

# Full-field model for dynamic precipitation

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This is the documentation for the crystal plasticity constitutive law for DAMASK including dynamic precipitation available in:

<https://github.com/LightForm-group/Damask-KWN>.

The installation instructions for this version of DAMASK are provided elsewhere:

[https://lightform-group.github.io/wiki/software\\_and\\_simulation/kwn-damask](https://lightform-group.github.io/wiki/software_and_simulation/kwn-damask).

The purpose of the document is to describe the KWN-crystal plasticity coupled model used in the constitutive law `kwnpowerlaw`. The program inputs corresponding to the variables mentioned in this document are shown in Section 5.

## 1 Kampmann-Wagner Model

The KWN model used in this full field model is the one described in [1]. A ternary system is considered, with main element C and solute atoms A and B. The precipitate composition - considered as constant - is given by  $A_x B_y C_z$ .

A number density distribution of precipitates,  $\phi(r, t)$ , is defined. The total number density of precipitates,  $N$ , is given by

$$N(t) = \int_0^\infty \phi(r, t) \, dr, \quad (1)$$

the average radius,  $R$ , is given by

$$R(t) = \frac{1}{N(t)} \int_0^\infty r \phi(r, t) \, dr, \quad (2)$$

and the volume fraction,  $f_v$ , is given by

$$f_v(t) = \int_0^\infty \frac{4}{3} \pi r^3 \phi(r, t) \, dr. \quad (3)$$

The conservation of precipitates is then represented by the following conservation law

$$\frac{\partial \phi}{\partial t} + \frac{\partial [v \phi]}{\partial r} = S \quad (4)$$

where,  $v(r, t)$  is the growth rate of the precipitates, and  $S(r, t)$  is the source term representing the nucleation of precipitates.

### 1.1 Initial distribution

A pre-existing precipitate distribution is considered. The initial distribution  $N(0)$  is assumed to follow a log-normal law, and is defined by its initial precipitate volume fraction  $f_v(0)$ , initial mean radius  $R(0)$  and standard deviation  $\sigma^1$  as [1]:

$$N(0) = f_v(0) \frac{\int_0^\infty f(r) dr}{\int_0^\infty f(r) \frac{4}{3} \pi r^3 dr} \quad (5)$$

where  $f$  is a function defined as:

$$f(r) = \frac{1}{\sqrt{2\pi}\sigma} \exp \left[ -\frac{R(0)^2}{2\sigma^2} \left( \ln\left(\frac{r}{R(0)}\right) + \frac{\sigma^2}{2R(0)^2} \right)^2 \right] \quad (6)$$

### 1.2 Precipitate growth

The growth rate,  $v(r, t)$  in Eq. (4) is assumed to be controlled by the slowest diffuser. Assuming the slowest diffuser is A,  $v(r, t)$  is given by:

$$v(r, t) = \frac{D}{r} \cdot \frac{\bar{x}_A - x_A^r}{x_A^{pr} - x_A^r} \quad (7)$$

where  $D$  is the diffusion coefficient, calculated as detailed later,  $\bar{x}_A$  is the average atomic fraction of A in the matrix,  $x_A^{pr}$  is the concentration in A in the precipitate ( $x_A^{pr} = \frac{x}{x+y+z}$ ), and  $x_A^r$  is the concentration in A at the interface between the matrix and a precipitate of radius  $r$ , calculated as the solution of the following system of equations [1, 2]:

$$\begin{cases} x \cdot (x_B^0 - x_B^r) = y \cdot (x_A^0 - x_A^r) \\ \exp \left[ \left( \frac{2\gamma v_m^{pr}}{rRT} + \frac{\Delta G_E \cdot v_m^{pr}}{RT} \right) \cdot (x + y + z) \right] \cdot (x_A^\infty)^x \cdot (x_B^\infty)^y = (x_A^r)^x \cdot (x_B^r)^y \end{cases} \quad (8)$$

where  $\Delta G_E$  is the elastic strain energy,  $x_A^0$  and  $x_B^0$  are the bulk concentrations of the alloy in species A and B, respectively;  $x_A^\infty$  and  $x_B^\infty$  are the equilibrium concentrations in A and B, respectively;  $x_B^r$  - which is the second variable of the system in Eq. (8) - is the concentration at the interface between the matrix and a precipitate of radius  $r$ , and  $v_m^{pr}$  is the molar volume of the precipitate.

### 1.3 Source term

Using classical nucleation theory, the source term in Eq. (4) can be expressed as

$$S = \delta(r - r_c) I \quad (9)$$

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<sup>1</sup>the log-normal distribution is sometimes defined by a dimensionless dispersion parameter  $s$  defined as  $s = \sigma/R(0)$

where  $r_c(t)$  is the critical nucleus size, and  $I(t)$  is the nucleation rate. The critical nucleus size  $r_c(t)$  is calculated as the radius such that:

$$v(r_c, t) = 0 \quad (10)$$

The nucleation rate is given by

$$I(t) = N_s Z \beta \exp \left[ -\frac{4\pi\gamma r_c^2}{3k_B T} \right] \quad (11)$$

where  $N_s$  is the number density of nucleation sites, calculated as

$$N_s = \frac{\bar{x}_A + \bar{x}_B}{V} \quad (12)$$

where  $V$  is the atomic volume of the matrix. In Eq. (11), the Zeldovich factor,  $Z$ , is given by

$$Z = \frac{V}{2\pi r_c^2} \sqrt{\frac{\gamma}{k_B T}}, \quad (13)$$

$\beta$  is given by

$$\beta = 4\pi D(\bar{x}_A + \bar{x}_B) \frac{r_c^2}{a^4}, \quad (14)$$

where  $a$  is the lattice parameter of the product phase.

## 1.4 Solute conservation

The precipitate solute concentration in element  $i$  (where  $i=A$  or  $B$ ),  $x_i^{\text{pr}}$ , is assumed to be constant, and the matrix solute concentration satisfies:

$$x_i^0 = f_v(t)x_i^{\text{pr}} + (1 - f_v(t))\bar{x}_i(t), \quad (15)$$

## 2 Crystal plasticity model

A crystal plasticity constitutive model is used and the plastic strain rate is given by [3]:

$$\mathbf{L}_p = \sum_{\alpha=1}^{12} \dot{\gamma}^\alpha n^\alpha \otimes s^\alpha \quad (16)$$

where the indices  $\alpha$  refer to the 12 FCC slip systems, defined by the slip plane normal  $n^\alpha$  and the slip direction  $s^\alpha$ , and  $\dot{\gamma}^\alpha$  is the slip rate. The slip rates are given by an empirical power law [4]:

$$\dot{\gamma}^\alpha = \dot{\gamma}_0 \left( \frac{\tau^\alpha}{\tau_*^\alpha} \right)^n \quad (17)$$

where  $\dot{\gamma}_0$  and  $n$  are constants,  $\tau^\alpha$  is the resolved shear stress on the slip system, and the slip resistance,  $\tau_*^\alpha$ , is given by [5]:

$$\tau_*^\alpha = \tau_s + \sqrt{(\tau_p)^2 + (\tau_d^\alpha)^2} \quad (18)$$

where  $\tau_s$ ,  $\tau_p$  and  $\tau_d^\alpha$  are the contributions of the atoms in solid solution, of the reinforcing precipitates, and of the dislocations, respectively.  $\tau_s$  is given by:

$$\tau_s = k_s \cdot (\bar{x}_A + \bar{x}_B)^{2/3} \quad (19)$$

where  $k_s$  is an empirical constant [5]. The precipitation hardening contribution  $\tau_p$ , is given by [5]:

$$\tau_p = \mu \sqrt{\frac{3f_v}{2\pi}} \cdot \frac{k_p}{R \cdot \sqrt{r_t}} \left( \int_0^{r_t} r \cdot \phi(r, t) \, dr + \int_{r_t}^\infty r_t \cdot \phi(r, t) \, dr \right)^{3/2} \quad (20)$$

where  $\mu$  is the shear modulus,  $k_p$  is an empirical constant, and  $r_t$  is the transition radius between shearing and Orowan looping [6]. In Eq. (18),  $\tau_d^\alpha$  is the resistance to slip due to forest dislocations, and the hardening rate is given by [4]:

$$\dot{\tau}_d^\alpha = h_0 \sum_{\beta=1}^N q_{\alpha\beta} \dot{\gamma}^\beta \left[ 1 - \frac{\tau_d^\beta}{\tau_\infty} \right]^m, \quad (21)$$

where  $h_0$ ,  $\tau_\infty$  and  $m$  are hardening parameters, and  $q_{\alpha\beta}$  is the latent hardening between the slip systems  $\alpha$  and  $\beta$ . The initial value of  $\tau_d^\alpha$  (considered as identical for all slip systems and labelled  $\tau_d(0)$ ) is a model input.

### 3 Excess vacancy model

Plastic deformation produces vacancies which enhance diffusivity and accelerate precipitation kinetics. A phenomenological model [1, 7, 8] is used to calculate the excess vacancy concentration as a function of temperature and deformation conditions. The rate of generation and annihilation of vacancies is given by

$$\dot{c}_V = \chi \frac{V}{E_V} \sum_{\alpha=1}^N \dot{\gamma}^\alpha \tau^\alpha + 0.5 \cdot \frac{c_j V}{4b^3} \sum_{\alpha=1}^N \dot{\gamma}^\alpha - \frac{D_V}{L^2} c_V, \quad (22)$$

where  $\chi$  is the fraction of plastic work converted into vacancies;  $E_V$ , is the vacancy formation energy;  $D_V = D_{V0} \exp \left[ -\frac{M_V}{k_B T} \right]$  is the vacancy diffusion coefficient in terms of the vacancy migration energy,  $M_V$ ;  $b$  the burgers vector of screw dislocations; and  $c_j$  is the dislocation jog concentration, calculated as  $c_j = \exp \left[ \frac{-E_j}{k_B T} \right]$ , where  $E_j$  is the jog formation energy.  $L$  is the average spacing between vacancy sinks, and is given by the minimum of dislocation spacing and grain size  $d_g$

$$L = \min \left[ \frac{0.27\mu b}{\max(\tau_d^\alpha) \cdot \kappa^2}, d_g \right]. \quad (23)$$

where  $\kappa$  is a constant between 1 and 10 related to the dislocation arrangement ([1, 7]). The solute diffusion coefficient  $D$  in Eq. (14) and Eq. (7) is related to the vacancy concentration through

$$D = D_{th} \left( 1 + \frac{c_v - c_{eq}}{c_{eq}} \right) \quad (24)$$

where  $c_{\text{eq}}$  is the equilibrium vacancy concentration, calculated as  $c_{\text{eq}} = 23 \exp[-\frac{E_f}{k_B T}]$  [1], and  $D_{\text{th}}$  is the solute diffusion coefficient in the absence of any deformation, calculated as  $D_{\text{th}} = D_0 \exp(-E_s/k_B T)$ , where  $E_s$  is the solute migration energy. It is assumed that  $c_v(0) = c_{\text{eq}}$ .

## 4 To improve the model

In the future, the influence of dislocation density on the nucleation site density,  $N_s$  in Eq. (11), through the elastic misfit energy,  $\Delta G_E$ , and the solute diffusion coefficient,  $D$ , through pipe diffusion, might be worth looking into.

### 4.1 Mass Transport

The solute and vacancy balance in Eq. (15) and Eq. (22) respectively local. In addition long-range transport can be incorporated by including diffusional fluxes

$$\dot{c}_V = \nabla D_V \nabla c_V + \chi \frac{\Omega}{E_V} \sum_{\alpha=1}^N \dot{\gamma}^\alpha \tau^\alpha - \frac{D_V}{L^2} c_V \quad (25)$$

$$\dot{c}_M = \nabla D_M \nabla c_M + \frac{\dot{f}}{1-f} (c_M - c_P) \quad (26)$$

### 4.2 Decreasing calculation time

Currently, the equilibrium composition at the interface between matrix and precipitate is calculated at each time step for all bins of the distribution. As the equilibrium concentration at the interface does not depend on time, the code could be modified to calculate the equilibrium composition only at the beginning.

### 4.3 Improve readability

Readability might be improved by converting all units in the `phase_mechanical_plastic_kwnpowerlaw.f9` source file to international unit system (for now, several variables are normed).

## 5 Model inputs

Symbol	Variable	Unit
$m$	a_sl	-
$V$	atomic_volume	$\text{m}^3$
$b$	burgers_vector	m
$[x_A^0, x_B^0]$	c0_matrix	at
$[x_A^\infty, x_B^\infty]$	ceq_matrix	at
$[x_A^{pr}, x_B^{pr}]$	ceq_precipitate	at
$\kappa$	dislocation_arrangement	-
$\dot{\gamma}_0$	dot_gamma_0_sl	$\text{m}^{-1}$
$\gamma$	gamma_coherent	$\text{Jm}^{-2}$
$h_0$	h_0_sl_sl	Pa
$q_{\alpha\beta}$	h_sl_sl	-
$R(0)$	initial_mean_radius	m
$f_v(0)$	initial_volume_fraction	-
$E_j$	jog_formation_energy	eV
$a$	lattice_parameter	m
$\Delta G_E$	misfit_energy	$\text{Jm}^{-3}$
$v_m^{pr}$	molar_volume	$\text{m}^3\text{mol}^{-1}$
$n$	n_sl	-
$k_p$	precipitate_strength_constant	-
$\mu$	shear_modulus	Pa
$D_0$	solute_diffusion0	$\text{m}^2\text{s}^{-1}$
$E_s$	solute_migration_energy	eV
$k_s$	solute_strength	-
$\sigma$	standard_deviation	m
$[x,y,z]$	stoichiometry	-
$r_t$	transition_radius	m
$\tau_d(0)$	xi_0_sl	Pa
$\tau_\infty$	xi_inf_sl	Pa
$D_{V0}$	vacancy_diffusion0	$\text{m}^2\text{s}^{-1}$
$E_V$	vacancy_energy	eV
$\chi$	vacancy_generation	-
$M_V$	vacancy_migration_energy	eV
$d_g$	vacancy_sink_spacing	m

## References

- [1] Madeleine Bignon, Pratheek Shanthraj, and Joseph D. Robson. Modelling dynamic precipitation in pre-aged aluminium alloys under warm forming conditions. *Acta Mater.*, 234:118036, aug 2022. ISSN 1359-6454. doi: 10.1016/

J.ACTAMAT.2022.118036. URL <https://linkinghub.elsevier.com/retrieve/pii/S1359645422004177>.

- [2] Myriam Nicolas and Alexis Deschamps. Characterisation and modelling of precipitate evolution in an Al-Zn-Mg alloy during non-isothermal heat treatments. *Acta Mater.*, 51(20):6077–6094, 2003. ISSN 13596454. doi: 10.1016/S1359-6454(03)00429-4.
- [3] F. Roters, M. Diehl, P. Shanthraj, P. Eisenlohr, C. Reuber, S. L. Wong, T. Maiti, A. Ebrahimi, T. Hochrainer, H. O. Fabritius, S. Nikolov, M. Friák, N. Fujita, N. Grilli, K. G.F. Janssens, N. Jia, P. J.J. Kok, D. Ma, F. Meier, E. Werner, M. Stricker, D. Weygand, and D. Raabe. DAMASK – The Düsseldorf Advanced Material Simulation Kit for modeling multi-physics crystal plasticity, thermal, and damage phenomena from the single crystal up to the component scale. *Comput. Mater. Sci.*, 158 (April 2018):420–478, 2019. ISSN 09270256. doi: 10.1016/j.commatsci.2018.04.030.
- [4] F. Roters, P. Eisenlohr, L. Hantcherli, D. D. Tjahjanto, T. R. Bieler, and D. Raabe. Overview of constitutive laws, kinematics, homogenization and multiscale methods in crystal plasticity finite-element modeling: Theory, experiments, applications. *Acta Mater.*, 58(4):1152–1211, 2010. ISSN 13596454. doi: 10.1016/j.actamat.2009.10.058. URL <http://dx.doi.org/10.1016/j.actamat.2009.10.058>.
- [5] A. Deschamps and Y. Brechet. Influence of predeformation and ageing of an Al-Zn-Mg Alloy-II. Modeling of precipitation kinetics and yield stress. *Acta Mater.*, 47(1): 293–305, 1998. ISSN 13596454. doi: 10.1016/S1359-6454(98)00296-1.
- [6] Derek Hull and D. J. Bacon. Introduction to Dislocations. *Introd. to Dislocations*, 2011. doi: 10.1016/C2009-0-64358-0.
- [7] J. D. Robson. Deformation Enhanced Diffusion in Aluminium Alloys. *Metall. Mater. Trans. A Phys. Metall. Mater. Sci.*, 51(10):5401–5413, 2020. ISSN 10735623. doi: 10.1007/s11661-020-05960-5. URL <https://doi.org/10.1007/s11661-020-05960-5>.
- [8] M. Militzer, W. P. Sun, and J. J. Jonas. Modelling the effect of deformation-induced vacancies on segregation and precipitation. *Acta Metall. Mater.*, 42(1):133–141, jan 1994. ISSN 0956-7151. doi: 10.1016/0956-7151(94)90056-6.
- [9] K. Linga Murty, K. Detemple, O. Kanert, and J. Th M. Dehossion. In-situ nuclear magnetic resonance investigation of strain, temperature, and strain-rate variations of deformation-induced vacancy concentration in aluminum. *Metall. Mater. Trans. A Phys. Metall. Mater. Sci.*, 29(1):153–159, 1998. ISSN 10735623. doi: 10.1007/s11661-998-0168-0.