

QC-Module

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Contents

1	DSE	1
2	BSE	2

1 DSE

The Dyson-Schwinger equation (DSE) in the Euclidean space takes the following form

$$S^{-1}(p) = Z_2 (i\gamma \cdot p + m_f(\Lambda^2)) + \int_q^\Lambda g^2 D_{\mu\nu}(p-q) \frac{\lambda^a}{2} \gamma_\mu S(q) \Gamma_\nu^a(q, p). \quad (1)$$

The dressed-quark propagator takes the form

$$S^{-1}(p) = \frac{1}{Z(p^2, \zeta^2)} (i\gamma \cdot p + M(p^2, \zeta^2)) = i\gamma \cdot p A(p^2, \zeta^2) + B(p^2, \zeta^2). \quad (2)$$

Then we apply the rainbow approximation

$$\Gamma_\nu^a = \gamma_\nu \frac{\lambda^a}{2}, \quad (3)$$

and the Qin-Chang form

$$g^2 D_{\mu\nu}(k) = \mathcal{G}(k^2) T_{\mu\nu}(k), \quad (4)$$

where

$$k^2 T_{\mu\nu}(k) = k^2 \delta_{\mu\nu} - k_\mu k_\nu \quad (5)$$

and

$$\frac{1}{Z_2^2} \mathcal{G}(k^2) = \frac{8\pi^2 D}{\omega^4} e^{-\frac{k^2}{\omega^2}} + \frac{8\pi^2 \gamma_m \mathcal{F}(k^2)}{\ln[\tau + (1 + \frac{k^2}{\Lambda_{QCD}^2})^2]}. \quad (6)$$

The parameters are listed below:

$$\begin{aligned} D &= 1.024 \text{GeV}^2, \\ \omega &= 0.5 \text{GeV}, \\ m_t &= 0.5 \text{GeV}, \\ \tau &= e^2 - 1, \\ n_f &= 4, \\ \Lambda_{QCD} &= 0.234 \text{GeV}, \\ \gamma_m &= 12/(33 - 2n_f). \end{aligned} \quad (7)$$

The results are shown in Fig.1.

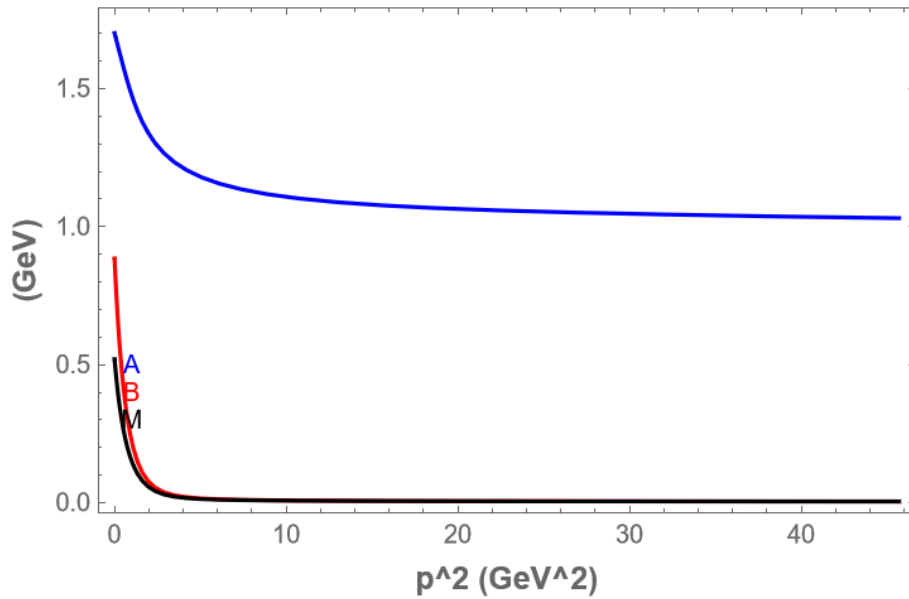


Figure 1: The function A, B and the quark mass function M.

2 BSE

The quark-antiquark bound state amplitude can be obtained by solving a homogeneous BSE. Employing the ladder truncation becomes

$$\Gamma^{f\tilde{g}}(p, P) = -Z_2^2 \int_q^\Lambda \mathcal{G}(k^2) T_{\mu\nu}(k) \frac{\lambda^a}{2} \gamma_\mu S^f(q_+) \Gamma^{f\tilde{g}}(q; P) S^g(q_-) \frac{\lambda^a}{2} \gamma_\nu. \quad (8)$$

Solving the eigenvalue equation

$$\lambda(P^2) \vec{F}_X(P^2) = \mathcal{K}_X(P^2) \vec{F}_X(P^2), \quad (9)$$

The solution at $\lambda(P^2 = -M^2) = 1$ corresponds to meson bound state. The results are shown in Fig.2.

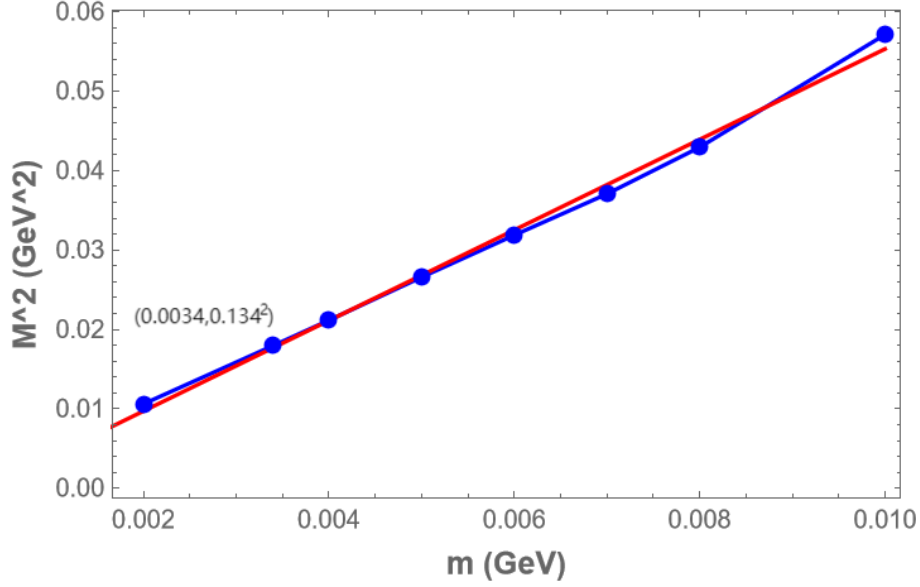


Figure 2: The blue points are directly calculated points.

This figure shows the Gell-Mann-Oakes-Renner relation.