

PHAS2423 - Self-Study - Cartesian Tensors - Problems and Solutions

(1) The summation convention.

(a) Express the following using the summation convention.

$$(a.1) x'_i = \sum_{j=1}^3 a_{ij} x_j; \quad (a.2) T'_{kl} = \sum_{i=1}^3 \sum_{j=1}^3 a_{ki} a_{lj} T_{ij}; \quad (a.3) B'_{pqr} = \sum_{i=1}^3 \sum_{j=1}^3 \sum_{k=1}^3 L_{pi} L_{qj} L_{rk} B_{ijk}.$$

(b) Write the following using explicit summation.

$$(b.1) C_{jknm} A_n B_{jk}; \quad (b.2) (\delta_{ij} \delta_{kl} + \delta_{il} \delta_{kj}) A_{ik}.$$

(c) Show that

$$(c.1) \delta_{ij} \delta_{jk} = \delta_{ik}; \quad (c.2) \delta_{ii} = N.$$

(d) Evaluate

$$(d.1) \delta_{ij} \delta_{jk} \delta_{km} \delta_{im}; \quad (d.2) \epsilon_{jk2} \epsilon_{k2j}; \quad (d.3) \epsilon_{23i} \epsilon_{2i3}.$$

Solutions

$$(a.1) x'_i = a_{ij} x_j \quad (a.2) T'_{kl} = a_{ki} a_{lj} T_{ij} \quad (a.3) B'_{pqr} = L_{pi} L_{qj} L_{rk} B_{ijk}$$

(b.1)

$$C_{jknm} A_n B_{jk} = \sum_{j=1}^3 \sum_{k=1}^3 \sum_{n=1}^3 C_{jknm} A_n B_{jk}$$

(b.2)

$$(\delta_{ij} \delta_{kl} + \delta_{il} \delta_{kj}) A_{ik} = \sum_{i=1}^3 \sum_{k=1}^3 (\delta_{ij} \delta_{kl} A_{ik} + \delta_{il} \delta_{kj} A_{ik}) = \sum_{i=1}^3 (\delta_{ij} A_{il} + \delta_{il} A_{ij}) = A_{jl} + A_{lj}$$

(c.1)

$$\delta_{ij}\delta_{jk} = \sum_{j=1}^3 \delta_{ij}\delta_{jk} = \delta_{i1}\delta_{1k} + \delta_{i2}\delta_{2k} + \delta_{i3}\delta_{3k} = \begin{cases} 1 & (\text{if } i = k) \\ 0 & (\text{if } i \neq k) \end{cases} = \delta_{ik}$$

(c.2)

$$\delta_{ii} = \delta_{11} + \delta_{22} + \dots \delta_{NN} = N$$

(d.1) Use the results of (c.1) and (c.2):

$$\delta_{ij}\delta_{jk}\delta_{km}\delta_{im} = \delta_{ik}\delta_{ki} = \delta_{ii} = 3 \quad (\text{or } N \text{ in the case of } N\text{-dimensional space}).$$

(d.2) It is sufficient to consider all combinations in which $j \neq k$ (all other contributions are equal to zero), and recall that $\epsilon_{ijk} = -\epsilon_{jik}$. Then,

$$\epsilon_{jk2}\epsilon_{k2j} = -\epsilon_{jk2}\epsilon_{kj2} = \epsilon_{jk2}\epsilon_{jk2} = 2.$$

(d.3) It is sufficient to consider all combinations in which $i=1$:

$$\epsilon_{23i}\epsilon_{2i3} = -\epsilon_{23i}\epsilon_{23i} = -(\epsilon_{23i})^2 = -(\epsilon_{231})^2 = -1.$$

(2) Rotation.(a) Orthogonality. An orthogonal matrix L has components L_{ij} . Evaluate the following:

$$(a.1) L_{ij}L_{jk}; \quad (a.2) L_{ji}L_{kj}; \quad (a.3) L_{ij}L_{ik}; \quad (a.4) L_{ij}L_{kj}.$$

(b) Rotation. Show that the transformation matrix L for a rotation of the coordinate system by an angle θ about \mathbf{e}_3 axis is

$$L = \begin{pmatrix} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

(c) Consecutive rotations. Show that two consecutive rotations of the coordinate system by an angle θ about \mathbf{e}_3 axis is also a rotation about the same axis with the value of the rotation angle of 2θ .**Solutions**

(a.1)

$$L_{ij}L_{jk} = (L^2)_{ik} = M_{ik} \quad (\text{matrix } M = L^2 \text{ is also orthogonal})$$

(a.2)

$$L_{ji}L_{kj} = L_{kj}L_{ji} = (L^2)_{ki} = M_{ki} = (M^T)_{ik} = (M^{-1})_{ik}$$

Alternatively,

$$L_{ji}L_{kj} = (L^T)_{ij}(L^T)_{jk} = (L^{-1})_{ij}(L^{-1})_{jk} = \left((L^{-1})^2\right)_{ik} = \left((L^2)^{-1}\right)_{ik} = (M^{-1})_{ik}$$

(a.3)

$$L_{ij}L_{ik} = \delta_{jk} \quad (\text{scalar product of columns } j \text{ and } k)$$

(a.4)

$$(a.4) \quad L_{ij}L_{kj} = \delta_{ik} \quad (\text{scalar product of rows } i \text{ and } k)$$

(b) See Fig. 2.2 of the lecture notes:

$$\begin{aligned} x'_1 &= x_1 \cos \theta + x_2 \sin \theta \\ x'_2 &= -x_1 \sin \theta + x_2 \cos \theta, \quad \text{i.e.,} \\ x'_3 &= x_3 \end{aligned} \quad \left(\begin{array}{c} x'_1 \\ x'_2 \\ x'_3 \end{array} \right) = \left(\begin{array}{ccc} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{array} \right) \left(\begin{array}{c} x_1 \\ x_2 \\ x_3 \end{array} \right)$$

(c)

$$L^2 = \left(\begin{array}{ccc} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{array} \right) \left(\begin{array}{ccc} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{array} \right) = \left(\begin{array}{ccc} \cos^2 \theta - \sin^2 \theta & 2 \sin \theta \cos \theta & 0 \\ -2 \sin \theta \cos \theta & \cos^2 \theta - \sin^2 \theta & 0 \\ 0 & 0 & 1 \end{array} \right).$$

Thus,

$$L^2 = \left(\begin{array}{ccc} \cos 2\theta & \sin 2\theta & 0 \\ -\sin 2\theta & \cos 2\theta & 0 \\ 0 & 0 & 1 \end{array} \right)$$

(3) Transformation of tensors.

(a) Contraction. Given that T_{ijk} and V_n are components of the 3rd order and 1st order tensors, respectively,

(a.1) Show that T_{ijj} is a 1st-order tensor.

(a.2) Show that $T_{ijk}V_k$ is a 2nd-order tensor.

(b) Outer product. If quantities A_{ij} and B_{kl} are components of 2nd order tensors, show that quantities T_{ijkl} formed by $T_{ijkl} = A_{ij}B_{kl}$ is a 4th-order tensor.

(c) Vectors. For the case of a two-dimensional space

(c.1) Show that $\mathbf{v} = (x_2, -x_1)$ transforms as a vector under rotation of the coordinate system.

(c.2) Show that $\mathbf{v} = (x_2, x_1)$ is not a vector.

(d) Scalars.

(d.1) Show that the scalar product of vectors \mathbf{a} and \mathbf{b} is, indeed, a scalar.

(d.2) Show that $\nabla \cdot \mathbf{v}$ is a scalar (assume that \mathbf{v} is a vector).

(e) Higher order tensors. Demonstrate that matrix T represents a 2^{nd} order tensor:

$$\mathbf{T} = \begin{pmatrix} x_2^2 & -x_1x_2 \\ -x_1x_2 & x_1^2 \end{pmatrix}.$$

Solutions

(a.1) Let $P_j = T_{ij}$. Then,

$$P'_j = T'_{ij} = L_{ip}L_{iq}L_{jr}T_{pqr} = \delta_{pq}L_{jr}T_{pqr} = L_{jr}T_{ppr} = L_{jr}P_r.$$

(a.2) Let $P_{ij} = T_{ijk}V_k$. Then,

$$P'_{ij} = T'_{ijk}V'_k = (L_{ip}L_{jq}L_{kr}T_{pqr})(L_{ks}V_s) = L_{ip}L_{jq}\delta_{rs}T_{pqr}V_s = L_{ip}L_{jq}(T_{pqs}V_s) = L_{ip}L_{jq}P_{pq}.$$

(b)

$$T'_{ijkl} = A'_{ij}B'_{kl} = (L_{ip}L_{jq}A_{pq})(L_{kr}L_{ls}B_{rs}) = L_{ip}L_{jq}L_{kr}L_{ls}(A_{pq}B_{rs}) = L_{ip}L_{jq}L_{kr}L_{ls}T_{pqrs}.$$

(c.1)

$$\begin{aligned} v'_1 &= x'_2 = -x_1 \sin \theta + x_2 \cos \theta = v_2 \sin \theta + v_1 \cos \theta \\ v'_2 &= -x'_1 = -(x_1 \cos \theta + x_2 \sin \theta) = v_2 \cos \theta - v_1 \sin \theta. \end{aligned}$$

Thus,

$$\begin{pmatrix} v'_1 \\ v'_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix},$$

i.e., upon rotation of the coordinate system components of \mathbf{v} transform in the same way as coordinates. Hence, \mathbf{v} is a vector.

(c.2)

$$\begin{aligned} v'_1 &= x'_2 = -x_1 \sin \theta + x_2 \cos \theta = -v_2 \sin \theta + v_1 \cos \theta \\ v'_2 &= x'_1 = x_1 \cos \theta + x_2 \sin \theta = v_2 \cos \theta + v_1 \sin \theta. \end{aligned}$$

Thus,

$$\begin{pmatrix} v'_1 \\ v'_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix},$$

i.e., components of \mathbf{v} do not transform as coordinates. Hence, \mathbf{v} is not a vector.

(d.1) Let $s = \mathbf{a} \cdot \mathbf{b}$. Then,

$$s' = \mathbf{a}' \cdot \mathbf{b}' = a'_i b'_i = (L_{ik}a_k)(L_{in}b_n) = (L_{ik}L_{in})a_k b_n = \delta_{kn}a_k b_n = a_n b_n = s.$$

(d.2)

$$\nabla' \cdot \mathbf{v}' = \frac{\partial}{\partial x'_i} v'_i = \left(\frac{\partial x_j}{\partial x'_i} \frac{\partial}{\partial x_j} \right) (L_{ik} v_k) = L_{ij} L_{ik} \frac{\partial}{\partial x_j} v_k = \delta_{jk} \frac{\partial}{\partial x_j} v_k = \frac{\partial}{\partial x_k} v_k = \nabla \cdot \mathbf{v}$$

(e) In order to check whether T_{ij} are components of a tensor, one has to show that

$$T'_{ij} = L_{ip} L_{jq} T_{pq}, \quad \text{where} \quad T' = \begin{pmatrix} x'_2 x'_2 & -x'_1 x'_2 \\ -x'_1 x'_2 & x'_1 x'_1 \end{pmatrix} \quad \text{and} \quad L = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix}.$$

First, express T'_{ij} in terms of x_1 and x_2 :

$$\begin{aligned} T'_{11} &= (x'_2)^2 \\ &= (-x_1 \sin \theta + x_2 \cos \theta)^2 \\ &= x_1^2 \sin^2 \theta + x_2^2 \cos^2 \theta - 2x_1 x_2 \sin \theta \cos \theta \\ \\ T'_{22} &= (x'_1)^2 \\ &= (x_1 \cos \theta + x_2 \sin \theta)^2 \\ &= x_1^2 \cos^2 \theta + x_2^2 \sin^2 \theta + 2x_1 x_2 \sin \theta \cos \theta \\ \\ T'_{12} &= -x'_1 x'_2 \\ &= -(x_1 \cos \theta + x_2 \sin \theta)(-x_1 \sin \theta + x_2 \cos \theta) \\ &= x_1^2 \cos \theta \sin \theta - x_2^2 \sin \theta \cos \theta - x_1 x_2 (\cos^2 \theta - \sin^2 \theta) \\ \\ T'_{21} &= T'_{12} \end{aligned}$$

Now calculate components of the transformed tensor T :

$$\begin{aligned} T'_{11} &= \sum_{i=1}^2 \sum_{j=1}^2 L_{1i} L_{1j} T_{ij} & T'_{22} &= \sum_{i=1}^2 \sum_{j=1}^2 L_{2i} L_{2j} T_{ij} \\ \\ T'_{12} &= \sum_{i=1}^2 \sum_{j=1}^2 L_{1i} L_{2j} T_{ij} & T'_{21} &= \sum_{i=1}^2 \sum_{j=1}^2 L_{2i} L_{1j} T_{ij} \end{aligned}$$

Then,

$$\begin{aligned} T'_{11} &= (\cos \theta)^2(x_2^2) + \cos \theta \sin \theta(-x_1x_2) + \sin \theta \cos \theta(-x_1x_2) + (\sin \theta)^2(x_1^2) \\ &= x_1^2 \sin^2 \theta + x_2^2 \cos^2 \theta - 2x_1x_2 \sin \theta \cos \theta \end{aligned}$$

$$\begin{aligned} T'_{22} &= (-\sin \theta)^2(x_2^2) + (-\sin \theta) \cos \theta(-x_1x_2) + \cos \theta(-\sin \theta)(-x_1x_2) + (\cos \theta)^2(x_1^2) \\ &= x_1^2 \cos^2 \theta + x_2^2 \sin^2 \theta + 2x_1x_2 \sin \theta \cos \theta \end{aligned}$$

$$\begin{aligned} T'_{12} &= \cos \theta(-\sin \theta)(x_2^2) + \cos \theta \cos \theta(-x_1x_2) + \sin \theta(-\sin \theta)(-x_1x_2) + \sin \theta \cos \theta(x_1^2) \\ &= x_1^2 \sin \theta \cos \theta - x_2^2 \sin \theta \cos \theta - x_1x_2(\cos^2 \theta - \sin^2 \theta) \end{aligned}$$

$$T'_{21} = T'_{12}$$

Clearly, both methods give the same expression for the values of T'_{kn} . Thus, \mathbf{T} is a tensor.

(4) Quotient theorem. Given that \mathbf{A} is an arbitrary tensor and \mathbf{B} is a non-zero tensor, prove the quotient theorem for the following cases:

$$(a) X_i A_{ij} = B_j$$

$$(b) X_{ij} A_k = B_{ijk}$$

Solutions

$$(a) X_i A_{ij} = B_j.$$

$$\begin{aligned} X'_i A'_{ij} &= B'_j = L_{jp} B_p = L_{jp} (X_q A_{qp}) \\ &= L_{jp} X_q (L_{mq} L_{np} A'_{mn}) = (L_{jp} L_{np}) L_{mq} X_q A'_{mn} = (\delta_{jn}) L_{mq} X_q A'_{mn} \\ &= L_{mq} X_q A'_{mj}. \end{aligned}$$

Replace the dummy index m by i and subtract the right hand side from the left hand side:

$$(X'_i - L_{iq} X_q) A'_{ij} = 0.$$

Since \mathbf{A} is an arbitrary tensor,

$$X'_i = L_{iq} X_q,$$

i.e., \mathbf{X} is a 1st order tensor.

(b) $X_{ij}A_k = B_{ijk}$.

$$\begin{aligned} X'_{ij}A'_k &= B'_{ijk} = L_{ip}L_{jq}L_{kr}B_{pqr} = L_{ip}L_{jq}L_{kr}X_{pq}A_r \\ &= L_{ip}L_{jq}L_{kr}X_{pq}(L_{mr}A'_m) = L_{ip}L_{jq}(L_{kr}L_{mr})X_{pq}A'_m \\ &= L_{ip}L_{jq}(\delta_{km})X_{pq}A'_m = L_{ip}L_{jq}X_{pq}A'_k. \end{aligned}$$

Subtract the right hand side from the left hand side:

$$(X'_{ij} - L_{ip}L_{jq}X_{pq})A'_k = 0.$$

Since \mathbf{A} is an arbitrary tensor,

$$X'_{ij} = L_{ip}L_{jq}X_{pq},$$

i.e., \mathbf{X} is a 2nd order tensor.

(5) Application of tensors ϵ_{ijk} and δ_{ij} Use properties of the Levi-Civita and Kronecker tensors to prove the following identities for vectors \mathbf{a} , \mathbf{b} , \mathbf{c} , and \mathbf{d} :

(a)

$$(\mathbf{a} \times \mathbf{b}) \cdot (\mathbf{c} \times \mathbf{d}) = (\mathbf{a} \cdot \mathbf{c})(\mathbf{b} \cdot \mathbf{d}) - (\mathbf{a} \cdot \mathbf{d})(\mathbf{b} \cdot \mathbf{c}).$$

(b)

$$(\mathbf{a} \times \mathbf{b}) \times (\mathbf{c} \times \mathbf{d}) = [(\mathbf{a} \times \mathbf{b}) \cdot \mathbf{d}]\mathbf{c} - [(\mathbf{a} \times \mathbf{b}) \cdot \mathbf{c}]\mathbf{d}$$

(c) Find an explicit expression for the i^{th} component of vector $\nabla \times (\nabla \times \mathbf{a})$.

Solutions

(a) Using the expression for the k^{th} component of the vector product $\mathbf{u} = \mathbf{a} \times \mathbf{b}$:

$$u_k = \epsilon_{kmn}a_mb_n,$$

we obtain

$$\begin{aligned}
(\mathbf{a} \times \mathbf{b}) \cdot (\mathbf{c} \times \mathbf{d}) &= (\mathbf{a} \times \mathbf{b})_i (\mathbf{c} \times \mathbf{d})_i \\
&= (\epsilon_{ijk} a_j b_k) (\epsilon_{ilm} c_l d_m) = (\epsilon_{ijk} \epsilon_{ilm}) a_j b_k c_l d_m \\
&= (\delta_{jl} \delta_{km} - \delta_{jm} \delta_{kl}) a_j b_k c_l d_m = \delta_{jl} \delta_{km} a_j b_k c_l d_m - \delta_{jm} \delta_{kl} a_j b_k c_l d_m . \\
&= a_j b_k c_j d_k - a_j b_k c_k d_j = (a_j c_j) (b_k d_k) - (a_j d_j) (b_k c_k) \\
&= (\mathbf{a} \cdot \mathbf{c}) (\mathbf{b} \cdot \mathbf{d}) - (\mathbf{a} \cdot \mathbf{d}) (\mathbf{b} \cdot \mathbf{c})
\end{aligned}$$

(b) The i^{th} component of the vector product $(\mathbf{a} \times \mathbf{b}) \times (\mathbf{c} \times \mathbf{d})$ is

$$\begin{aligned}
[(\mathbf{a} \times \mathbf{b}) \times (\mathbf{c} \times \mathbf{d})]_i &= \epsilon_{ijk} (\mathbf{a} \times \mathbf{b})_j (\mathbf{c} \times \mathbf{d})_k \\
&= \epsilon_{ijk} (\epsilon_{j pq} a_p b_q) (\epsilon_{k rs} c_r d_s) \\
&= \epsilon_{j pq} (\epsilon_{k ij} \epsilon_{k rs}) a_p b_q c_r d_s = \epsilon_{j pq} (\delta_{ir} \delta_{js} - \delta_{is} \delta_{jr}) a_p b_q c_r d_s \\
&= \epsilon_{j pq} \delta_{ir} \delta_{js} a_p b_q c_r d_s - \epsilon_{j pq} \delta_{is} \delta_{jr} a_p b_q c_r d_s \\
&= \epsilon_{j pq} a_p b_q c_i d_j - \epsilon_{j pq} a_p b_q c_j d_i = [(\epsilon_{j pq} a_p b_q) d_j] c_i - [(\epsilon_{j pq} a_p b_q) c_j] d_i \\
&= [(\mathbf{a} \times \mathbf{b})_j d_j] c_i - [(\mathbf{a} \times \mathbf{b})_j c_j] d_i \\
&= [(\mathbf{a} \times \mathbf{b}) \cdot \mathbf{d}] c_i - [(\mathbf{a} \times \mathbf{b}) \cdot \mathbf{c}] d_i.
\end{aligned}$$

Hence, in the vector form:

$$(\mathbf{a} \times \mathbf{b}) \times (\mathbf{c} \times \mathbf{d}) = [(\mathbf{a} \times \mathbf{b}) \cdot \mathbf{d}] \mathbf{c} - [(\mathbf{a} \times \mathbf{b}) \cdot \mathbf{c}] \mathbf{d}.$$

(c) The i^{th} component of vector $\nabla \times (\nabla \times \mathbf{a})$ is

$$(\nabla \times (\nabla \times \mathbf{a}))_i = \epsilon_{ijk} (\nabla)_j (\nabla \times \mathbf{a})_k = \epsilon_{ijk} \epsilon_{kmn} (\nabla)_j (\nabla)_m a_n = \epsilon_{ijk} \epsilon_{kmn} \frac{\partial^2 a_n}{\partial x_j \partial x_m}.$$

Use

$$\epsilon_{ijk} \epsilon_{kmn} = \epsilon_{kij} \epsilon_{kmn} = \delta_{im} \delta_{jn} - \delta_{in} \delta_{jm}.$$

Then,

$$(\nabla \times (\nabla \times \mathbf{a}))_i = (\delta_{im} \delta_{jn} - \delta_{in} \delta_{jm}) \frac{\partial^2 a_n}{\partial x_j \partial x_m} = \frac{\partial^2 a_n}{\partial x_i \partial x_n} - \frac{\partial^2 a_i}{\partial x_m \partial x_m} = \frac{\partial}{\partial x_i} (\nabla \cdot \mathbf{a}) - \frac{\partial^2}{\partial x_m^2} a_i.$$

Finally,

$$(\nabla \times (\nabla \times \mathbf{a}))_i = [\nabla(\nabla \cdot \mathbf{a})]_i - \nabla^2 a_i,$$

i.e.,

$$\nabla \times (\nabla \times \mathbf{a}) = \nabla(\nabla \cdot \mathbf{a}) - \nabla^2 \mathbf{a}.$$

(6) A rigid body consists of eight particles, each of mass m , held together by light rods. In a certain coordinate system the particles are at positions

$$\pm a(3, 1, -1) \quad \pm a(1, -1, 3) \quad \pm a(1, 3, -1) \quad \pm a(-1, 1, 3).$$

The body rotates about an axis passing through the origin. Show that, if the angular velocity and angular momentum vectors are parallel, then their ratio must be $40ma^2$, $64ma^2$, or $72ma^2$.

Solution. If a body rotates about a fixed axis, the vector of the angular velocity $\boldsymbol{\omega}$, vector of the orbital momentum \mathbf{L} , and the moment of inertia are related as

$$\mathbf{L} = I\boldsymbol{\omega},$$

where scalar I is defined for that rotation axis. If the rotation axis is not fixed, the vectors $\boldsymbol{\omega}$ and \mathbf{L} may not be parallel and the relation between them becomes

$$L_i = I_{ij}\omega_j,$$

where I_{ij} is a tensor of inertia. Hence, vectors $\boldsymbol{\omega}$ and \mathbf{L} become parallel if the rotation axis coincides with a principal axis of the tensor of inertia and, therefore, the ratio between \mathbf{L} and $\boldsymbol{\omega}$ is equal to the principal moment of the tensor of inertia.

$$I_{ij} = \sum m(r^2\delta_{ij} - x_i x_j)$$

$$\begin{aligned} I_{11} &= 2ma^2(4 \times 11 - 9 - 1 - 1 - 1) = 64ma^2 \\ I_{22} &= 2ma^2(4 \times 11 - 1 - 1 - 9 - 1) = 64ma^2 \\ I_{33} &= 2ma^2(4 \times 11 - 1 - 9 - 1 - 9) = 48ma^2 \\ I_{12} &= -2ma^2(3 - 1 + 3 - 1) = -8ma^2 \\ I_{13} &= -2ma^2(-3 + 3 - 1 - 3) = 8ma^2 \\ I_{23} &= -2ma^2(-1 - 3 - 3 + 3) = 8ma^2 \end{aligned}$$

Hence,

$$I = 8ma^2 \begin{pmatrix} 8 & -1 & 1 \\ -1 & 8 & 1 \\ 1 & 1 & 6 \end{pmatrix}$$

Solve characteristic equation $\det(\mathbf{I} - \lambda \mathbf{E}) = 0$:

$$\begin{vmatrix} 8 - \lambda & -1 & 1 \\ -1 & 8 - \lambda & 1 \\ 1 & 1 & 6 - \lambda \end{vmatrix} = (8 - \lambda)(\lambda^2 - 14\lambda + 47 - 2) = (8 - \lambda)(\lambda - 9)(\lambda - 5) = 0.$$

Thus, principal moments are $40ma^2$, $64ma^2$, and $72ma^2$.