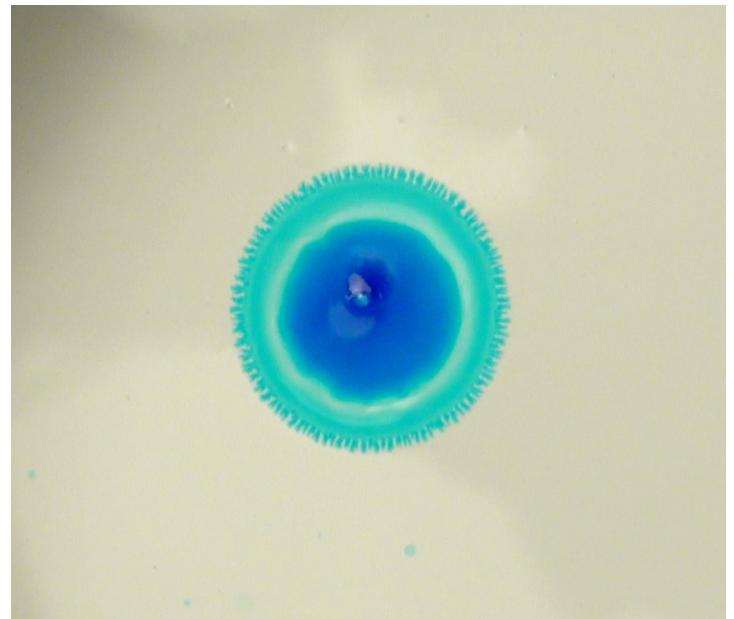




# 1. Fractal Fingers

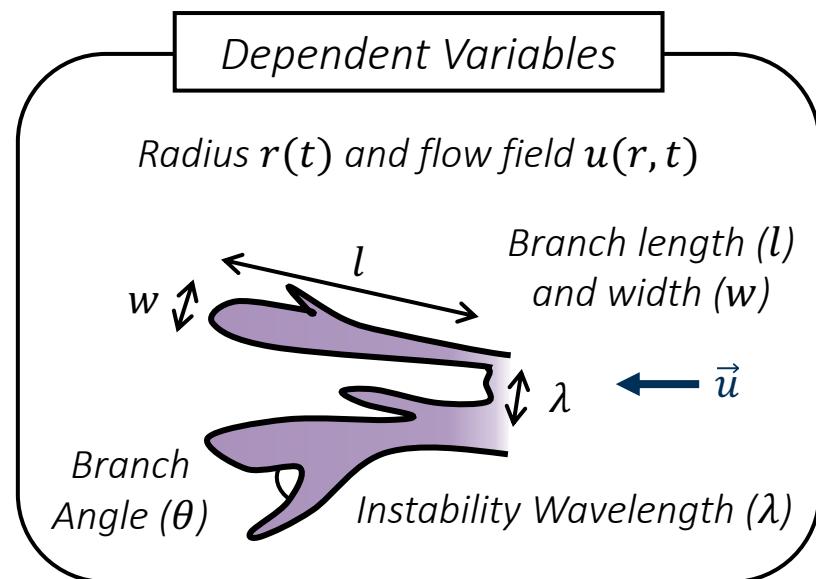
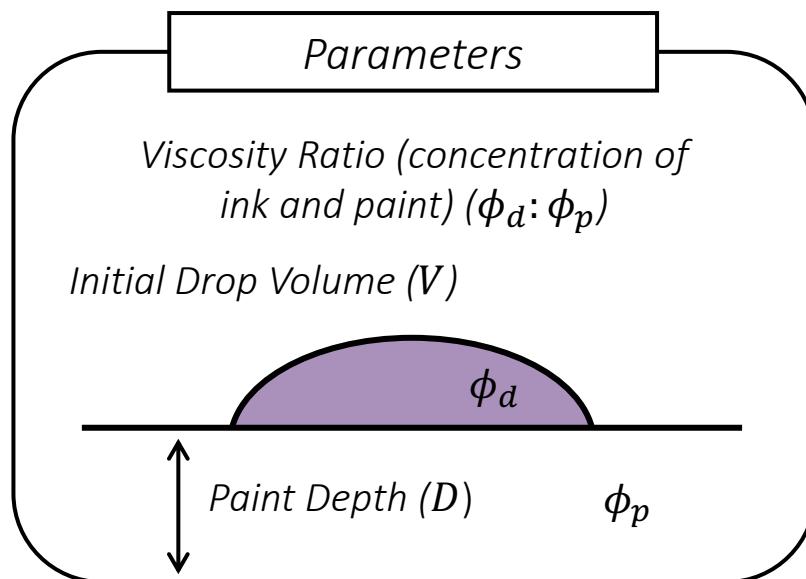
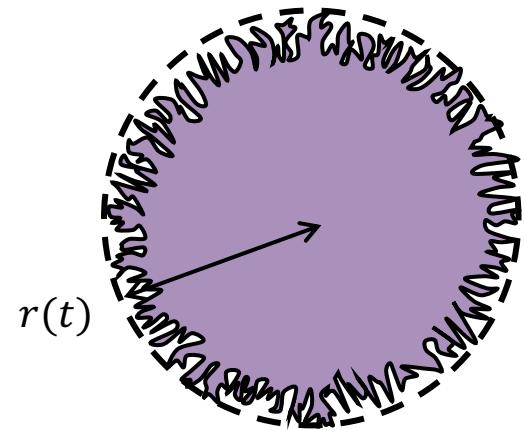
*“The effect of **fractal fingering** can be observed if a droplet of an **ink-alcohol mixture** is deposited onto diluted acrylic paint. How are the **geometry** and **dynamics** of the fingers influenced by **relevant parameters?**”*



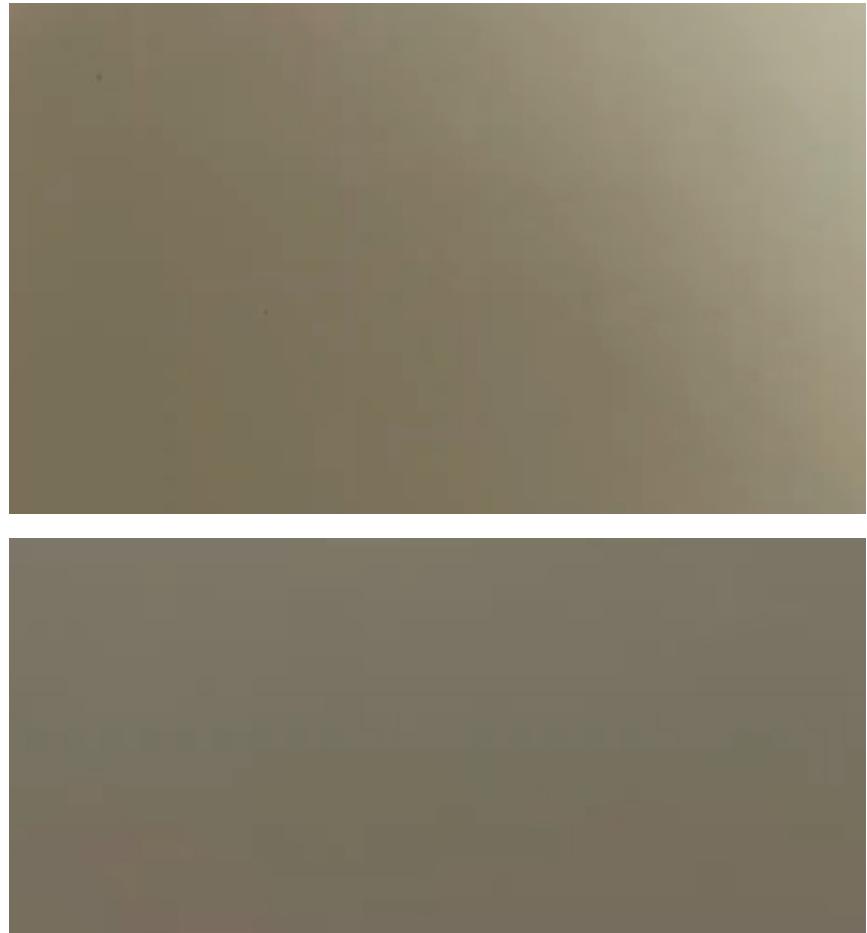
# Tasks

"The effect of **fractal fingering** can be observed if a droplet of an **ink-alcohol mixture** is deposited onto diluted acrylic **paint**. How are the **geometry** and **dynamics** of the fingers influenced by **relevant parameters**?"

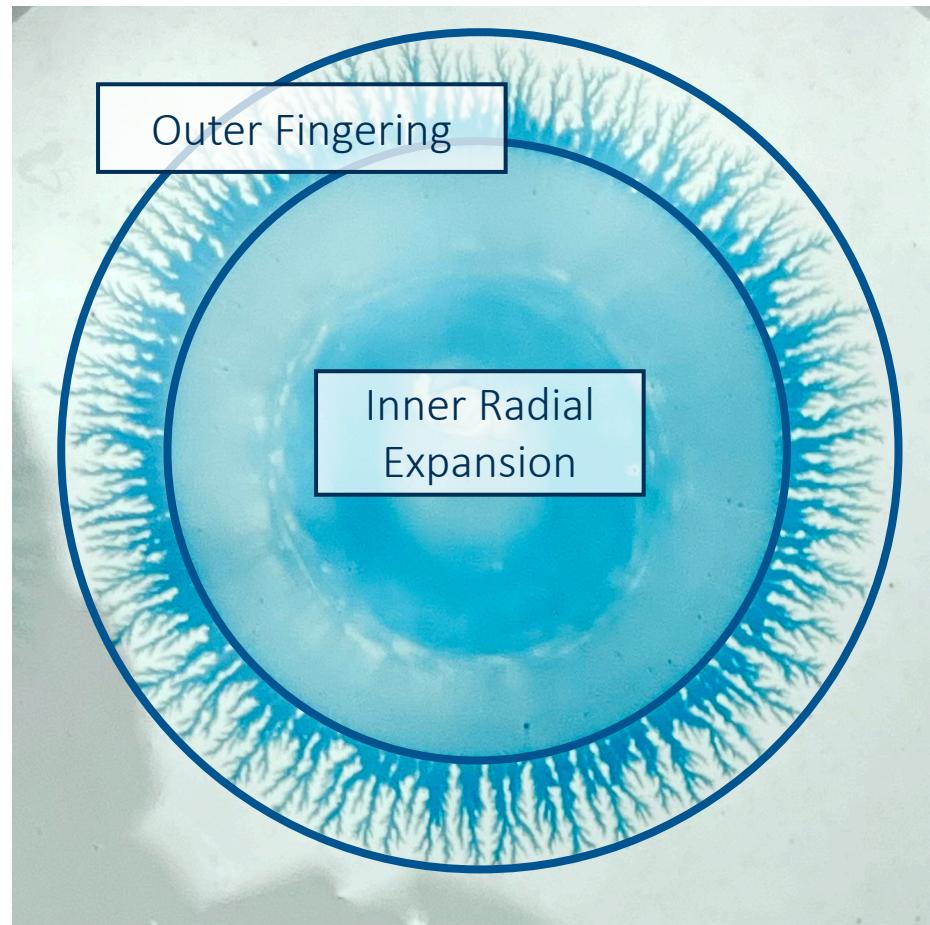
- 1 Explore the **geometry** and **dynamics** of the fractal
- 2 Quantify **conditions** for branching
- 3 Create a **reproducible** experiment



# Phenomenon



# Phenomenon



# Overview

1

## Experimental Setup

*General Setup, Viscosity Measurements*

2

## Droplet Dynamics

*Mass Conservation, Evaporation, Velocity Vector Field*

3

## Branching Geometry

*Fractal Dimension, Saffman-Taylor Instability*

4

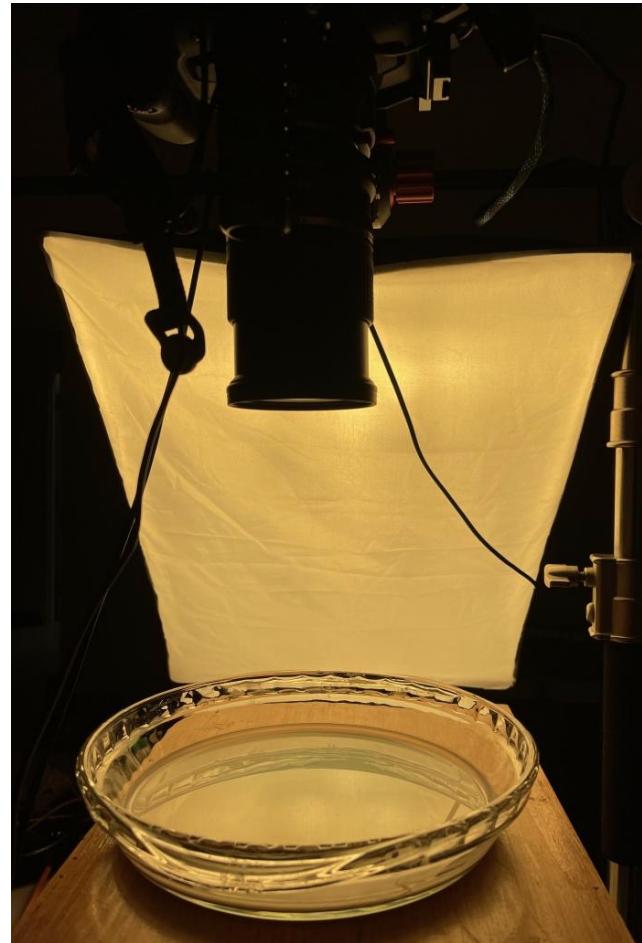
## Insights

*Key Parameters, Further Investigations, Conclusion*

# Experimental Setup

# Experimental Setup

Camera



Lighting

Paint

# Experimental Measurement



Digital Gram Scale  
 $(\pm 0.005\text{g})$



Micropipette  
 $(\pm 0.6\%)$



200 mL graduated cylinder  
 $(\pm 2\text{mL})$

# Viscometry Setup



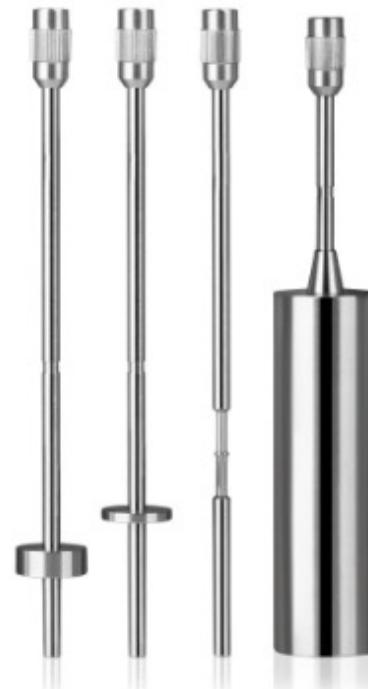
Experimental Setup

Droplet Dynamics

Branching Geometry

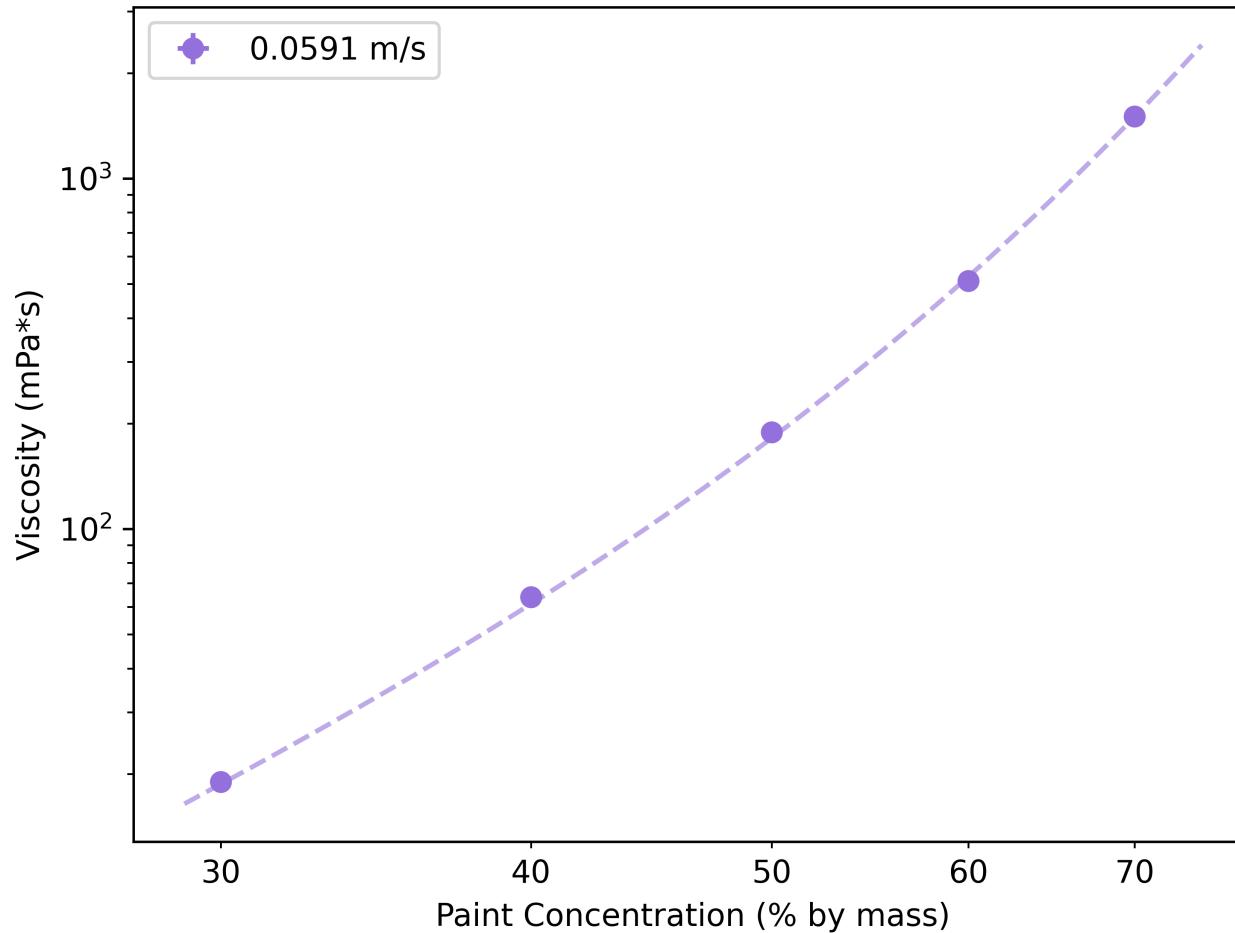
Insights

*Varying Rotor Sizes*



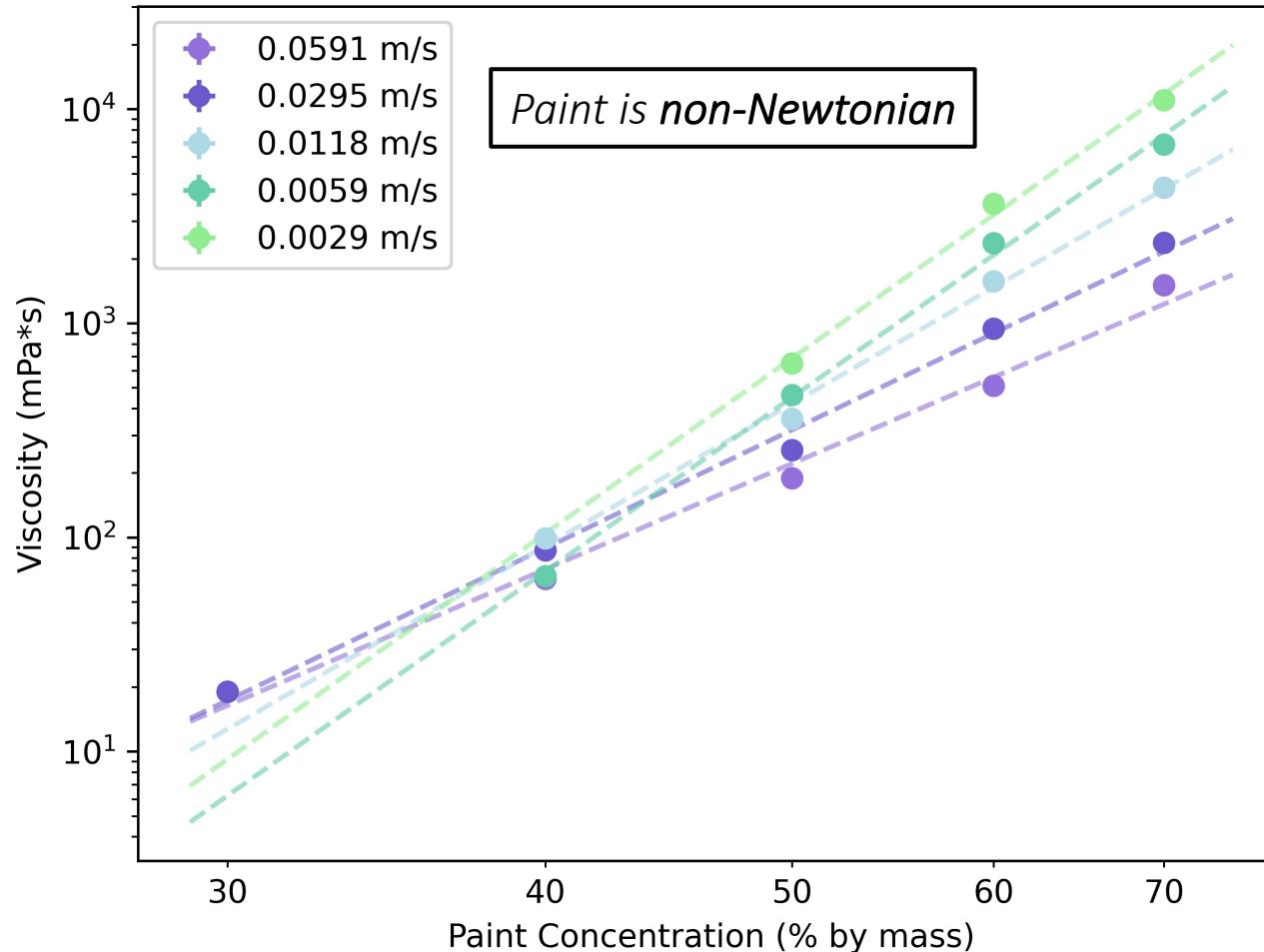
# Viscosity of Paint

Viscosity vs. Paint Concentration for Varying Velocities



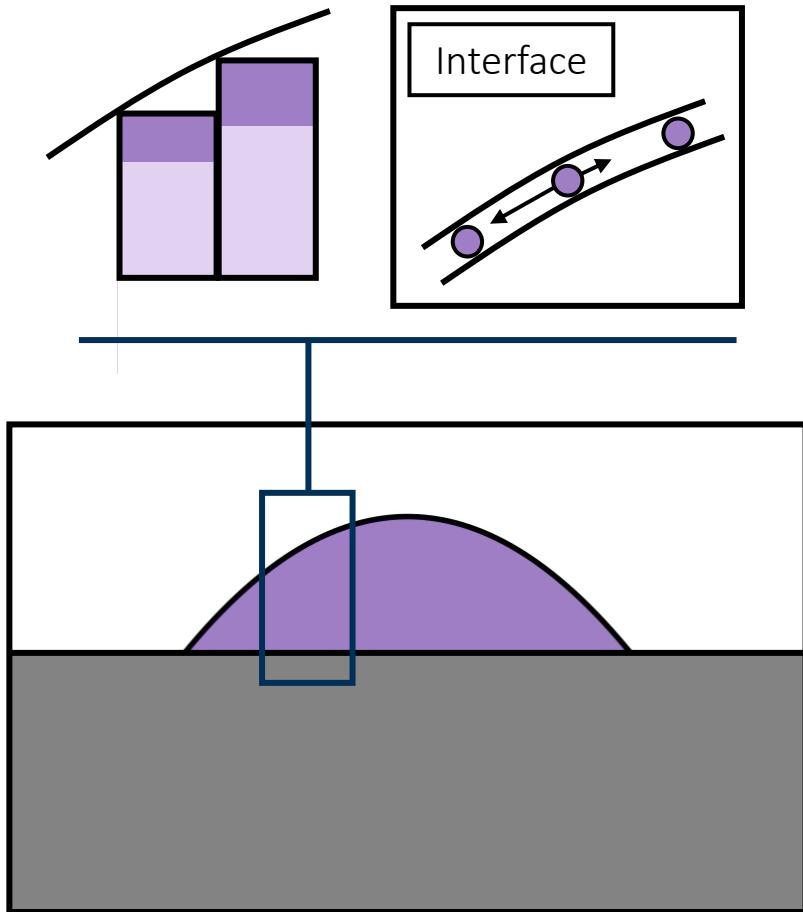
# Viscosity of Paint

Viscosity vs. Paint Concentration for Varying Velocities



# Droplet Dynamics

# Qualitative Explanation



Experimental Setup

Droplet Dynamics

Branching Geometry

Insights

1 *Uneven evaporation along droplet profile*

2 *Radial gradients*

3 *Marangoni flow and radial expansion*

# Assumptions



*Evaporation of water is negligible*



*Ambient room temperature and humidity remains constant*

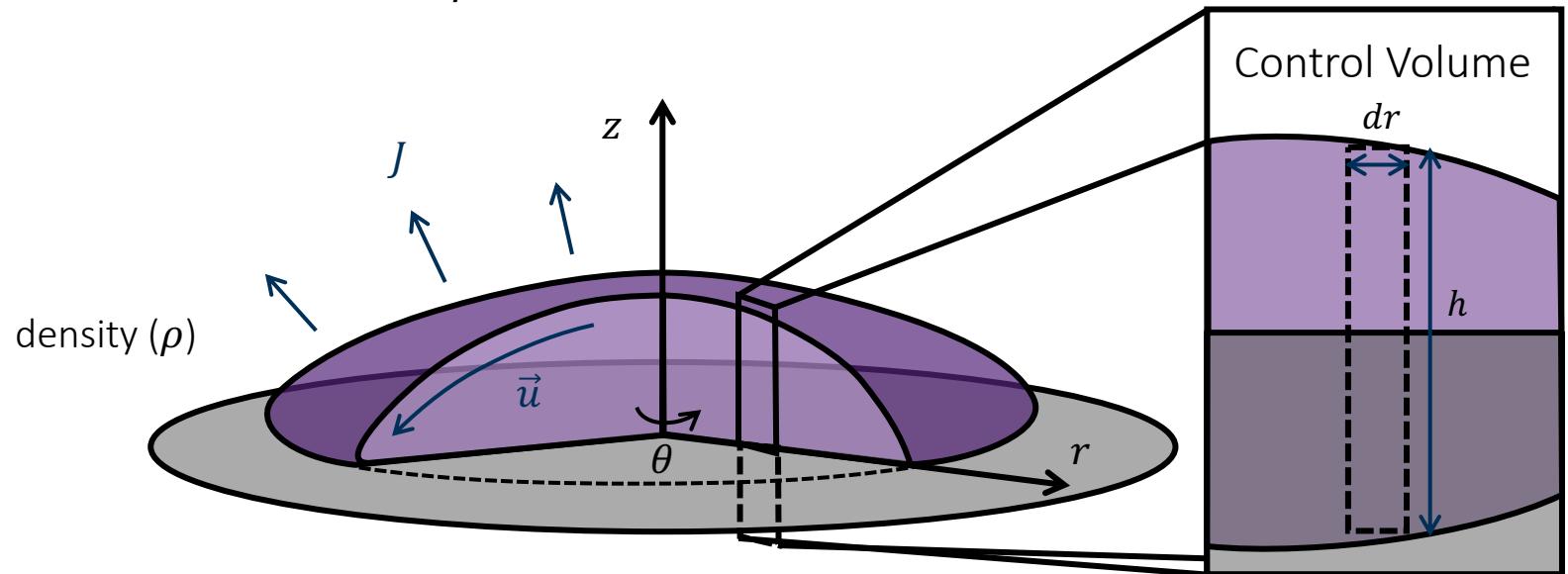


*The evaporation during the dropping process is negligible*

# Continuity

For axisymmetric flow, by mass conservation,

$$\frac{\partial h}{\partial t} = -\frac{1}{\rho} J + \frac{1}{r} \frac{\partial Q}{\partial r} \quad \text{where} \quad \vec{Q} = h \vec{u}$$



# Volumetric Flow



*Air is shear-free*

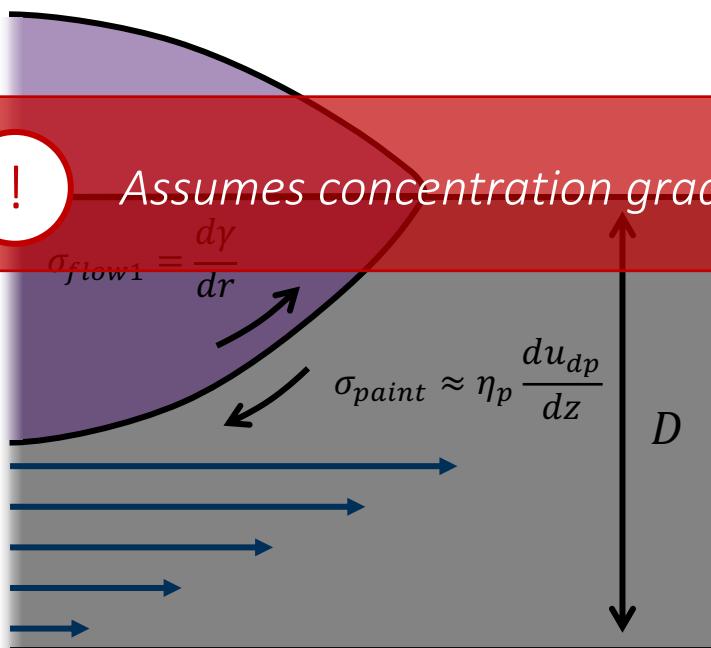
*Equating shearing at drop-paint interface:*

$$\frac{\partial \gamma}{\partial r} = \eta_p \frac{\partial u_{dp}}{\partial z}$$



*Assumes concentration gradient is linear over time and fluids are immiscible*

*No slip condition:  $u(z) = 0 \Big|_{z=-D}$  so,  $\frac{\partial u}{\partial z} \Big|_{z=-D}$*



$$u_{dp} = \frac{D}{\eta_p} \frac{\partial \gamma}{\partial r}$$

# Stage I: Marangoni Stress

$\Delta\sigma$  = surface tension difference

$\sigma_a$  = surface tension of diluted acrylic paint

$\sigma_d$  = surface tension of droplet

$S$  = mixing length

$R$  = radius of droplet film

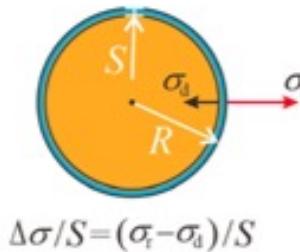
$k$  = constant of proportionality

$\mu_a$  = dynamic viscosity of diluted acrylic paint

$\rho_a$  = density of diluted acrylic paint

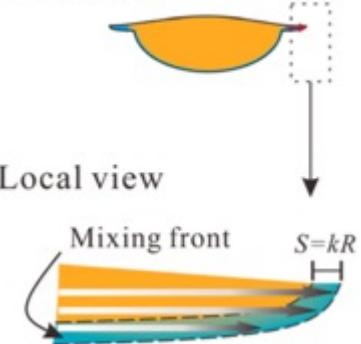
$t$  = time

Top view



$$\Delta\sigma/S = (\sigma_f - \sigma_d)/S$$

Front view



When the droplet is first deposited onto the diluted acrylic paint bath,

$$\frac{\Delta\sigma}{S} = \frac{\sigma_a - \sigma_d}{kR} \sim \frac{\mu_a^{1/2} \rho_a^{1/2} R}{t^{3/2}}$$

The first stage power-law relation can be derived as,

$$R \sim k^{1/2} (\sigma_a - \sigma_d)^{1/2} \mu_a^{-1/4} \rho_a^{-1/4} t^{3/4}$$

# Stage II: Normal Diffusion

$\Delta\sigma$  = surface tension difference

$\sigma_a$  = surface tension of diluted acrylic paint

$\sigma_d$  = surface tension of droplet

$S$  = mixing length

$R$  = radius of droplet film

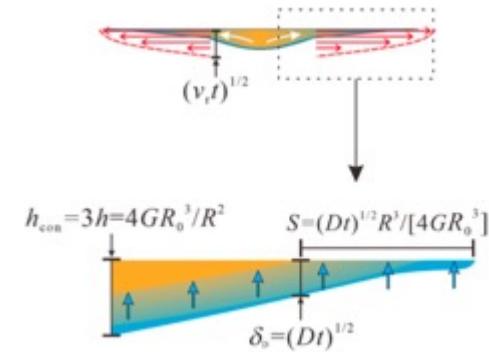
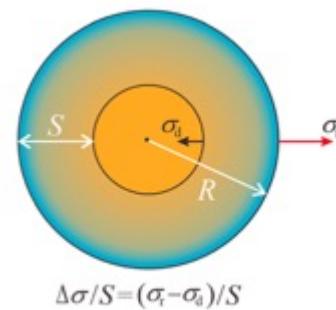
$k$  = constant of proportionality

$h_{cone}$  = height of cone

$\delta_D$  = mass diffusion layer thickness

$D$  = binary–liquid mutual diffusion coefficient

$G$  = volume fraction of droplet



Once the normal diffusion dominates,

$$\frac{R}{S} = \frac{h_{cone}}{(Dt)^{1/2}} = \frac{4GR_0^3}{R^2(Dt)^{1/2}}$$

So,

$$R \sim \frac{(4(\sigma_a - \sigma_d)GR_0^3)^{1/4}}{(\mu_a \rho_a D)^{1/8}} t^{1/4}$$

# Stage III: Tangential Stretch

$\Delta\sigma$  = surface tension difference

$\sigma_a$  = surface tension of diluted acrylic paint

$\sigma_d$  = surface tension of droplet

$S$  = mixing length

$R$  = radius of droplet film

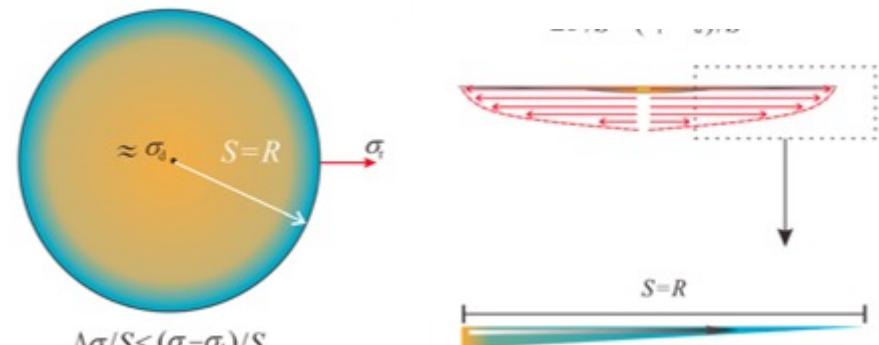
$k$  = constant of proportionality

$h_{cone}$  = height of cone

$\delta_D$  = mass diffusion layer thickness

$D$  = binary–liquid mutual diffusion coefficient

$G$  = volume fraction of droplet



Once the tangential stretch dominates,

$$\frac{\Delta\sigma}{S} = \frac{(\sigma_a - \sigma_d)}{R} = \frac{4GR_0^3}{R^2(Dt)^{1/2}}$$

So,

$$R \sim \mu_a^{1/2} \rho_a^{1/2} t^{3/4}$$

# Three-Stage Flow

*Stage I: Marangoni Flow*

$$R^* \sim k^{-1/2} t_1^{*1/4}$$

$$t_1^* = \frac{t}{\mu_a^3 (\sigma_a - \sigma_d)^{-2} \rho_a^{-1}}$$

*Stage II: Normal Diffusion*

$$R^* \sim t_2^{*1/4}$$

$$t_2^* = \frac{t}{(\mu_a^{\frac{17}{2}} D^{\frac{1}{2}} (\sigma_a - \sigma_d)^{-5} \rho_a^{-\frac{7}{2}} R_0^{-3})}$$

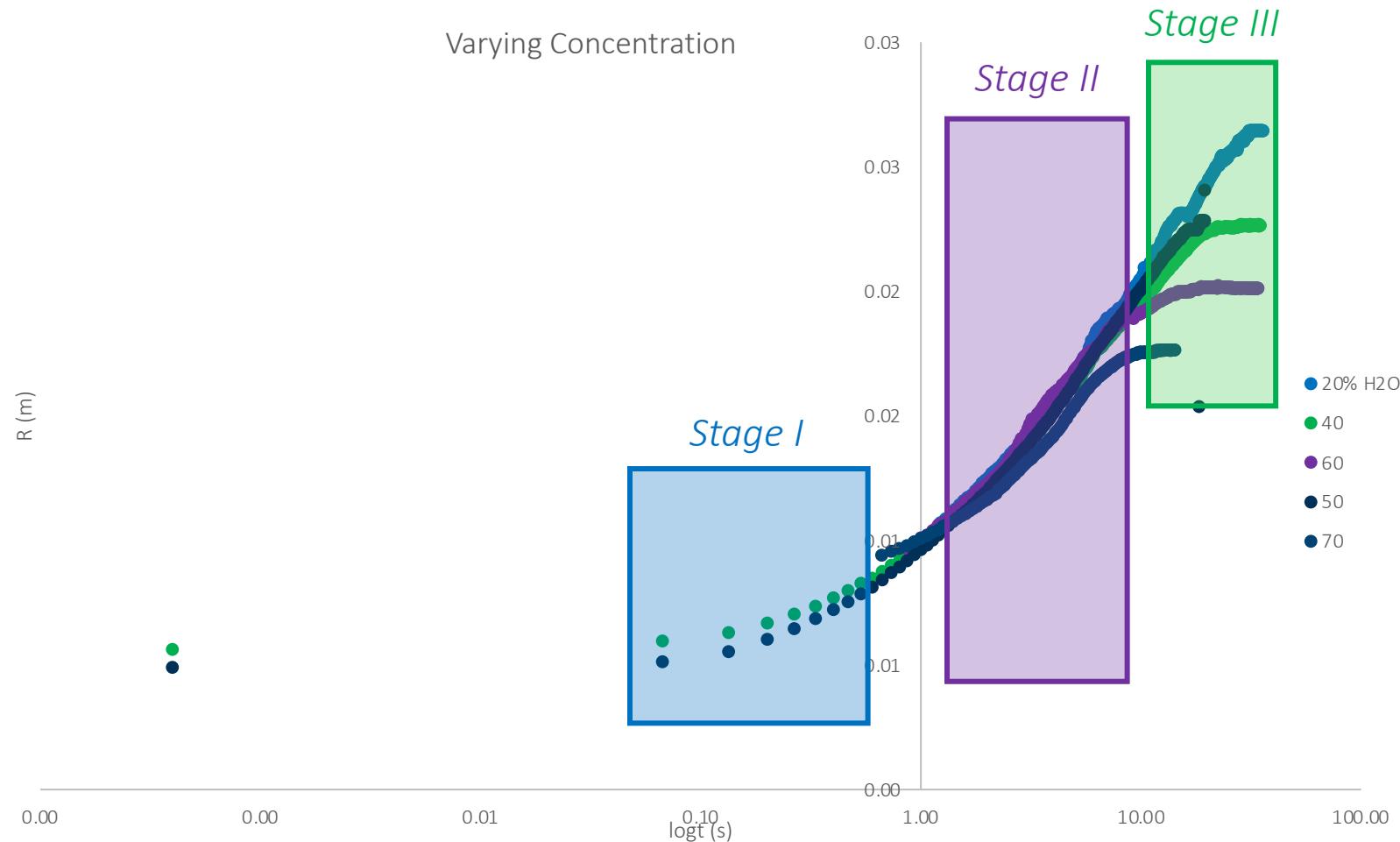
*Stage III: Tangential Stretch*

$$R^* \sim t_3^{*3/4}$$

$$t_3^* = \frac{t}{\mu_a^3 (\sigma_a - \sigma_d)^{-2} \rho_a^{-1}}$$

(Jia et al., 2022)

# Experimental Verification



Experimental Setup

Droplet Dynamics

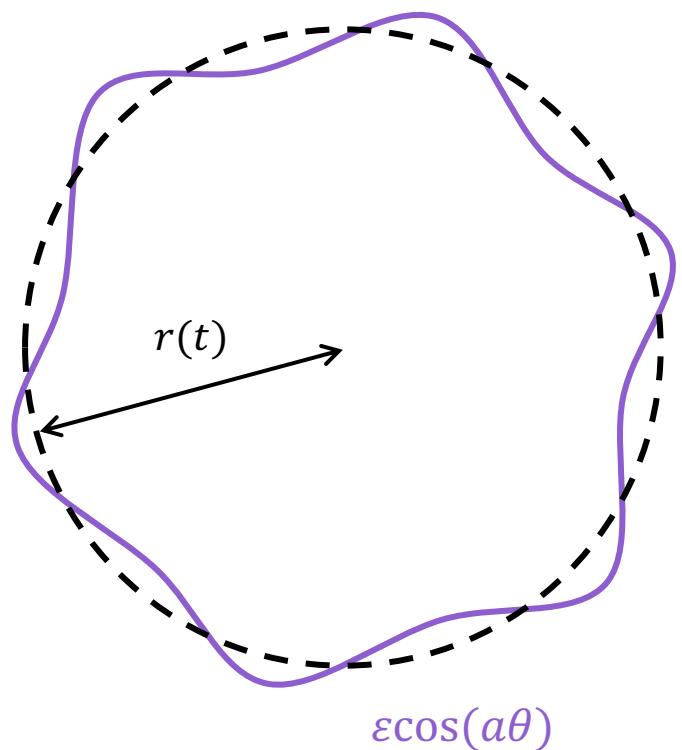
Branching Geometry

Insights

# Branching Geometry

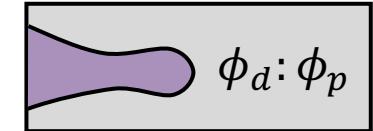
# Saffman-Taylor Instability

*Wavelike perturbation of magnitude  $\varepsilon$*



*Viscosity Difference*

*Drives linear instability*



1

*Small sinusoidal perturbation in flow*

2

*Perturbations of different wavelengths grow/shrink depending on parameters*

3

*Self similarity of perturbations that grow fastest result in fractals*

# Experimental Observations

80% Isopropanol



30% Isopropanol

Concentrated  
Fingers at Edge



*Weaker Marangoni Flow results in inward fractalling*

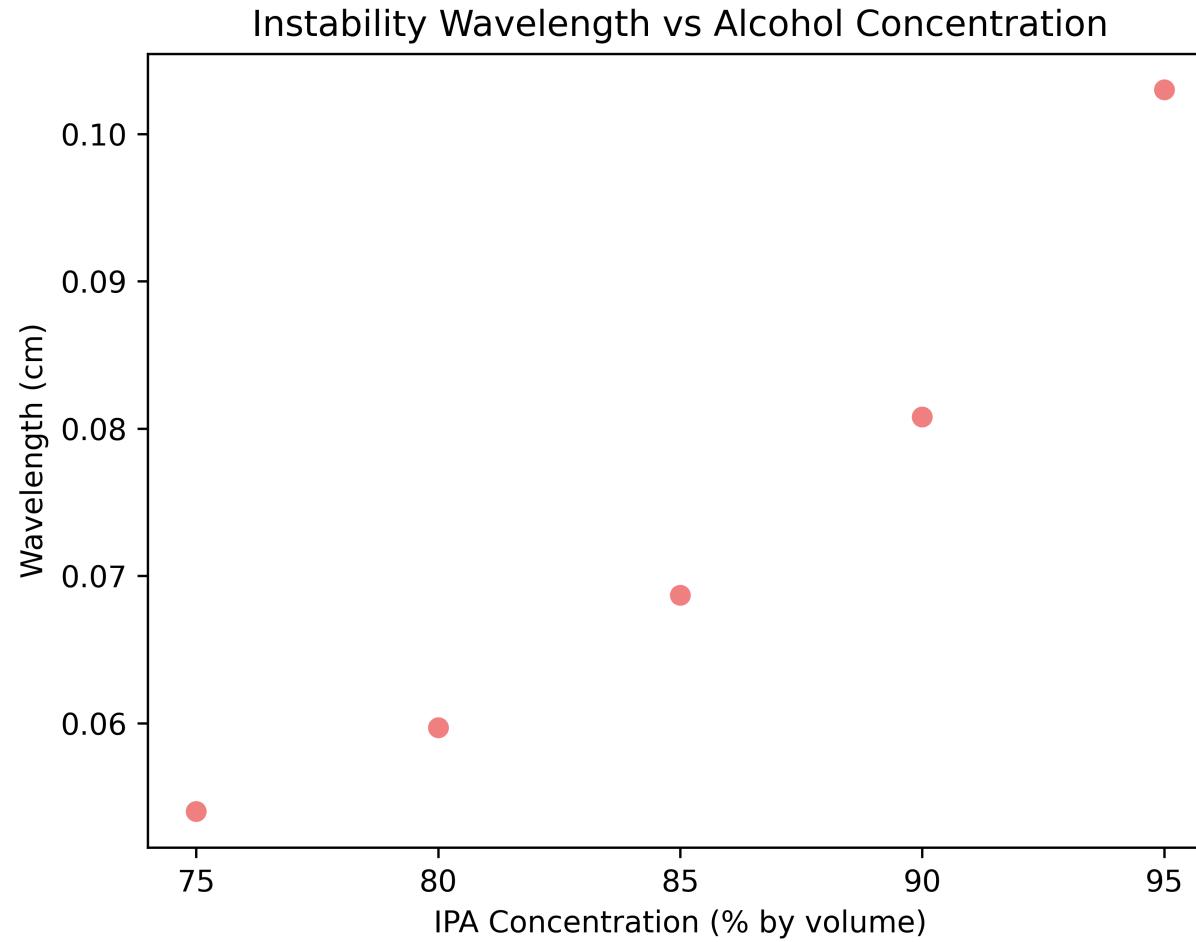
Outer Expansion



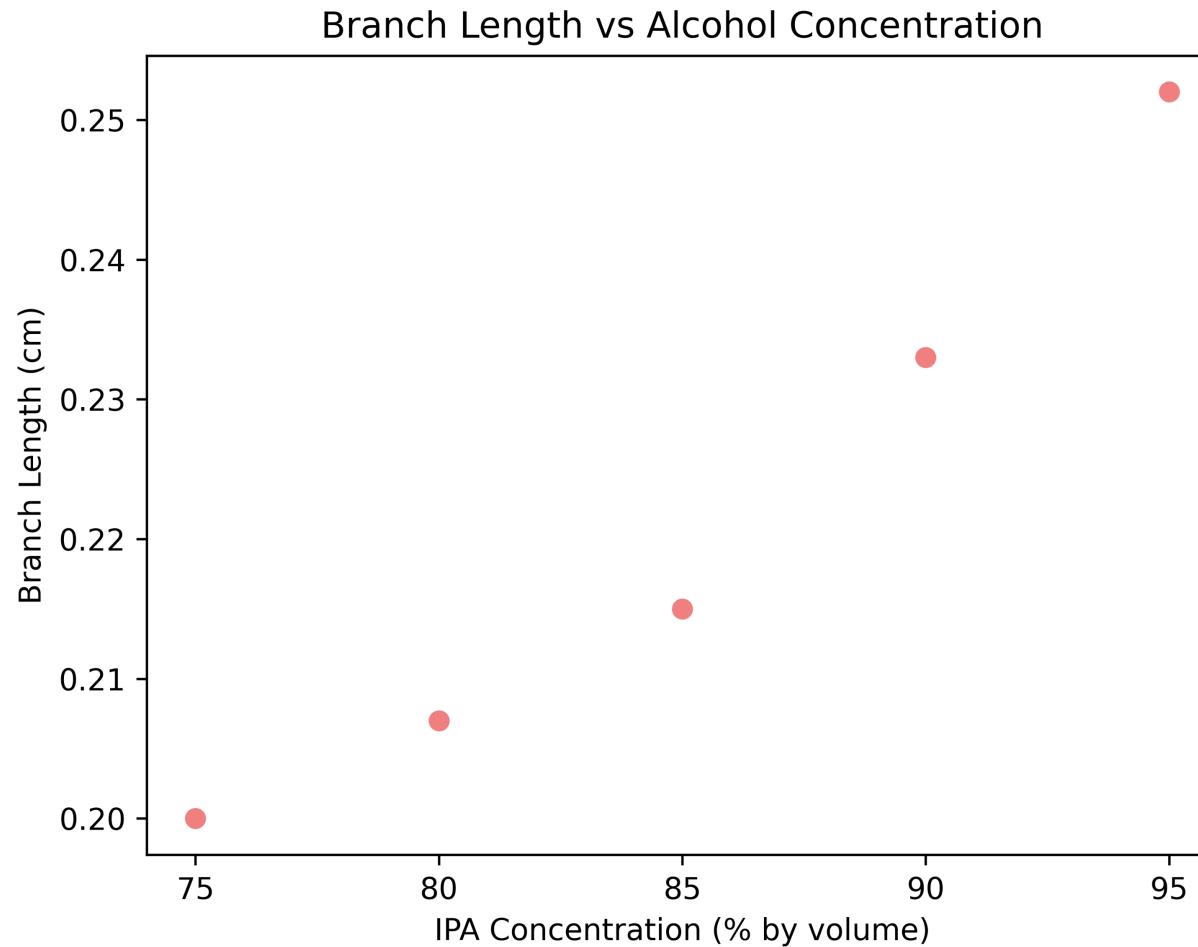
Inward Fractal



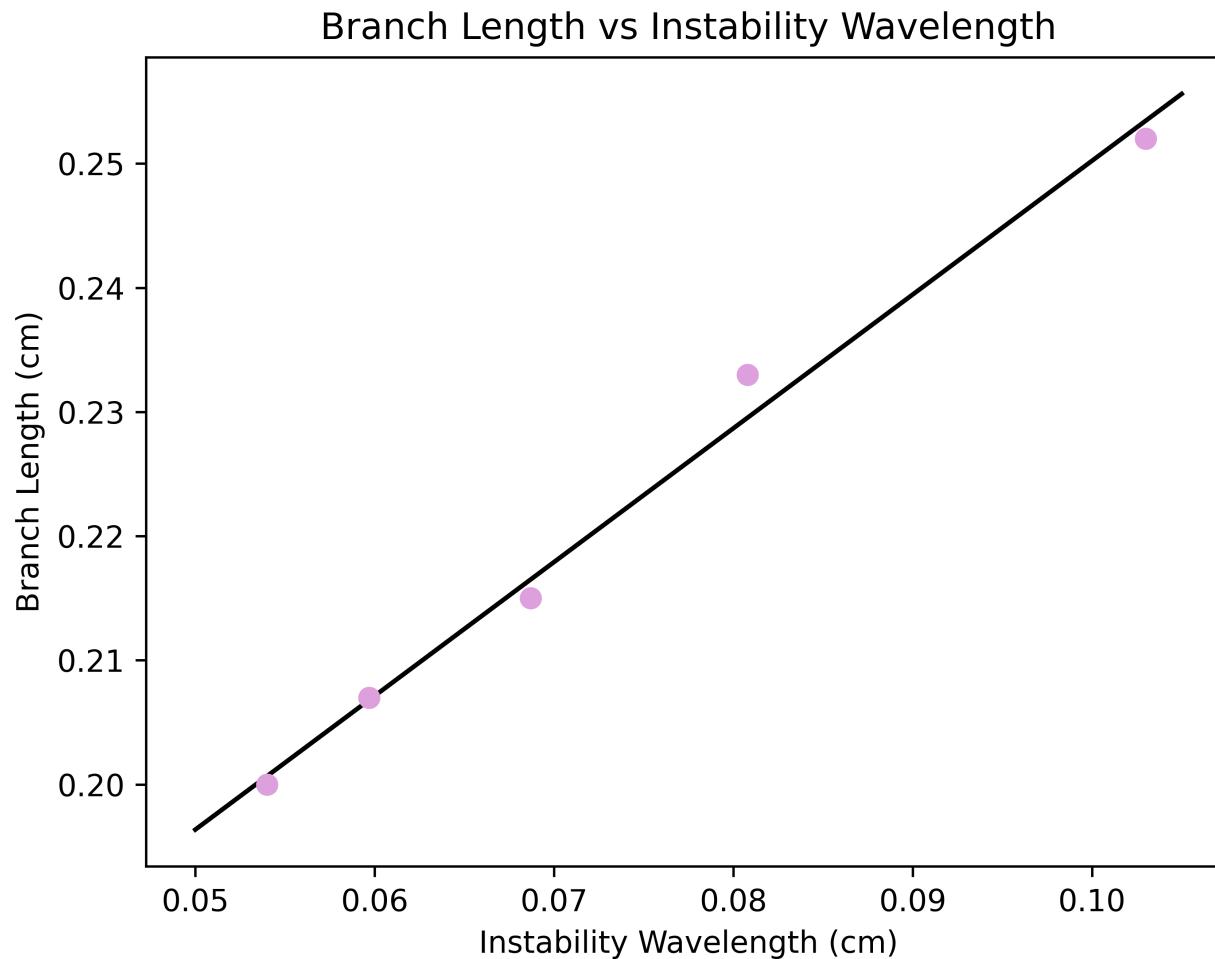
# Wavelength of Instability



# Branch Length

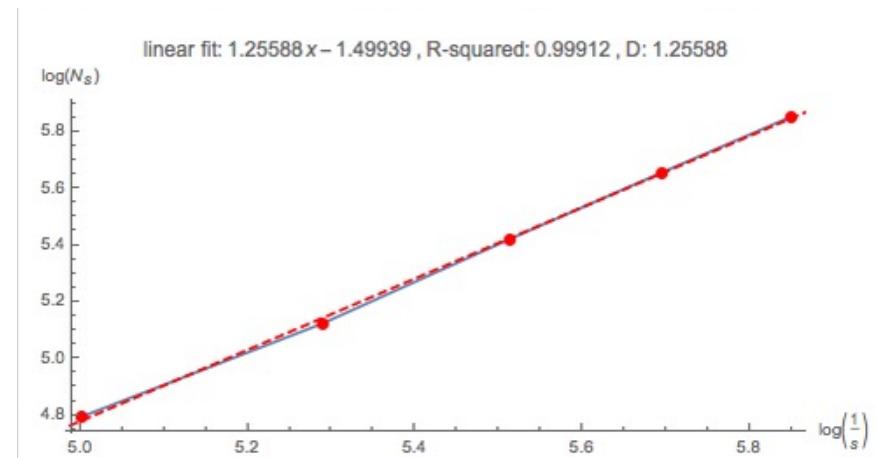
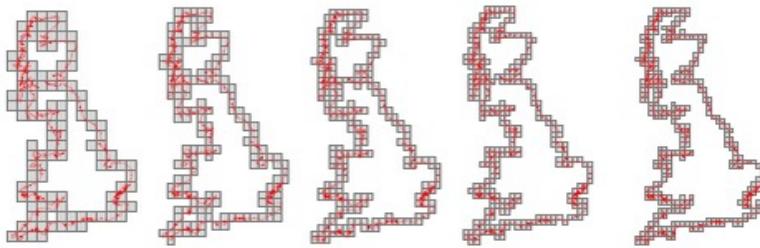


# Linear Instability



# Fractal Dimension

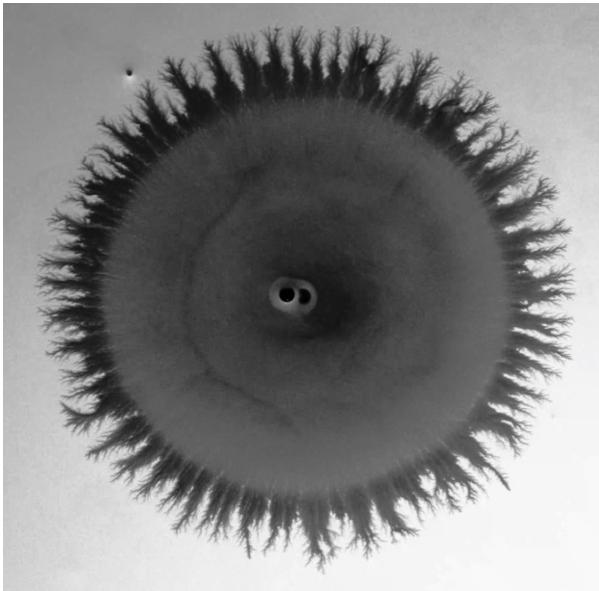
- Fractals are defined shapes with a **Hausdorff Dimension** that is higher than the integer **topological dimension**
- Hausdorff Dimension can be calculated using a variety of different methods
  - Box Counting Dimension
- Dimension =  $\lim_{\varepsilon \rightarrow 0} \frac{\log N(\varepsilon)}{\log(1/\varepsilon)}$



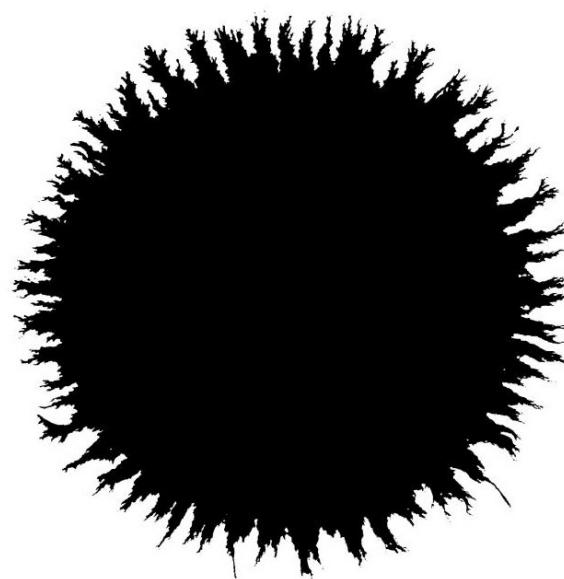
Examples from Wolfram Alpha's [Demonstration](#)

# Fractal Dimension

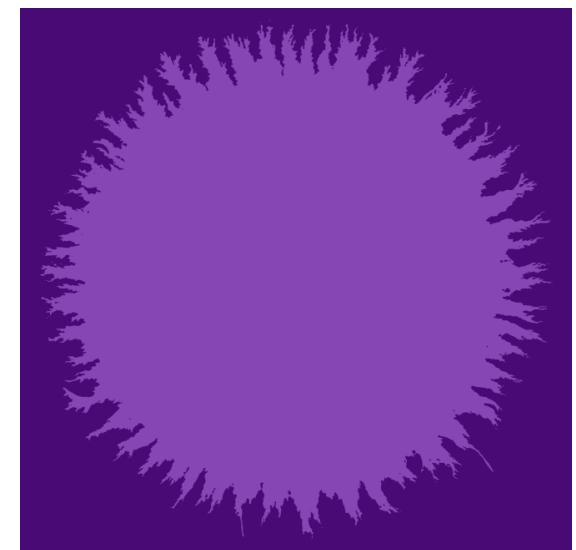
*Raw Video*



*Grayscale*

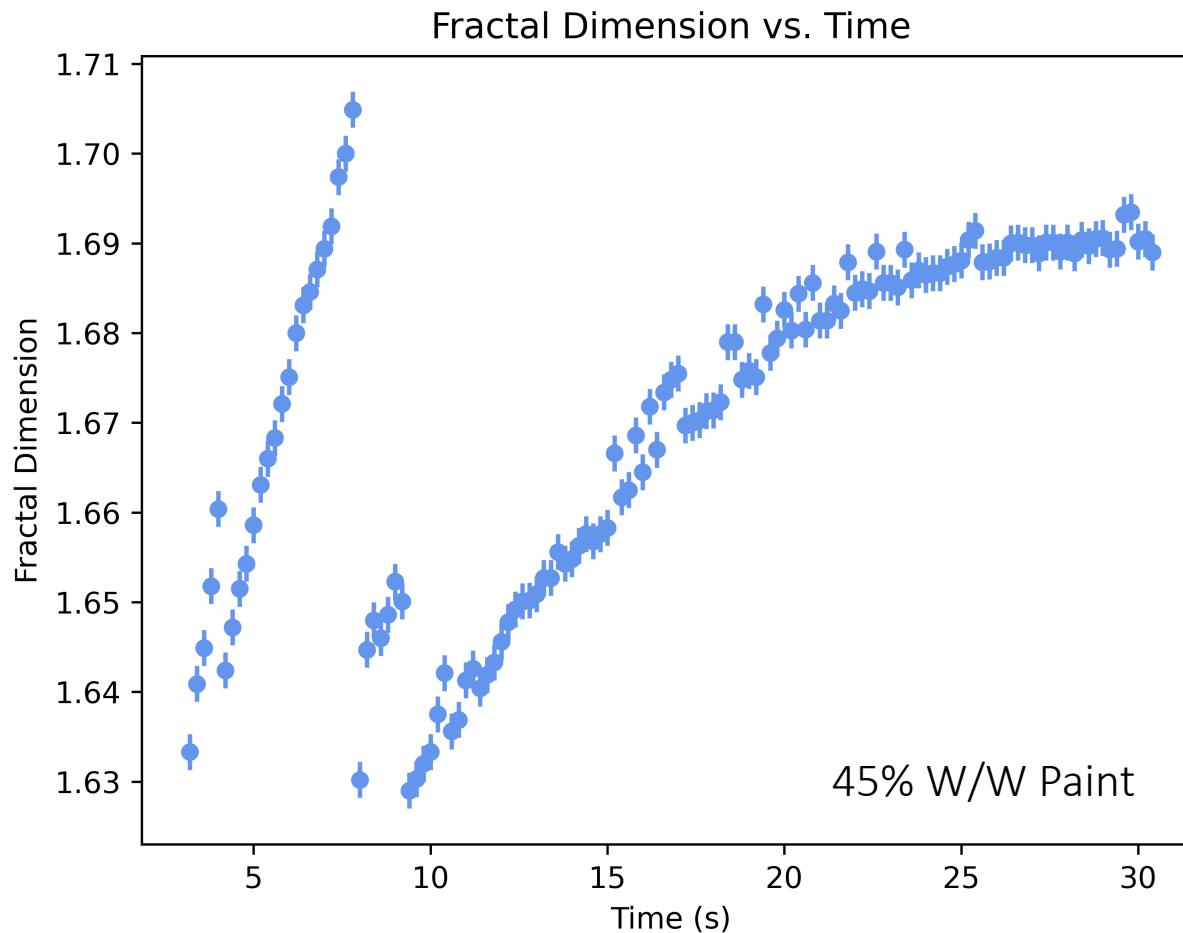


*FracLac Box Counting*

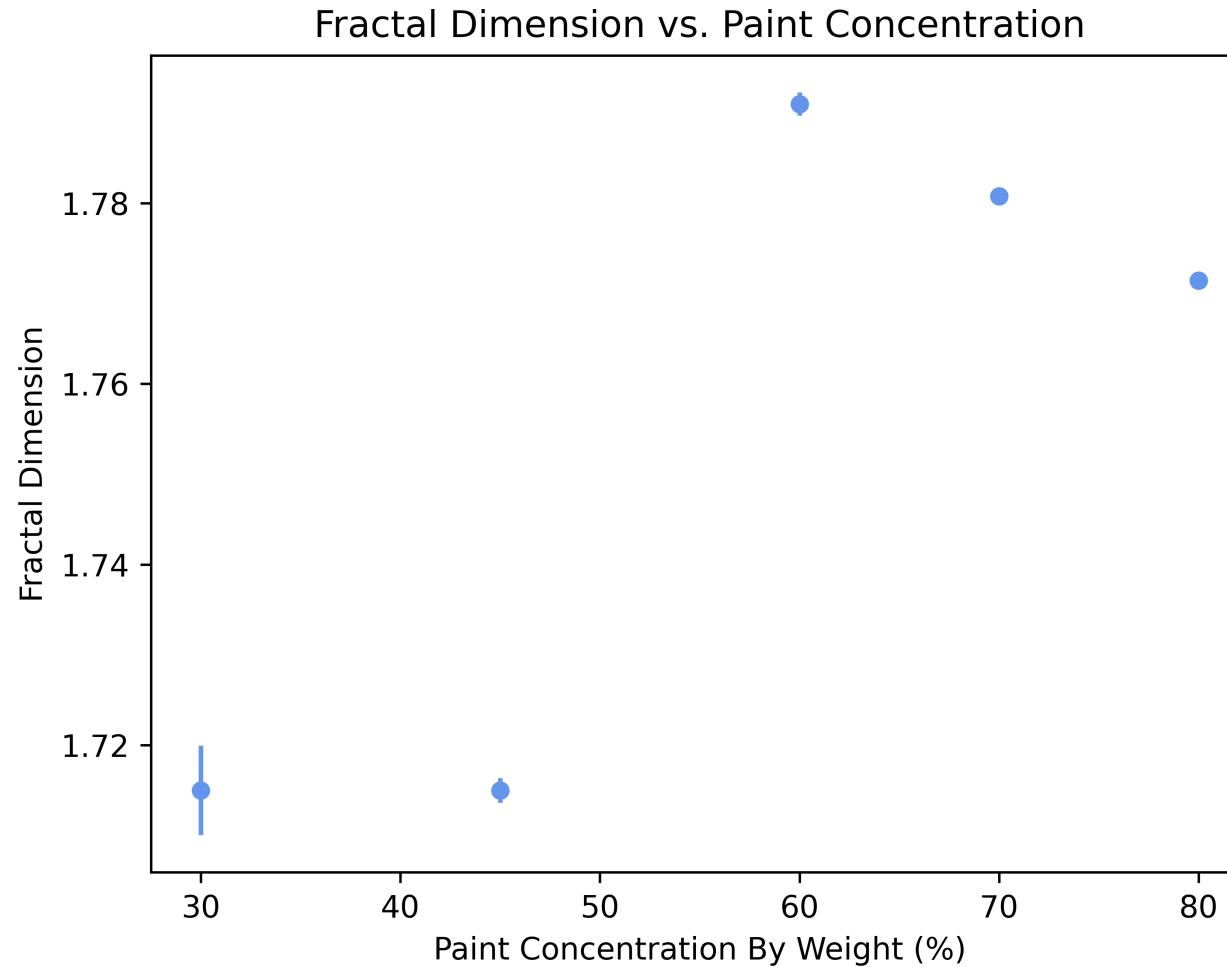


*Able to accurately and efficiently track fractal dimension over time*

# Experimental Data



# Experimental Data



# Insights

# Varying Key Parameters

Flow Velocity

Negative Correlation

Paint Concentration

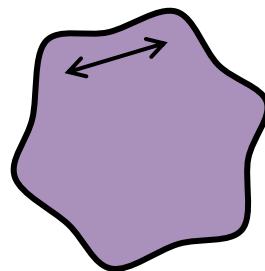
Positive Correlation

Ink Concentration

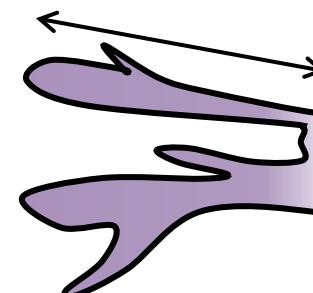
Instability Wavelength

Branch Length

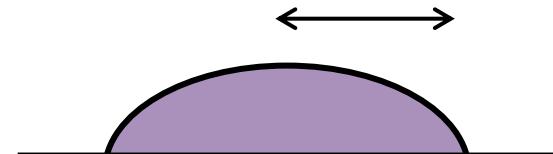
Maximum Radius



Experimental Setup



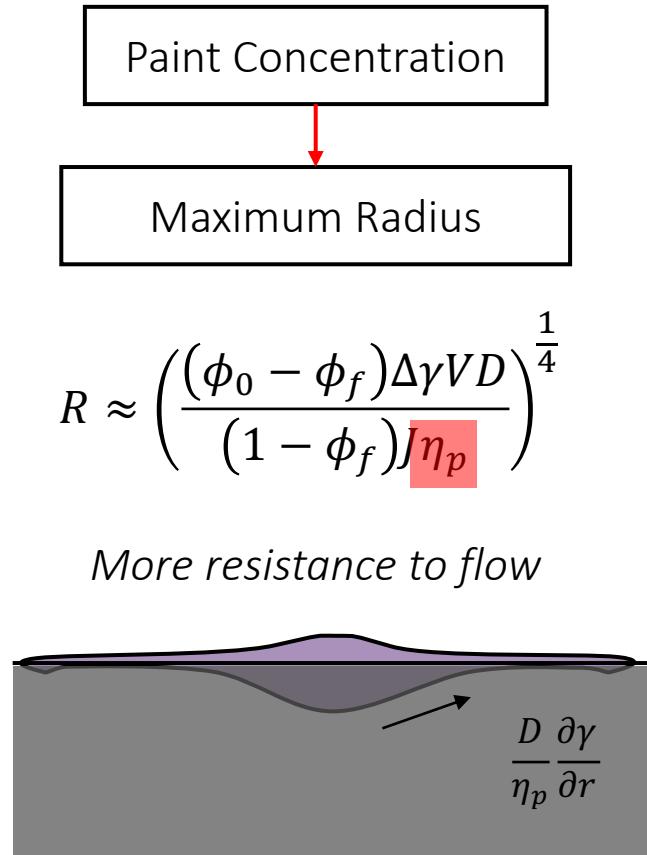
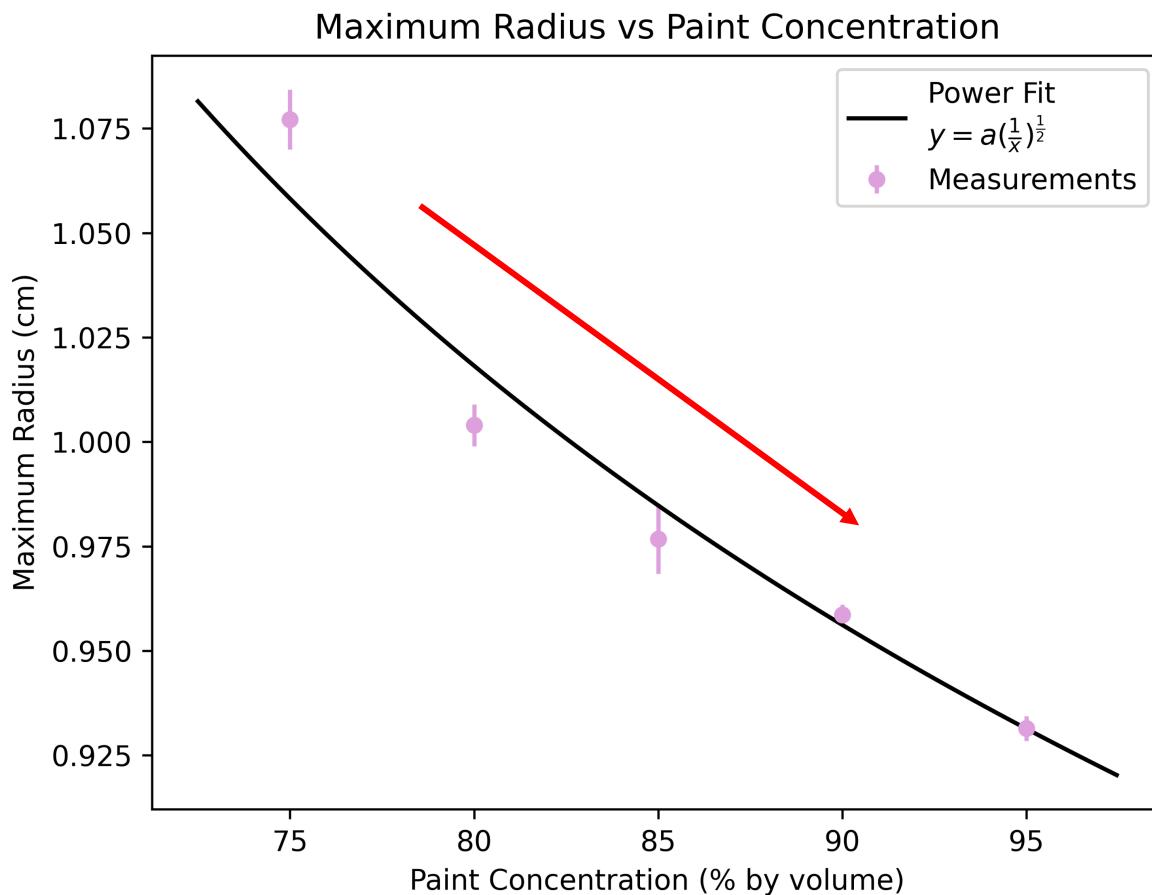
Droplet Dynamics



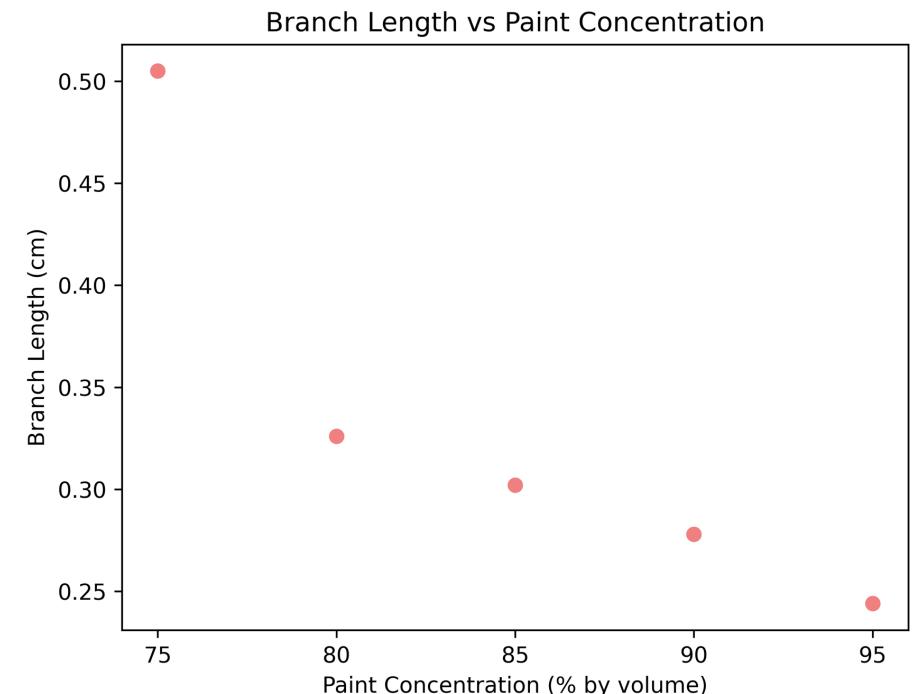
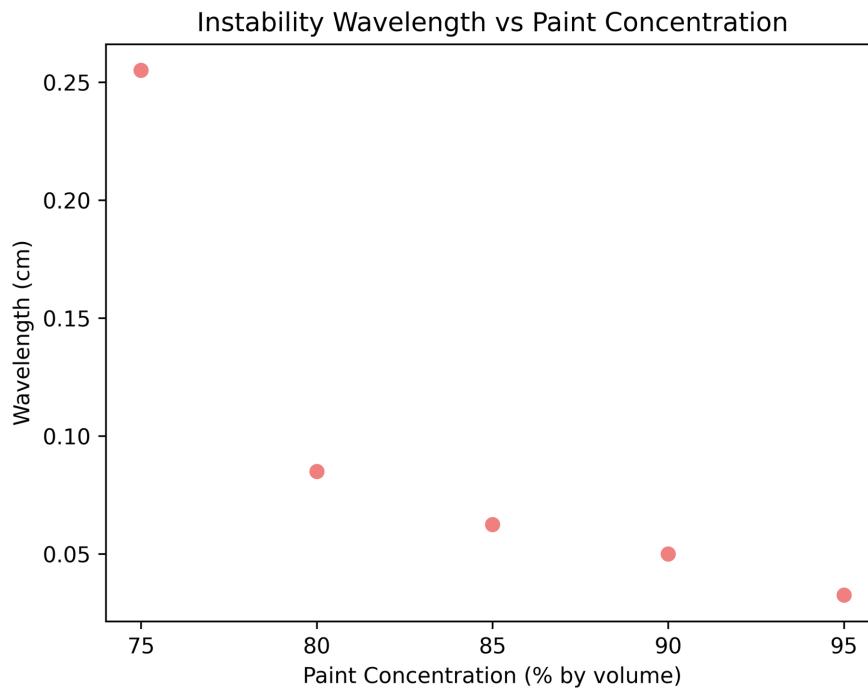
Branching Geometry

Insights

# Varying Key Parameters

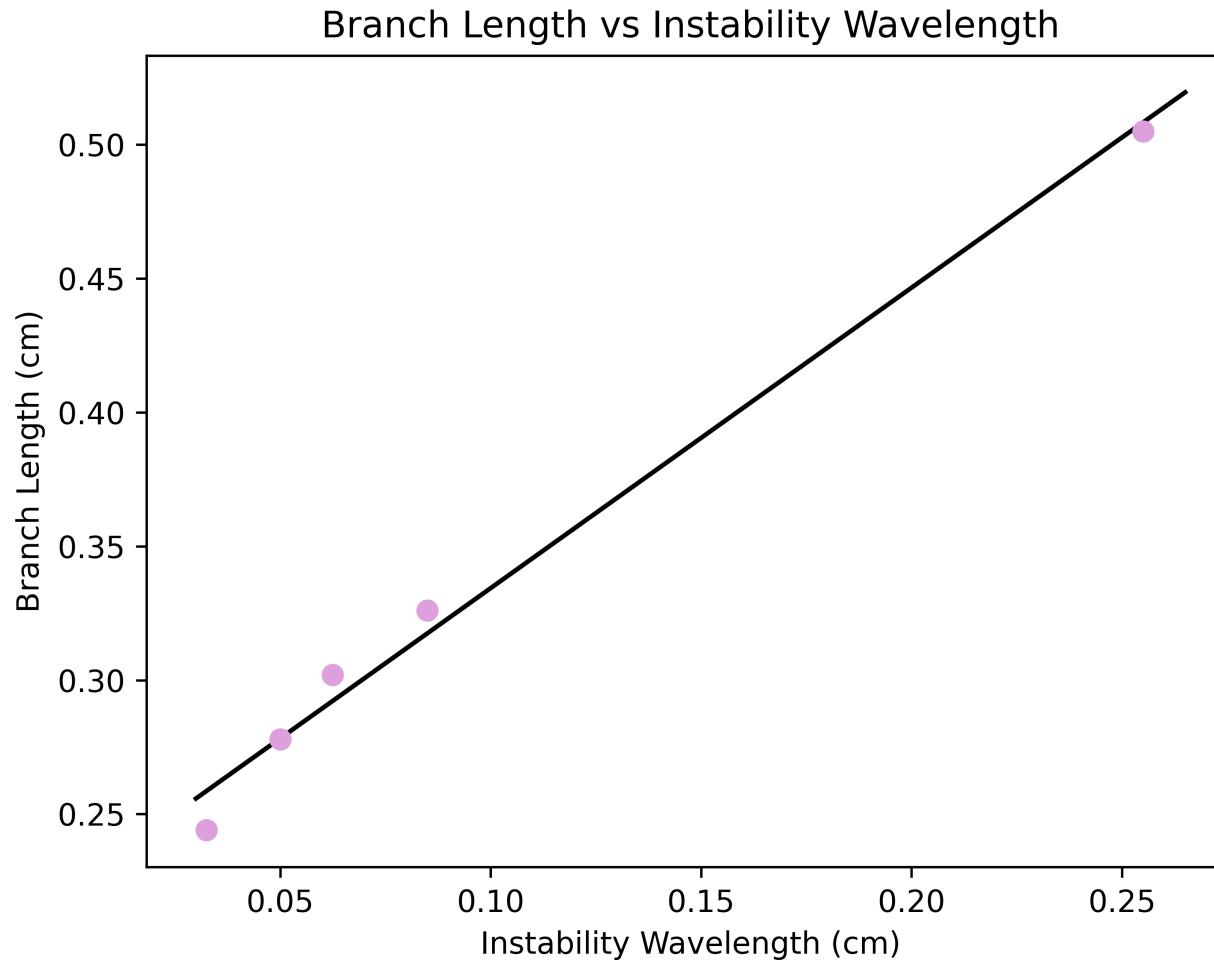


# Varying Key Parameters

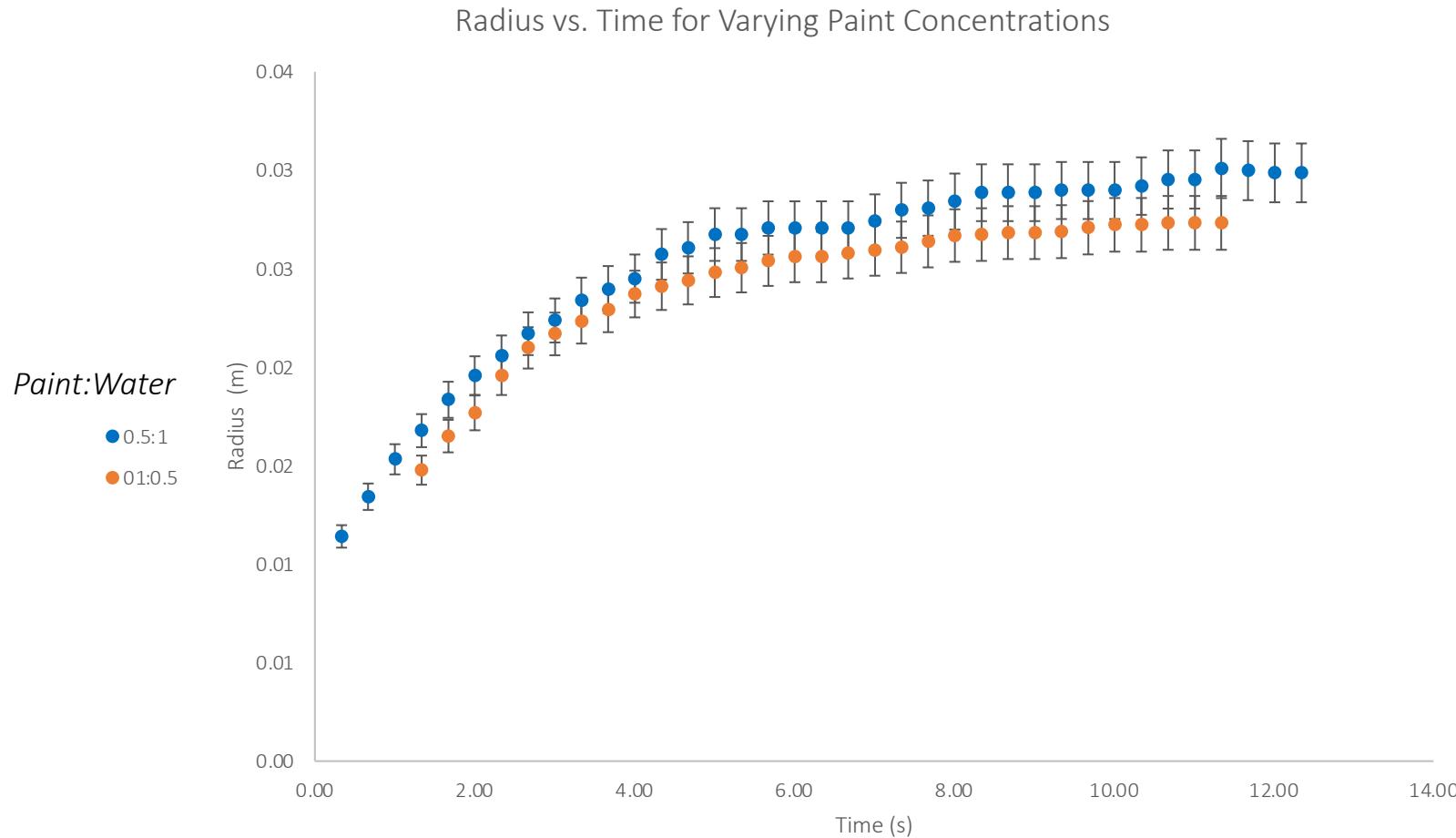


*As flow velocity increases, the instability amplitude and wavelength increase*

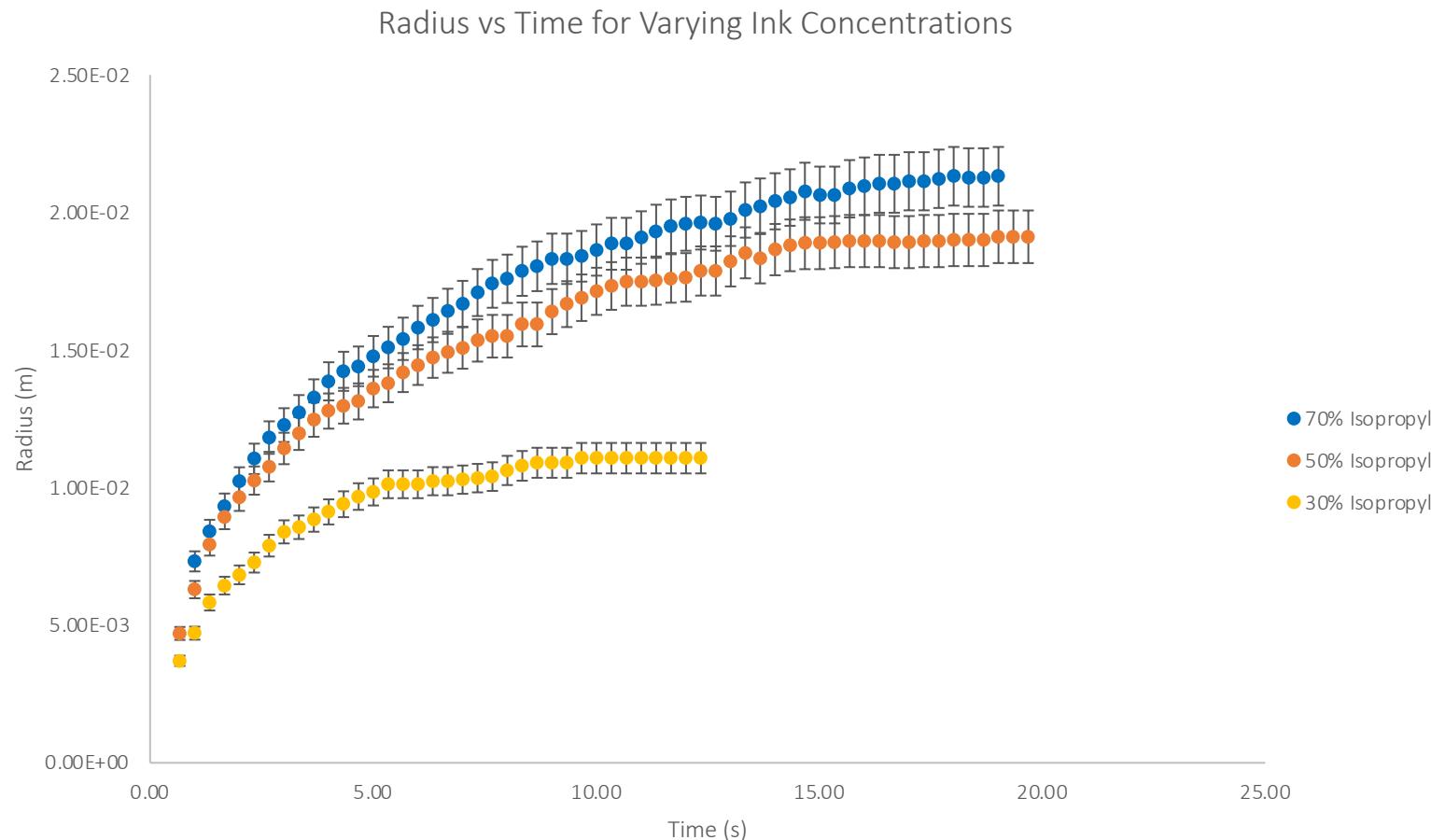
# Varying Key Parameters



# Varying Key Parameters



# Varying Key Parameters



# References

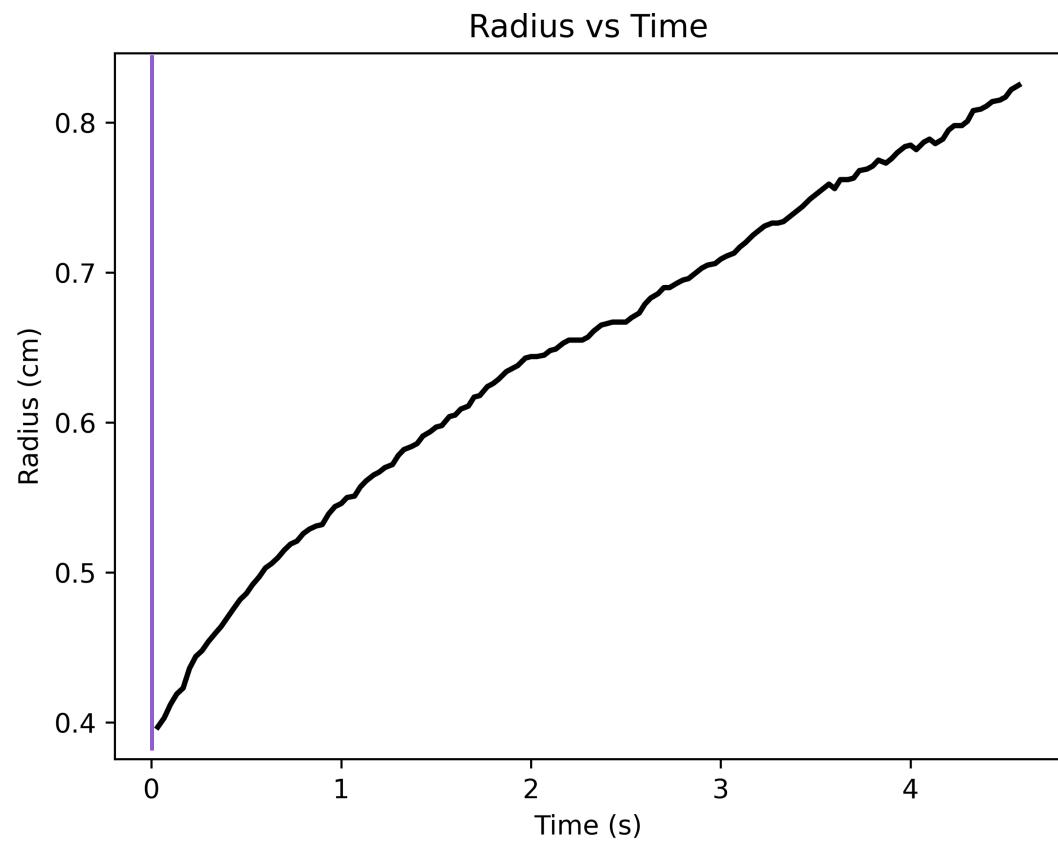
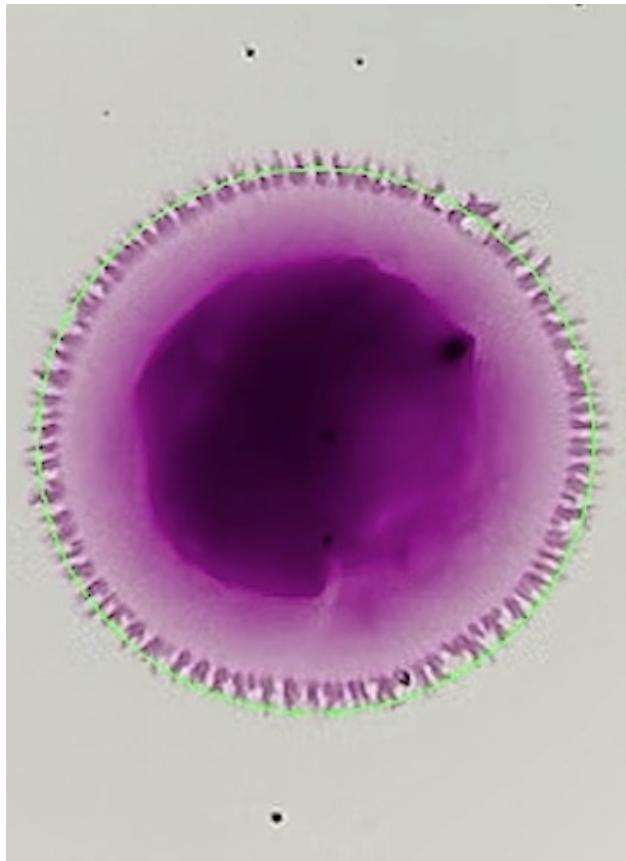
- Gelderblom, H., Diddens, C., & Marin, A. (2022). *Evaporation-driven liquid flow in sessile droplets*. *Soft Matter*, 18(45), 8535–8553.  
<https://doi.org/10.1039/d2sm00931e>
- Jia, F., Wang, T., Peng, X., & Sun, K. (2022). *Three stages of Marangoni-driven film spreading for miscible fluids*. *Physics of Fluids*, 34(12), 121705–121705.  
<https://doi.org/10.1063/5.0132216>
- Keiser, L., Bense, H., Colinet, P., Bico, J., & Reyssat, E. (2017). *Marangoni Bursting: Evaporation-Induced Emulsification of Binary Mixtures on a Liquid Layer*. *Physical Review Letters*, 118(7). <https://doi.org/10.1103/physrevlett.118.074504>
- Kondic. (n.d.). *Saffman-Taylor Instability Motivation*. Retrieved January 20, 2023, from <https://web.njit.edu/~kondic/capstone/2015/saffman-taylor.pdf>
- Lancaster, J. F. (1999). *Introductory Metallurgy of Welding*, 1–39.  
<https://doi.org/10.1533/9781845694869.1>

Thank you for listening!

# Appendix

Appendix

# Radius Tracking



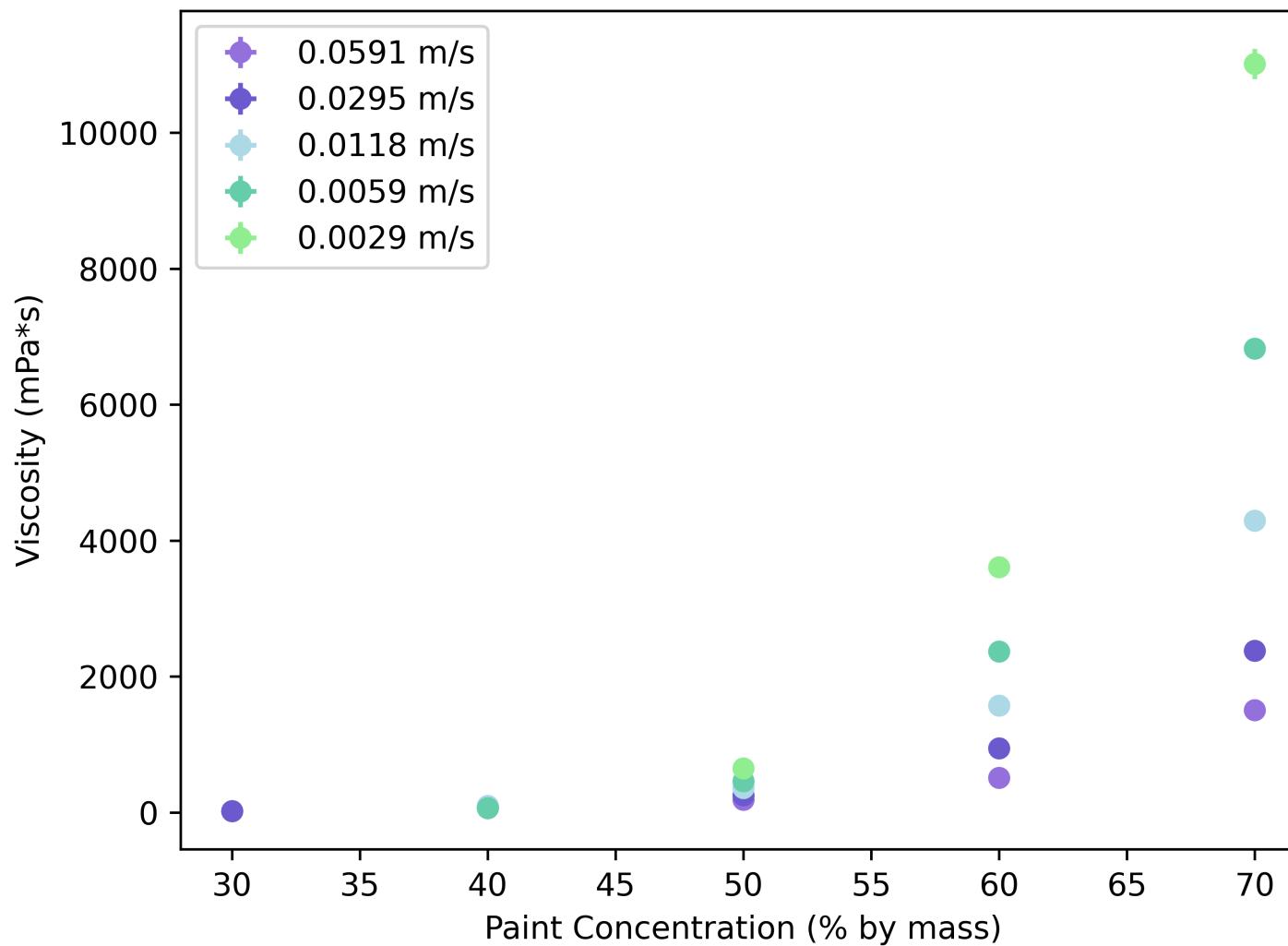
Experimental Setup

Droplet Dynamics

Branching Geometry

Insights

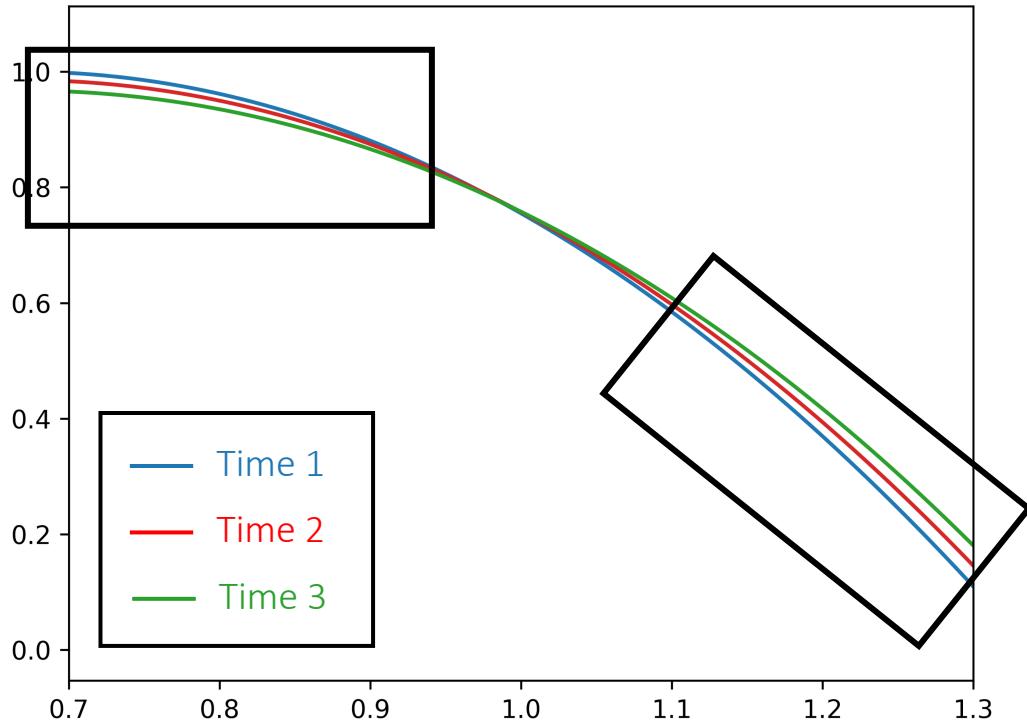
### Viscosity vs. Paint Concentration for Varying Velocities



# Numerical Solution

$$-\frac{\partial h}{\partial t} = \frac{1}{\rho} J - \frac{1}{r} \frac{\partial}{\partial r} \left( k \frac{(h^2 + 2Dh)}{\eta_p \eta_d} \frac{\partial}{\partial r} \left( -\frac{Jt}{h\rho} \right) \right)$$

where  $-\frac{Jt}{h\rho} = 0 \quad \text{if} \quad Jt \geq \phi_d h$

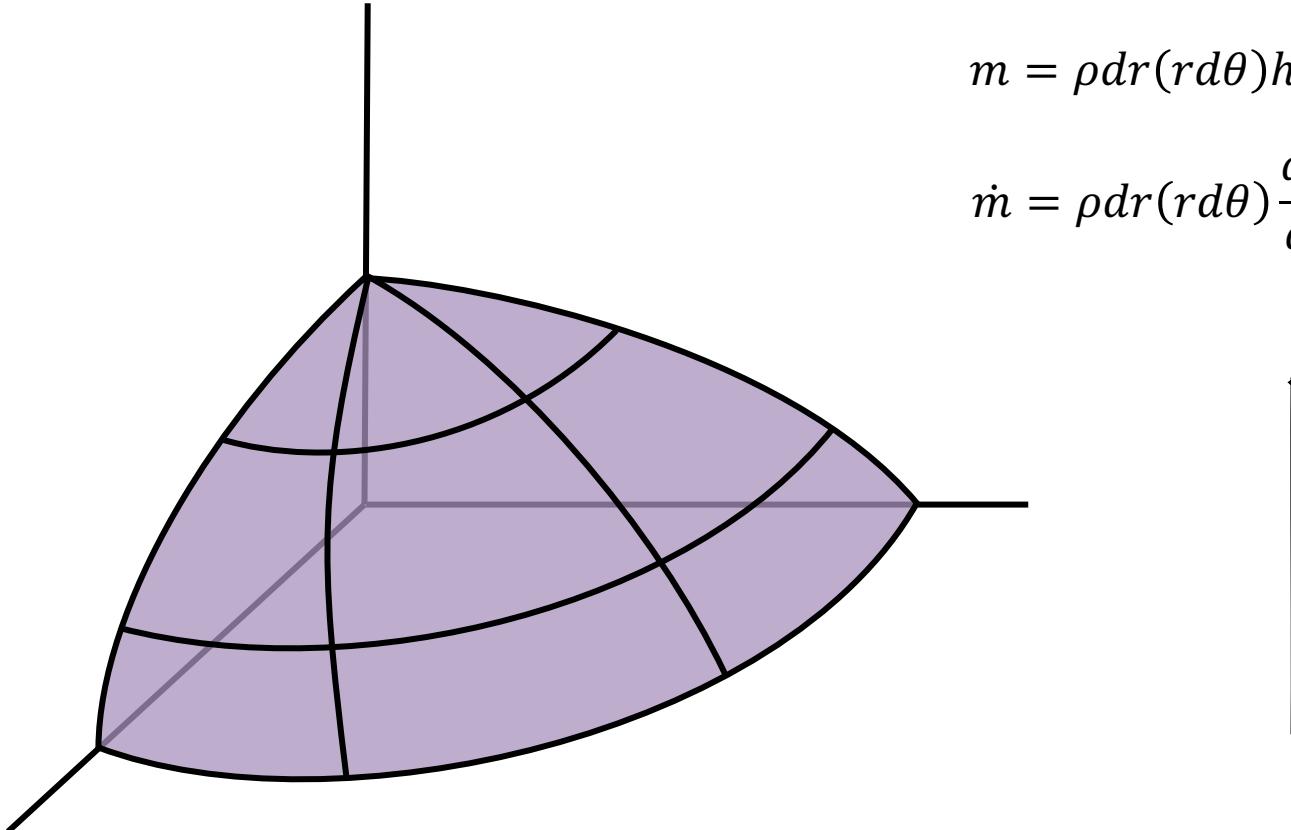


Height decrease at small  $r$

Height increase at large  $r$

✓ Flow as predicted by qualitative

# Continuity Derivation



$$m = \rho dr(r d\theta) h$$

$$\dot{m} = \rho dr(r d\theta) \frac{dh}{dt} = \sum \dot{m} = m_u + m_e$$

$$\rho dr(r d\theta) \frac{dh}{dt} = -J r d\theta dr + \rho h(r d\theta) \left( u_r \Big|_{r+dr} - u_r \Big|_r \right) + (\rho h dr) \left( u_\theta \Big|_{\theta+d\theta} - u_\theta \Big|_\theta \right)$$

Experimental Setup

Droplet Dynamics

Branching Geometry

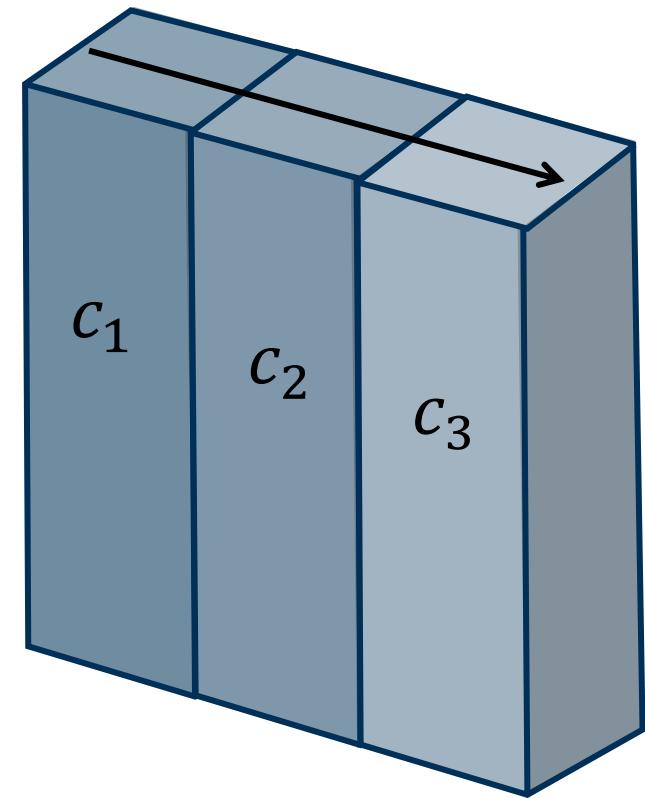
Insights

# Concentration Gradient

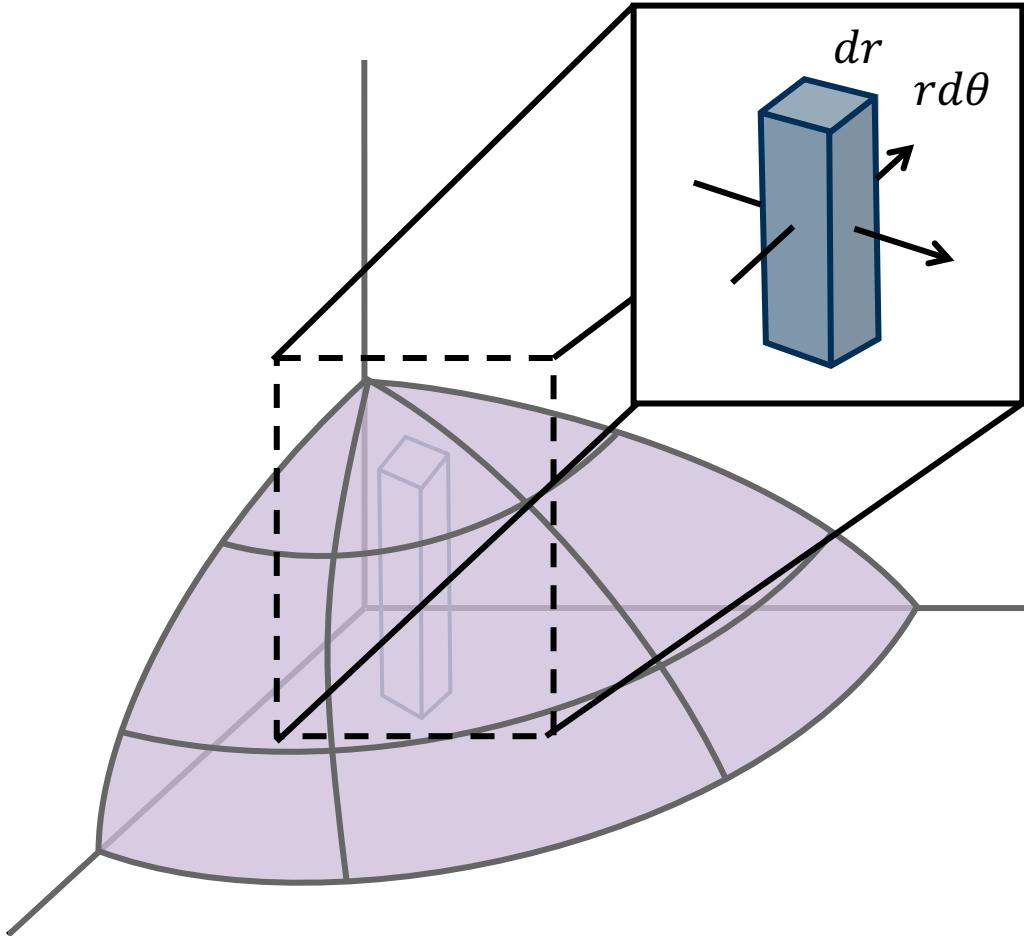
*Previous Assumption*

$$c(r, t) \approx \frac{\phi_d h - \frac{J}{\rho} t}{h}$$

*Concentration over time is nonlinear*



# Three-Dimensional Model



Experimental Setup

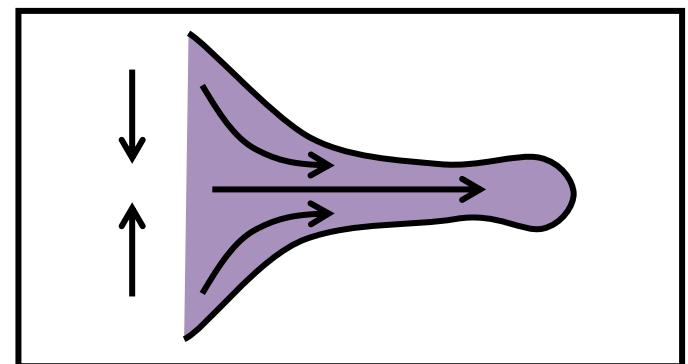
Droplet Dynamics

Branching Geometry

Insights

$$\frac{dh}{dt} = -\frac{1}{\rho} J + \frac{\partial Q_r}{\partial r} + \frac{1}{r} \frac{\partial Q_\theta}{\partial \theta}$$
$$\left\{ \begin{array}{l} u_r = \frac{h^2 + 2Dh}{\eta_p \eta_d} \frac{\partial \gamma}{\partial r} \\ u_\theta = \frac{h^2 + 2Dh}{\eta_p \eta_d} \frac{\partial \gamma}{\partial \theta} \end{array} \right.$$

*Accounting for azimuthal flow*



# Governing Equations

*Continuity for drop and alcohol with a restriction such that  $Q = 0$  if  $h_{alcohol} = 0$*

$$\left\{ \begin{array}{l} \frac{dh_{total}}{dt} = -\frac{1}{\rho}J + \frac{\partial Q_r}{\partial r} + \frac{1}{r}\frac{\partial Q_\theta}{\partial \theta} \\ \frac{dh_{alcohol}}{dt} = -\frac{1}{\rho}J + \frac{\partial Q_r}{\partial r} + \frac{1}{r}\frac{\partial Q_\theta}{\partial \theta} \end{array} \right.$$

*Flow derived from shearing of drop interfaces*

$$\left\{ \begin{array}{l} u_r = \frac{h^2 + 2Dh}{\eta_p \eta_d} \frac{\partial \gamma}{\partial r} \\ u_\theta = \frac{h^2 + 2Dh}{\eta_p \eta_d} \frac{\partial \gamma}{\partial \theta} \end{array} \right.$$

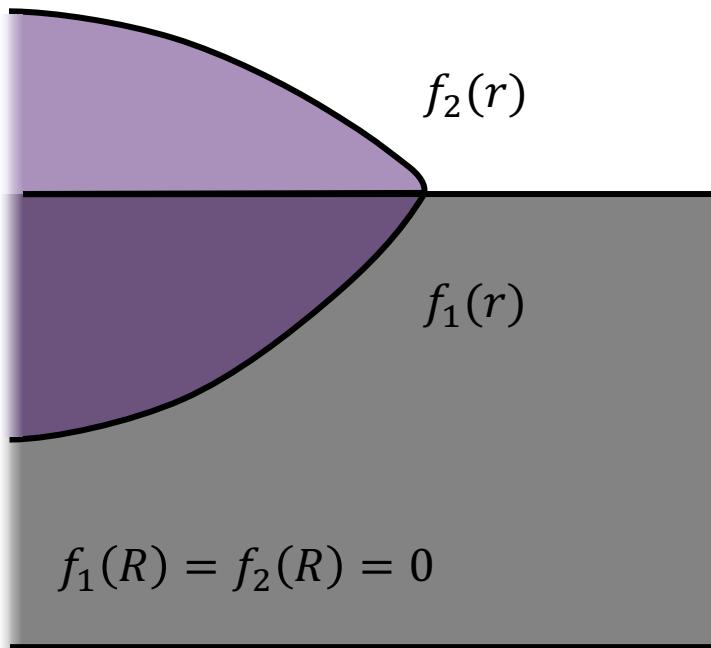
*Surface tension gradient is proportional to concentration*

$$\left\{ \begin{array}{l} \frac{\partial \gamma}{\partial r} = k \frac{h_{alcohol}}{h_{total}} \end{array} \right.$$

# Initial Conditions

*Minimize Total Energy of Drop:*

$$U_d + U_p + E_{ad} + E_{dp}$$



$$U_d = \rho_d g \pi \int_0^R r ([f_2(r)]^2 - [f_1(r)]^2) dr$$

$$U_p = \rho_p g \pi \int_0^R r ([f_1(r)]^2) dr$$

$$E_{ad} = \gamma_{ad} 2\pi \int_0^R r \sqrt{1 + \left(\frac{df_2}{dr}\right)^2} dr$$

$$E_{dp} = \gamma_{dp} 2\pi \int_0^R r \sqrt{1 + \left(\frac{df_1}{dr}\right)^2} dr$$

$$V = 2\pi \int_0^R r (f_2(r) - f_1(r)) dr$$

Find  $f_2(r) - f_1(r)$

# Continuity Derivation

$$\rho dr(r d\theta) \frac{dh}{dt} = -J r d\theta dr + \rho h(r d\theta) \left( u_r \Big|_{r+dr} - u_r \Big|_r \right) + (\rho h dr) \left( u_\theta \Big|_{\theta+d\theta} - u_\theta \Big|_\theta \right)$$

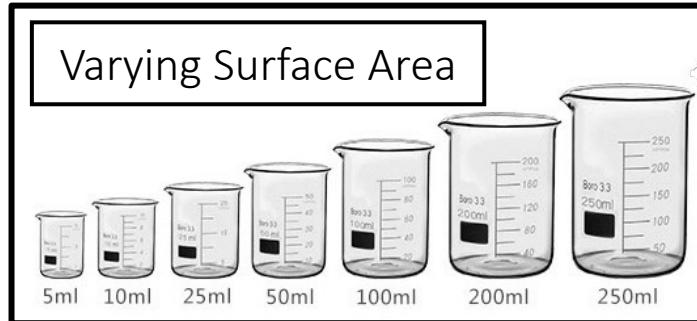
$$\frac{dh}{dt} = -\frac{J}{\rho} + h \frac{\partial u_r}{\partial r} + \frac{h}{r} \frac{\partial u_\theta}{\partial \theta}$$

$$\boxed{\frac{dh}{dt} = -\frac{J}{\rho} + \nabla \cdot (h \vec{u}) = -\frac{J}{\rho} + \nabla \cdot \vec{Q}}$$

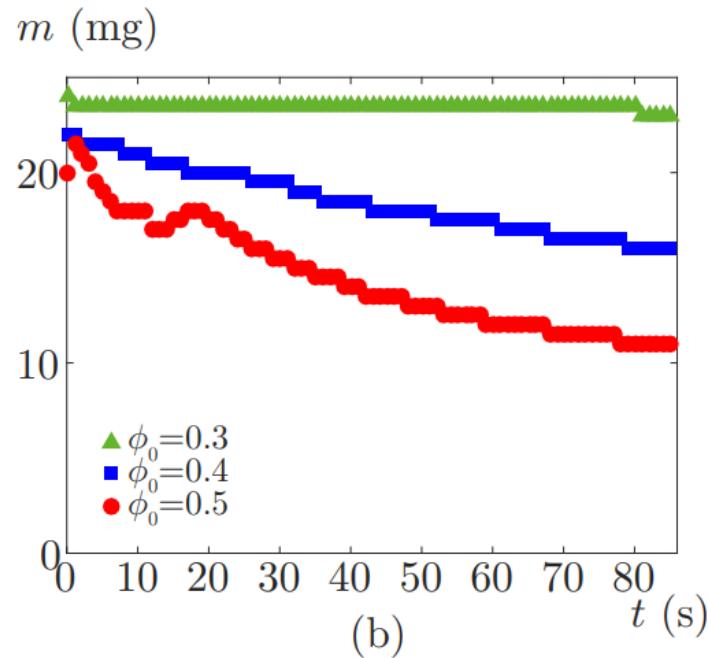
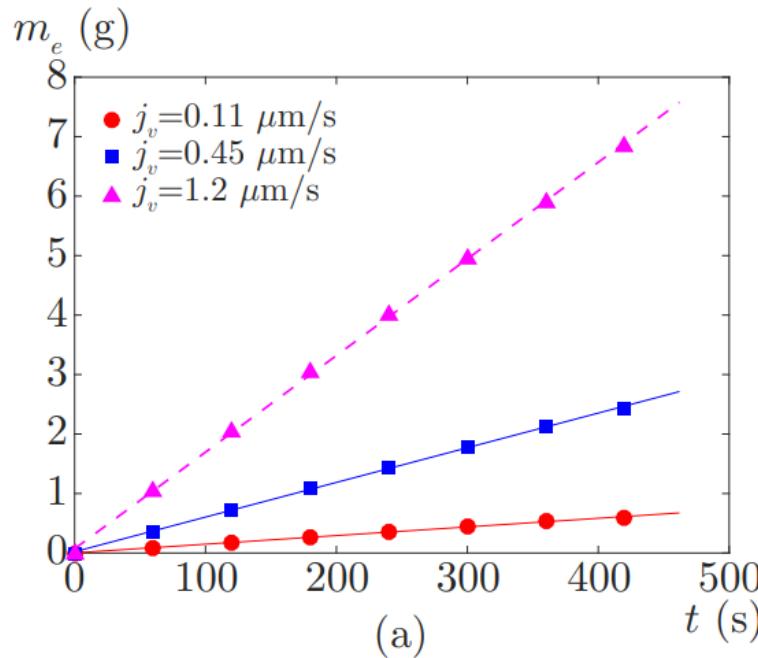
# Evaporative Flux



Change in volume is too small to observe, analytical balance required to analyze mass change over time



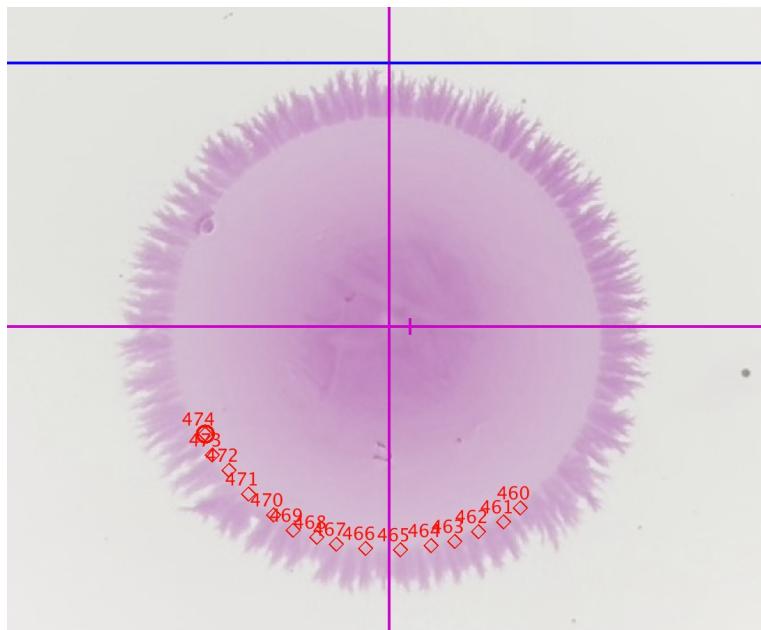
# Concentration Dependence On $J$



Evaporation rate of alcohol does not vary greatly with concentration

Keiser, L., et al. "Marangoni Bursting: Evaporation-Induced Emulsification of Binary Mixtures on a Liquid Layer." *Physical Review Letters*, vol. 118, no. 7, 17 Feb. 2017, 10.1103/physrevlett.118.074504.

# Radius Measurement: Tracker



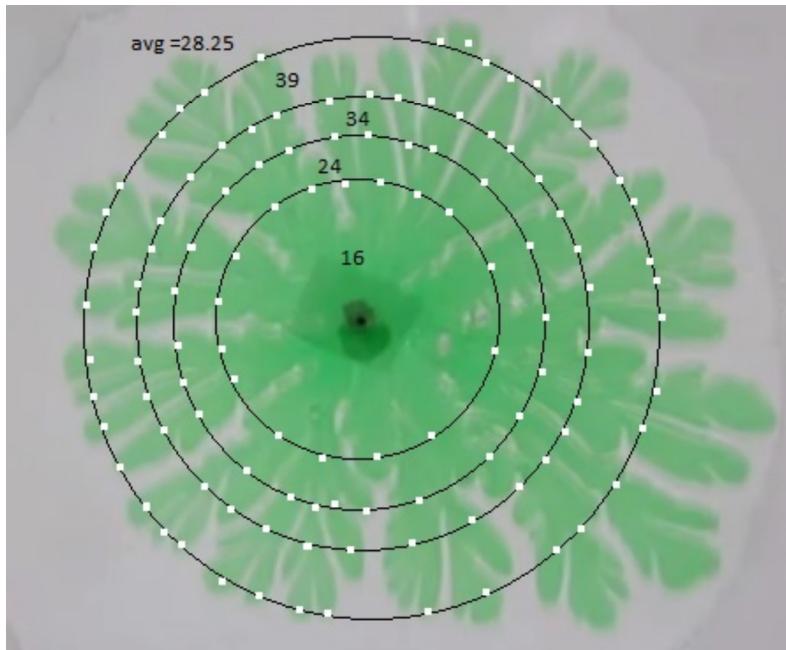
1 Measure *inner* and *outer* radius around the fractal

2 Determine the *average* of these values to find *radius* as defined (in the report)

3 Error bars correspond to the *standard error* of all tracked values

4 Scale the error by the derivative of the radius equation when scaling graph

# Instability Wavelength Measurement



Kondic, L., Lastra, Y., Moulton, L., & Anazco, R. (2014). (rep.). *The Hele-Shaw Cell/Saffman Taylor Instability: Theoretical and Experimental Comparison of Newtonian Fluids.*

1 Determine **number of branches** touching circles of different radii

2 Calculate **arc length** between each branch

3 Divide arc length by radius to determine **angular distance** between branches

4 Divide total # of fingers by # of circles to get **average number of fingers**

# Other Slides

# Literature

## Viscous Stress

Equating the stresses,

$$\frac{\Delta\gamma}{R} \approx \frac{u\eta_p}{H} \quad u \approx \frac{D\Delta\gamma}{R\eta_p}$$

$$\sigma_{drop} \approx \frac{\Delta\gamma}{R}$$

$$\sigma_{paint} \approx \frac{u\eta_p}{D}$$

$$T \approx \frac{R}{u} = \frac{\eta_p R^2}{D\Delta\gamma}$$

$$R \approx \left( \frac{(\phi_0 - \phi_f)\Delta\gamma V D}{(1 - \phi_f)J\eta_p} \right)^{\frac{1}{4}}$$

## Alcohol Volume

$$\phi_f V_f - \phi_0 V = JAT$$

$$\phi_f V_f - \phi_0 V \approx J(R)^2 T$$

## Ink Volume

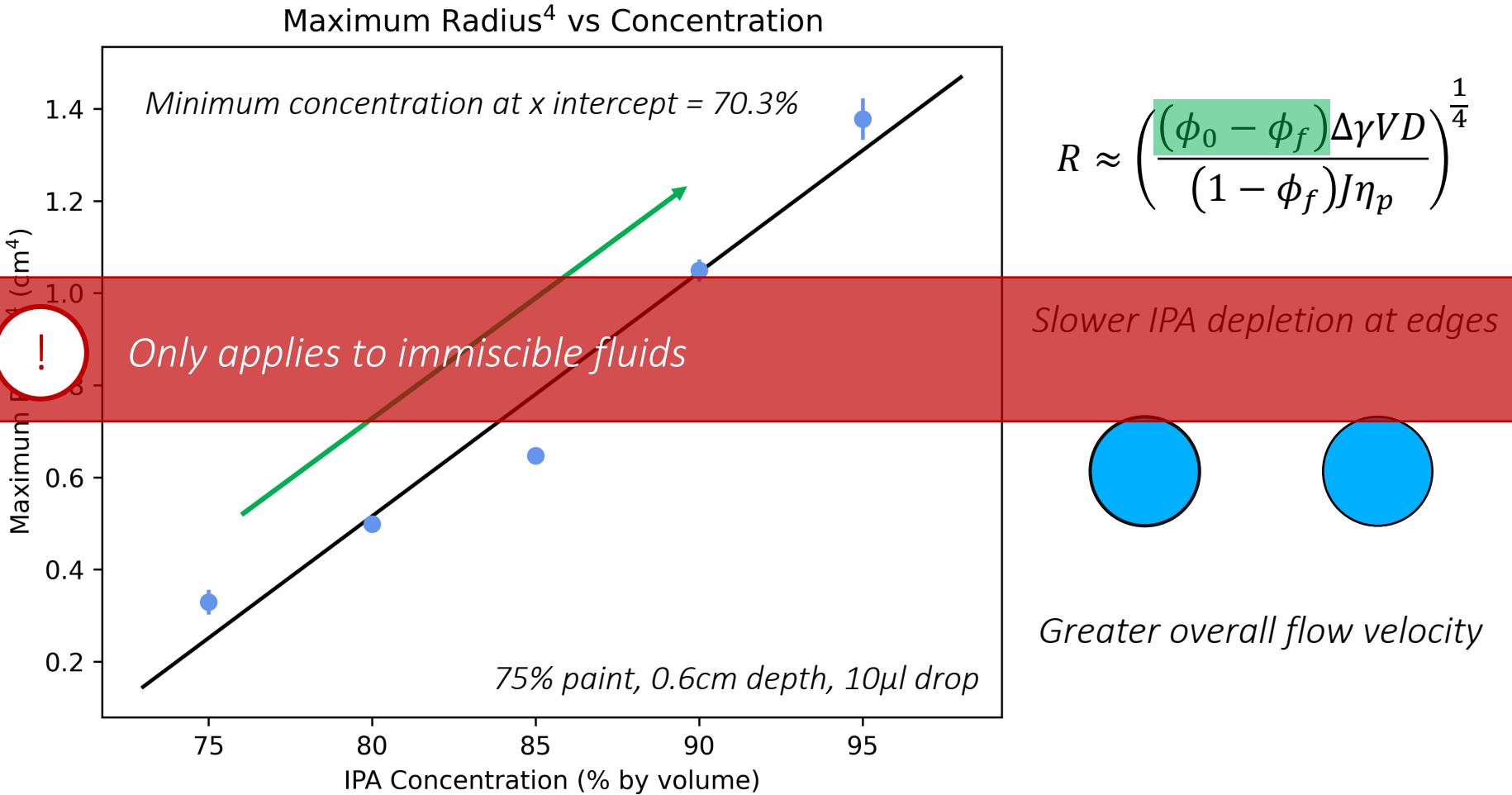
Conserved;

$$(1 - \phi_0)V = (1 - \phi_f)V_f$$

$$T \approx \left( \frac{(\phi_0 - \phi_f)V\eta_p}{(1 - \phi_f)J\Delta\gamma D} \right)^{\frac{1}{2}}$$

(Keiser et al., 2017)

# Literature Experimental Verification



# Miscible Viscous Fingering

*Velocity Induced Fractal Fingering*

$$\lambda_c = \pi b \sqrt{\frac{\sigma}{\Delta \mu V}}$$

$\sigma$  = surface tension

$\mu$  = viscosity of diluted acrylic paint

$b$  = space between parallel plates

$V$  = interfacial velocity imposed by an external pressure gradient

*Qualitatively, the Marangoni flow and alcohol evaporation induces the velocity in this problem, increasing the instability.*

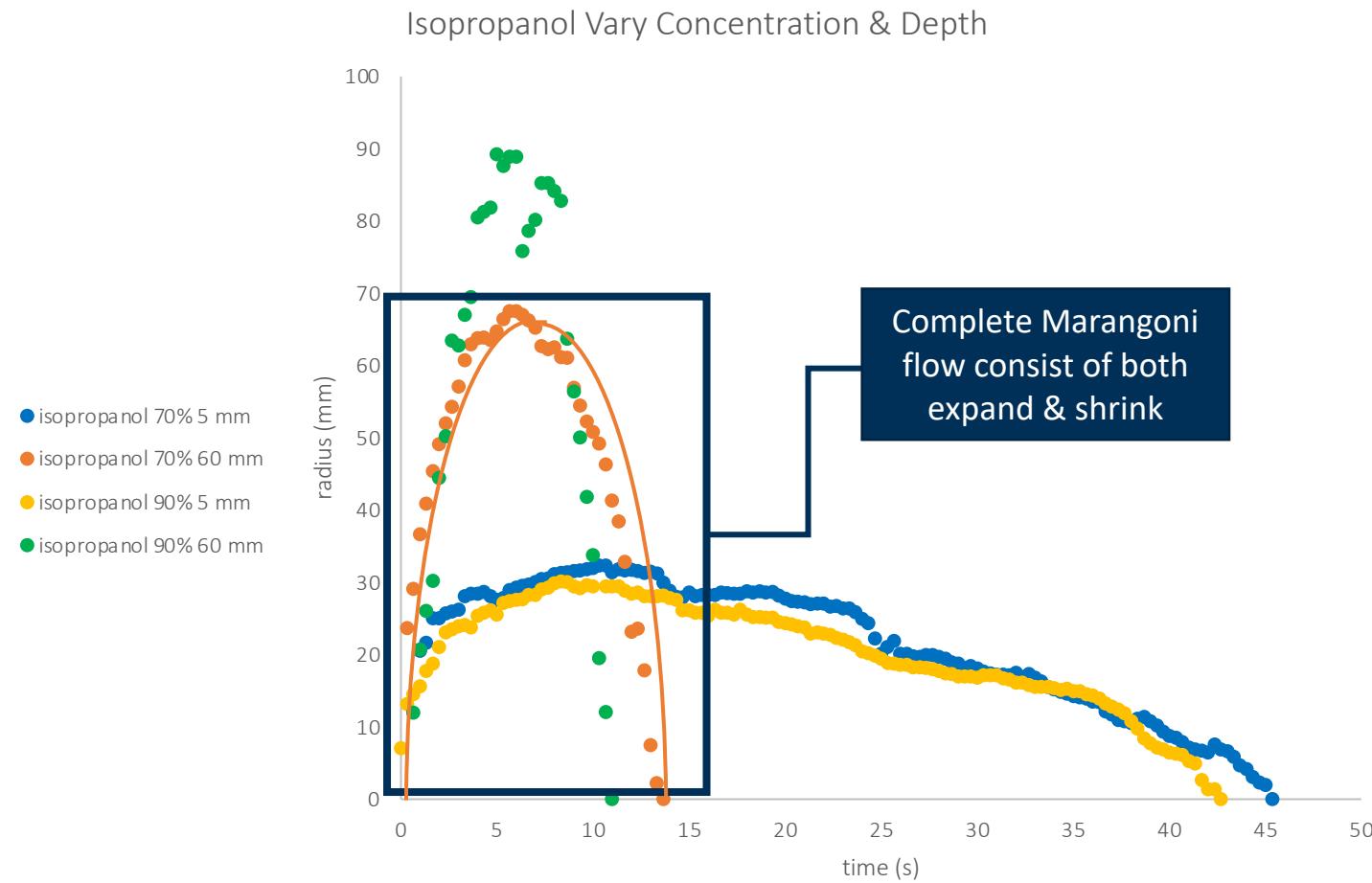
*Viscosity Ratio Function*

$$A \equiv \frac{\mu_a - \mu_d}{\mu_a + \mu_d}$$

$A$  = dimensionless parameter

As  $A$  approaches 1,  $\frac{Rf}{Ri}$  increases rapidly within a very small range of  $A$ .

# Experimental Verification



Three Stage Marangoni Flow

Finger Growth

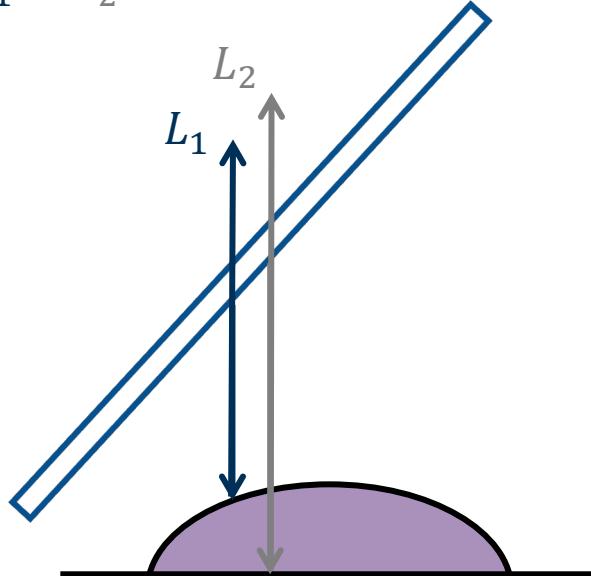
Geometry Analysis

# Non log log plot

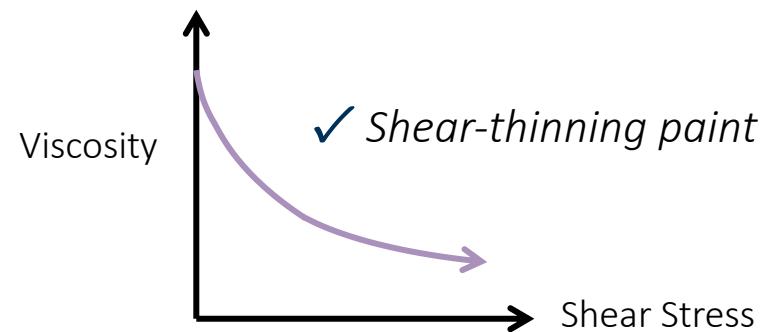
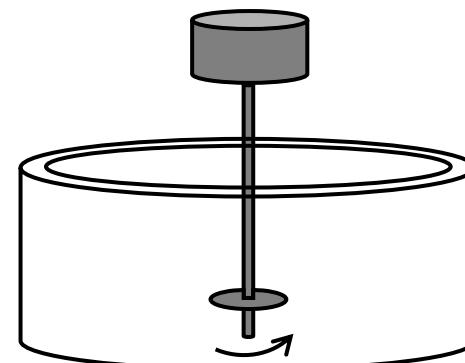
# Further Insights

## Interferometry

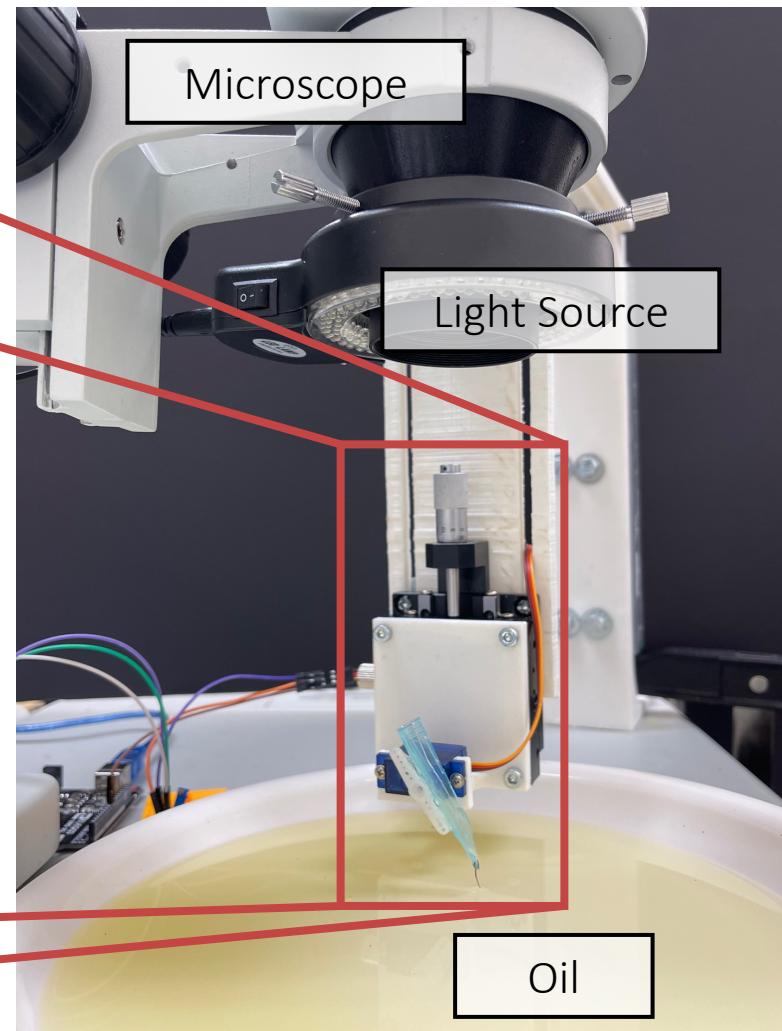
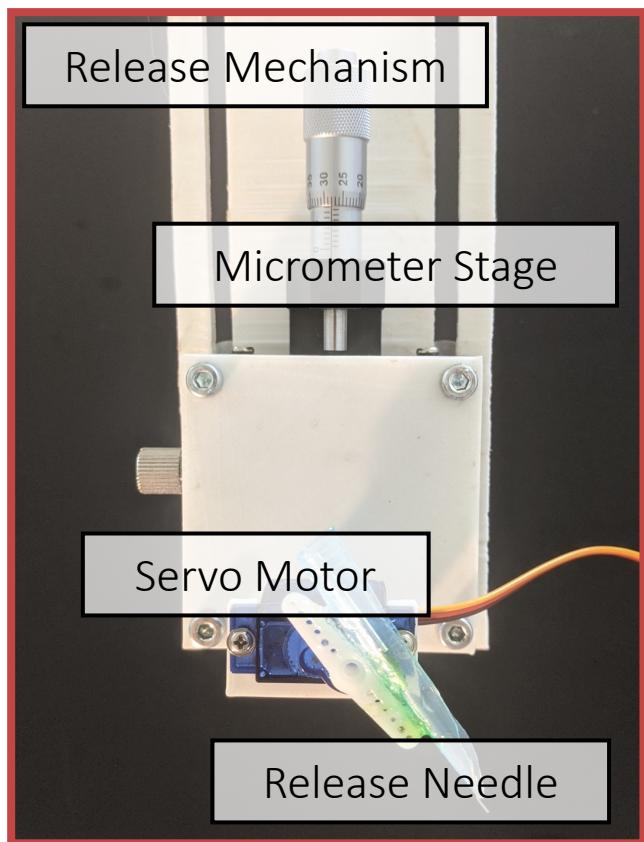
$$L_1 - L_2 = 2h$$



## Viscometry



# Microscopy Setup



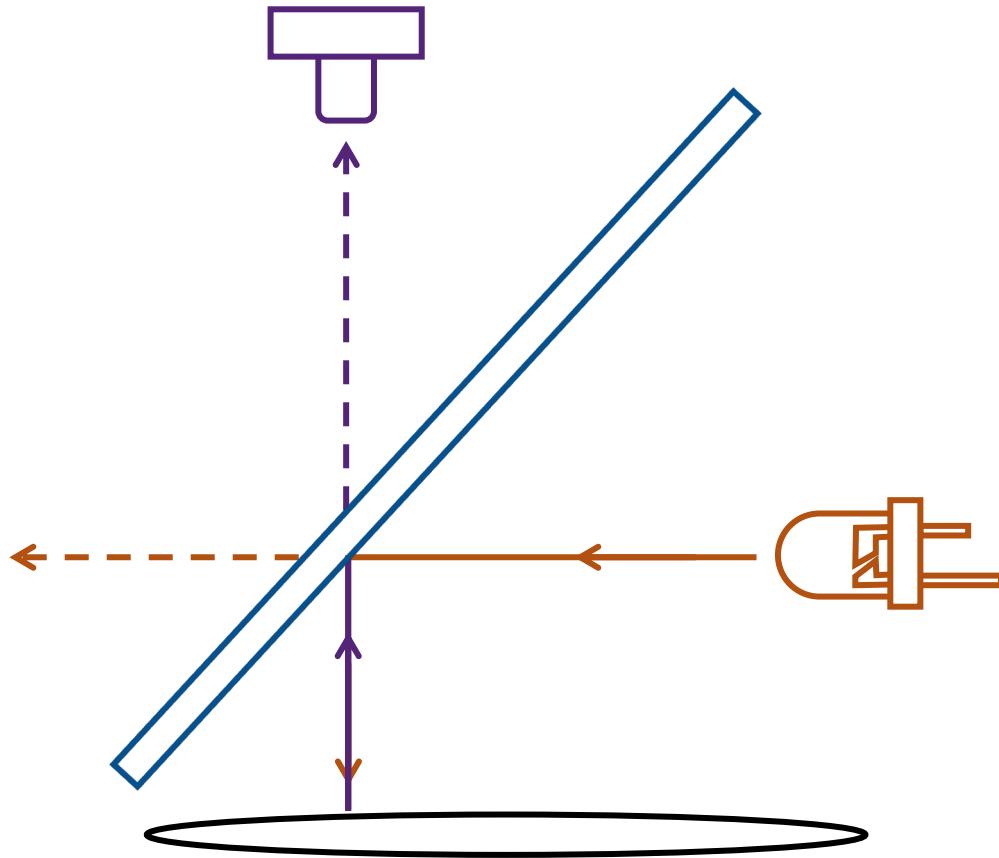
Experimental Setup

Droplet Dynamics

Branching Geometry

Insights

# Lighting

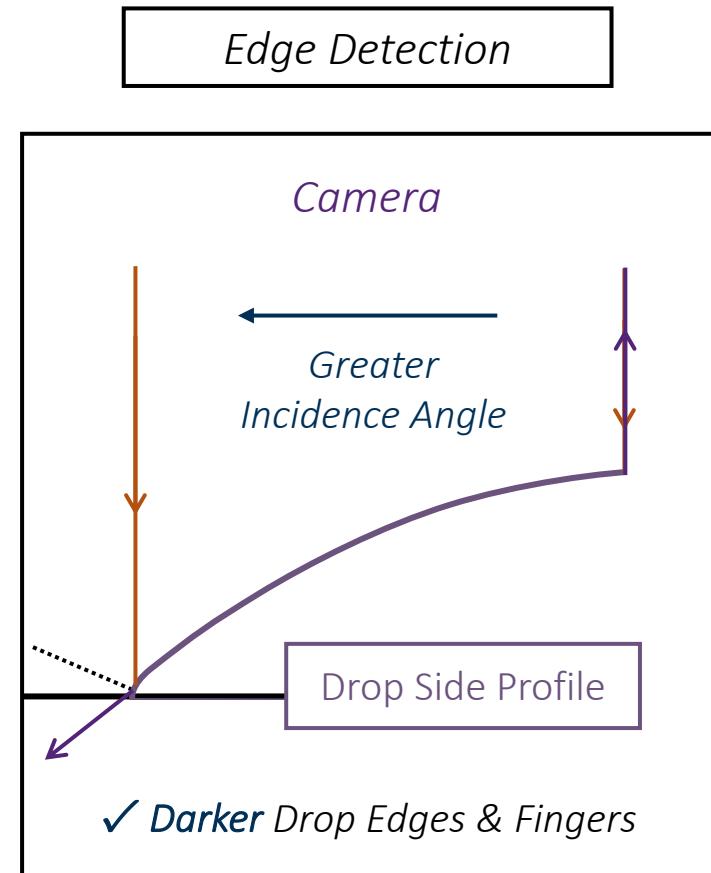


Experimental Setup

Droplet Dynamics

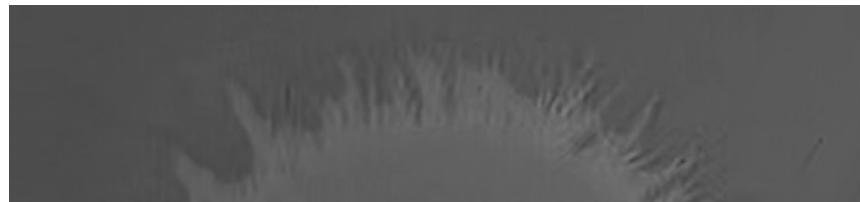
Branching Geometry

Insights



# Lighting

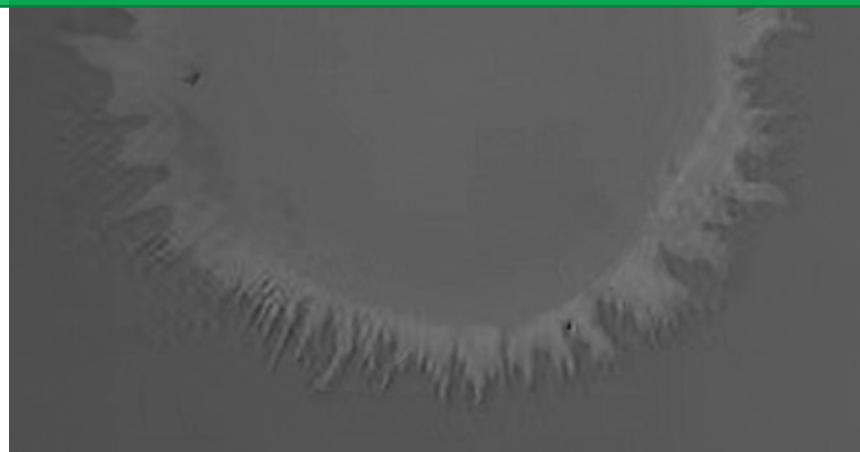
Without Lighting



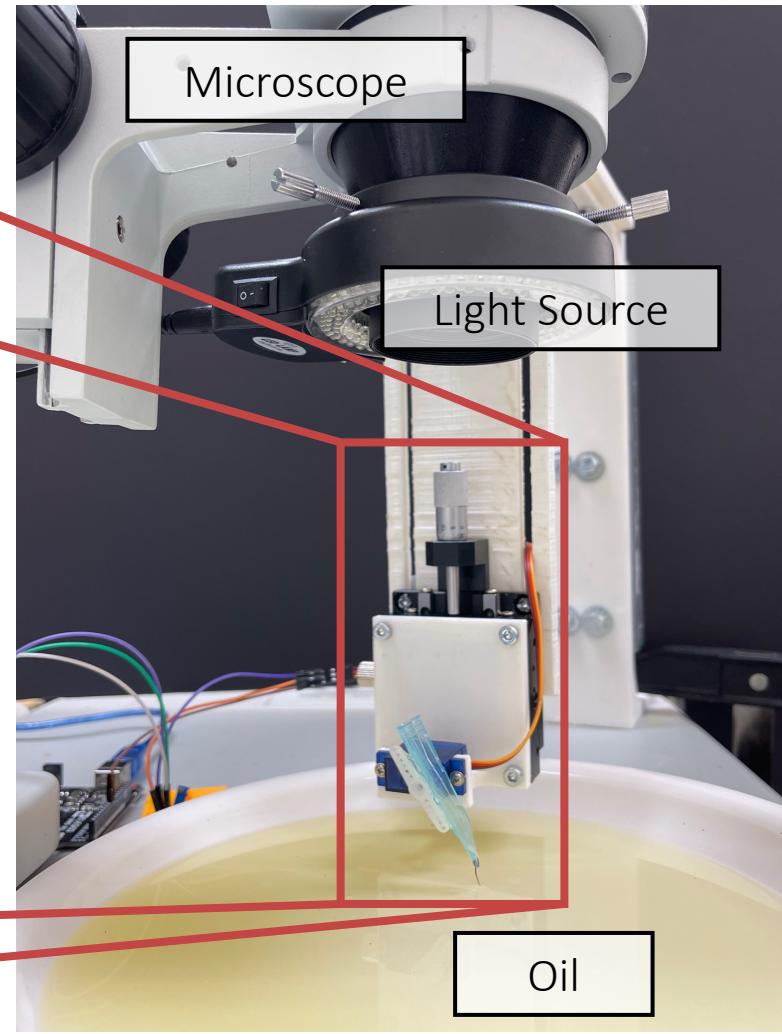
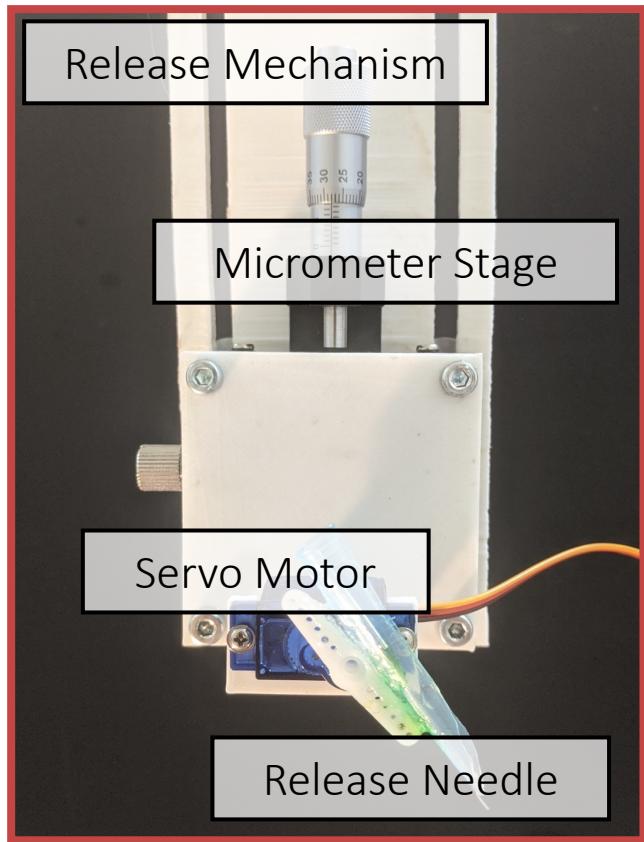
With Lighting



*Edge detection can distinguish between dark rings, fingers, and inner radius*



# Microscopy Setup



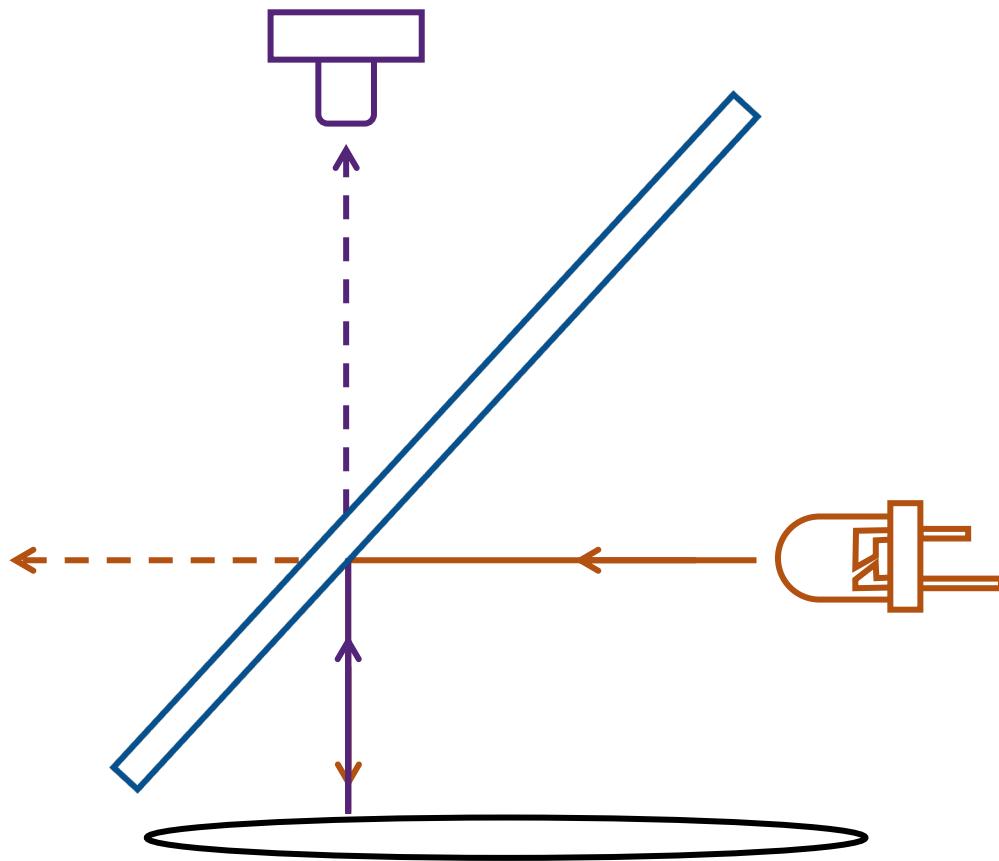
Experimental Setup

Droplet Dynamics

Branching Geometry

Insights

# Lighting

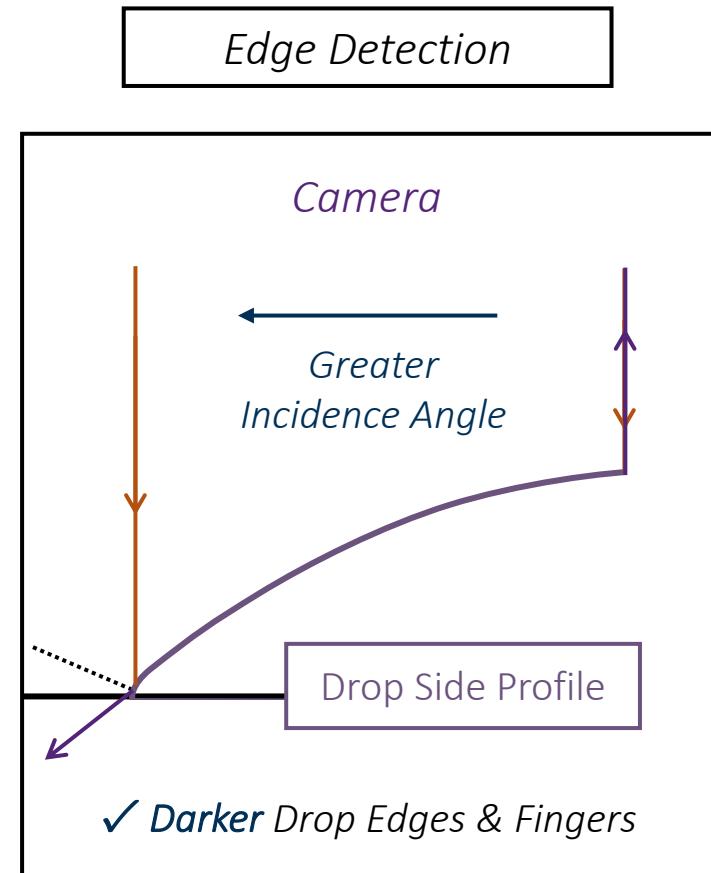


Experimental Setup

Droplet Dynamics

Branching Geometry

Insights



Edge Detection

Camera

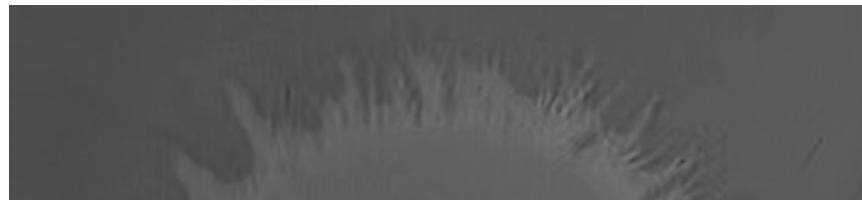
Greater  
Incidence Angle

Drop Side Profile

✓ Darker Drop Edges & Fingers

# Lighting

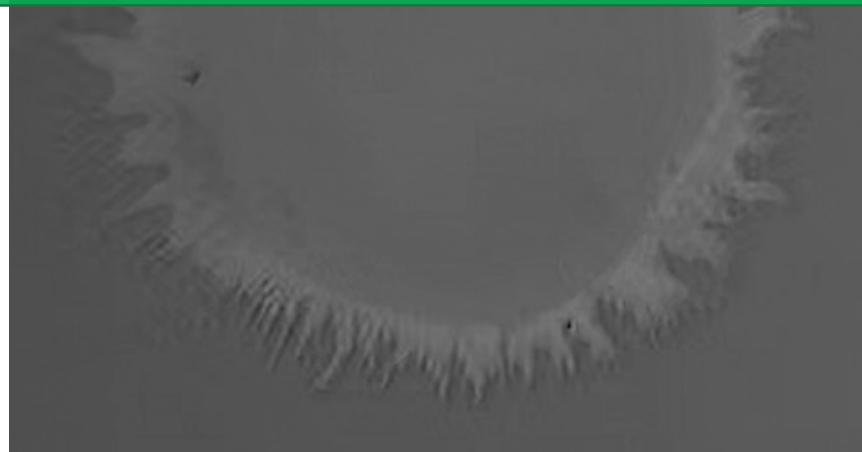
Without Lighting



With Lighting



*Edge detection can distinguish between dark rings, fingers, and inner radius*



[2]

# Marangoni and VS Stress

$$\text{Marangoni Stress} \sim \frac{\partial \gamma}{\partial r}$$



$$\text{Viscous Shear Stress} \sim \frac{\mu u}{l}$$



Setting them both as equal:

$$\frac{\partial \gamma}{\partial r} \approx \mu(u/l) \rightarrow u \approx \frac{r^{\frac{1}{3}}}{(\mu p)^{\frac{1}{3}}} \left( \frac{\partial \gamma}{\partial r} \right)$$

$$u \approx \frac{(\gamma_+ - \gamma_-)^{2/3}}{(\mu \rho r)^{1/3}} \quad u = \frac{\partial r}{\partial t} = \frac{(\gamma_+ - \gamma_-)^{2/3}}{(\mu \rho r)^{1/3}}$$

$r$  = radius

$u$  = expansion speed at  
 $r$

$l$  = depth into acrylic

$\gamma_-$  = surface tension of  
alcohol-ink

$\gamma_+$  = surface tension of  
acrylic

$\mu$  = viscosity of acrylic

$\rho$  = density of diluted  
acrylic

For  $\gamma_-(t)$ , surface tension of alcohol-ink,  
dependent on alcohol concentration varying through time

# Marangoni Flow



$$\text{Marangoni Stress} \sim \frac{\partial \gamma}{\partial r}$$

$$\text{Viscous Shear Stress} \sim \frac{\mu u}{l}$$

Equating Marangoni stress with opposing viscous stresses,

$$\frac{\partial \gamma}{\partial r} \approx \mu(u/l) \rightarrow u \approx \frac{r^{\frac{1}{3}}}{(\mu p)^{\frac{1}{3}}} \left( \frac{\partial \gamma}{\partial r} \right)$$

$$u \approx \frac{(\gamma_+ - \gamma_-)^{2/3}}{(\mu \rho r)^{1/3}} \quad u = \frac{\partial r}{\partial t} = \frac{(\gamma_+ - \gamma_-)^{2/3}}{(\mu \rho r)^{1/3}}$$

$r$  = radius

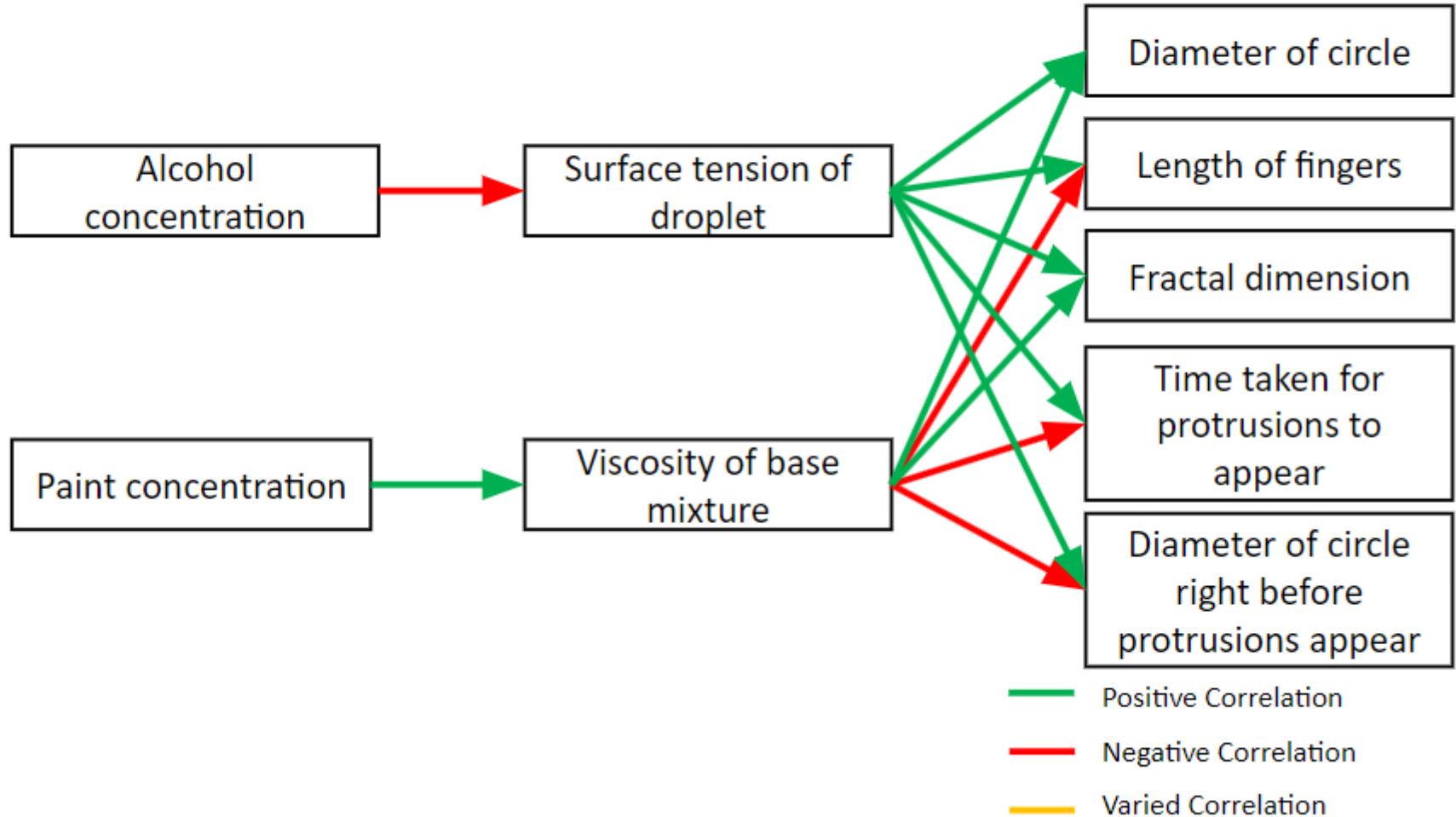
$u$  = expansion speed at  
 $r$

$l$  = depth into acrylic

$\gamma_-$  = surface tension of  
alcohol-ink

$\gamma_+$  = surface tension of  
acrylic

# Key Parameters



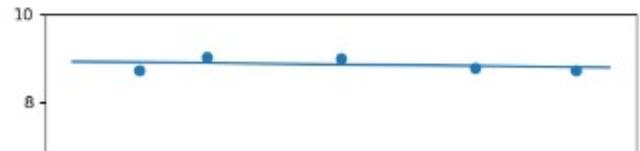
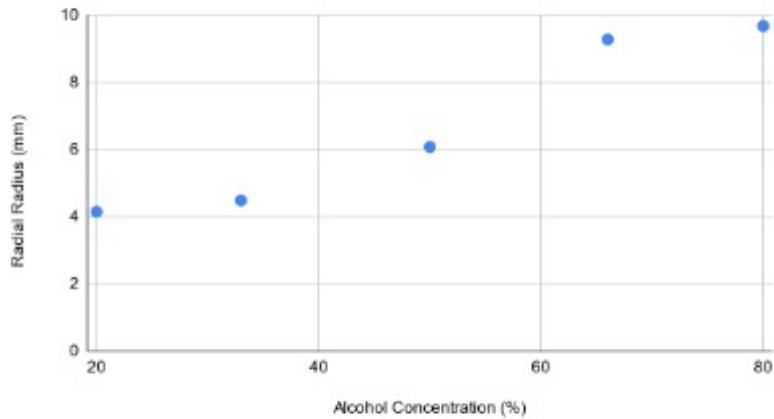
# Fingering Results

Alcohol Concentration  $\phi$  has obvious positive correlation with radius, supported by Marangoni Flow Theory

$$r = \left( \frac{4}{3} \int \frac{(\gamma_+ - \gamma_-)^{2/3}}{(\mu\rho)^{1/3}} dt \right)^{3/4}$$

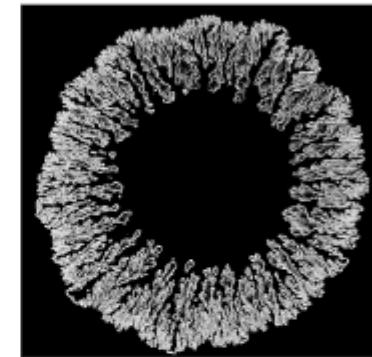
Concentration  $\Phi$  increase implies  $\gamma_-$

Maximum Radius vs Alcohol Concentration



# Fractal Dimension

The Minkowski Dimension, also called the box-counting Dimension, is a quantitative index for the (loosely speaking) *roughness* or *complexity* of a fractal. The definition of a fractal is that its detail stays constant at different zoom levels. If you zoom in a lot, it will still look as rough and complex as it originally did.



$$\dim_{\text{box}}(S) = \lim_{\epsilon \rightarrow 0} \frac{\log N(\epsilon)}{\log 1/\epsilon}$$

The box-counting dimension qualitatively measures this by seeing how much grid boxes on the plane the shape coincides with at a certain degree of zoom level, divided by the zoom level. The dimension is the limit of this ratio as the grid boxes become infinitely finer. The upper and lower box dimensions, counting the maximum and minimum box count at local ranges of zoom, are

# Fractal Dimension

## Numerical Dimension Approximation

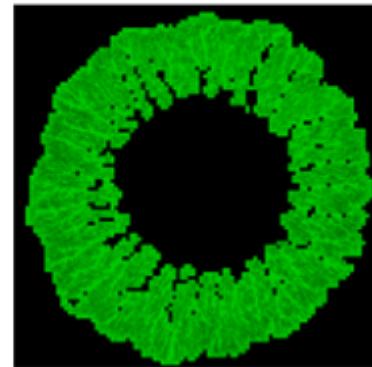
As our images don't have infinite resolution,  
We computationally approximate the box dimension  
And verify the constant roughness by counting the  
grid boxes taken up at different box lengths (or zoom  
levels).

We plot their logarithms. The resulting  
graph being linear means that the shape  
instead has fractal properties, and the  
(negated) slope is thus the dimension.

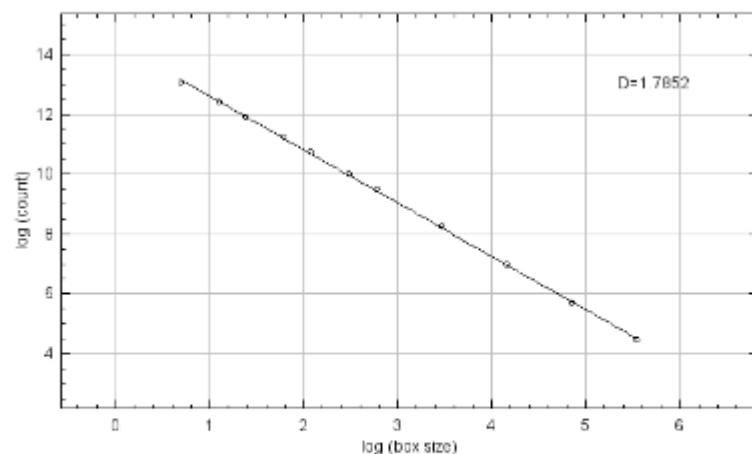
$$\frac{\log N(l)}{\log l_0/l}$$

$l_0$  = Base Box Length  
 $l$  = Box Length  
 $N(l)$  = Minimal Box Count

Green boxes show grid boxes  
being taken up, here we show  
a relatively small box length



$\log(\text{box\_count})$  vs  $\log(\text{box\_size})$  of the above  
image, note the linearity of the graph and the  
slope (negating slope gives the dimension)

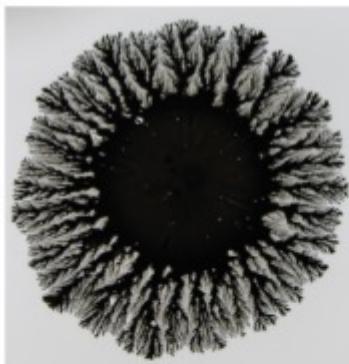


Plot Shows Linear Relation, Verified Fractality Of Fingering

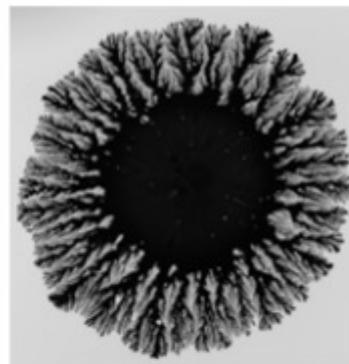
# Processing Pipeline



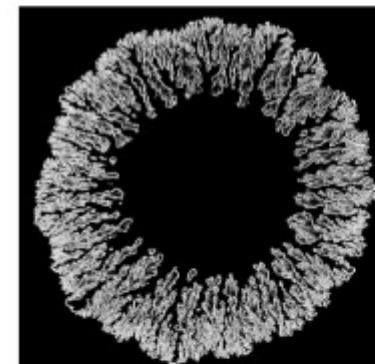
Julia Image Processing Pipeline. The original image is fed in, and grayscale edge detection is performed. Dimension analysis is run on the resulting contour.



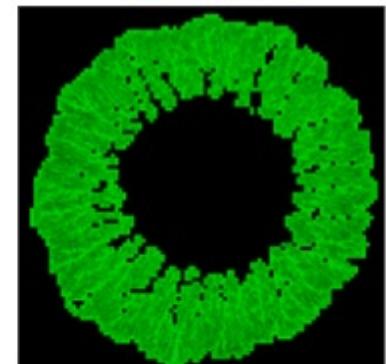
Original



Grayscale



Edge Detection



Box Counting