## PHAS1202 - Atoms, Stars and The Universe Problem Solving Tutorial Sheet 2 Model Answers

### **Objectives:**

- 1. Perform a dimensional analysis on Planck's constant.
- 2. Gain practise with normalising wave-functions and calculating their properties, including expectation values.
- 3. Practise using the TISE to calculate energies for an unfamiliar potential.
- 4. Get practise with the integration and differentiation which arises very often in problems of this kind.
- 5. Gain practise in solving the TISE for different trial solutions and showing that certain functions are not solutions.

#### **Useful definitions**

Planck's constant h is  $6.6 \times 10^{-34}$  Js (2 s.f.).

The time-independent Schrödinger equation (TISE) for a particle in one-dimensional potential V(x) with mass m, energy E with wave-function  $\psi(x)$  is

$$-\frac{\hbar^2}{2m}\frac{d^2\psi(x)}{dx^2} + V(x)\psi(x) = E\psi(x).$$

### 1: The dimensions of h

In Physics, studying the dimensions of quantities can be a very important way of discovering relationships (or potential relationships) between them. Here we will consider the dimensions arising in Planck-Einstein's photon energy law and Bohr's atomic model.

- 1.a ) Write down the dimensions of Planck's constant h in terms of the fundamental quantities of mass M, length L and time T, taking the definition of h to be through Planck and Einsteins relationship E=hf. Why does  $\hbar$  have the same dimensions as h?
- 1.b ) Find the dimensions (in terms of mass M, length L and time T) of angular momentum and show that these are the same as those of Plancks constant h. Hence show that Bohrs quantisation assumption  $L = n\hbar$  (where L is the angular momentum and n is an integer) is dimensionally consistent.

We shall use square brackets to denote units, e.g. [h] refers to the units of h.

1.a ) From Planck-Einstein's formula:

$$[h] = \frac{[E]}{[f]}$$

The units of frequency are Hz or  $s^-1$ , hence

$$[f] = [T]^{-1}$$

To express the units of energy in terms of mass [M], length [L] and time [T] we can refer to any equation for energy which contains only these quantities. Kinetic energy will do:

$$K.E = \frac{1}{2}mv^2$$

and hence

$$[E] = [M][L]^2[T]^{-2}$$

Thus

$$[h] = \frac{[E]}{[f]} = [M][L]^2[T]^{-1}.$$

 $\hbar$  has the same dimensions as h since  $\hbar = h/2\pi$  and  $2\pi$  is dimensionless.

1.b ) To find the dimensions of angular momentum we can use

$$l = rp$$

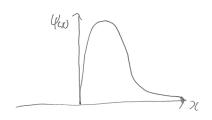
where p is momentum. The dimensions of momentum are  $[M][L][T]^{-1}$  from p = mv.

Hence

$$[l] = [L][M][L][T]^{-1} = [M][L]^2[T]^{-1}$$

Therefore the left-handside and the right-handside of Bohr's angular momentum assumption are both equal, and  $\hbar$  has the same dimension as angular momentum.

# 2: Calculating properties of wave-functions



(a) Clearly can focus just on x > 0. Normalisation:

$$\int_0^\infty |\psi(x)|^2 dx = A^2 \int_0^\infty \left( e^{-x} - e^{-2x} \right)^2 dx = 1$$

$$= A^2 \int_0^\infty \left( e^{-2x} - 2e^{-3x} + e^{-4x} \right) dx = 1$$

$$= A^2 \left[ -\frac{e^{-2x}}{2} + 2\frac{e^{-3x}}{3} - \frac{e^{-4x}}{4} \right]_0^\infty = 1$$

$$A^2 \left( \frac{1}{2} - \frac{2}{3} + \frac{1}{4} \right) = 1$$

$$A = \sqrt{12}$$

(b)

$$\int_0^\infty x |\psi(x)|^2 dx = 12 \int_0^\infty \left( xe^{-2x} - 2xe^{-3x} + xe^{-4x} \right) dx$$

Integration by parts:

$$\int_0^\infty xe^{-ax}dx = \left[xe^{-ax}\right]_0^\infty + \int_0^\infty \frac{e^{-ax}}{a}dx$$
$$\int_0^\infty xe^{-ax}dx = \left[-\frac{e^{-ax}}{a^2}\right]_0^\infty = \frac{1}{a^2}$$

substituting into above

$$\int_0^\infty x |\psi(x)|^2 dx = 12 \left(\frac{1}{4} - \frac{2}{9} + \frac{1}{16}\right)$$
$$\int_0^\infty x |\psi(x)|^2 dx = 12 \left(\frac{5}{16} - \frac{2}{9}\right)$$
$$\int_0^\infty x |\psi(x)|^2 dx = 12 \left(\frac{45 - 32}{144}\right)$$
$$\int_0^\infty x |\psi(x)|^2 dx = \frac{13}{12}$$

(c)

Maximum probability corresponds to  $d\psi/dx=0$ 

$$\frac{d\left|\psi(x)\right|^2}{dx} = 0$$

$$\frac{d|\psi(x)|^2}{dx} = 12\frac{d[e^{-2x} - 2e^{-3x} + e^{-4x}]}{dx}$$

$$0 = 12 \left[ -2e^{-2x} + 6e^{-3x} - 4e^{-4x} \right]$$

Removing common terms

$$0 = \left[ -1 + 3e^{-x} - 2e^{-2x} \right]$$

factorises to

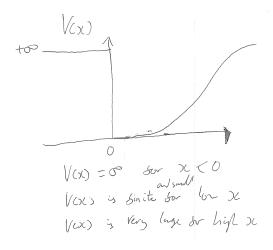
$$0 = (2e^{-x} - 1)(1 - e^{-x})$$

the right term is only zero at x = 0, which is outside the allowed range. Which leaves:

$$x = \ln 2$$

This is not the same as the expectation value. Mean and peak values are not the same for skewed distributions .

(d)



## 3: Which functions are valid wave-functions? (seen)

To represent a physical particle, wave-functions must be continuous and normalisable. Consider the following functions. Can they represent a wave-function for a physical system? If not, explain why.

3a)  $f_1(x) = \cos(px/\hbar)$ 

As we saw in lectures,

$$\int_{-\infty}^{\infty} \sin^2(px/\hbar) dx$$

diverges to infinity, the same is true of

$$\int_{-\infty}^{\infty} \cos^2(px/\hbar) dx$$

and therefore this function cannot be normalised and cannot represent a physical state.

3b)  $f_2(x) = \begin{cases} \cos(4\pi x) - \cos(2\pi x) & \text{if } 0 \le x \le 1\\ 0 & \text{otherwise} \end{cases}$ 

Yes - this is a physical wavefunction. It is continuous, and normalisable, and so satisfies the criteria in the question.

Give yourself a bonus mark if you noted that this wavefunction is also continuous in its first derivative. Since in the middle region,

$$\frac{df_3(x)}{dx} = -4\pi \sin(4\pi x) + 2\pi \sin(2\pi x)$$

which at x=0 is equal to 0 and at x=1 is also equal to 0, matching the zero first derivative of the constant wavefunction outside this region.

3c)  $f_3(x) = \begin{cases} \cos(\pi x) & x \le 0\\ \sin(\pi x) & x > 0 \end{cases}$ 

This is not a valid wave-function since it is *not continuous*. sin(0) = 0 but cos(0) = 1 hence there is a discontinuity at x = 0.

$$f_4(x) = e^{-x}$$

This wave-function blows up to infinity as x gets large and negative, and thus the normalisation integral also blows up to infinity. This wavefunction cannot be normalised and is unphysical.

### 4: Ground state of a chemical bond

A chemical bond can be approximated by a simple spring between two masses (which represent the atoms), the potential energy of which depends on their separation x, according to:

$$V(x) = \frac{1}{2}kx^2$$

where k is the spring constant. If we solve the TISE for such a potential, the lowest energy state has a wave function of the form:

$$\psi(x) = A \exp[-\alpha^2 x^2/2]$$

where  $\alpha = \sqrt{m\omega/\hbar}$  and  $\omega = \sqrt{k/m}$ .

4a) Use the TISE to compute the energy of wavefunction  $\psi(x)$ .

The TISE for this potential is

$$-\frac{\hbar^2}{2m}\frac{d^2\psi(x)}{dx^2} + \frac{1}{2}kx^2\psi(x) = E\psi(x)$$

We know that  $\psi(x)$  is a solution to this, so by substituting it into the TISE we can compute the energy. We need to compute

$$\frac{d^2\psi(x)}{dx^2} = \frac{d}{dx}A\left(\frac{-2x\alpha^2}{2}\right)\exp[-\alpha^2x^2/2] = -A\alpha^2\frac{d}{dx}\left(x\exp[-\alpha^2x^2/2]\right)$$
$$= -A\alpha^2\left(\exp[-\alpha^2x^2/2] + \frac{-2\alpha^2}{2}x^2\exp[-\alpha^2x^2/2]\right) = A(\alpha^4x^2 - \alpha^2)\exp[-\alpha^2x^2/2]$$

where we used the chain rule and the product rule.

Hence we obtain:

$$-\frac{\hbar^2}{2m}A(\alpha^4x^2-\alpha^2)\exp[-\alpha^2x^2/2] + \frac{1}{2}kx^2A\exp[-\alpha^2x^2/2] = EA\exp[-\alpha^2x^2/2]$$

and after cancellations

$$-\frac{\hbar^2}{2m}(\alpha^4 x^2 - \alpha^2) + \frac{1}{2}kx^2 = E$$
$$\left(-\frac{\hbar^2}{2m}\alpha^4 + \frac{1}{2}k\right)x^2 = (E - \frac{\hbar^2}{2m}\alpha^2)$$

When we substitute for the values of  $\alpha$  and  $\omega$  given above the LHS cancels to zero. Hence

$$(E - \frac{\hbar^2}{2m}\alpha^2) = 0$$

$$E = \frac{\hbar^2 \alpha^2}{2m} = \frac{\hbar^2 m\omega}{2m\hbar} = \frac{\hbar\omega}{2}$$

- 4b) Use the symmetry of the wave-function to determine the expectation value for separation x.
- $\langle x \rangle = 0$ . The wavefunction  $\psi(x)$  is an even function. Therefore the probability density  $|\psi(x)|^2$  is also even. The integral:

$$\langle x \rangle = \int_{-\infty}^{\infty} x |\psi(x)|^2 dx.$$

is an integral over the odd function  $x|\psi(x)|^2$  which evaluates to zero.

4c) Find the value of A which normalises this wavefunction, and calculate the expectation value of separation x.

You may find the following integral useful

$$\int_{-\infty}^{\infty} e^{-y^2} dy = \sqrt{\pi}$$

To find A we solve:

$$\int_{-\infty}^{\infty} |\psi(x)|^2 dx = 1$$

Hence

$$\int_{-\infty}^{\infty} |\psi(x)|^2 dx = A^2 \int_{-\infty}^{\infty} \exp[-\alpha^2 x^2] dx$$

To use the given integral I need to change variables:  $y = \alpha x$  and hence  $dy = \alpha dx$ .

$$A^{2} \int_{-\infty}^{\infty} \exp[-\alpha^{2} x^{2}] dx = \frac{A^{2}}{\alpha} \int_{-\infty}^{\infty} \exp[-y^{2}] dy = \frac{A^{2}}{\alpha} \sqrt{\pi}.$$

Hence

$$A = \frac{\sqrt{\alpha}}{\pi^{1/4}}.$$

Now we need to compute  $\langle x \rangle$  via:

$$\langle x \rangle = \int_{-\infty}^{\infty} x |\psi(x)|^2 dx = A^2 \int_{-\infty}^{\infty} x \exp[-\alpha^2 x^2] dx.$$

You are not given an identity to solve this equation, instead you can solve it from its structure. x is an odd function, while  $\exp[-\alpha^2 x^2]$  is even. Hence the integrand is odd, and thus the integral is zero. Hence  $\langle x \rangle = 0$ .

## 5: Solving the Time Independent Schrödinger Equation for a free particle (unseen)

In lectures we saw that a general sinusoidal function:

$$\psi(x) = A \sin\left(\frac{2\pi(x-\phi)}{\lambda}\right)$$

was a solution of the TISE for a free particle, and we calculated the associated energy. In this question, you are going to attempt to solve the TISE for a free particle yourselves, using a variety of trial functions.

The TISE for a free particle is:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi(x)}{dx^2} = E\psi(x)$$

5a) In lectures you heard that the only classes of functions which are proportional to their second derivative have sine-like and exponential-like forms. Let's look at some other functions and see how they fail. First let's consider

$$\psi(x) = x^2$$

Show that this function is *not* a solution to the free particle TISE.

$$\frac{d\psi(x)}{dx} = 2x$$
$$\frac{d^2\psi(x)}{dx^2} = 2$$

2 is a constant and is not proportional to  $x^2$ . Therefore  $x^2$  cannot be a solution to the TISE.

5b) Is there any (finite) value of n for which

$$\psi(x) = x^n$$

could be a solution to the TISE for a free particle?

It is the same argument but more general.

$$\frac{d\psi(x)}{dx} = nx^{n-1}$$
$$\frac{d^2\psi(x)}{dx^2} = n(n-1)x^{n-2}$$

 $n(n-1)x^{n-2}$  is not proportional to  $x^n$  (question to the students - why?). Therefore  $x^n$  cannot be a solution to the TISE for any n.

It is not a coincidence that solutions to the TISE are functions such as sine and cos which have infinitely many terms in their power series expansion.

5c) Now consider the following function:

$$\psi(x) = A\sin(ax) + B\cos(bx)$$

Find an expression for the ratio

$$\frac{d^2\psi(x)}{dx^2} \div \psi(x).$$

Show that if b=a this ratio is a constant, and that the wavefunction is solution to the TISE and calculate the energy of the particle as a function of parameters a, A and B.

$$\frac{d\psi(x)}{dx} = Aa\cos(ax) - Bb\sin(bx)$$
$$\frac{d^2\psi(x)}{dx^2} = -Aa^2\sin(ax) - Bb^2\cos(bx)$$

Let's rewrite the TISE in a more convenient form:

$$\frac{d^2\psi(x)}{dx^2} = -\frac{2mE}{\hbar^2}\psi(x)$$

We need to check whether  $\frac{d^2\psi(x)}{dx^2}$  can be proportional to (with a negative constant of proportionality)  $\psi(x)$ .

To simplify notation let  $\kappa = \frac{2mE}{\hbar^2}$ . Hence we want our wavefunction to satisfy:

$$\frac{d^2\psi(x)}{dx^2} = -\kappa\psi(x)$$

$$-Aa^{2}\sin(ax) - Bb^{2}\cos(bx) = -(Aa^{2}\sin(ax) + Bb^{2}\cos(bx)) = -\kappa(A\sin(ax) + B\cos(bx))$$

we can rewrite this:

$$\kappa = \frac{Aa^2 \sin(ax) + Bb^2 \cos(bx)}{A \sin(ax) + B \cos(bx)}$$

If b = a

$$\kappa = \frac{Aa^2 \sin(ax) + Ba^2 \cos(ax)}{A \sin(ax) + B \cos(ax)} = \frac{a^2 (A \sin(ax) + B \cos(ax))}{A \sin(ax) + B \cos(ax)} = a^2$$

To calculate the energy, recall that:

$$\kappa = \frac{2mE}{\hbar^2} = a^2$$

Hence

$$E = \frac{\hbar^2 a^2}{2m}$$

NB the energy does not depend on A or B.