} |0^n\rangle \xrightarrow{i \text{ is a string } \{0, 1\}^n} \frac{1}{\sqrt{2^n}} \sum_{j \in \{0, 1\}^n} (-1)^{\sum_{i,j} i_j} |j\rangle$$

Verify this with $n = 3$.

Measurement

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If we have a qubit $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$

Measuring this qubit in computational basis gives ~~it's~~ collapses the state $|0\rangle$ into $|0\rangle$ with probability $|\alpha|^2$.

And collapses $|\psi\rangle$ into $|1\rangle$ with probability $|\beta|^2$.

Recall that qubits have norm 1. Therefore the measurement outcome will always be either $|0\rangle$ or $|1\rangle$.

Similarly, if we had a 2 qubit state

$$|\psi\rangle = \alpha_1|100\rangle + \alpha_2|010\rangle + \alpha_3|101\rangle + \alpha_4|110\rangle$$

then measuring the 2 qubits ^{in computational basis} will give collapse

$|\psi\rangle$ into $|100\rangle$ with probability $|\alpha_1|^2$

$|010\rangle$ with probability $|\alpha_2|^2$

$|101\rangle$ with probability $|\alpha_3|^2$

and

$|110\rangle$ with probability $|\alpha_4|^2$.

But instead of measuring both the qubits ^{say} but instead you measure, the first one then

$$\begin{aligned} \text{the } |\psi\rangle &= |0\rangle (\alpha_1|10\rangle + \alpha_2|11\rangle) + |1\rangle (\alpha_3|10\rangle + \alpha_4|11\rangle) \\ &= \sqrt{(|\alpha_1|^2 + |\alpha_2|^2)} |0\rangle \left(\frac{\alpha_1|10\rangle + \alpha_2|11\rangle}{\sqrt{|\alpha_1|^2 + |\alpha_2|^2}} \right) + \sqrt{(|\alpha_3|^2 + |\alpha_4|^2)} |1\rangle \left(\frac{\alpha_3|10\rangle + \alpha_4|11\rangle}{\sqrt{|\alpha_3|^2 + |\alpha_4|^2}} \right) \end{aligned}$$

With probability $|\alpha_1|^2 + |\alpha_2|^2$ we get outcome $|0\rangle$ and state collapses to $|0\rangle \left(\frac{\alpha_1|10\rangle + \alpha_2|11\rangle}{\sqrt{|\alpha_1|^2 + |\alpha_2|^2}} \right)$

With probability $|\alpha_3|^2 + |\alpha_4|^2$ we get outcome $|1\rangle$ and state collapses to $|1\rangle \left(\frac{\alpha_3|10\rangle + \alpha_4|11\rangle}{\sqrt{|\alpha_3|^2 + |\alpha_4|^2}} \right)$

Teleportation Protocol

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Suppose we have two parties Alice & Bob.

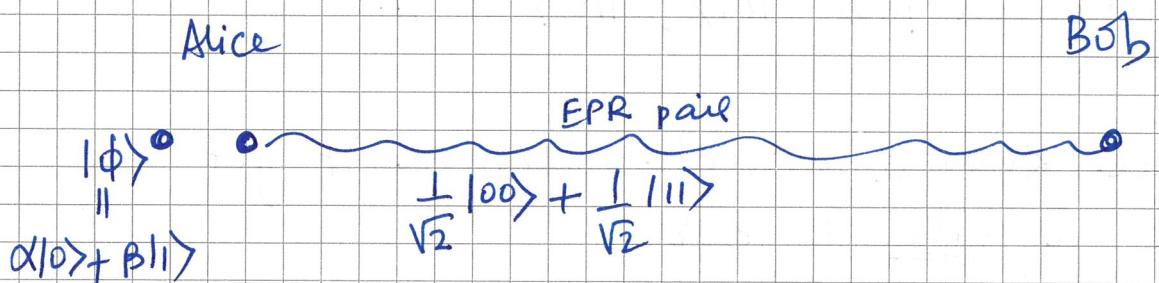
Alice has a qubit $| \phi \rangle = \alpha | 0 \rangle + \beta | 1 \rangle$ for some $\alpha, \beta \in \mathbb{C}$ but Alice doesn't know the values.

Alice wants to send $| \phi \rangle$ to Bob using classical bits

Additionally Alice & Bob share a EPR pair

$$\frac{1}{\sqrt{2}} | 00 \rangle + \frac{1}{\sqrt{2}} | 11 \rangle$$

Recall that this is an entangled state



Joint state of Alice and Bob

$$\frac{\alpha}{\sqrt{2}} \underset{A \otimes B}{| 000 \rangle} + * \frac{\alpha}{\sqrt{2}} \underset{A \otimes B}{| 011 \rangle} + \frac{\beta}{\sqrt{2}} \underset{A \otimes B}{| 100 \rangle} + \frac{\beta}{\sqrt{2}} \underset{A \otimes B}{| 111 \rangle}$$

↓ Alice performs CNOT

$$\frac{\alpha}{\sqrt{2}} | 000 \rangle + \frac{\alpha}{\sqrt{2}} | 011 \rangle + \frac{\beta}{\sqrt{2}} | 110 \rangle + \frac{\beta}{\sqrt{2}} | 101 \rangle$$

↓ Alice performs H on first qubit

$$\begin{aligned} & \frac{\alpha}{\sqrt{2}} \left(\frac{1}{\sqrt{2}} | 0 \rangle + \frac{1}{\sqrt{2}} | 1 \rangle \right) | 00 \rangle + \frac{\alpha}{\sqrt{2}} \frac{1}{\sqrt{2}} \left(| 0 \rangle + | 1 \rangle \right) | 11 \rangle + \frac{\beta}{\sqrt{2}} \frac{1}{\sqrt{2}} \left(| 0 \rangle - | 1 \rangle \right) | 10 \rangle \\ & \quad + \frac{\beta}{\sqrt{2}} \frac{1}{\sqrt{2}} \left(| 0 \rangle - | 1 \rangle \right) | 01 \rangle \end{aligned}$$

$$= \frac{\alpha}{2} (| 000 \rangle + | 100 \rangle + | 011 \rangle + | 111 \rangle) + \frac{\beta}{2} (| 101 \rangle - | 110 \rangle + | 001 \rangle - | 111 \rangle)$$

Rearranging the terms we get

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$$\begin{aligned} & \frac{1}{2} |00\rangle (\alpha|0\rangle + \beta|1\rangle) \\ & + \frac{1}{2} |01\rangle (\alpha|1\rangle + \beta|0\rangle) \\ & + \frac{1}{2} |10\rangle (\alpha|0\rangle - \beta|1\rangle) \\ & + \frac{1}{2} |11\rangle (\alpha|1\rangle - \beta|0\rangle) \end{aligned}$$

Now Alice Measures in computational basis with probability $\frac{1}{4}$ Alice gets one of the four outcomes $|00\rangle$, $|01\rangle$, $|10\rangle$ or $|11\rangle$

Once Alice knows which outcome it got it sends 2 bits to Bob.

Then Bob knows what to do.

For e.g. If Alice gets $|00\rangle$ then Bob has to do nothing. because Bob has $\alpha|0\rangle + \beta|1\rangle$ now.

If Alice gets $|01\rangle$, then Bob applies X on its qubit. and tells that to Bob

If Alice gets $|10\rangle$ and communicates the same to Bob then Bob also applies Z gate on its qubit.

If Alice gets $|11\rangle$ then Bob after knowing that applies X and Z .

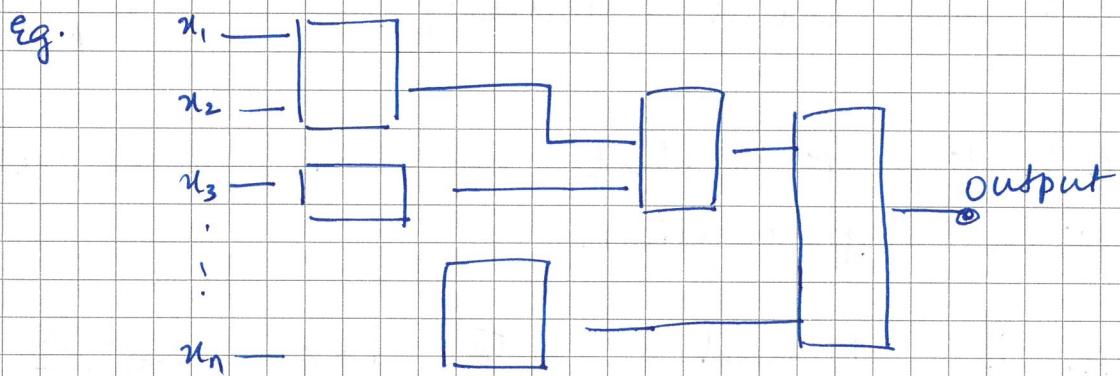
Ultimately Bob is able to get state $|0\rangle$.

Classical Circuits

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A Boolean circuit is a directed acyclic graph with AND, OR & NOT gates as internal nodes. It has n input nodes and some output nodes. Let us stick to 1 output.

We say a Boolean circuit C computes a Boolean function $f: \{0, 1\}^n \rightarrow \{0, 1\}$ if $C(x) = f(x) \ \forall x \in \{0, 1\}^n$.



Quantum Circuit also a directed acyclic graph contains ~~m~~ m wires but elementary quantum gates as we discussed earlier.

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

$$\text{CNOT} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} \quad I_2 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

But a slightly different model.

Instead of inputs in the wire as we see in the classical circuit we access the input through an 'Oracle' which is a unitary operation.

Bit Oracle

$$O_n |i, 0\rangle \rightarrow |i, u_i\rangle$$

more generally

$$O_n |i, b\rangle \rightarrow |i, b \oplus u_i\rangle$$

Exercise: Show O_n is a unitary.

where

my input is
 n and

$$u = u_1 u_2 \dots u_N$$

are N bits
of my
input.

The best part is that we can access the input in superposition.

for e.g. if I had a state $|\Psi\rangle = \frac{1}{\sqrt{2^n}} \sum_{i \in \{0,1\}^n} |i\rangle$

$$|\Psi\rangle = \frac{1}{\sqrt{2^n}} \sum_i |i\rangle$$

$$|\Psi\rangle = \frac{1}{\sqrt{2^n}} \sum_{i \in \{0,1\}^n} |i\rangle$$

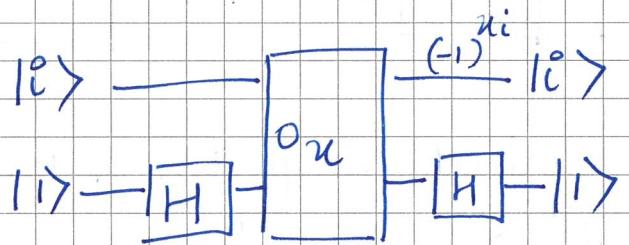
$$\text{Then } O_n |\Psi\rangle |b\rangle \rightarrow \frac{1}{\sqrt{2^n}} \sum_{i \in \{0,1\}^n} O_n |i\rangle |b\rangle$$

$$= \frac{1}{\sqrt{2^n}} \sum_{i \in \{0,1\}^n} |i\rangle |b \oplus u_i\rangle$$

Phase Oracle:

$$O_{n,i} |i\rangle \rightarrow (-1)^{x_i} |i\rangle$$

Let's implement a phase oracle from this bit oracle



$$|i\rangle |j\rangle \xrightarrow{\frac{I \otimes H}{\sqrt{2}}} \frac{|i\rangle (|0\rangle - |1\rangle)}{\sqrt{2}}$$

$\downarrow O_n$

$$\frac{O_n |i\rangle |0\rangle}{\sqrt{2}} \neq - \frac{O_n |i\rangle |1\rangle}{\sqrt{2}}$$

$$= \frac{1}{\sqrt{2}} (|i\rangle |x_i\rangle - |i\rangle |1 \oplus x_i\rangle)$$

$$= \frac{1}{\sqrt{2}} |i\rangle (|0\rangle - |1\rangle) \quad \text{if } x_i = 0$$

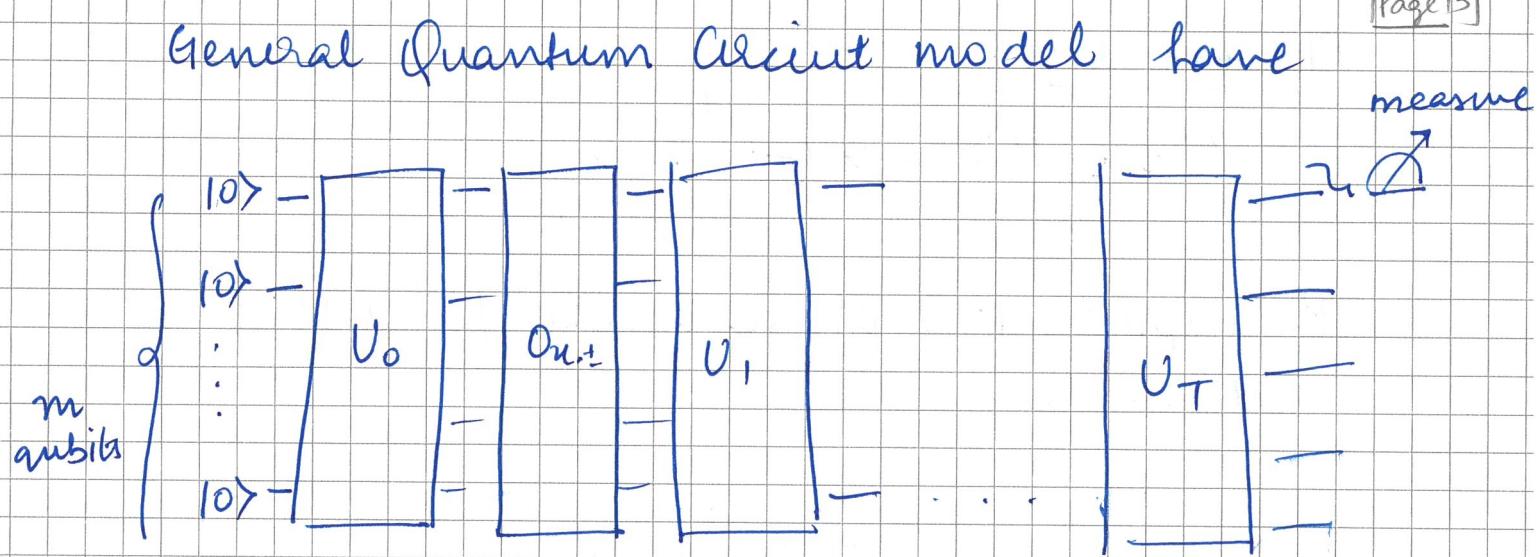
else

$$\frac{1}{\sqrt{2}} |i\rangle (-|0\rangle + |1\rangle) \quad \text{if } x_i = 1$$

$$= \frac{1}{\sqrt{2}} (-1)^{x_i} |i\rangle (|0\rangle - |1\rangle)$$

$\downarrow I \otimes H$

$$\frac{1}{\sqrt{2}} (-1)^{x_i} |i\rangle |j\rangle$$



If We say a circuit computes a Boolean function $f: \{0,1\}^n \rightarrow \{0,1\}$ iff

$$C(x) = f(x) \text{ with probability } \geq \frac{2}{3} \quad \forall x \in \{0,1\}^n$$

Our inputs are given via O_n , the oracle.

Deutsch - Jozsa Problem

Suppose $N = 2^n$, we are given an $x \in \{0, 1\}^N$ such that either

1. all x_i have same value ("constant" 0 or 1)
2. $\frac{N}{2}$ of the x_i have 0 and rest are 1 ("balanced").

We want to find out if x is constant or balanced.

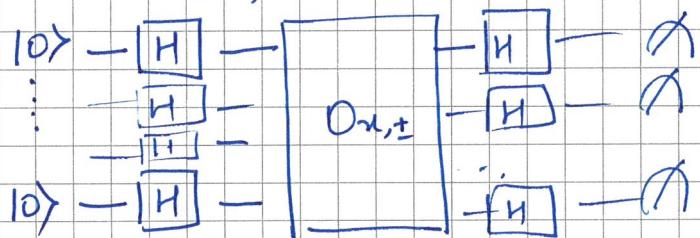
For example let's take $n=2 \Rightarrow N=4$

The possible inputs are { 0000, 1111, 0011, 0101, 1100, 1010, 0110, 1001 }

Q. Algorithm

We start with an all zero quantum state.

$$|0\rangle^{\otimes n} = \underbrace{|0\rangle |0\rangle \dots |0\rangle}_{n \text{ of them}}$$



Let's look at this algorithm

$$\begin{aligned}
 |0\rangle^{\otimes n} &\xrightarrow{H^{\otimes n}} \frac{1}{\sqrt{2^n}} \sum_{i \in \{0,1\}^n} |i\rangle \\
 &\xrightarrow{O_{n,\pm}} \frac{1}{\sqrt{2^n}} \sum_{i \in \{0,1\}^n} (-1)^{\sum_j u_j i_j} |i\rangle \\
 &\xrightarrow{H^{\otimes n}} \frac{1}{\sqrt{2^n}} \sum_{i \in \{0,1\}^n} (-1)^{\sum_{i,j} u_j i_j} |i\rangle
 \end{aligned}$$

$$\frac{1}{2^n} \sum_{i \in \{0,1\}^n} (-1)^{x_i} = \begin{cases} 1 & \text{if } x_i = 0 \text{ for all } i \\ -1 & \text{if } x_i = 1 \text{ for all } i \\ 0 & \text{if } x \text{ is balanced.} \end{cases}$$

which means by looking at the final state on n qubits we can ~~it~~ conclude with certainty if our given x was constant or balanced.

How many queries did we use? Just one!
 (But in superposition).

Homework: Verify the algorithm for $n=3$ case.

A classical deterministic algorithm needs about $\frac{N}{2} + 1$ queries