## Local Boundedness — De Giorgi and Moser

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#### Setup

In the following we consider a weak sub-solution  $u \in H_0^1(B_1)$  of the PDE

$$Lu + cu < f$$

where

$$Lu = \sum_{i,j} D_i(a_{ij}D_j).$$

In other words, for any non-negative  $\phi \in H_0^1(B_1)$ ,

$$\int a_{ij} D_i u D_j \phi + c u \phi \le \int f \phi \tag{1}$$

In the following, we assume  $\{a_{ij}(x)\}_{ij}$  is bounded and satisfies an ellipticity condition,

$$a_{ij}(x)\xi_i\xi_j \ge \lambda |\xi|^2 > 0,$$

and  $c, f \in L^q$  for  $q > \frac{n}{2}$ . We also take  $\Lambda$  so that

$$|a_{ij}|_{\infty} + ||c||_q \leq \Lambda.$$

**Remark.** The requirement that  $q > \frac{n}{2}$  is a bare-minimum to guarantee the integrals in 1 are finite. Indeed, by Hölder,

$$||cu\phi|| \le ||c||_q ||u\phi||_{1-\frac{1}{q}} \le ||f||_q ||u||_{2-\frac{2}{q}} ||\phi||_{2-\frac{2}{q}}$$

If we only know  $u, \phi \in H^1$ , we need  $2(1 - \frac{1}{q}) < \frac{2n}{n-2}$  to make use of the Sobolev inequality, or  $q > \frac{n}{2}$ .

Theorem 1 (De Giorgi [1], Nash [4], Moser [3]). With the above conditions, assume u satisfies the above weak form. Write  $u^+$  as the positive part of u. Then for some constant C depending only on  $n, \lambda, \Lambda, p$ ,

$$\sup_{B_{\theta}} u^{+} \leq C(n, \lambda, \Lambda, p) \left\{ \frac{1}{(1-\theta)^{n/p}} \|u\|_{L^{p}(B_{1})} + \|f\|_{q} \right\}.$$

This note will sketch the proof of this classical result for the case p=2,  $\theta=1/2$  using Moser's and De Giorgi's methods. For the extension to all  $p \in (0, \infty)$  and  $\theta \in (0, 1)$ , and for many more (important and instructive) details, see section 4 of [2].

# Moser's Approach

## The Homogeneous Case

Let's simplify things and assume f = 0 and c = 0,

$$\int a_{ij} D_i u D_j \phi \le 0 \tag{2}$$

This was the original equation studied in the papers referenced above.

Suppose  $u \geq 0$ ; otherwise replace u with its positive part. Plugging in  $\phi = \eta^2 u$  and applying both boundedness and ellipticity of  $a_{ij}$ , we obtain the estimate

$$\int |D(u\eta)|^2 \le C \left\{ \int |D\eta|^2 u^2 \right\}.$$

With 0 < r < R < 1, choose  $\eta$  supported on  $B_R$  with  $\eta \equiv 1$  on  $B_r$ . Then by the Sobolev inequality,

$$||u||_{L^{2\chi}(B_r)} \le C \frac{1}{R-r} ||u||_{L^2(B_R)}$$

where  $\chi=\frac{n}{n-2}>1$  so that  $2\chi=2^*$ . We have successfully bounded a higher  $L^p$  norm of u by its  $L^2$  norm on a larger set, at the cost of a factor  $\frac{1}{R-r}$  and a smaller domain, a price we'll gladly pay. If the same were true for this higher power of u, we could iterate,  $u\to u^\chi\to u^{\chi^i}$  and obtain a chain of estimates leading back to  $\|u\|_{L^2(B_R)}$ . As  $\chi^i\to\infty$ , this gives control over the  $L^\infty$  norm of u, albeit on a smaller set than the initial ball  $B_R$ . So long as the decrements of the radius, R-r, decrease fast enough, we can get an estimate of  $\sup_{x\in B_{\frac{1}{2}}}u=\lim_{i\to\infty}\|u\|_{L^2\chi^i(B_{r_i})}$  in terms of  $\|u\|_{L^2(B_R)}$  with R=1.

Lets see the details. Notice that since  $x \mapsto x^{\chi}$  is a convex function when  $\chi > 1$ , and has positive derivative when x > 0,  $u^{\chi}$  is also a sub-solution to 2. A-fortiori, we iterate the bound in the previous paragraph; set  $r_i = \frac{1}{2} + \frac{1}{2^{i+1}}$ . Then

$$||u||_{L^{2\chi^{i}}(B_{r_{i}})} \le C^{\frac{1}{\chi^{i}}} 2^{-\frac{i}{\chi^{i}}} ||u||_{L^{2\chi^{i-1}}B(r_{i-1})}$$

Repeatedly applying this bound to the left-hand side,

$$||u||_{L^{2\chi^{i}}(B_{r_{i}})} \le C^{\sum_{j \le i} \frac{1}{\chi^{j}}} 2^{-\sum_{j \le i} \frac{j}{\chi^{j}}} ||u||_{L^{2}B(r_{0})}$$

Letting  $i \to \infty$  on both sides gives the result: the left-hand side becomes  $||u||_{L^{\infty}(B_{1/2})}$ , while the constants appearing on the right have convergent sums as exponents:

$$\sup_{B_{1/2}} u \le C \|u\|_{L^2(B_1)}$$

### The General Case

We return to the general case of  $c, f \in L^q(B_1)$  for  $q > \frac{n}{2}$ ,

$$\int a_{ij} D_i u D_j \phi + c u \phi \le \int f \phi. \tag{3}$$

With care, Moser's approach will work once again. Set  $\tilde{u} = u^+ + k$  with  $k \geq 0$  to be determined in order to handle  $||f||_q$ . Plug in  $\phi = \eta^2 \tilde{u}$ , again applying ellipticity, to obtain

$$\int |D(\tilde{u}\eta)^2| \le C \left\{ \int |D\eta|^2 \tilde{u}^2 + \int |c|\eta^2 \tilde{u}^2 + \int f\eta^2 \tilde{u} \right\}$$

Note that  $\tilde{u} \geq k$ , so  $f\tilde{u} \geq \frac{f}{k}\tilde{u}^2$ . Choosing  $k = ||f||_q$ , we can group the |c| and |f| terms, apply Holder's inequality and use the condition that q > n/2 to obtain

$$\int |D(\tilde{u}\eta)^2| \le C \left\{ \int |D\eta|^2 \tilde{u}^2 + \int \eta^2 \tilde{u}^2 \right\}$$

Choosing  $\eta$  in the same way as above and applying the Sobolev inequality

$$\|\tilde{u}\|_{L^{2\chi}(B_r)} \le C \frac{1}{R-r} \|\tilde{u}\|_{L^2(B_R)}$$

This is once again the starting point of the Moser iteration scheme. But there is a problem: we don't know whether  $\tilde{u}^{\chi}$  satisfies the same type of estimate as  $\tilde{u}$ . Before, convexity saved us. But now we need to check by hand. Luckily,  $x \mapsto x^{\chi}$  is a nice enough function for things to work; we can still apply Holder's inequality to products involving  $u^{\chi}$ . For the details, see section 4.2 of [2]. At the end of the day, we found the right bound,

$$\sup_{B_{1/2}} u \le \sup_{B_{1/2}} \tilde{u} \le C \|\tilde{u}\|_{L^2(B_1)} \le C \left\{ \|u\|_{L^2(B_1)} + \|f\|_q \right\}.$$

## De Giorgi's Approach

Starting from 3, set  $\phi = \eta^2 v$  where  $v = (u - k)^+$ , with k to be chosen later (but not for the purpose of removing  $||f||_q$ ). Our goal is to show, for k large enough,

$$\int_{B_{1/2}} ((u-k)^+)^2 = 0$$

It immediately follows that

$$\sup_{B_{1/2}} u^+ \le k$$

Hopefully, k is of the form in the theorem (it will be, don't worry). This is a decidedly different method from Moser's approach: instead of using the Sobolev inequality to directly lower-bound  $\int |D(\eta \tilde{u})|^2$ , in this approach, we want the  $L^2$  of v on the *left-hand side* rather than the Sobolev norm. To make progress, apply Holder's inequality:

$$||v\eta||_2 \le ||v\eta||_{2^*}^2 |\{v\eta > 0\}|^{1 - \frac{2}{2^*}}$$
  
$$\le ||D(v\eta)||_2^2 |\{v\eta > 0\}|^{\frac{2}{n}}$$

This  $\frac{2}{n}$  will be crucial later. If we apply the usual ellipticity conditions to 3 with our chosen  $\phi$ , we find

$$\int D(v\eta)^2 \le C \left\{ \int |D\eta|^2 v^2 + \int |c|uv\eta^2 + \int |f|v\eta^2 \right\}$$

To make progress, we apply Holder's inequality to the c and f terms. For example,

$$\begin{split} \int |f|v\eta^2 &\leq \|f\|_q \|\eta v\|_{2^*} |\{v\eta>0\}|^{1-\frac{1}{2^*}-\frac{1}{q}} \\ &\leq \|f\|_q \|D(\eta v)\|_2 |\{v\eta>0\}|^{\frac{1}{2}+\frac{1}{n}-\frac{1}{q}} \\ &\leq \frac{1}{2\varepsilon} \|f\|_q |\{v\eta>0\}|^{1+\frac{2}{n}-\frac{2}{q}} + \frac{\varepsilon}{2} \|D(\eta v)\|_2^2 \end{split}$$

We can  $\varepsilon$  as small as we like to absorb the second term into our constant C. Since q > n/2, we can replace the measure term by

$$|\{v\eta > 0\}|^{1 + (\frac{2}{n} - \frac{1}{q}) - \frac{1}{q}} = C(n, q)|\{v\eta > 0\}|^{1 - \frac{1}{q}}$$

Doing the same for the c term and applying the above inequality to  $\int |D(\eta v)|^2$ , we obtain the start of an iteration scheme different from Moser's.

Define  $A(k,r) = \{u > k\} \cap B_r$ . Using the same  $\eta$  function as in the Moser section

$$||v||_{L^{2}(B_{r})}^{2} \leq C \left\{ \frac{1}{(R-r)^{2}} |A(k,R)|^{\varepsilon} ||v||_{L^{2}(B_{R})} + (k+||f||_{q})^{2} |A(k,R)|^{1+\varepsilon} \right\}$$

where  $\varepsilon = \frac{2}{n} - \frac{1}{q} > 0$ . Without this  $\varepsilon$  of room, there is no hope of iteration. Remember that  $v = (u - k)^+$ . In order to iterate, when going from right to left, we need to

REFERENCES REFERENCES

decrease the size of the domain and *increase* the cutoff to some h > k. We can then chain estimates to end with  $\|(u-h)^+\|_{L^2(B_{1/2})}$  on the left-hand side.

To see this in action, we need bounds on |A(k,R)|. We follow [2] closely. By Markov's inequality,  $|A(k,R)| \leq \frac{1}{k} \|u^+\|_{L^2}$ , so the above inequality holds for  $k_0 = C\|u^+\|_{L^2}$  with C large enough. Now note that  $A(k,r) \subset A(k,R)$ , and if h > k, then  $A(k,r) \supset A(h,r)$ . We can apply Markov's inequality along with these inclusions to obtain

$$|A(k,r)| \le \frac{1}{(h-k)^2} \int_{A(k,R)} (u-k)^2$$

Set  $r_i = \frac{1}{2} + \frac{1}{2^{i+1}}$  and  $k_i = k_0 + k(1 - \frac{1}{2^i})$ . Writing  $\phi(k, r) = \|(u^+ - k)^+\|_{L^2(B_r)}$ , we have a chain of inequalities,

$$\phi(k_i, r_i) \le C2^i \phi(k_{i-1}, r_{i-1})^{1+\varepsilon}$$

We have the freedom to choose our end-point, k, as large as we like. In particular, we can make it so that, for some constant  $\gamma > 1$ ,

$$\phi(k_i, r_i) \le \frac{\phi(k_0, r_0)}{\gamma^i} \tag{4}$$

From this we can show  $\Phi(k_{\infty}, r_{\infty}) = \|(u-k_{\infty})^+\|_{L^2(B_{1/2})}^2 = 0$ . Since  $k_{\infty} = k$ , whichever value of k we choose to make 4 hold will yield our desired bound. For details, once again see [2]. The important point is that the  $1 + \varepsilon$  power of the right-hand side allows the inequality chain to accumulate powers of  $\phi(k_0, r_0)$ . If this balances with the accumulating powers of  $2^i$  and C, the argument goes through.

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