

博士論文

Doctoral Dissertation

Time Dependent Charge-Parity Violation in $B^0 \rightarrow K_s^0 K_s^0 K_s^0$ in Belle

II early operation

(Belle II 初期データを使った $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ 崩壊の時間に依存する
荷電・パリティ非保存の研究)

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**Time Dependent Charge-Parity Violation in $B^0 \rightarrow K_s^0 K_s^0 K_s^0$ in
Belle II early operation**

by

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Abstract

Belle II experiment is a next-generation super B-factory experiment. The targeted instantaneous luminosity is $8 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$ and the expected integrated luminosity is 50 ab^{-1} by 2030 with the majority of data collected at the $\Upsilon(4S)$ resonance using SuperKEKB accelerator.

The thesis is based on the time-dependent CP violation study of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ decay to precisely measure the CP parameters \mathcal{S} and \mathcal{A} in penguin-dominated $b \rightarrow s$ transition, which is sensitive to New Physics effects. Such a precise measurement mainly depends on determination of the distance between two vertices of two neutral B mesons. The blind analysis and fit by a unbinned maximum likelihood method are performed using about 62.8 fb^{-1} recorded experiment data from Belle II detector 2019 and 2020 (spring and summer) operation. The measurement results: $\mathcal{S} = -\sin(2\phi_1) = -0.82 \pm 0.85 \text{ (stat)} \pm 0.07 \text{ (syst)}$ and $\mathcal{A} = -0.21 \pm 0.28 \text{ (stat)} \pm 0.06 \text{ (syst)}$ are obtained.

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Chapter 1

Introduction

1.1 The Standard Model

The Standard Model (SM) was built in the late 70th of 20th century to describe the matter compositions and interactions using a group of fundamental particles - fermions and bosons. In the Standard Model, there are three generations of quarks and leptons, along with their anti-particles, which are all fermions. On the other hand, the bosons in the Standard Model consist of gluons, photons, W^\pm and Z^0 bosons that are all gauge bosons and one Higgs boson that is a scalar boson. This group of particles are summarized in Figure 1-1. The Standard Model depicts the interactions between elementary particles as the exchange of the bosons. The strong interaction requires the exchange of gluons. Photons, W^\pm and Z^0 bosons carry the electromagnetic force and weak force, which are unified as electroweak interaction in the Standard Model. Higgs boson is responsible for the generation of masses for the gauge bosons through electroweak symmetry breaking[1]. The Standard Model has been proved to be an excellent theoretical model that can be used to explain many experimental observations, but sadly not all of them. For instance, neutrino mass is expected to be zero in the Standard Model but the flavor oscillation indicates non-zero mass of neutrinos. The observation of Charge-Parity (CP) asymmetry in universe presented by the absence of antimatter can not be fully explained by the CP violation sources within the Standard Model. These experimental observations

require further researches beyond the Standard Model, which is called New Physics (NP) studies.

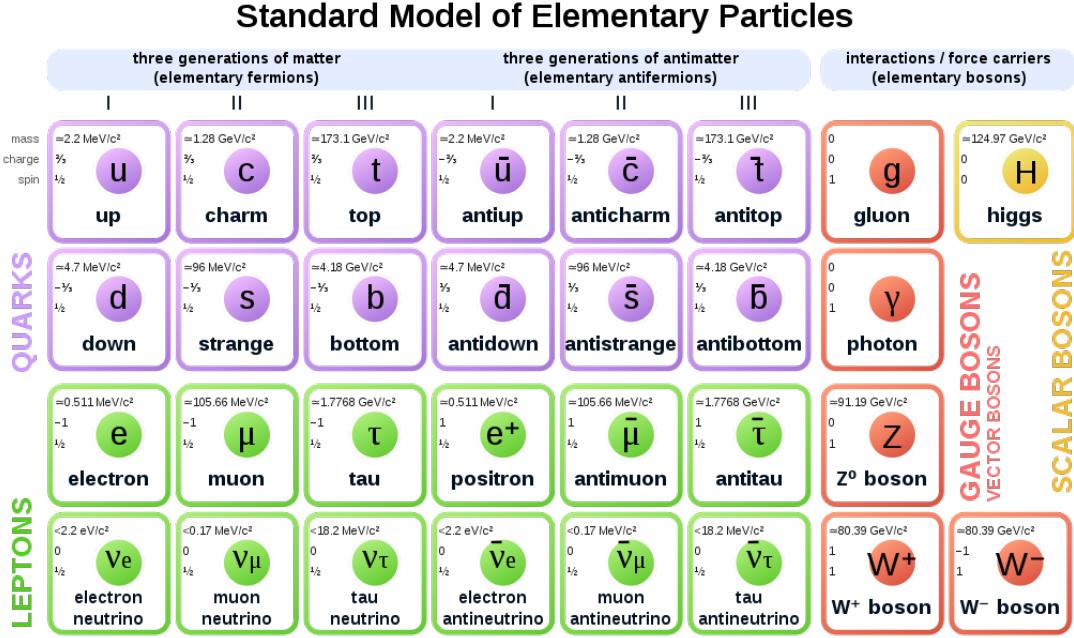


Figure 1-1: Elementary particles in the Standard Model.[2]

1.2 Symmetry Violation

Symmetry violation has been one of the focuses in NP studies due to the internal link between symmetries and conservation laws, which makes it a good probe for possible NP theories beyond the SM. When a known symmetry is found to be broken, it usually leads to the discovery of a new theory.

There are three types of discrete symmetric operations which play important roles in particle physics. Charge-conjugation C is the operation that turns particle to its anti-particle. Parity transformation P is the one that puts a negative sign before all the spatial related vector such as $\vec{r} \rightarrow -\vec{r}$. The time-reversing operation T is to reversely proceed a physical process backward time. Physicists were convinced that each of these three symmetric operations makes no change to any physics system. However, in 1950s, Lee and Yang [3] first questioned that parity symmetry might be

broken in weak interactions. They offered a few possible ways to test it and then by Wu [4], an observation on the β decay of ^{60}Co was presented that the electrons emitted from ^{60}Co decay prefer the direction of nuclear spin that can be controlled by the external magnetic field. The violation of P symmetry was discovered by this clear evidence.

The first evidence of CP violation was discovered in neutral K^0 system by Cronin and Fitch's experiment[5]. The neutral K^0 mesons can be observed as two states that have significantly different lifetime (called as " K_S^0 " and " K_L^0 " for short and long lifetime particles) with opposite CP . The experiment measured the decay products at 57 foot of a neutral K^0 beamline assuming all the particle at the end of the beam should be long lifetime K_L^0 , nearly no K_S^0 . But 0.002% of K_L^0 were found to decay into $\pi^+\pi^-$ which is the main decay process of K_S^0 . ($CP = 1$ in $\pi^+\pi^-$ final states, while K_L has $CP = -1$). Given that the expected distance to have 0.002% of K_S^0 at about speed of light is no more than 1 meter in the beamline, such a deviation at 57 foot is an obvious evidence that $K_L^0 \rightarrow \pi^+\pi^-$ exists and therefore CP symmetry is violated in the neutral K^0 system.

In 1973, Kobayashi and Maskawa introduced a quark mixing matrix called CKM matrix for three or more generations of quarks before the discovery of the third generation of the quark family[6]. The theory naturally explained an irreducible complex phase in CKM matrix and it accounts for the origin of CP asymmetries of weak interactions in the Standard Model. The experimental evidence of CP violation in B meson system was observed in 2001 by Belle and Babar experiments[7][8]. They measured the time-dependent decay time difference of B and \bar{B} in the decay of $B \rightarrow J/\psi K_S^0$. This channel provided a good clearness in theoretical prediction and has relatively large branching fraction, thus it's called the "golden mode"[9]. In 2008, Kobayashi and Maskawa were rewarded the Nobel Prize to highly value their contribution to CP violation mechanism in the SM, to which Belle experiment contributes greatly. Later in 2010, the upgrade of Belle, Belle II and the upgrade of KEK accelerator, SuperKEKB, were approved to further push the understanding of CP violation along with other topics in New Physics researches.

1.3 CKM mechanism

$$\Phi = \begin{pmatrix} \phi^+ \\ \nu + \frac{H+i\chi}{\sqrt{2}} \end{pmatrix} \quad (1.1)$$

Equation 1.1 is the Higgs potential doublets in the SM, where the value of H is 174 GeV as the expected Higgs potential for vacuum[10]. The ϕ and χ are the psuedo-Goldstone fields which are appearing when introducing Higgs field ϕ without breaking the gauge symmetry. The Lagrangian for Yukawa interaction of the quark fields[11] can be presented by Equation 1.2.

$$\mathcal{L}_{Yuk}^q = -Q^\dagger Y^d \Phi d'_R - Q^\dagger Y^u \epsilon \Phi^* u'_R + h.c. \quad (1.2)$$

where the primed fields stand for the weak eigenstates of quarks. The ϵ is a 2×2 matrix and Q^\dagger is the left-handed doublets that stand for weak eigenstates of up and down types quarks, see Equation 1.3 and 1.4.

$$\epsilon = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \quad (1.3)$$

$$Q = \begin{pmatrix} u' & d' \\ c' & s' \\ t' & b' \end{pmatrix}_L \quad (1.4)$$

Yukawa matrix is an arbitrary 3×3 complex matrix $Y^{u,d}$ which gives the rise of up and down type massive quark field $M^{u,d} = Y^{u,d} \nu$ according to Equation 1.2. The representation of the quark fields using weak eigenstates can be transformed to mass eigenstates by Equation 1.5 and 1.6.

$$S_{L,R}^u \begin{pmatrix} u' \\ c' \\ t' \end{pmatrix}_{L,R} = \begin{pmatrix} u \\ c \\ t \end{pmatrix}_{L,R} \quad (1.5)$$

$$S_{L,R}^d \begin{pmatrix} d' \\ s' \\ b' \end{pmatrix}_{L,R} = \begin{pmatrix} d \\ s \\ b \end{pmatrix}_{L,R} \quad (1.6)$$

In Equation 1.5 and 1.6, $S_{L,R}^{u,d}$ are all unitary matrices since they are generated by the normalized eigenstate states of Yukawa matrix. The mass item in the Lagrangian can be presented as Equation 1.7

$$\mathcal{L}_m = - \sum_{q=u,c,t,d,s,b} M_q q^\dagger q \quad (1.7)$$

where the $q = (q_L + q_R)$ is four-component Dirac field, and $q_L^\dagger q_L = q_R^\dagger q_R = 0$. As a result of diagonalizing $Y^{u,d}$, the charged current W^\pm interactions couple to the physical quarks and the Lagrangian is written as Equation 1.8, where $V_{CKM} \equiv S_L^u S_L^{d\dagger}$.

$$\mathcal{L}_W^q = \frac{g}{\sqrt{2}} \left[\begin{pmatrix} \bar{u} & \bar{c} & \bar{t} \end{pmatrix}_L \gamma^\mu W_\mu^+ V_{CKM} \begin{pmatrix} d \\ s \\ b \end{pmatrix}_L + \begin{pmatrix} \bar{d} & \bar{s} & \bar{b} \end{pmatrix}_L \gamma^\mu W_\mu^- V_{CKM}^\dagger \begin{pmatrix} u \\ c \\ t \end{pmatrix}_L \right] \quad (1.8)$$

The Lagrangian hereby clearly declares the transition of different charged quarks through the coupling of charged current W^\pm , where such a coupling only applies for the left-handed quarks. For example, a left-handed charm quark only transits to left-handed strange quark by a W boson. By only applying C or P conjugation, the Lagrangian is not invariant, indicating the non-conservation of C or P individually. However, if the CP conjugation is applied, the Equation 1.8 transits as Equation 1.9 shows.

$$\begin{pmatrix} \bar{u} & \bar{c} & \bar{t} \end{pmatrix}_L \gamma^\mu W_\mu^+ V_{CKM} \begin{pmatrix} d \\ s \\ b \end{pmatrix}_L \Leftrightarrow \begin{pmatrix} u & c & t \end{pmatrix}_L \gamma^\mu W_\mu^- V_{CKM} \begin{pmatrix} \bar{d} \\ \bar{s} \\ \bar{b} \end{pmatrix}_L \quad (1.9)$$

Comparing Equation 1.9 and 1.8, the CP symmetry requires the invariance before and after CP conjugation, meaning that Equation 1.10 is expected.

$$u_L^i V_{ij} \bar{d}_L^j \gamma^\mu W_\mu^- = u_L^n V_{nm}^* \bar{d}_L^m \gamma^\mu W_\mu^- \quad (1.10)$$

The same indices ij and nm are summed over on both side. This is equivalent to Equation 1.11:

$$V_{ij} = V_{ij}^* \quad (1.11)$$

On the one hand, if the CKM matrix is real, CP will be conserved in the weak interaction in the SM due to the natural hold of Equation 1.11. On the other hand, from Equation 1.10, it's still possible to make Lagrangian invariant even if V_{CKM} is not real, which can be achieved by introducing non-physical phases for each quark field $u_L^k e^{(i\phi_{uk})}$ and $d_L^j e^{(i\phi_{dj})}$, the Equation 1.11 becomes Equation 1.12.

$$V_{kj} e^{i(\phi_{dj} - \phi_{uk})} = V_{kj}^* e^{i(\phi_{uk} - \phi_{dj})} \quad (1.12)$$

Assuming the complex phase of the kj -th element in CKM matrix is θ_{kj} , it's obviously required Equation 1.13 to hold.

$$\theta_{kj} = \phi_{uk} - \phi_{dj} \quad (1.13)$$

If the number of generations in quark family is 3 or more, the non-physical phases can not render proper values to ensure the hold of Equation 1.13, and there will always be one irreducible complex phase parameter in the CKM matrix in the existence of three generations of quarks, which means CP symmetry is no longer conserved in the weak interactions.

The 3×3 unitary CKM matrix can be written as Equation 1.14 based on the

quark fields it connects using Equation 1.5 and 1.6.

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \quad (1.14)$$

It can be parameterized into the form of Equation 1.15.

$$V_{CKM} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{-i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{-i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}s_{23}s_{13}e^{-i\delta} & -c_{12}c_{23} - s_{12}s_{23}s_{13}e^{-i\delta} & c_{23}c_{13} \end{pmatrix} \quad (1.15)$$

where the $c_{jk} = \cos(\theta_{jk})$ and $s_{jk} = \sin(\theta_{jk})$, and δ is the irreducible complex phase. By measuring the relative branching ratio of $b \rightarrow c$, $s \rightarrow u$ and $b \rightarrow u$ in tree level transitions as shown in Equation 1.16.

$$|V_{ub}| \ll |V_{cb}| \ll |V_{us}| \quad (1.16)$$

The relations in Equation 1.17 are often used to simplify CKM matrix presentation.

$$s_{13} = \lambda, s_{23} = A\lambda^2, s_{13}e^{i\delta} = A\lambda^3(\rho - i\eta) \quad (1.17)$$

By using Equation 1.17, CKM matrix is parameterized as Equation 1.18.

$$V_{CKM} = \begin{pmatrix} 1 - 1/2\lambda^2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - 1/2\lambda^2 & A\lambda^2 \\ A\lambda^3(\rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4) \quad (1.18)$$

Using the unitary condition, the Equation 1.19 is obtained.

$$1 + \frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*} + \frac{V_{td}V_{tb}^*}{V_{cd}V_{cb}^*} = 0 \quad (1.19)$$

Using Equation 1.19, 1.20 and 1.21, the shape of CKM triangle can be defined on the

complex plane in Figure 1-2.

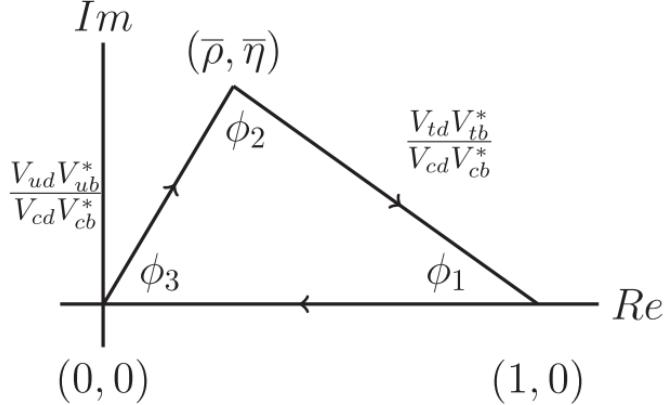


Figure 1-2: The unitary triangles of CKM[12].

$$\bar{\rho} + i\bar{\eta} = -\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*} \quad (1.20)$$

$$1 - (\bar{\rho} + i\bar{\eta}) = -\frac{V_{td}V_{tb}^*}{V_{cd}V_{cb}^*} \quad (1.21)$$

These angles are obtained by drawing the $(\bar{\rho}, \bar{\eta})$ on the complex coordinates, and they are also well-known in the names as: $\phi_1 = \beta, \phi_2 = \alpha, \phi_3 = \gamma$. The results presenting the measurement of CKM angles or $(\bar{\rho}, \bar{\eta})$ in 2019 are shown in Figure 1-3.

The measurement of ϕ_1 and ϕ_2 are mainly obtained from the time-dependent CP violations (TDCPV) measurement. The ϕ_1 in the tree-level dominated decays has been precisely measured due to the small hadronic uncertainties. Flavor-Changing-Neutral-Current (FCNC) processes can rise through the $B_d^0 - \bar{B}_d^0$ mixing in box diagram, and it's believed that potential NP processes might contribute to the difference in between results of CKM angles measured from experiments, such as ϕ_1 value in tree-dominated processes and penguin-dominated processes, where both involve $b \rightarrow s$ transition. It requires the precise measurements on multiple decay channels to search for the potential NP effects. The prospective large Belle II data and improved detector performance will be much useful to help the discovery of NP in future.

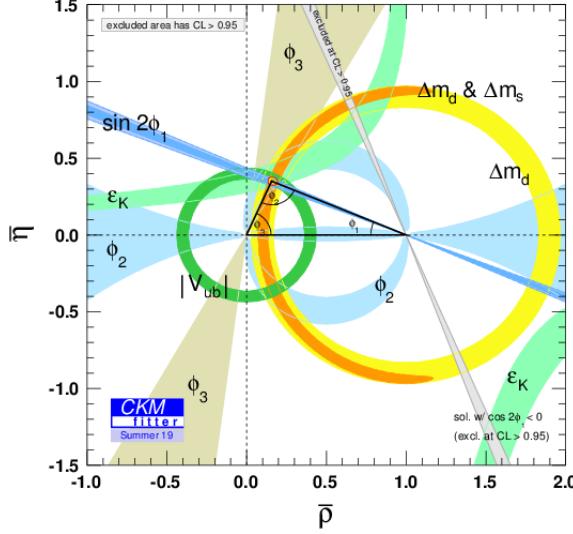


Figure 1-3: The CKM triangle fit in the complex plane of $\bar{\rho} - \bar{\eta}$.[12]

1.4 Time Dependent CP violation

1.4.1 CP violation in neutral B system

The ϕ_1 , ϕ_2 and ϕ_3 are essentially measuring the CKM CP violating phase since there's only one complex phase in the CKM matrix and it can be determined by these three angles. For determining the value of ϕ_1 , TDCPV measurements provide a good experimental environment. From Figure 1-2, one can obtain ϕ_1 and ϕ_2 by Equation 1.22 and 1.23.

$$\phi_1 = \text{Arg}\left(-\frac{V_{td}V_{tb}^*}{V_{cd}V_{cb}^*}\right) \quad (1.22)$$

$$\phi_2 = \text{Arg}\left(-\frac{V_{td}V_{tb}^*}{V_{ud}V_{ub}^*}\right) \quad (1.23)$$

The time-dependent CP violation comes from the interference of neutral B mixing phase and the weak phase in the decay amplitude. The mass eigenstates which are driving the propagation of neutral B meson states with mixing are: $|B\rangle_{H,L} = p|B\rangle \pm q|\bar{B}\rangle$, where H and L stand for the heavier and lighter mass eigenvalues. The $|B\rangle$ and $|\bar{B}\rangle$ present the flavor eigenstates of neutral B mesons. The Hamiltonian

matrix can be written using flavor eigenstates as shown in Equation 1.24.

$$M_\Gamma = \begin{bmatrix} m - i/2\Gamma & M_{12} - i/2\Gamma_{12} \\ M_{12}^* - i/2\Gamma_{12}^* & m - i/2\Gamma \end{bmatrix} \quad (1.24)$$

Considering the time evolution of mass eigenstates, the time-dependent states can be shown as Equation 1.25 and 1.26 by using the notation of $B_{H,L}$ as physical states at $t = 0$.

$$B_H(t) = e^{-im_H t} e^{-\Gamma_H t/2} B_H \quad (1.25)$$

$$B_L(t) = e^{-im_L t} e^{-\Gamma_L t/2} B_L \quad (1.26)$$

where $M_{H,L}$ and $\Gamma_{H,L}$ are the masses and decay width of two mass eigenstates. By expanding the mass eigenstates using flavor eigenstates, which are shown in Equation 1.27 and 1.28.

$$B(t) = (1/2p)e^{-im_H t} e^{-\Gamma_H t/2}(pB + q\bar{B}) + (1/2p)e^{-im_L t} e^{-\Gamma_L t/2}(pB - q\bar{B}) \quad (1.27)$$

$$\bar{B}(t) = (1/2q)e^{-im_H t} e^{-\Gamma_H t/2}(pB + q\bar{B}) - (1/2q)e^{-im_L t} e^{-\Gamma_L t/2}(pB - q\bar{B}) \quad (1.28)$$

Replacing $g_\pm(t) = \frac{1}{2}(e^{-im_H t - \Gamma_H/2t} \pm e^{-im_L t - \Gamma_L/2t})$, Equation 1.27 and 1.28 become Equation 1.29 and 1.30.

$$B(t) = g_+(t)B + \frac{q}{p}g_-(t)\bar{B} \quad (1.29)$$

$$\bar{B}(t) = g_+(t)\bar{B} + \frac{p}{q}g_-(t)B \quad (1.30)$$

Considering all the phase-spaces of the decay from flavor eigenstates to final states $f(\bar{f})$ are included in the amplitudes $\mathcal{A}_f(\bar{\mathcal{A}}_{\bar{f}})$, one needs to expand the flavor eigenstates using the final states amplitudes to have the differential decay rate $\Gamma(B \rightarrow f, t)$. From $B(t) \propto \mathcal{A}_f \psi_f + h.c$ and $(\bar{B}(t) \propto \bar{\mathcal{A}}_{\bar{f}} \psi_{\bar{f}} + h.c)$, combined with Equation 1.29 and

1.30, the decay rate can be shown in Equation 1.31 and 1.32.

$$\Gamma(B \rightarrow f, t) = |\mathcal{A}_f|(|g_+(t)|^2 + |\lambda_f|^2 |g_-(t)|^2 + 2\text{Re}(\lambda_f g_+^*(t) g_-(t))) \quad (1.31)$$

$$\Gamma(\bar{B} \rightarrow \bar{f}, t) = |\bar{\mathcal{A}}_f|(|g_+(t)|^2 + |\bar{\lambda}_{\bar{f}}|^2 |g_-(t)|^2 + 2\text{Re}(\bar{\lambda}_{\bar{f}} g_+^*(t) g_-(t))) \quad (1.32)$$

where the parameter λ_f and $\bar{\lambda}_{\bar{f}}$ can be defined as Equation 1.33 and 1.34.

$$\lambda_f \equiv (q/p)(\bar{\mathcal{A}}_f / \mathcal{A}_f) \quad (1.33)$$

$$\bar{\lambda}_{\bar{f}} \equiv (q/p)(\mathcal{A}_{\bar{f}} / \bar{\mathcal{A}}_{\bar{f}}) \quad (1.34)$$

The q/p is introduced by the coefficient of mass eigenstates from weak eigenstates. Using the Hamiltonian matrix, q/p can be presented using Equation 1.35

$$q/p = \frac{\Delta M - i/2\Delta\Gamma}{2(M_{12} - i/2\Gamma_{12})} \quad (1.35)$$

where the M_{12} and Γ_{12} stands for the contribution of non-diagnosed term in the Hamiltonian matirx. $\Delta M = M_H - M_L$ and $\Delta\Gamma = \Gamma_H - \Gamma_L$ are the difference of mass and decay width for two mass eigenstates, respectively. It's obvious that if $|\mathcal{A}_f| \neq |\bar{\mathcal{A}}_{\bar{f}}|$, direct CP violation will occur. The time-dependent decay rate difference is defined as Equation 1.36.

$$\begin{aligned} A_{CP}(t) &\equiv \frac{\Gamma(B \rightarrow f, t) - \Gamma(\bar{B} \rightarrow \bar{f}, t)}{\Gamma(B \rightarrow f, t) + \Gamma(\bar{B} \rightarrow \bar{f}, t)} \\ &= \frac{\mathcal{S}\sin(\Delta Mt) - \mathcal{A}\cos(\Delta Mt)}{\cosh(\Delta\Gamma t/2) + A_{\Delta\Gamma}^f \sinh(\Delta\Gamma t/2)} \end{aligned} \quad (1.36)$$

where

$$\mathcal{S} = \frac{2\text{Im}(\lambda_f)}{1 + |\lambda_f|^2} \quad (1.37)$$

$$\mathcal{A} = \frac{1 - |\lambda_f|^2}{1 + |\lambda_f|^2} \quad (1.38)$$

$$A_{\Delta\Gamma}^f = -\frac{2\text{Re}(\lambda_f)}{1 + |\lambda_f|^2} \quad (1.39)$$

From Equation 1.37 and ??, the time-dependent CP violation parameters \mathcal{S} and \mathcal{A} are dependent on the parameter λ_f , which are determined by Equation 1.33.

1.4.2 ϕ_1 from $B^0 \rightarrow J/\psi K_S^0$

If final states are CP eigenstates, the amplitude is obtained by $\mathcal{A}_f \equiv \langle f | H | B \rangle$ and $\bar{\mathcal{A}}_f \equiv \langle f | H | \bar{B} \rangle$. In $B_d^0 - \bar{B}_d^0$ mixing system, the q/p can be treated as $e^{i\phi_d}$ as pure phase term. This relative phase accounts the transition from b to up-type quarks to strange quark s in mixing, so it can be presented as $\phi_d = \text{Arg}(V_{td}^* V_{tb}) / (V_{tb}^* V_{td}) \approx 2\phi_1$ based on negligible correction to the SM. In mode $B^0 \rightarrow J/\psi K_S^0$, considering $\Delta\Gamma$ can be treated as zero in the SM in this case[13], Equation 1.36 can be reduced to Equation 1.40.

$$A_{CP}(t) = \mathcal{S}\sin(\Delta M t) - \mathcal{A}\cos(\Delta M t) \quad (1.40)$$

which receives contributions from tree-level and loop-level processes shown in Figure 1-4 ,

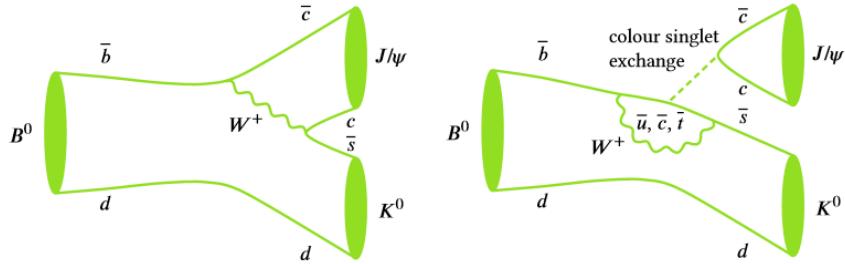


Figure 1-4: The dominated tree-level (left) and the suppressed loop-level (right) of $B \rightarrow J/\psi K^0$, in which K^0 particles are detected K_S^0 [14].

Using the relation $|V_{ub}| \ll |V_{cb}| \ll |V_{us}| < |V_{cs}|$, it's obvious that $V_{ub}^* V_{us} \ll V_{cb}^* V_{cs}$, so the penguin-mode is suppressed in the Standard Model. Defining η_f as the CP eigenvalue,

Given $\eta_f = 1$ and $|\lambda_f| = 1$ in $B^0 \rightarrow J/\psi K_S^0$, from 1.37, CP violation parameters

can be presented as shown in Equation 1.41.

$$\mathcal{S} = \text{Im}(\lambda_f) = -\sin(\phi_d)\eta_f = -\sin(2\phi_1); \mathcal{A} = 0 \quad (1.41)$$

From Equation 1.41, ϕ_1 can be obtained precisely in the measurement of time-dependent CP violation in $B^0 \rightarrow J/\psi K_S^0$.

1.4.3 ϕ_1 from penguin-dominated mode $b \rightarrow q\bar{q}s$

Compared to $B^0 \rightarrow J/\psi K_S^0$ channel, the measurement of \mathcal{S} and \mathcal{A} from penguin-dominated channels through $b \rightarrow q\bar{q}s$ where q is u, d, s can be different due to the varied tree-to-penguin amplitude ratio. Furthermore, they are quite sensitive to NP effects for the following reasons[15]. First, they can probe $B^0 - \bar{B}^0$ through different short-distance vertices than the tree-level dominated decays. Second, the tree-level decay amplitude is suppressed and penguin-level amplitude is dominated, where the overall non-NP amplitude is relatively small so NP effects may show up easier. Last but not least, they comprise a large number of different final states, which can help disentangling non-perturbation long-distance physics from short-distance information, such as ϕ_1 or NP contributions to the weak Hamiltonian.

Considering possible New Physics contribution besides the tree-level and penguin-level processes as A_f^{NP} , the decay amplitude can be rendered as Equation 1.42

$$\mathcal{A}_f = \lambda_u^s T_f + \lambda_c^s P_f + A_f^{NP} \quad (1.42)$$

where T_f and P_f are tree-level and penguin-level amplitudes. The coefficients λ_u^s and λ_c^s are determined from CKM matrix elements by $\lambda_i^q \equiv V_{ib}^* V_{iq}$. Note that compared to the $B^0 \rightarrow J/\psi K_S^0$, the tree level amplitude T_f is suppressed and penguin amplitude P_f is dominated in $b \rightarrow q\bar{q}s$. It is also worth noting that T_f contains tree-level W^\pm exchange, QCD and electroweak penguin contributions. These carry the combination of CKM matrix elements $\lambda_t^s = V_{ts} V_{tb}^* = -(1 + \epsilon_{uc})\lambda_c^s$ where $\epsilon_{uc} \equiv \lambda_u^s / \lambda_c^s = \mathcal{O}(\lambda^2)$. In the SM with neglected ϵ , $b \rightarrow q\bar{q}s$ modes are pure penguin with the same weak phase

as $B^0 \rightarrow J/\psi K_S^0$ has. Thus, direct CP violation vanishes and time-dependent CP violation reflects \mathcal{S} in the same way as $B^0 \rightarrow J/\psi K_S^0$ does.

Departures from this limit, non-neglected tree amplitude T_f (often called “tree pollution”), as well as possible NP effects, could give different results on ϕ_1 . Introducing the tree-penguin ratio $r_f^T = T_f/P_f$, NP-to-SM ratio $r_f^{NP} = \mathcal{A}_f^{NP}/(\lambda_c^s P_f)$, the following statements are usually used[15]:

- Branching ratios are affected at $\mathcal{O}(|\epsilon_{uc} r_f^T|, |r_f^{NP}|)$
- Direct CPV in the SM are of $\mathcal{O}(\epsilon_{uc} \text{Im}(r_f^T))$
- $-n_f^{CP} \mathcal{S} = \sin(2\phi_1) + \Delta \mathcal{S}$, where $\Delta \mathcal{S} = 2\cos 2\phi_1 \sin \phi_3 |\epsilon_{uc}| \text{Re}(r_f^t) + \Delta \mathcal{S}^{NP}$

1.4.4 ϕ_1 from $B^0 \rightarrow K_S^0 K_S^0 K_S^0$

Since Belle reported the time-dependent CP analysis on various $b \rightarrow q\bar{q}s$ which experimentally showed that the difference on ϕ_1 has a margin for NP effects[16], the improved measurements with a larger data collection is popularly discussed in order to reduce the impact of uncertainties and clear the tension between results. The decay channel $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ is one of the most promising modes for this purpose. The CP eigenvalue of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ is positive ($CP = +1$). Since there’s no up-quark shown in the final states, the potential contribution of $b \rightarrow u\bar{u}s$ re-scattered into $b \rightarrow s\bar{s}s$ is almost of absence. This makes $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ a cleaner channel than $B^0 \rightarrow K^+ K^- K_S^0$ which has a different weak phase contribution [17]. Any NP effects expected in the $B^0 \rightarrow \phi K_S^0$ should also affect $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ and the absence of NP effects will lead the same CP violation just as $J/\psi K_S^0$ [17]. Currently there’s no specific theoretical calculation on the $\Delta \mathcal{S}$ for three-body $B^0 \rightarrow K_S^0 K_S^0 K_S^0$. However, due to the same weak phase of this decay as $\eta' K_S^0$ and ϕK_S^0 , the theoretical prediction on $\Delta \mathcal{S}$ is usually applied to $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ as well[17]. The expected range on $\Delta \mathcal{S}$ is typically at level of ~ 0.05 [18], which requires the expected precision improvement for both statistical and systematic uncertainty in future data. The result of ϕ_1 from $B^0 \rightarrow J/\psi K_S^0$ using full Belle data is presented as $\mathcal{S}_{J/\psi K_S^0} = +0.670 \pm 0.029(\text{stat}) \pm 0.013(\text{syst})$ [15]. In the meantime, the latest result of ϕ_1 from $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ using full Belle data[19] is presented as: $\mathcal{S}_{3K_S^0} = -0.71 \pm 0.23(\text{stat}) \pm 0.05(\text{syst})$, and the

result from BaBar [20] is: $\mathcal{S}_{3K_S^0} = -0.94^{+0.21}_{-0.24}(\text{stat}) \pm 0.06(\text{syst})$. Both results have shown a small deviation from the result in $B^0 \rightarrow J/\psi K_S^0$ for NP effects while the statistical uncertainties are dominated which prevents the claim whether NP effects are spotted. For $\Delta\mathcal{S}$ from $B^0 \rightarrow K_S^0 K_S^0 K_S^0$, the experimental sensitivity of $\Delta\mathcal{S}$ will be dominated by $\mathcal{S}_{3K_S^0}$ because the total uncertainty from $J/\psi K_S^0$ will be reduced to about 0.005 at 50 ab^{-1} Belle II data[15]. The Figure 1-5 shows the scaled $\Delta\mathcal{S}$ uncertainty regarding the luminosity in Belle II prospective[21], which only includes the statistical uncertainty of Table 1.1. If the conservative estimation of systematic uncertainty from Belle is considered, the red arrow shows the approximate luminosity where the experimental sensitivity becomes comparable with theoretical prediction at ~ 0.05 . If the future result is different from that in $J/\psi K_S^0$ with 5 times of the uncertainty at 0.05, then it could be an evidence for NP effects. Of course, smaller the total uncertainty is, easier identifying the existence of NP effects will be.

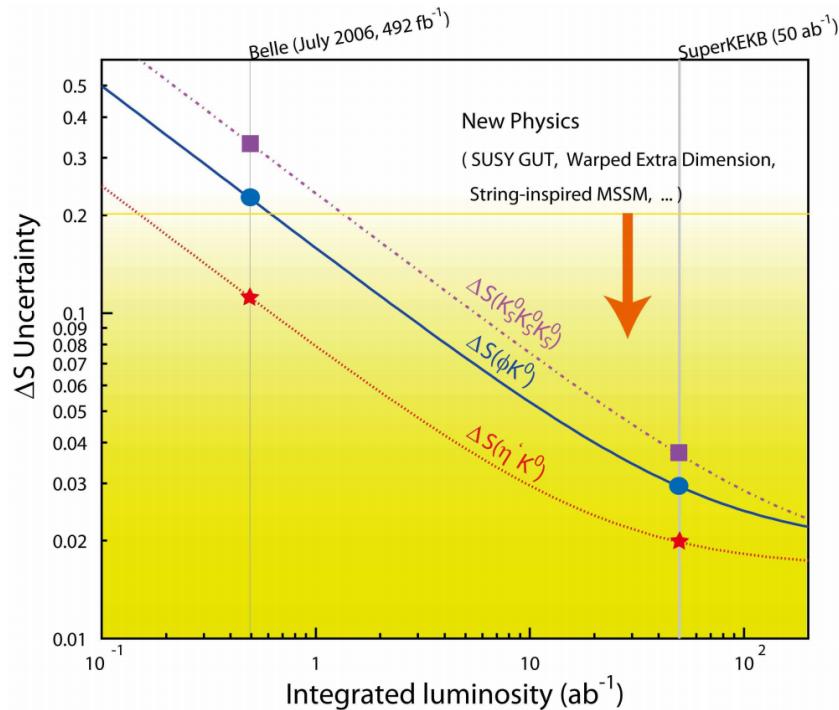


Figure 1-5: Expected sensitivity of $\Delta\mathcal{S}$ regarding the integrated luminosity of Belle II prospective.[21]

Table 1.1: $\Delta\mathcal{S}$ estimated statistical uncertainty with respect to the integral luminosity in Belle II technical design report[21].

Observable	Belle (0.5 ab $^{-1}$)	Belle II (5 ab $^{-1}$)	Belle II (50 ab $^{-1}$)
$\Delta\mathcal{S}_{\phi K_S^0}$	0.22	0.073	0.029
$\Delta\mathcal{S}_{\eta' K_S^0}$	0.11	0.038	0.020
$\Delta\mathcal{S}_{K_S^0 K_S^0 K_S^0}$	0.33	0.105	0.037

Chapter 2

Belle II experiment

2.1 Belle II and SuperKEKB overview

The goal of the Belle II experiment is to search for evidence of New Physics, and the expected operation period is from 2019 to the end of 2030. The facilities are located in KEK, Tsukuba City, around 70 km in the north of Tokyo, Japan. The SuperKEKB accelerator enables electron-positron collision at the center-of-mass energy on the region of $\Upsilon(4S)$ resonances which is about the mass of two B mesons. The electron and positron beams are designed at 7 GeV and 4 GeV, respectively, with boost factor of 0.28, providing an environment for measuring time-dependent CP violation by displacing the decay vertices of a B meson pair in a measurable distance along the boosted direction. The SuperKEKB has a targeted luminosity of $8 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$, a factor of 40 times higher than its predecessor, the KEKB. Some key parameters of the SuperKEKB are listed in Table 2.1. The overview of SuperKEKB and Belle II are shown in Figure 2-1.

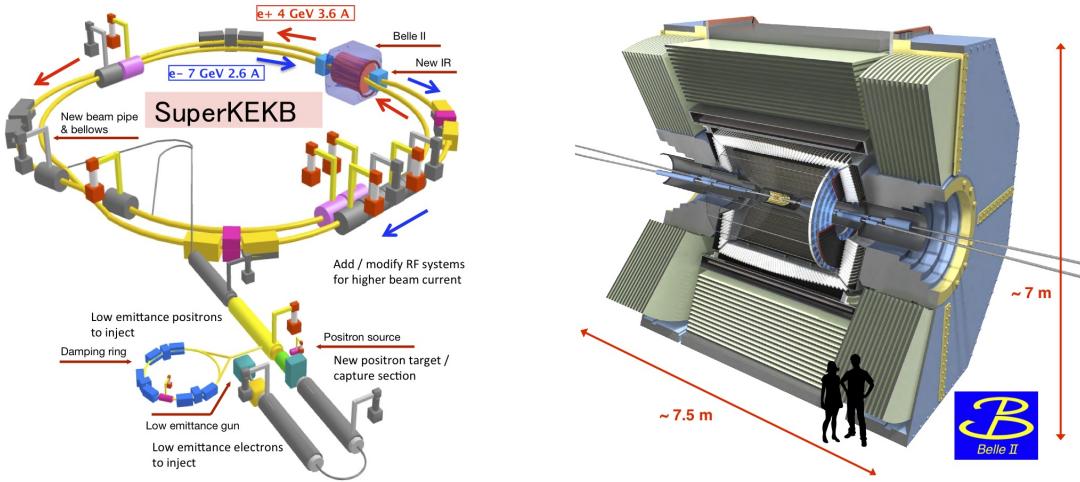


Figure 2-1: Overview of SuperKEKB and Belle II detector [21].

Table 2.1: SuperKEKB parameters for low energy (LER) and high energy (HER) rings.[15]

Parameters	LER(e^+)	HER(e^-)	Unit
Energy	4.0	7.0	GeV
Half crossing angle	41.5		mrad
Horizontal emittance	3.2	4.6	nm
Emittance ratio	0.27	0.25	%
Beta functions at IP (x/y)	32/0.27	25/0.30	mm
Beam currents	3.6	2.6	A
Beam-beam parameter	0.0881	0.0807	
Luminosity	8×10^{35}		$\text{cm}^{-2}\text{s}^{-1}$
Perimeter of ring	3		km

The Belle II detector has a close size as the Belle detector so that it is placed in the same shell, but all sub-detectors and electronic systems have been either newly built or considerably upgraded. The advantage of the SuperKEKB requires that the

Belle II has to be able to stably operate at a 40 times higher events rates as well as 10 to 20 times higher beam background compared to Belle. The mitigation of the effects caused by such high beam background is essential to the success of Belle II. Higher background level leads to higher occupancy and radiation damage to the detectors, along with more fake hits in the vertex detectors and central drift chamber, pile-up backgrounds in electromagnetic calorimeter and neutron-induced hits in muon detector. Data-acquisition system (DAQ) and trigger are also upgraded not only to adapt to higher luminosity but also for a better low-multiplicity event sensitivity. Belle II detector in the top view is shown in Figure 2-2.

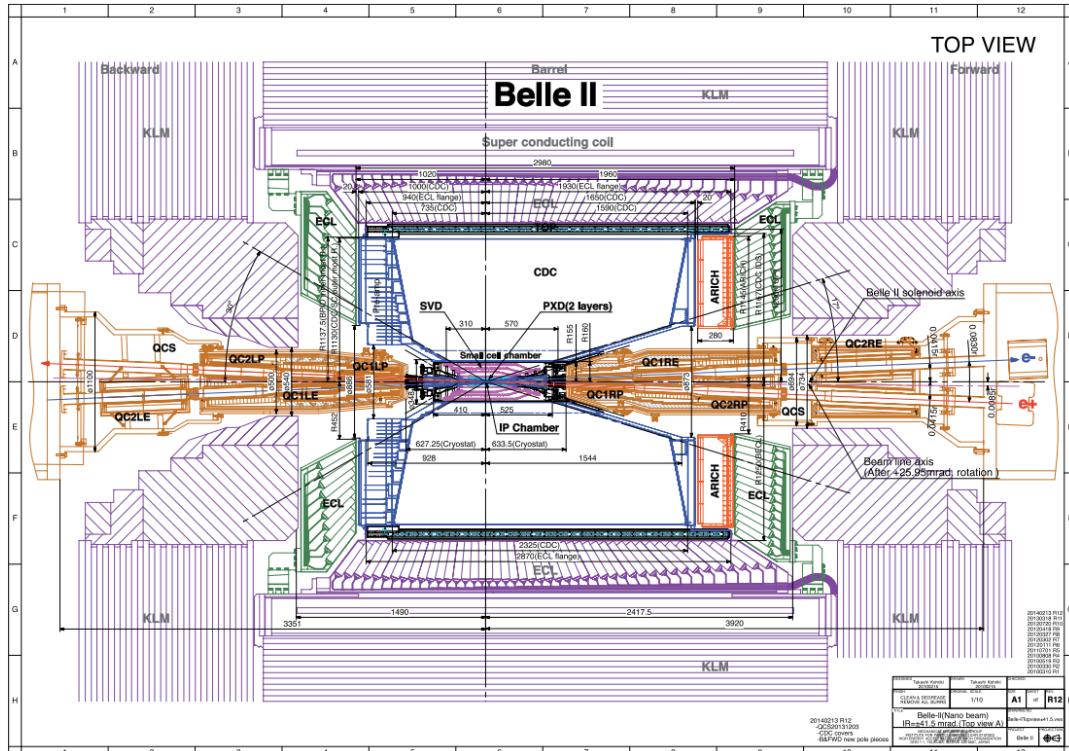


Figure 2-2: Belle II detector top view [15].

The success of Belle II detector depends on the complex of sub-detectors which each of them is design for specific purposes. The critical components and features are explained in the following sections.

2.2 Vertex detector (VXD)

The vertex detector is composed of two components, the silicon based pixel detector (PXD) and silicon based vertex detector (SVD), where total 6 layers are placed in the inner-most region from interaction point (IP). The geometry of VXD is shown in Figure 2-3. The PXD is placed at a radii of $r = 14$ mm and $r = 22$ mm with DEPFET[21] type pixel sensors, which is designed to provide two dimensional hit position information. The inner layer leaves a sufficient space for possible variations of the beampipe layout. The size of two layers are determined by the required acceptance angle, which is 17 degrees (forward) to 150 degrees (backward). The pixel sensor is a monolithic structure with current-digitizing electronics at the end of the sensor which makes the a very thin a layer at about 50 microns. The schematic view of sensors on PXD is shown in Figure 2-4. As the very close range the PXD is, the sensors

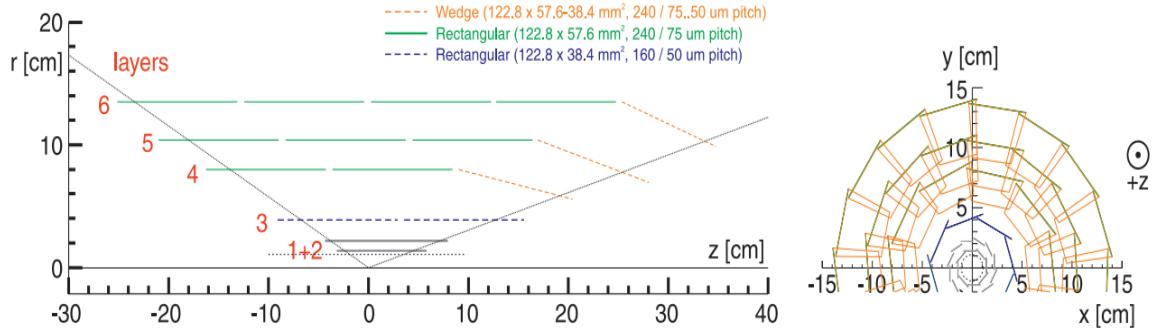


Figure 2-3: A schematic view of PXD (2 layers in gray) and SVD (4 layers in green and orange). [21].

are exposed to a very high event rate and very high beam background environment. The large data flow from PXD without any data reduction scheme is problematic for data acquisition system. In order to reduce the data that is not interested by physics analysis such as beam backgrounds, a fast online tracking system is built up for searching a “region of interest” (ROI) on the PXD sensors. To be specific, the data from PXD will be first readout to a system called “ONSEN” which can store large size temporary data up to 5 seconds. In this timing window, a fast online tracking system will perform a track fitting using vertex detector and central drift chamber to

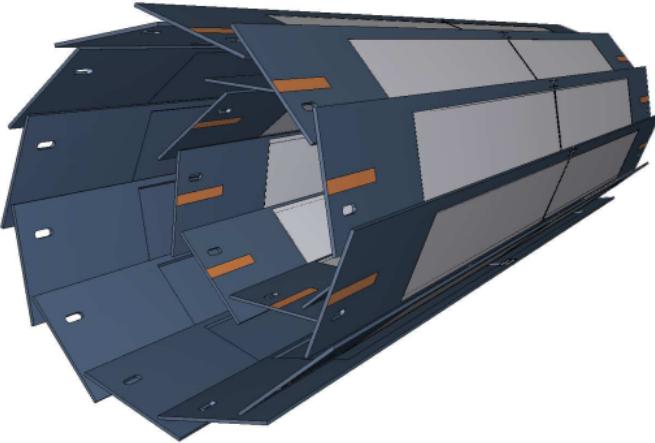


Figure 2-4: The geometry of sensors on PXD where the light grey surfaces are DEPFET sensors with a thinness of 50 microns. The full length including the out modules is 174 mm[21].

extrapolate the fitted tracks backward to PXD plane so the ROI on the PXD sensors can be defined. The data from PXD outside of the ROI is not read out to external tapes where offline data is written.

SVD detector consists of 4 layers of detectors called “double-sided silicon detectors” (DSSDs) at 39 mm, 80 mm, 104 mm, and 135 mm away from IP, respectively. The two sides of the sensors are called *p*-side and *n*-side, where the former is for strips on $r - \phi$ direction and the latter is for strips on the z direction. To suppress the background hits, a readout chip with a fast shaping time of $\mathcal{O}(50\text{ ns})$ is indispensable. The APV25 chip[22] is chosen as the readout chip that was originally developed for CMS silicon tracker, with total 128 identical channels of low-noise preamplifiers followed by a 50 ns peaking time shaper stage. The polar angular acceptance ranges from 17 degrees to 150 degrees, which is asymmetric to account for the forward boost of the center-of-mass frame. The combination between sensors, electronics and the supporting structure uses so-called “Origami” concept which stands for a Chip-on-Sensor design. In the Origami scheme, the readout chips APV25 are placed on a single flexible circuit mounted on the *n*-side of the sensors. The channels of *p*-side are attached by small flexible fan-outs wrapped around the edge of the sensors. All connections between flex pieces, sensor, and APV25 chips are made by wire bonds.

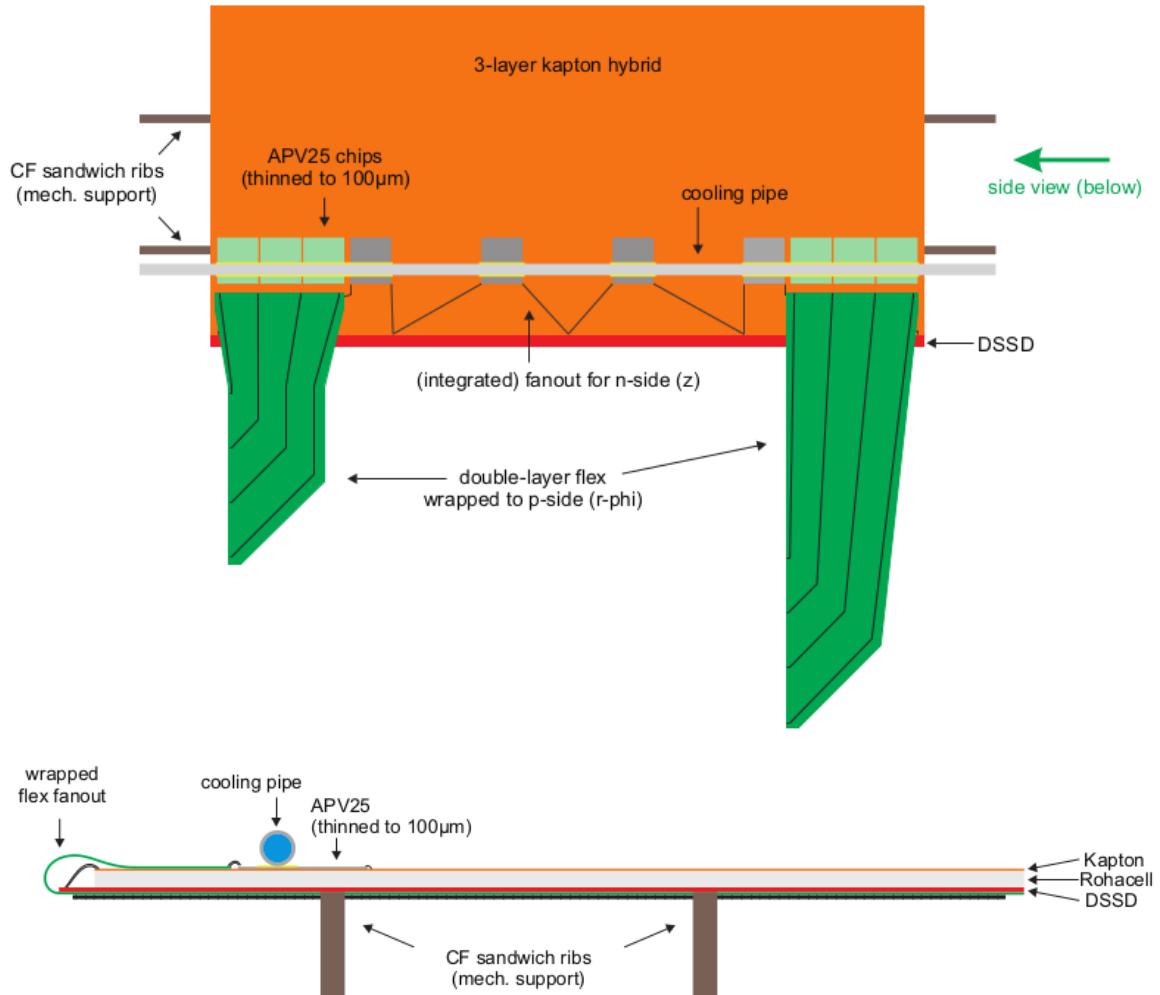


Figure 2-5: The top and side views of Origami Chip-on-Sensor design for DSSDs of SVD. Top: the APV25 chips in grey read out the same side sensors channel while chips in green read out the sensors on the opposite side using wrapped-around flex pieces. Bottom: side view of the Origami design shows the location of wrapped flex which connects the strips of the bottom sides which are placed at the left edge.[21]

2.3 Central drift chamber (CDC)

The Central drift chamber (CDC) is the core component of spectrometer in Belle II, which consists of a fairly big drift chamber made of many small drift cells filled with gas. The chamber gas is comprised of a He–C₂H₆ 50%:50% mixture with an average drift velocity of 3.3 cm μs^{-1} and a maximum drift time of about 350 ns for a 17 mm cell size. The out radius of CDC has been extended to 1130 mm from 880 mm of Belle thanks to the much thinner layers in barrel region. The whole CDC contains 14336 sense wires in 56 layers, placed in the axial direction and the stereo direction[21][15]. Such a design can utilize the information from axial and stereo wires to construct a full 3 dimensional hits which reflects helix tracks in CDC volume. Thus, CDC is one of the key components for measuring the helix parameters for tracking, providing precise information on the charged tracks momentum. Also, it provides particle identification information using measurements of energy loss within its gas volume. Low-momentum tracks, which do not reach the particle identification device, can be identified using the CDC alone. Finally, it provides efficient and reliable trigger signals for charged particles.

Belle II CDC is expected to handle higher trigger rates with less dead time. The front-end electronics are located near the backward end-plate and send digital signals to the electronics hut through optical fibers. Due to the higher radiation and higher beam background in Belle II, also to create more space for SVD volume, the inner radius of CDC in Belle II is 160 mm and the outer radius is 1130 mm. Within this volume, CDC can create three dimensional trigger information. A dedicated trigger type called z trigger[21] based on the 3D tracking achieved by an FPGA using axial and stereo wires is implemented.

The structure of CDC consists of three main components which are a thin carbon-fiber reinforced plastic (CFRP) inner cylinder, two aluminum endplates, and a CFRP outer cylinder, as shown in Figure 2-6. The outer cylinder is a thickness of 5 mm structure supporting most of the wire tension of 4 tonnes. The inner cylinder is as thin as 0.5 mm to minimize the material and support small cell chamber such as the

layers in the inner most region.

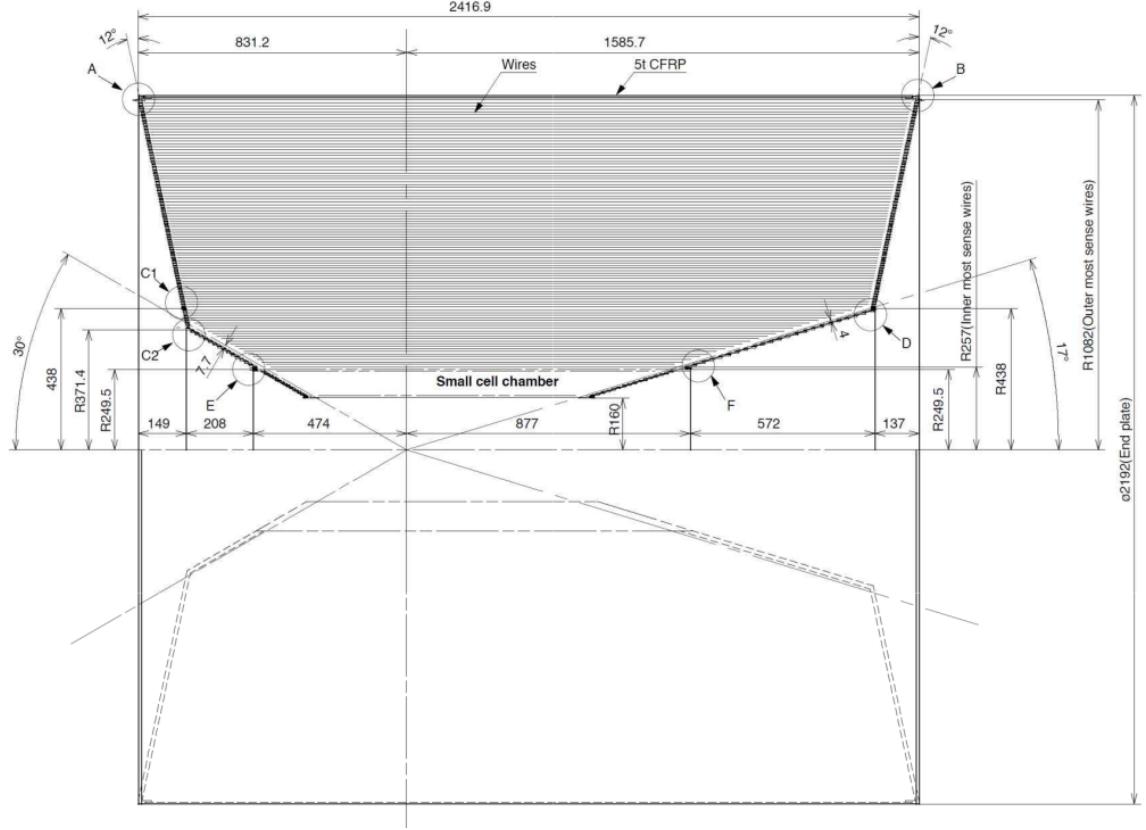


Figure 2-6: CDC structure schematic view[21].

2.4 TOP and ARICH detectors

The particle identification system of Belle II mainly consists of two parts, time-of-propagation counter (TOP) and aerogel based Cherenkov radiation imaging ring (ARICH).

TOP is the specialized detector that can reconstruct Cherenkov radiation's time of arrival and generated position by a photon detector placed at the end of a 2.6 cm quartz bar. The TOP is placed at the barrel region of the spectrometer, as shown in Figure 2-2. The conceptional view and the working principle of TOP counter are shown in Figure 2-7. In this counter, the time of propagation of the Cherenkov photons internally reflected inside a quartz radiator is measured. The quartz radiator is composed of three components: a long bar for radiating Cherenkov light and propa-

gating this light via total internal reflection to the bar end, where the MCP-PMTs are mounted; a spherical mirror mounted on the forward end of the bar for focusing the light; and, for the 1-bar design, a prism that attaches to the backward end of the bar and allows the Cherenkov ring image to expand before the photons are recorded by the PMTs. A 3-dimentional information with $x - y$ position and a timing information are obtained by micro-channel plate (MCP) PMTs at the end surfaces of the quartz bar. The resolution of starting time is achieved about 50 ps[21]. As the key component of the photon detector, the squared shape MCP PMTs, donated as SL-10[23], have been developed with a 4×4 anode array, a multi-alkali photocathode, two MCP plates with $10 \mu\text{m}$ pore size, and an aluminum layer on the second MCP to protect against ion feedback. The image of a SL-10 MCP PMT and an anode schematic view are shown in Figure 2-8.

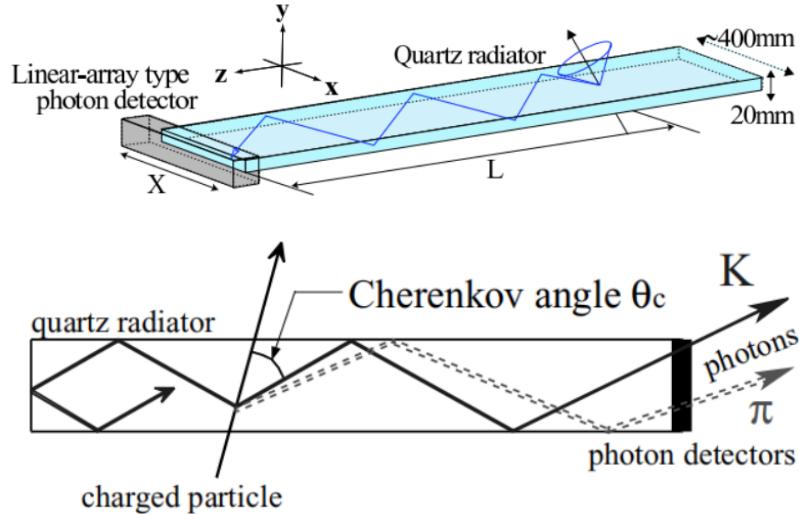


Figure 2-7: Conceptional view of TOP counter (up) and its imaging process of K^\pm and π^\pm (down)[21] for PID purpose.

Aerogel Ring-Imaging Cherenkov detector (ARICH) is located at the forward end-cap in Figure 2-2 to separate charged particles in a momentum range from 0.5 GeV to 4 GeV, which requires single-photon-sensitive high-granularity sensor to reconstruct the Cherenkov angle with small photon yield. Japan has developed a hybrid avalanche photon detector (HAPD) to meet the requirements. Each sensor is $73 \times 73 \text{ mm}^2$ em-

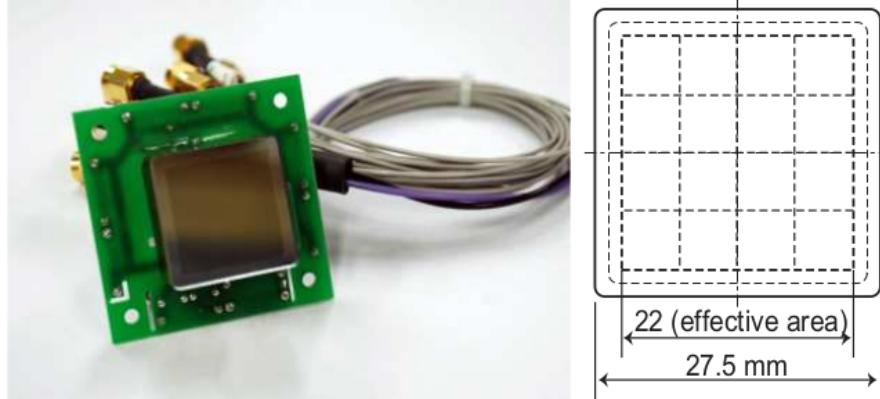


Figure 2-8: SL-10 MCP PMT (left) and the schematic view of 4×4 anode (right)[21]

bedded with 144 channels to accelerate emitted electrons in a 8 kV field. Avalanche photo-diodes (APD) are used for the detection of electrons at the end of electron acceleration, see Figure 2-9. The ARICH detector outlook and the ring image of cosmic muon on the HAPD sensors are shown in Figure 2-10.

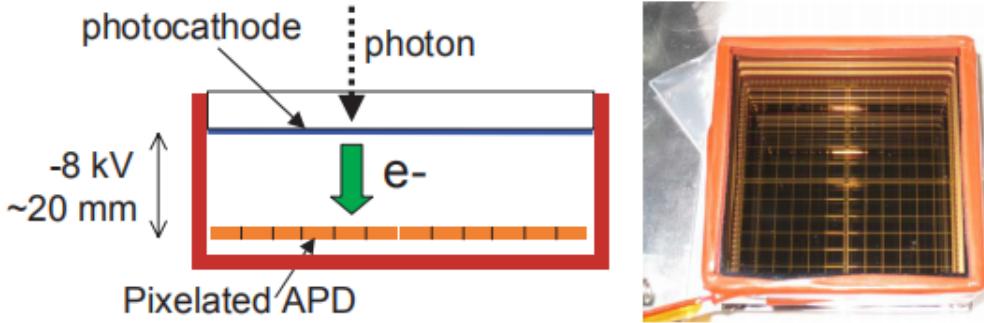


Figure 2-9: Photon-electrons acceleration (left) and pixelated APD (right) at the end[21].

2.5 Electromagnetic calorimeter (ECL)

The electromagnetic calorimeter (ECL) in Belle II is mainly responsible for the detection of γ radiation and electrons, providing energy deposition information for trigger, particle reconstruction and PID. ECL consists of three sections as shown in Figure 2-2: a 3 m long barrel section with an inner radius 1.25 m, and two annular endcaps

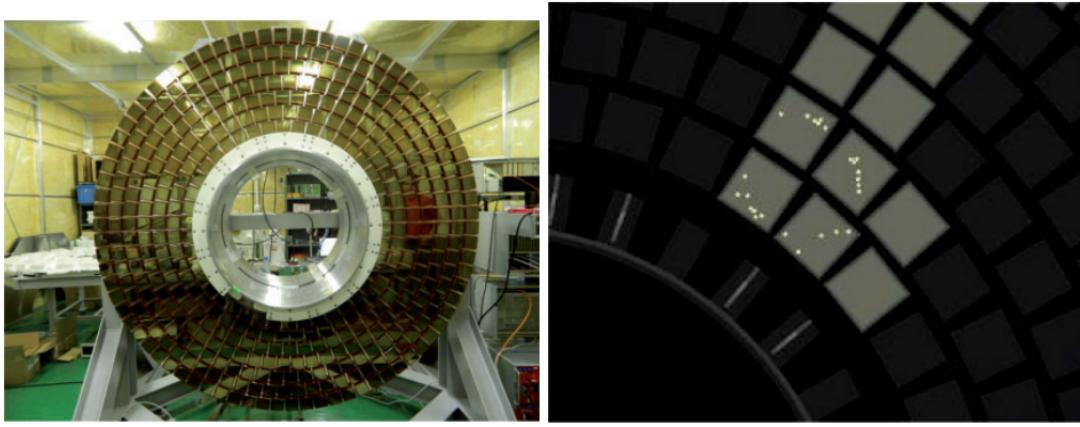


Figure 2-10: ARICH detector (left) and the ring image of cosmic muon on the HAPD sensors[15]

at $z = 1.96$ m (forward) and $z = -1.02$ m (backward) from the IP. The barrel section contains 6624 CsI(Tl) crystals of 29 distinct shapes and each crystal is a pyramid shape with about 6×6 in cross section and 30 cm in length. The endcaps section contains 2112 CsI crystals of 69 shapes and the total number of crystals is 8736, with a total mass of about 43 tons[21].

As the basic component of ECL, the thallium doped caesium iodide CsI(Tl) crystals are assembled tightly in end-caps and barrel sections. Compared to the previous ECL in Belle, the pre-amplifiers and the structures remain unchanged, while the readout electronics have been upgraded. The estimated background level in Belle II ECL will cause the much longer decay time in the scintillation of CsI(Tl). This will lead to the pile-up effect of readout noise. To compensate this effect, wave-form sampling electronics are embedded with the photon detectors (PMT). Especially in the forward direction of the electron beamline, where the level of beam background is much higher, the effect of pile-up noise becomes even worse and the performance of ECL will be of trouble if no special measure taken. Therefore, the pure CsI crystal is considered to be chosen as the material of detector to achieve a fast wave-shaping time and higher radiation tolerance compared to the dosed CsI(Tl), which is an backup option for the future upgrade. ECL is the most important detector for providing trigger information for low multiplicity events, since the main feature of these events

is one or two energetic photon(s) emitted from IP region while the charged tracks are missing.

2.6 K_L^0 muon detector (KLM)

The K_L^0 and muon detector (KLM) system of Belle II consists of a sandwich stacked iron plates at outside of the superconducting solenoid. The iron plates serve as the interaction materials with > 3.9 times the interacting length of material compared to the ECL, allowing K_L^0 particles to shower through. The octagonal barrel covers the polar angle range from 45 degrees to 125 degrees, while the endcaps extend this coverage from 20 degrees to 155 degrees. There are 15 detector layers and 14 iron plates in the barrel and 14 detector layers and 14 iron plates in each endcap. The side view of KLM is shown in Figure 2-11. The Belle KLM material uses the glass-electrode resistivity plate chambers (RPC) which is not suitable for Belle II due to high background level. Neutrons dose is significantly larger due to the much more electromagnetic radiation reaction on detector materials. The long dead time of RPC under such dose rate will reduce the efficiency of KLM. To mitigate this problem, the RPCs are replaced by the layers of scintillator strips with wavelength-shifting fibers, read out by silicon photomultipliers (called “SiPMs”, Geiger mode operated APDs) as light sensors, which is proven to be able to reliably operate by setting up the discrimination threshold [15].

2.7 Trigger and DAQ system

The interesting topics in Belle II physics analysis highly depend on the trigger system. The Belle II trigger system is composed of two levels: a hardware-based, low-level trigger called “L1” trigger, and a software-based high-level trigger (HLT). The L1 trigger has a latency of $\sim 5\mu\text{s}$ and the maximum trigger output rate is 30 kHz, which is limited by the read-in rate of data acquisition system (DAQ). Considered the high event rate and background level from future Belle II luminosity, a series of upgrades

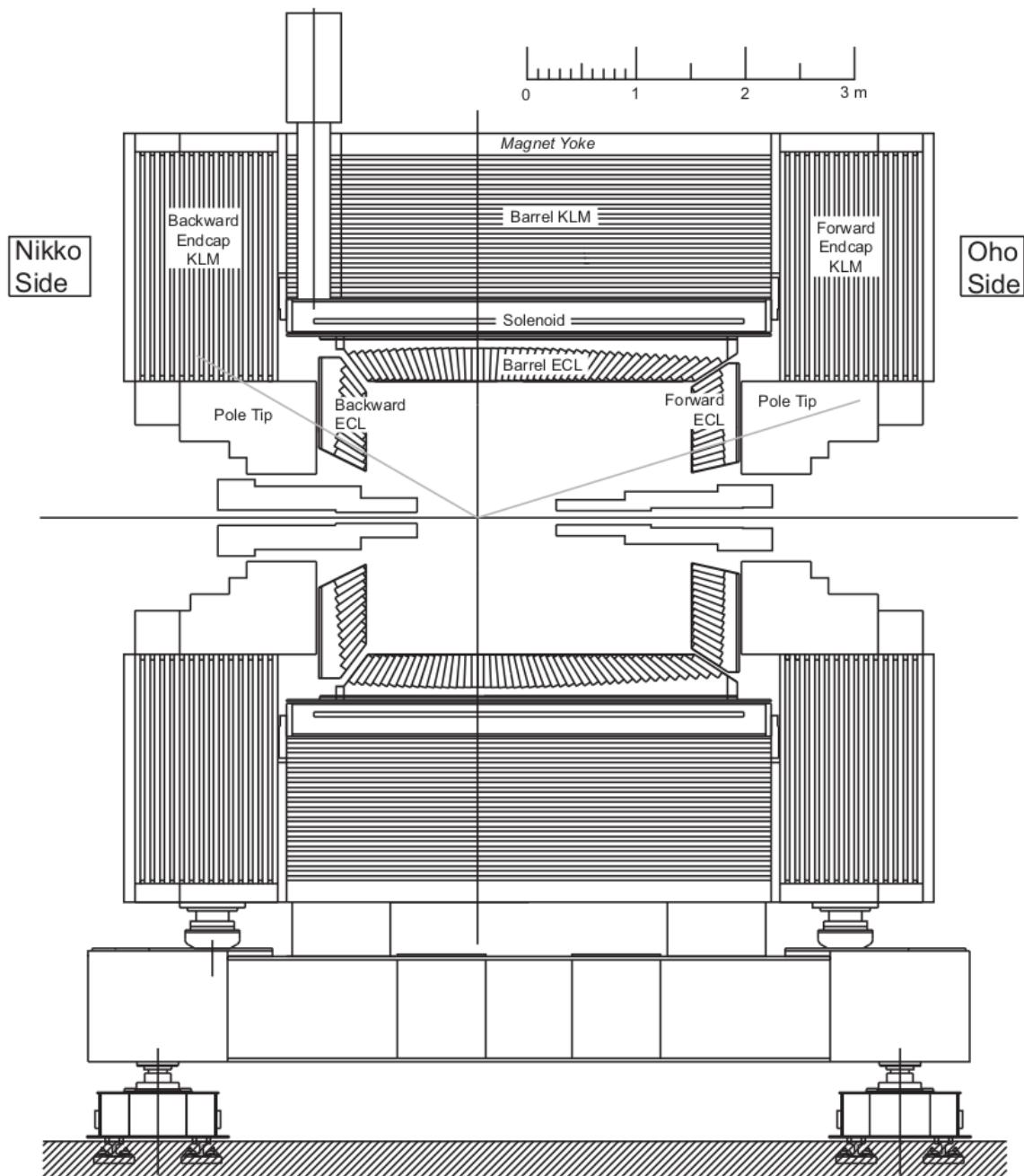


Figure 2-11: The side view of KLM in between the ECL and the solenoid, which the grey lines presents the nominal acceptance angle of the Belle II[21].

have been implemented for L1 trigger. The key improvements of L1 come from the firmware-based reconstruction algorithm and trigger logic.

HLT, as the second level of Belle II trigger systema, plays an important role in DAQ. As discussed in the section of PXD, the data size in PXD is huge at high luminosity and the ROI selection must be applied to reduce the data size. The event rate reduction relies on the tracking extrapolation on PXD plane, which is the key function of HLT. HLT can suppress the event rate down to 15 kHz using the information from the CDC tracking and ECL reconstruction. The information from the detectors other than PXD are fed into the first event builder for full event reconstruction. The event rate is further reduced to below 10 kHz by using full reconstruction information to find track-associated hits on PXD, as introduced as ROI before. The workflow of DAQ with HLT is demonstrated in Figure 2-12. The reduced event rate by applying HLT on PXD and other detector read-out systems are combined into the second event builder and eventually written to the offline storage.

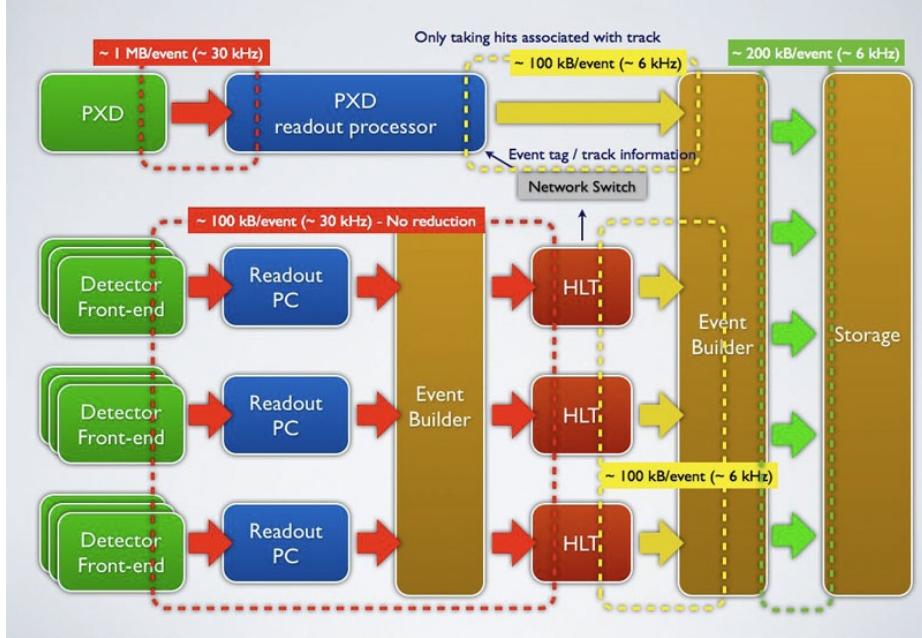


Figure 2-12: The Belle II DAQ workflow with HLT between two event builder to reduce the original 30 kHz event rate down to about 6 kHz for offline storage.

Since the primary goal of Belle II is focusing on B physics studies, it is natural that the trigger system should be able to operate over all of the interesting B physics

conditions, with normally 3 or more CDC tracks and large energy deposition in ECL. By offline reconstructing the simulated events and studying the efficiency, close to 100% B decays are recorded by Belle II trigger system. However, the extensive capabilities of studying a large range of physics not only in b sector brings a challenge to Belle II trigger and DAQ system. The control of beam background becomes essential. The main sources of beam background are beam-gas scattering, synchrotron radiation, the radioactive Bhabha scattering, the two-photon process, beam-beam effects, and Touschek effect. Their impacts depend on many factors such as beam current, luminosity and vacuum conditions, etc. One of the featured topology of these beam background events is the combination of two charged tracks in CDC and one or two clusters in ECL. The sources of the main beam backgrounds and their event rates in simulation is listed in Table 2.2.

Table 2.2: Simulated beam background rate[15]

Type	Source	Rate (MHz)
Radiative Bhabha	HER	1320
Radiative Bhabha	LER	1294
Radiative Bhabha(wide angle)	HER	40
Radiative Bhabha (wide angle)	LER	85
Touschek scattering	HER	31
Touschek scattering	LER	83
Beam–gas interactions	HER	1
Beam–gas interactions	LER	156
Two-photon QED	-	206

The improvements on both L1, HLT and the DAQ system are dedicated to serve the operation of the Belle II detector with much higher event rate in future. The further upgrades are also considered such as the replacement of the current DAQ module (COPPER board[21] on the read-out PCs) with the latest PCIe-40 platform[24] to enlarge the bandwidth between detector front-end electronics and the event builder in Figure 2-12.

2.8 Analysis software framework

The data acquired by the Belle II experiment or simulation can be processed by Belle II Analysis Software Framework, called BASF2. It has a good capability to handle multiple tasks for the Belle II data analysis, from the simulated data production to physics events reconstruction. BASF2 takes the advantage of good efficiency and reliability of C++ as the programming languages, but the use of Python is also allowed when it shows clear advantages, such as steering the analysis workflow.

2.8.1 BASF2 Core Structure

The core structure of BASF2 contains three major parts: the analysis packages required by the needs of analyzing the Belle II data such as finding tracks and combining particles, the external libraries as the third-party such as ROOT, and the tools for configuring and installing BASF2 which are mostly Python and shell scripts. Physics analysis is supported by providing a series of modules belonged to BASF2 for appropriate reconstruction based on their specific needs. To realize this, a modular analysis workflow, where each module can handle the event data through an unified method such as ROOT I/O based object persistency, is desired. Other processes, such as data summary table (DST) processing, simulation of each sub-detectors, and data skimming, are done with the packages built for sub-detectors.

The analysis packages are categorized based on the different levels of Belle II detector components, like the packages of base-level system control called “framework”, the package that provides the simulation of each sub-detectors like “svd”, the package for track reconstruction called “tracking”, and the one for post-reconstruction data analysis called “analysis”, etc. Users can work either with compiled binary version of BASF2 installed centrally on working servers, or build from the source based on their own need. Furthermore, the distributed computing is also supported by the installations of BASF2 through the management service provided by DIRAC system[25]. The detail information about the core structure of BASF2 can be found here[26].

2.8.2 Event processing workflow

The data from Belle II detector or from the simulation, are organized into a set of runs that are defined by either experimental conditions or simulation conditions. For instance, the simulation data from a certain detectors' condition are packed together, marked with the conditions' database index that is used during the simulation. Such data sample then is divided into different runs based on estimated luminosity from experiment, which can contain the different number of events in each run. This scheme is used for categorizing experimental data as well, so that users can easily know which experiment conditions are used. Thus, when BASF2 processes a data set, the functions are called for every event based on different configurations that are corresponding to the different experiment conditions. For example, in a data set where events are recorded with the different magnetic fields, BASF2 can automatically change the configurations of the magnetic fields event-by-event to provide a better track measurement. Based on this idea, all BASF2 functions (called “modules”) are developed based on a python module class which contains following embedded functions to be called at event-based level:

- initialize: called at the start of processing a event to properly set up constants needed for this run.
- beginRun: called at the start of calling this module, including setting up database conditions used in this run (run-dependent configurations) or event (event-based configurations).
- event: called for each event. This is the actual processing step, such as perform tracking or combining all daughters to find a mother particle.
- endRun: called at the end of a run, usually to register all processed information to the storage, such as physics variables from all reconstructed particles.
- terminate: called at the end of the processing of all events, release the buffered space and memory.

BASF2 executes a series of modules loaded dynamically to process the data set for analysis purposes, which is shown as Figure 2-13. The selection, configuration and

executed order of the modules are defined by a file called “steering file” written in Python. The modules parameters are attributes which can be set during the runtime using the steering file. For example, the “Path” object declared in a steering file stores the sequence of modules that will be executed, to which allow other modules such as “mdstInput” or “reconstructDecay” to be added. Users can use “boolean” type variable set in “event” function to create a conditional branch of a “Path” in case that one event needs to be processed with different modules at the same time. For instance, in the decay reconstruction package, if a decay chain is not fulfilled by missing one particle in the “event” functions, other back-up decay chains can be checked to see if a successful reconstruction is possible.

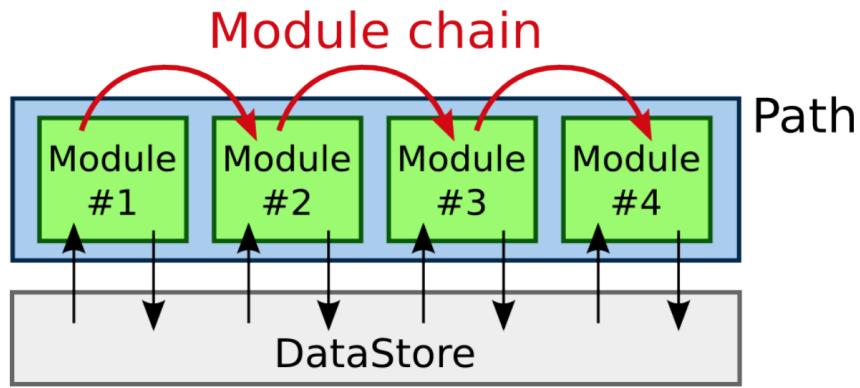


Figure 2-13: The module-based analysis workflow in BASF2.

The object that interacts with BASF2 I/O is called “DataStore”, as shown in Figure 2-13. This implementation doesn’t depend on the event data model. The only mandatory component is called “EventMetaData” which presents the experiment, run and event number of a event. “Unpacker” module converts the raw digits into digits-based object in BASF2. In simulation, digitization is done by module called “digitizer”. The digits-based objects are further processed to form hits or clusters depending on detector types. Higher level functions such as tracking and decay reconstructions are implemented based on these basic information by their packages. Eventually, BASF2 writes out the information based on users’ needs, like kinematics

variables, to ROOT[27] format files, or simply prints out processing statistics to the standard output. In practice, BASF2 starts running when it checks there is at least one module specifying the number of events to be processed in a “path” from the “steering file”, then it reads in the information from DataStore in the input ROOT file, execute all the requested modules in the “steering file” and return the time and number of events as information printed in standard output.

2.8.3 mDST structure

The output of BASF2 processing from the online data contains several detector-specific objects, which are restored as mini data summary table (mDST) type ROOT file. For a mDST level analysis, the goal is usually aimed to find particles from physics processes and reconstruct decay information. A output mDST ROOT file contains the reconstructed objects from each sub-detectors, and the following items are required for $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ analysis.

- Track: object presenting any charged particle trajectory. It’s linked to multiple track fit results using different nominal mass hypotheses as well as their track fit quality to help select good tracks.
- TrackFitResult: the fitting result of tracks with different mass hypotheses. It consists of five helix parameters, their covariance matrix and p-value from the fit. It also stores the information of hit pattern on VXD and CDC.
- V0: object for the relative long-lived neutral particles that fly out of interaction region but mostly decay or interact inside detector region. In Belle II, these are mostly K_S^0 , Λ and photon converted to electron pairs. V0 also stores their relation to the charged daughter tracks and track fit results for further selections.
- PIDLikelihood: it presents for the possibility of a charged track to be an electron, muon, charged kaon and pion, proton and deuteron provided by particle identification system.
- MCParticle: simulated particles and particle-detectors relations are created if simulated particles are correctly reconstructed as tracks or clusters.

2.8.4 Conditional Database

In addition to the physics data, analysis relies on various conditional data that are different calibration of detector, weight files for multi-variate analysis usage like PID and so on. This data is stored in a central database server called central Conditional Database (CDB)[28].

Conditions are made of payloads and each payload has its own “Intervals of Validity” (IoV). It defines in which runs the payload is valid. A set of payloads and IoVs are called a global tag (GT). Considered the GT that is required by the different analysis purposes may change even though the experiment condition is still same, GT is subjected to be updated once new calibrations of detectors or weight files for analysis tools are available.

Users instantiate one of the creation classes, add objects to them and commit them to the configured database with a user-supplied IoV. This includes the support for run dependency as well. The capability to use a local file-based database allows for easy preparation and validation of new payloads before they are uploaded to the CDB. Only the creator of the CDB objects has the right to add, recall, replace and remove the GT from CDB, which guarantees the stability.

2.9 Belle II simulation

This section briefly describes simulation (MC) used in the studies presented in this thesis. As for the focus of this analysis is in b sector which is mainly from $\Upsilon(4S)$ decay, the discussed simulation is based on the collisions with center-of-mass (CMS) energy at $\sqrt{s} = 10.58$ GeV.

In the previous section, it's shown that external packages and functionalities have been integrated with BASF2, including the core components of Belle II simulation in B decay: *evtgen* as event generator[29] and *GEANT4* as the simulator of detectors[30]. For the simulation and the reconstruction used in this analysis, the latest release of BASF2 was used. Thanks to the CDB management, BASF2 can utilize the same constants such as the magnetic field distribution for the consistence

between simulation and reconstruction.

All simulations start with at least one event generator that configures the physics processes. The *evtgen* requires a decay file that describes the decay chain from a certain mother particle, branching fraction for all processes and decay-related information such as flavor mixing or *CP* violation information. MC sample is centrally produced using Belle II grid computing service by DIRAC system and skimmed, of which the output is for physics analysis to create ROOT files. Each round of MC sample is packed and marked by their production index, such as *MC13*, which is the latest MC sample with improvements in PID. In the following content of this thesis, all MC samples are produced in *MC13* if not specifically stated.

For the analysis in this thesis, there are two MC samples, one is called *signal MC* and the other is called *generic MC*. *Signal MC*, as its name suggests, is the MC sample that describes the whole decay chain of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$. The mother particle of the decay chain is $\Upsilon(4S)$, then it decays into a pair of $B^0 - \bar{B}^0$ at branching fraction of 100%, with the model *EvtVSSMix*[29] describing the decay model. Then, one of the B meson is set to decay into three K_S^0 based on phase-space model (“PHSP”) at 100% branching fraction. The default configuration of *evtgen* can not handle multi-bodies charmless B decay with TDCPV. A modified decay model profile is under-development and not fully validated yet. Thus, MC sample of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ yields zero *CP* parameters by default. As for the other B meson, it decays into all possible final states that are described by Belle II generic decay file.

As for *generic MC*, all hadronic processes in a $\sqrt{s} = 10.58$ GeV collision are simulated. The total production cross section receives contributions from not only $\Upsilon(4S)$ (b -flavor decay dominated), but also u, d, s, c . Their relative branching fractions are taken from cross sections at $\sqrt{s} = 10.58$ GeV as shown in Table 2.3. *Generic MC* sample contains 6 types of MC samples due to this production arrangement, where $\Upsilon(4S)$ produces *mixed* (neutral) and *charged* B meson pairs and the rest are *uubar*, *ddbar*, *ccbar* and *ssbar*, respectively. In this thesis, the latter 4 types of MC samples are combined and called *qqbar* for simplicity. In the mixed MC sample, the branching fraction of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ is set at 6×10^{-6} and the branching fraction of $K_S^0 \rightarrow \pi^+ \pi^-$

Table 2.3: Production cross section for different hadronic flavors from collision at $\sqrt{s} = 10.58$ GeV used in Belle II *generic MC*[15].

Processes	$\Upsilon(4S)$	$u\bar{u}(\gamma)$	$d\bar{d}(\gamma)$	$s\bar{s}(\gamma)$	$c\bar{c}(\gamma)$
Cross section [nb]	1.110 ± 0.008	1.61	0.40	0.38	1.30

is set at 0.692. Both values are taken from Particle Data Group (PDG)[31]. Same as *signal MC*, CP violation is set to zero for signal events in *generic MC* since they use the same model at generator level.

In addition to the simulation of physics processes, MC data is produced with at least two beam background conditions, called *BG0* without beam background and *BG1* with one overlay of beam background. The components of them have been discussed briefly in section 2.7. The mixing of simulated beam background to simulated physics events is done by adding simulated hits on each sub-detector output. Possible pile-up of hits is therefore inherently included. The average number of background events of a given type to be added to a single simulated event is determined from the rate R_{BG} of beam background sample and the time window Δt in which the background is mixed shown in Equation 2.1:

$$\bar{N} = sR_{BG}\Delta t \quad (2.1)$$

where s is an optional scaling factor. The injected background events are based on a Poisson distribution with mean \bar{N} . Within the timing window, the background events are shifted randomly to simulate contributions from different bunches. To use real experiment background events (data-based beam background), the random triggered events are measured and added to simulated BG0 MC sample for a more precise background configuration. This method can give a more realistic description of actual beam background but with a possibility to introduce bias due to the pile-up effect of multiple background events in a short timing window. In the early stage of Belle II, the level of background is not high and the background pile-up effect is small.

In total, there are 2 million events generated in *signal MC*. Half of the *signal MC* (1

Table 2.4: MC samples with and without beam background used in $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ analysis.

Events number	BG0	BG1
<i>signal MC</i>	10^6	10^6
<i>generic MC</i>	None	1 ab^{-1}

million) is produced without beam background for cross-checking the reconstruction performance. For *generic MC*, 1 ab^{-1} sample including mixed, charged and $q\bar{q}$ events are produced with beam background at $\sqrt{s} = 10.58$ GeV. The MC sample used in this analysis is summarized in Table 2.4

2.10 Belle II data taking

The Belle II beam test operation started in 2016 which was focused on the commissioning and test of the SuperKEKB accelerator. Later in 2018, the commissioning of the Belle II detector was accomplished, with partial installation of PXD and full installation of SVD. From 2019 April, the Phase 3 operation that marked the beginning of official physics runs has started. By the end of 2020, Belle II has been operating in Phase III for 4 total run seasons. The integrated luminosity collected during this period of time is about 84.73 fb^{-1} , shown in Figure 2-14. The data processing is regularly performed along with the data taking. For the analysis reported in this thesis, the experimental data collection from experiment number 7, 8, 10 and 12 is used. The data of experiment 7, 8 and 10 were collected during 2019 and the data of experiment 12 and 14 were collected during 2020, as shown in Figure 2-14. Correspondingly, the integrated luminosity for offline reconstruction that is used for this thesis is about 62.8 fb^{-1} [32] which is from experiment 7, 8, 10 and 12. The experiment 14 is not used due to the unfinished processing of the latest experiment data.

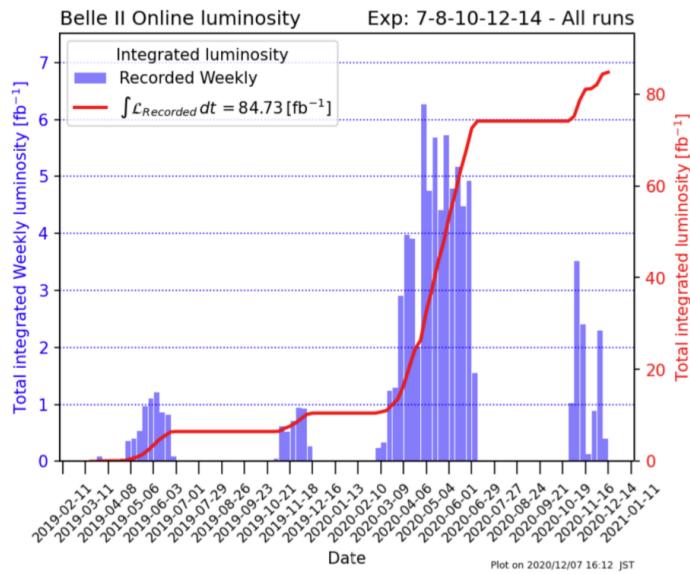


Figure 2-14: Belle II online luminosity from 2019 April to the end of 2020[32].

Chapter 3

K_S^0 reconstruction study

The final states of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ only depends on the decay of K_S^0 . The main decay channels of K_S^0 is to either $\pi^+ \pi^-$ at branching fraction of about 0.692, or to $\pi^0 \pi^0$ at branching fraction of 0.307, referenced from PDG[31]. The characteristics of these two decays are much different in terms of the response from the Belle II detector. The charged decay that yields $\pi^+ \pi^-$ leaves two tracks originating from VXD or CDC volumes with opposite charges. On the other hand, the π^0 main decay channel is $\pi^0 \rightarrow \gamma\gamma$ which typically results in the photon clusters on the ECL. There are mainly two reasons for not selecting π^0 as final states. First, $\pi^0 \rightarrow \gamma\gamma$ can yield a large fraction of fake K_S^0 . The reconstruction of two photons using ECL clusters provides no constrain on K_S^0 vertex so it's almost impossible to suppress the combinatorial background using vertexing quality in this case. The photons could be originating from many other resources, such as beam background and charged track radiation. Besides, the main viable selection is the mass of K_S^0 which is typically distributed around its nominal mass with a few hundred of keV. However, using the mass window of K_S^0 could not effectively reject the noticeable fraction of fake K_S^0 , especially when using photons. Second, B^0 that decays to one or more K_S^0 reconstructed from neutral pions have poorly reconstructed vertices. Even with $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ which only uses K_S^0 from charged pions in the final states, there is no direct charged tracks from IP, which leads to the worse resolution of vertex position compared to the channel like a $B^0 \rightarrow J/\psi K_S^0$, which has two direct charged tracks of $e^+ e^-$ or $\mu^+ \mu^-$ from J/ψ . If

one (or more) of K_S^0 has the poor vertexing quality from its decay products, it can further reduce the precision of vertex positions of B^0 , which eventually leads to a large uncertainties in defining the decay time of signal B^0 and the decay time difference as the key observables of TDCPV measurement. Therefore, only K_S^0 reconstructed using charged pions is considered to reconstructed B^0 in this analysis.

3.1 Cut-based K_S^0 Reconstruction

The K_S^0 has average life time at $(8.954 \pm 0.004) \times 10^{-11}$ s in PDG. Therefore, the flight length of K_S^0 is comparable with the scale of VXD size. In the Belle II energy scale, the flight length of K_S^0 is in a range from a few μm away from B vertex to more than 13.5 cm that is further than the outmost layer of SVD ladders, see Figure 3-1.

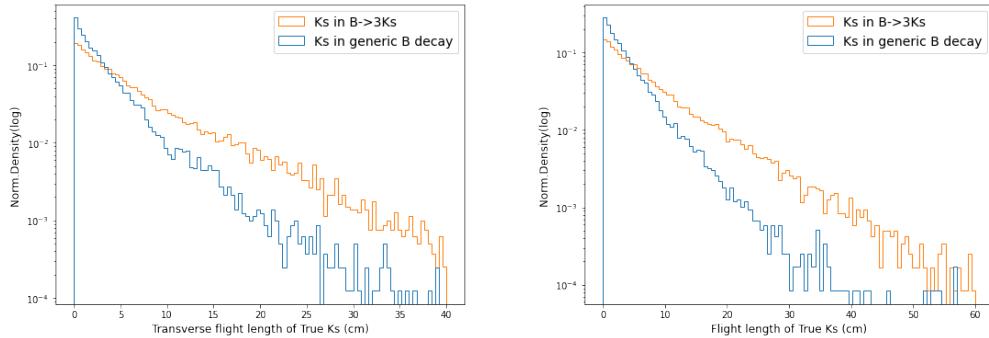


Figure 3-1: The left is the transverse flight length distribution and the right is the total flight length distribution from true K_S^0 . The blue is from *generic MC* and the orange is from *signal MC*. Both plots are normalized.

Due to the different topology of B^0 decay, the average momentum of K_S^0 in *generic MC* is different from the ones from the *signal MC*. In general, the cut-based reconstruction for K_S^0 is first performed by the selection of invariant mass from its decay products. After the selection on invariant mass is applied, a vertex fit for each K_S^0 using two reconstructed charged pions is done without IP constraint. This reconstruction is mainly achieved by using standard BASF2 particle list, in which two K_S^0 collections are first reconstructed and then merged. We first take all the V0 objects from BASF2 which use 2 online reconstructed charged tracks with opposite charges and a converged fitted vertex. In this step, charged tracks with mass hypothesis of π^\pm

are used, which the tracks and PID of charged pions are pre-selected by the criteria in Table 3.1. The K_S^0 candidates with invariant mass M between $0.45 < M < 0.55$ GeV are selected. In addition to these K_S^0 from V0 objects, another K_S^0 collection from offline reconstruction is also formed. The V0 based K_S^0 and offline reconstructed K_S^0 are merged and the vertex fit is performed using *TreeFit*[33]. The duplication of K_S^0 between two K_S^0 collections are possible so that the objects' index of two charged pions' tracks in BASF2 are compared, from which the identical combinations are removed to avoid duplication. The B^0 reconstruction efficiency is highly sensitive to the efficiency of charged pions because the final state particles are three identical K_S^0 decaying to six charged pions. That's why a very loose selection on π^\pm is applied. The selected K_S^0 collection using cut-based method contains many fake candidates from *signal MC* as shown in Figure 3-2.

Table 3.1: Pre-selection criteria of $\pi^+\pi^-$ for K_S^0 reconstruction.

Selection Criteria	θ	CDC Hits Number	PID
CDC acceptance		> 20	pionID > 0.1

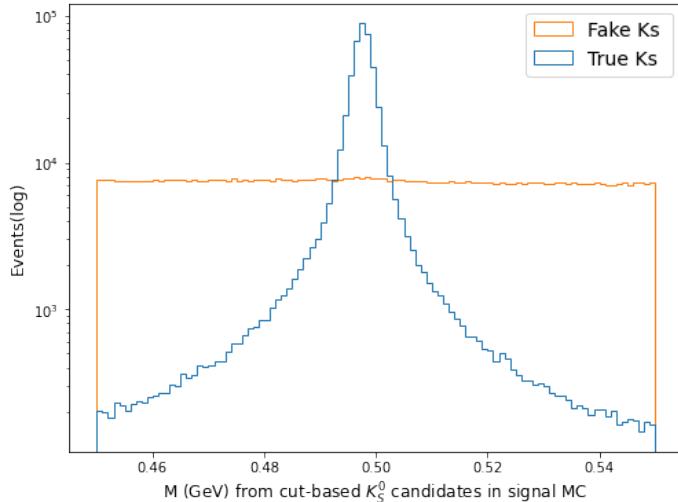


Figure 3-2: M of K_S^0 from cut-based selection in *signal MC*. The blue line is the true K_S^0 and the orange is the fake K_S^0 . 200000 candidates are used in total.

The reconstruction quality of K_S^0 also depends on the flight distance. K_S^0 that decay in the inner region of VXD yields more hits on the SVD layers from its charged

daughters, which is critical in performing a proper tracking. Belle II tracking efficiency gets poor due to higher beam background when performing track finding for tracks without inner detector hits such as SVD. For certain fraction of K_S^0 decaying outside of layer 5 of SVD, it's much likely that there is no SVD hits assigned to their daughters' tracks. This is due to the feature of SVD track finding, where a track candidate needs either at least 3 SVD hits to form a good SVD track, or 2 hits to form a hit double-lets to be used as a complete track. Single hit on layer 5 or layer 6 is filtered out to suppress the large fraction of beam background induced by random single hits. This effect is shown in Figure 3-3. K_S^0 are categorized based on how many SVD hits their daughters have, in which $SVD10$ and $SVD01$ stands for K_S^0 that only π^+ and π^- has non-zero SVD hits number, $SVD11$ and $SVD00$ stands for K_S^0 that both or neither charged pions have SVD hit non-zero SVD hits number. This is related to the track quality of K_S^0 where $SVD11$ K_S^0 has the best quality and $SVD00$ has the worst. Thus, the efficiency and purity of K_S^0 with long flight length is reduced. It's clear that $SVD00$ K_S^0 show up at about 11 cm where SVD layer 5 is placed. Most of $SVD10(SVD01)$ K_S^0 start to show up at the similar range. The geometric structure of PXD and SVD is shown in Figure 2-3 and the fraction for each types of K_S^0 in $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ is listed in Table 3.2.

K_S^0 type	$SVD11$	$SVD00$	$SVD10$	$SVD01$
% in Belle II	52%	39%	5%	5%

Table 3.2: The fraction of each category of K_S^0 based on pions SVD hits in $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ signal MC.

Fake K_S^0 candidates costs a large extra processing time and the number of combinatorial backgrounds in $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ becomes high, which largely reduces the signal significance and introduce bias to the CP parameters measurement. Thus, a multi-variate analysis (MVA) based K_S^0 classification package, *KsFinder*, is developed to further reject the fake K_S^0 from cut-based selected candidates.

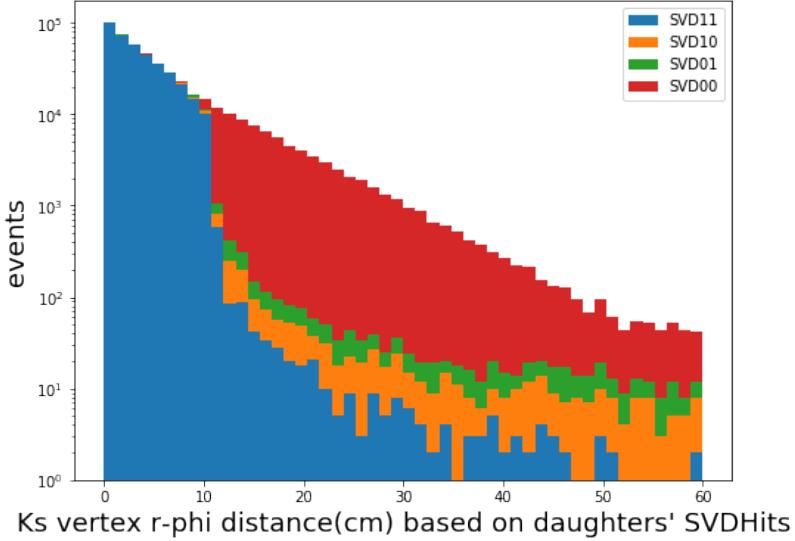


Figure 3-3: K_S^0 transverse flight length based on SVD hits of pions. $SVD11$: both pions have SVD hits, $SVD10(SVD01)$, positive(negative) pions have SVD hits, and $SVD00$: no SVD hits from pions. The result is from *signal MC* of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$.

3.2 MVA-based K_S^0 selection

3.2.1 Belle II K_S^0 classification

As mentioned in the last section, in order to improve the reconstruction performance of K_S^0 from cut-based selection, a MVA-based package called *KsFinder* has been developed. The reconstruction of K_S^0 can be treated as a typical classification problem. The input is a set of variables that describes the characteristics of $K_S^0 \rightarrow \pi^+ \pi^-$ decay. The training target is the true or fake flag from the MC truth-matching variable called *isSignal* where *isSignal* = 1 (0) stands for being a true (fake) K_S^0 . It aims to improve the limitations in the Belle K_S^0 MVA classification tool.

In Belle, the K_S^0 reconstruction was first done by using cut-based method to select primary candidates, then a MVA-based classifier was implemented by assigning two likelihood indicators to each K_S^0 candidates. The package used by Belle is called *nisKsFinder*[15] which outputs the two likelihood variables based on NeuroBayes algorithm[34]. The Belle tool defines the goodness of K_S^0 , called *nb_nolam* and *nb_vlike*, respectively. As their names suggest, *nb_nolam* is the likelihood of not being a Λ particle and *nb_vlike* is the likelihood of being a V0-like particle. A good

K_S^0 candidate from *nisKsFinder* is the one with a low likelihood of being Λ particle and a high likelihood of being a V0-like particle, assuming the major backgrounds for K_S^0 is the mis-identified Λ among V0-like particles. By putting cuts on these two variables, a purification of K_S^0 can be made, see Figure 3-4. It can effectively reduce fake K_S^0 from cut-based selected candidates, however, there are a few disadvantages about this method. First, NeuroBayes is a commercial product that was developed over 10 years ago. The official support and update is stopped nowadays, so it's not an ideal method for an experiment like Belle II that has a quite long prospective in operation. Second, the classification is based on a joint cut on two variables, which might make the cut values hard to choose, for example, two different cuts might have very close purity. Last but not least, the classification of K_S^0 is not the directly targeted output of the neuro-network. Instead, it classifies the V0-like particle and Λ . Besides, the computation speed of NeuroBayes is not optimized.

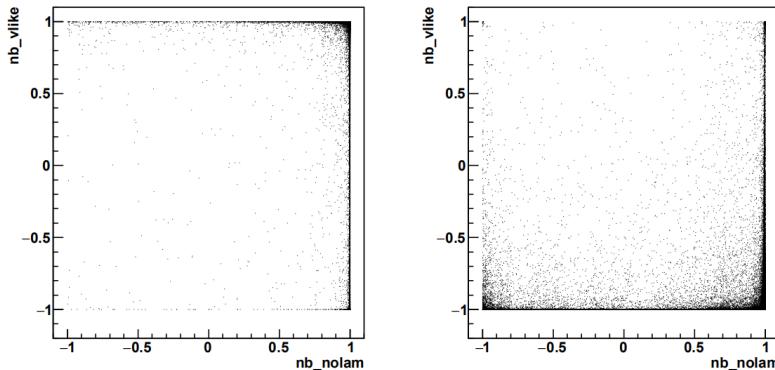


Figure 3-4: The distribution of two variables outputs: nb_nolam and nb_vlike for K_S^0 candidates from Belle *signal MC*. The left is from true K_S^0 and the right is from the fake K_S^0 . In Belle, the standard cuts for K_S^0 is $nb_vlike > 0.5$ and $nb_nolam > -0.4$ [19].

Such a dedicated K_S^0 classification tool is not implemented yet in BASF2 framework until 2019. Considered the limitation of NeuroBayes, the development of K_S^0 classifier demands another algorithm and structure. The *Boosted Decision Trees* (BDT) is widely employed for multivariate classification and regression tasks in high energy physics field. Particularly, a speed-optimized and cache-friendly implementation of such a method called FastBDT (FBDT) is popularly used[35]. Compared

to other popular classification algorithms such as TMVA[36], scikit-learn[37] and XGBoost[38], FastBDT method is proven to be one order of magnitude faster during the training and applying phases[35]. By using FastBDT algorithm, *KsFinder* in Belle II is expected to give a single output which directly presents the goodness of a candidate of being a true K_S^0 . Since the FastBDT algorithm depends on the variables that are different in signal and backgrounds, a set of training variables are selected based on K_S^0 decay topology. The K_S^0 variables used in the training of *KsFinder* might be differently distributed in different decay channels, therefore a *KsFinder* trained using MC sample from one channel may not be able to perform a good classification on the other. Thus, *KsFinder* is designed as a general package that provides a mode-dependent K_S^0 classification which mainly consists of four components: *KsFinderSampler*, *KsFinderTeacher*, *KsFinderApplier* and *KsFinderTest*. *KsFinderSampler* is a function that automatically generates training and/or testing sample from mDST files where the cut-based reconstruction is used as section 3.1. *KsFinderTeacher* is responsible for extracting variables to perform training of the FastBDT model and generate a weight file containing all the nodes information in ROOT format, which also provides a function to communicate with BASF2 CDB so that users can share or download others' weight file in their own analysis. *KsFinderApplier* can apply the weight file generated by *KsFinderTeacher* (or downloaded from BASF2 CDB) to the independent data sample and assign each K_S^0 candidate a goodness index used as a single cut value in the further analysis. *KsFinderTest* is the evaluation function that can use a test sample to check for over-training, efficiency, purity. By providing MC sample(s) from a certain decay mode(s), users can easily generate their own weight file(s) of K_S^0 classification that suits different decay modes despite K_S^0 variables distribution may be varied. Such a design largely improves the flexibility of *KsFinder* compared to Belle MVA tool which indirectly classify K_S^0 with two outputs.

3.2.2 Decay Topology of $K_S^0 \rightarrow \pi^+\pi^-$

As introduced in section 3.2.1, the first step for developing K_S^0 MVA classification is to determine the input variables for FastBDT algorithm that can represent the

decay features of K_S^0 against possible backgrounds. The remaining background of $K_S^0 \rightarrow \pi^+ \pi^-$ after the cut-based reconstruction comes from different sources, mainly including the false combination of tracks (including π^\pm misidentification), V0-like particle misidentification and self-looped tracks. For instance, a D^0/D^* from a B decaying to $K\pi$ with K misidentified as π , could give a false combination of tracks. On the other hand, it's also possible that both of two tracks are correctly identified as π^\pm but they are not from the same mother particle, or the mother is not a K_S^0 particle due to the missing of other daughters, such as $D^+ \rightarrow K_S^0(\rightarrow \pi^+ \pi^-) \pi^+$. The decay shape resembled the above cases are illustrated in Figure 3-5.

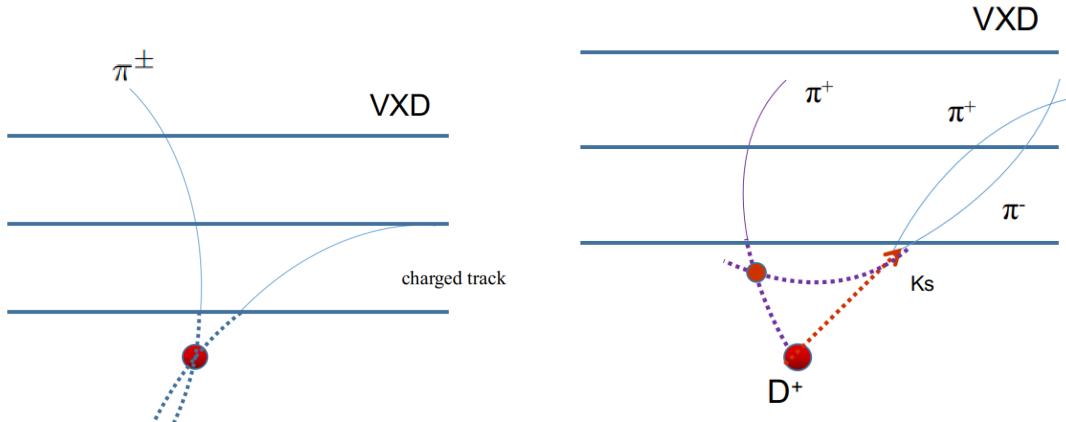


Figure 3-5: The left shows the case when a charged track (not π^\pm) combined with a charged pion to form a fake K_S^0 , the right shows the case when two daughters are correctly reconstructed as pions but not from the correct mother particle, which is falsely taken as a K_S^0 .

The V0-like particles mainly refer to K_S^0 , Λ and γ . $\gamma \rightarrow e^+ e^-$ yield is significantly lower than the other two types and the mass difference between pion and electron is very large, so the PID values can be used to well-distinguish them. As for the contribution of $\Lambda \rightarrow p^+ \pi^-$, it happens when the positive charged tracks (proton track) is wrongly identified as π^+ , see Figure 3-6 left. The key observable to distinguish this background is the invariant mass of mother particle, which is 1.115 Λ GeV, much larger than the K_S^0 . The number of left-over Λ after the cut-based reconstruction in section 3.1 is small, and can be further reduced by rejecting the candidates whose positive charged daughter has $\text{PID}(\pi^\pm)$ smaller than $\text{PID}(p)$.

When a charged pion only carries a minimal of its mother's transverse momentum p_T , the curvature of its track may form a self-loop of which radius is comparable with the size of Belle II detector (mainly VXD and CDC). In this case, one charge pion could leave two charged tracks candidates with the opposite charge and similar p_T , with a possibility to form a converged vertex to form a fake K_S^0 , see Figure 3-6 right.

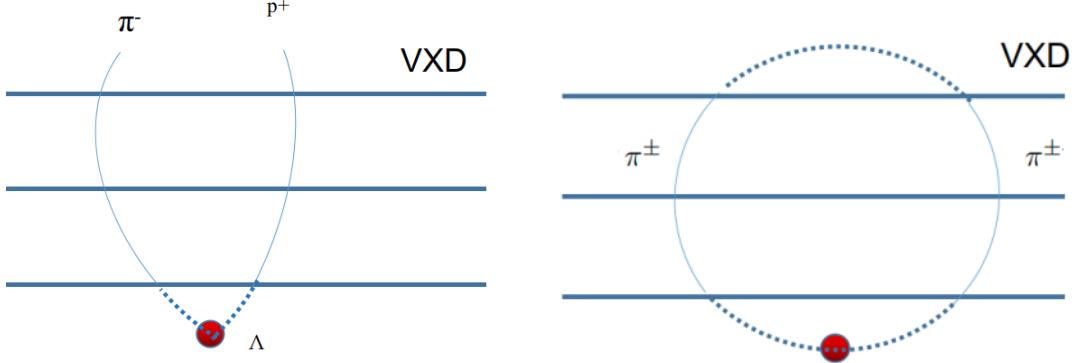


Figure 3-6: The left shows the $\Lambda \rightarrow p^+\pi^-$ decay shape that can be treated as K_S^0 , the right shows a self-loop formed by a low p_T charged pion reconstructed as two separated tracks with a vertex.

3.2.3 Determination of training observables from K_S^0 decay

Given the characteristics of $K_S^0 \rightarrow \pi^+\pi^-$ discussed in the previous section, a set of variables as training features of *KsFinder* can be selected. The set includes variables related to K_S^0 kinematics, decay shape parameters, particle identifications and detector hits information. The summarized information of training variables is listed in Table 3.3.

The cosine between K_S^0 vertex and momentum direction (named *cosVertexMomentum*) is of the most importance because it demonstrates the best separation between a true and a fake K_S^0 . For instance, if a falsely reconstructed K_S^0 is made of two tracks, it's likely that the momentum direction of the fake K_S^0 is not aligned with the its vertex direction from IP. So the projection of vertex position of K_S^0 on the reconstructed momentum direction could be negative value for fake K_S^0 . While in case of a true K_S^0 , such projection is almost always a positive value, shown in Figure

K_S^0 variables	Meaning
cosVertexMomentum	cosine between K_S^0 vertex and momentum direction (lab)
flight distance	K_S^0 flight distance projected on its momentum direction
significanceOfDistance	relative error of flight length from IP
cosHelicityAngleMomentum	cosine between π^\pm and K_S^0 (lab)
ImpactXY	Impact parameters in transverse plane for K_S^0
x, y, z, px, py, pz	K_S^0 vertex position and momentum
p_D1(D2)	momentum magnitude for $\pi^+(\pi^-)$
pionID, muonID	PID values of π^+
decayAngle_D1(D2)	angle between $\pi^+(\pi^-)$ and K_S^0 (K_S^0 CMS)
daughterAngle2body	angle between π^\pm (lab)
daughtersDeltaZ	Z-direction distance of two tracks helix
nSVDHits_D1(D2)	SVD detector hits of $\pi^+(\pi^-)$
nPXDHits_D1(D2)	PXD detector hits of $\pi^+(\pi^-)$
M, InvM	K_S^0 invariant mass before(after) vertex fit

Table 3.3: Summary of *KsFinder* input variables, where “lab” means angles in lab frame and “ K_S^0 CMS” means in K_S^0 rest frame. Other variables are calculated in lab frame by default.

3-7. This often happens when the two tracks taken as π^\pm are accidentally crossed, or due to the misidentified track(s). The abbreviations and importance rank of input variables from *KsFinderTest* function is shown in Table 3.4.

As a FastBDT method relies on the distribution of variables to calculate signal and background separation, there are a few points to be checked before feeding the training data to the algorithm or applying the classification. First, the distribution of the observables should be different in true K_S^0 and the fake ones, so the FastBDT classifier can effectively separate the true and the fake K_S^0 at each node to maximize the separation gain. Second, there will a correlation among the training observables and they should also be different in signal and background. The boosting step will create a sequence of shallow DTs whose structures are not same. Different correlations helps improve the performance of DTs in tuning of structure. For instance, a true K_S^0 flights longer due to larger momentum in general, so its daughters’ detector hits number becomes fewer. Then these two observables have negative correlations in true K_S^0 . In case a fake K_S^0 , the flight length could be a deep outside of VXD but daughters may have full hits on SVD, without strong correlation, see Figure 3-8 . At last, one should also avoid using many observables with too strong correlations, since in this

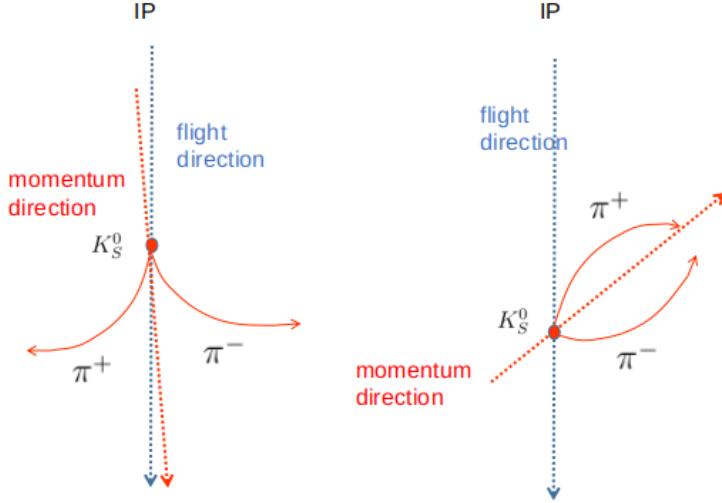


Figure 3-7: The left shows a true K_S^0 decay shape where the cosine angle of K_S^0 vertex position (blue dashed arrow) against reconstructed momentum direction (red dashed arrow) is positive. While the right shows a fake K_S^0 decay shape where cosine angle of K_S^0 vertex position against reconstructed momentum direction can be negative.

case, many DTs might have a potentially equivalent structure in the boosting step. Therefore, the separation power of many DTs doesn't gain any improvement and the collection of observables might be redundant. The correlation between variables are shown in Figure 3-8.

3.2.4 Training, Applying and Testing of *KsFinder*

The variables are internally registered inside the *KsFinder* so it can automatically retrieve their values from a mDST file in BASF2. The first step of using *KsFinder* is to call *KsFinderSampler* on a MC sample to generate training and testing data sample. To show the flexibility and stability of *KsFinder* on different modes, *KsFinderSampler* extracts MC data points from both *signal MC* and *generic MC* (see MC definition in section 2.9), respectively. *KsFinder* configures that the depth of each DT is 3 and boosting steps is 200. In both MC samples, the ratio of true and fake K_S^0 is set to 1:1 and each component contains 200000 data points. The distribution of input variables in *signal MC* is shown in Appendix A.

To train the *KsFinder*, *KsFinderTeacher* function is called for the training samples

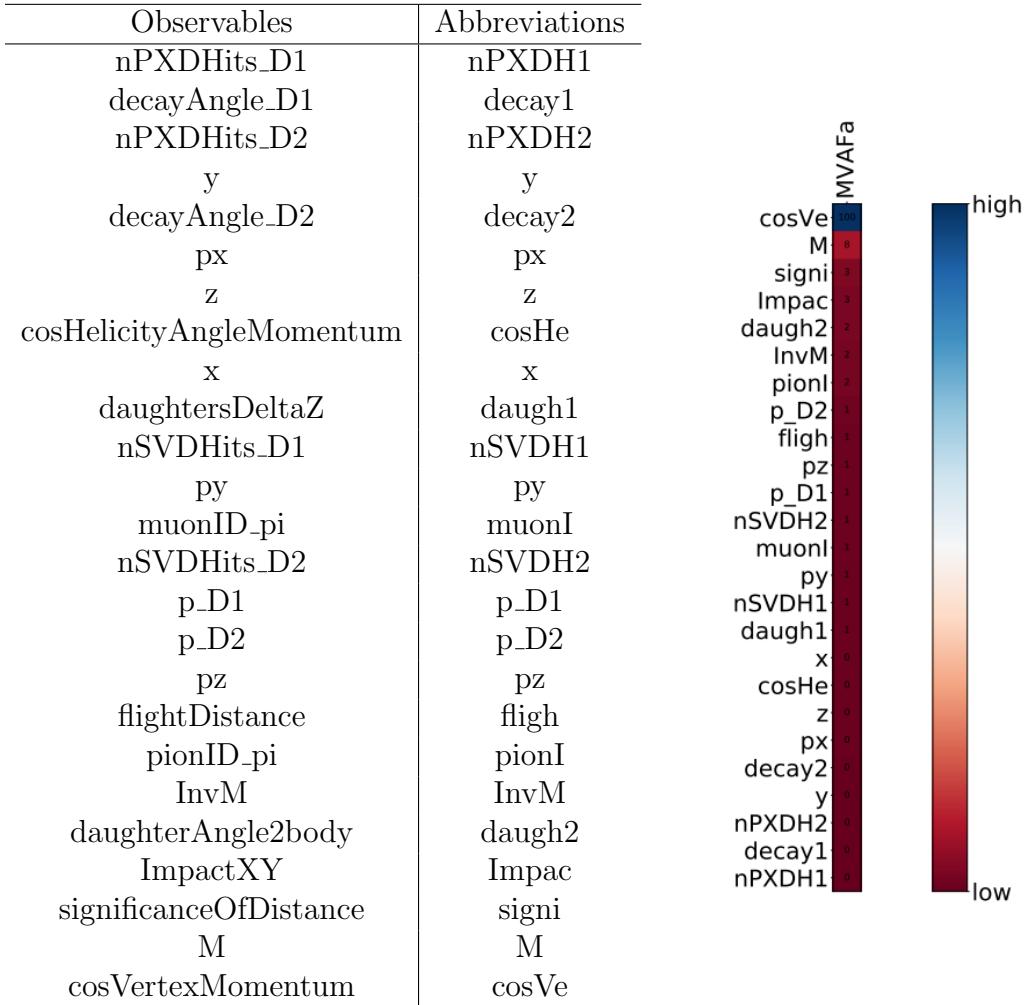


Table 3.4: The abbreviations (left) and importance rank (right) of input variables from *KsFinderTest*, where the most important variable is *cosVertexMomentum*.

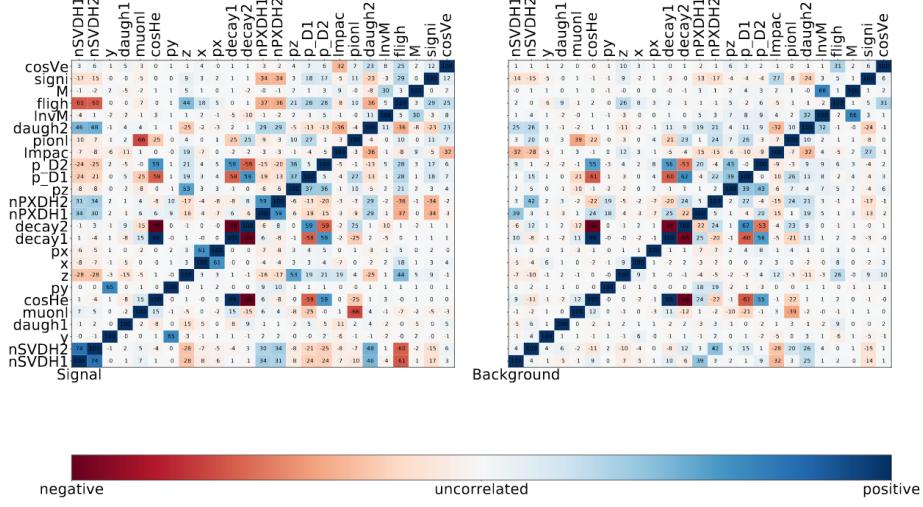


Figure 3-8: The correlation between input variables for *KsFinder*. As the given example, flight length has negative correlation with SVD hits in signal while uncorrelated in background.

from MC and weight files are saved. To apply the classification of K_S^0 , *KsFinderApplier* reads in the testing samples from MC and calculate output using saved weight files, so that each K_S^0 candidate is assigned with a goodness index named *FBDT_Ks*. It ranges from 0 to 1 where 1 stands for the best goodness. After the applying of *KsFinder* on the testing samples, *KsFinderTest* is called to check the performance and over-training of *KsFinder* on the testing samples, which will be discussed in the next section.

3.2.5 The Performance and Over-fitting check

To evaluate the performance of *KsFinder* on both signal and *generic MC* samples, signal efficiency and background rejection are calculated by cutting on the different values on *FBDT_Ks*, as defined in Equation 3.1 and 3.2.

$$\text{signal efficiency} = \frac{\text{Number of true } K_S^0 \text{ with } \text{FBDT_Ks} > \text{cut value}}{\text{Number of all true } K_S^0} \quad (3.1)$$

$$\text{background rejection} = \frac{\text{Number of fake } K_S^0 \text{ with } \text{FBDT_Ks} < \text{cut value}}{\text{Number of fake true } K_S^0} \quad (3.2)$$

The ROC (receiver operating characteristics) curve is usually taken as an indicator of the performance where the curve shows the dependence of rejection power with respect to the signal purity. The larger area under a ROC curve means that the better performance is achieved since background rejection drops slower when increasing the cut. The ROC curves and efficiency/background rejections are shown in Figure 3-9 and 3-10, where the former is for *signal MC* and the latter is for *generic MC*.

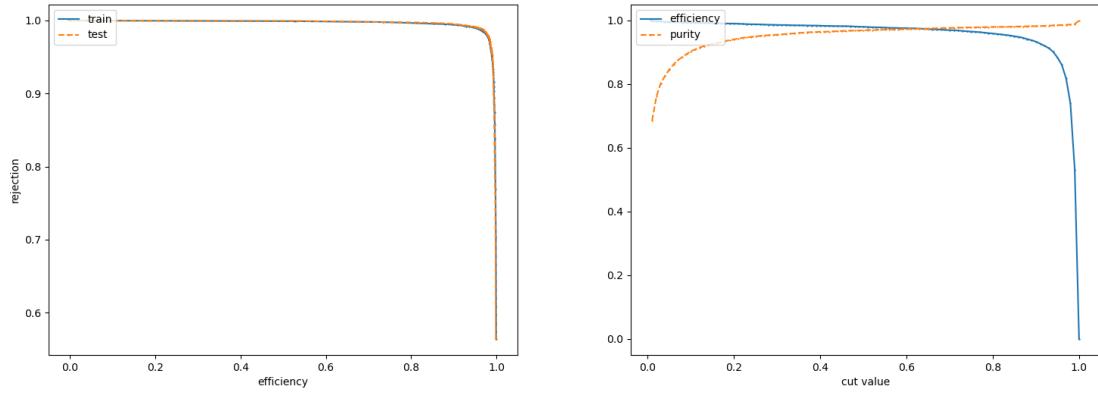


Figure 3-9: The left is ROC curve(blue for training and orange for testing) and the right is efficiency and purity (blue for efficiency and orange for purity) depending on cut of *KsFinderoutput*. Results are from *signal MC* sample.

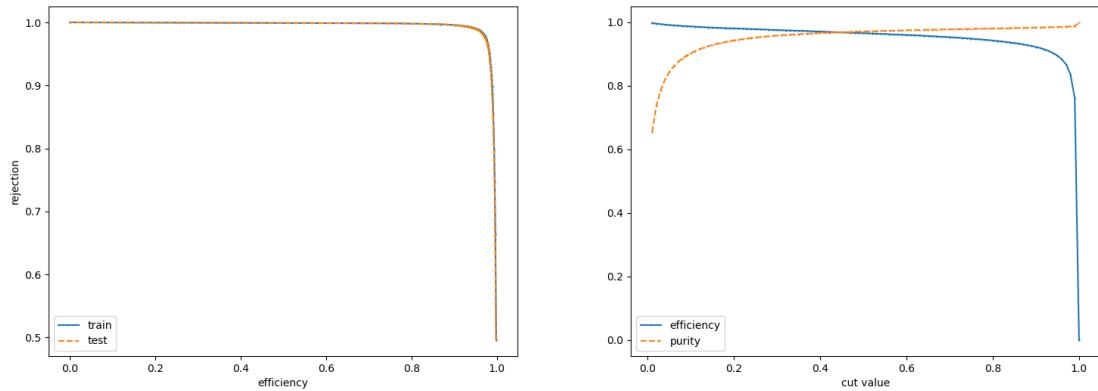
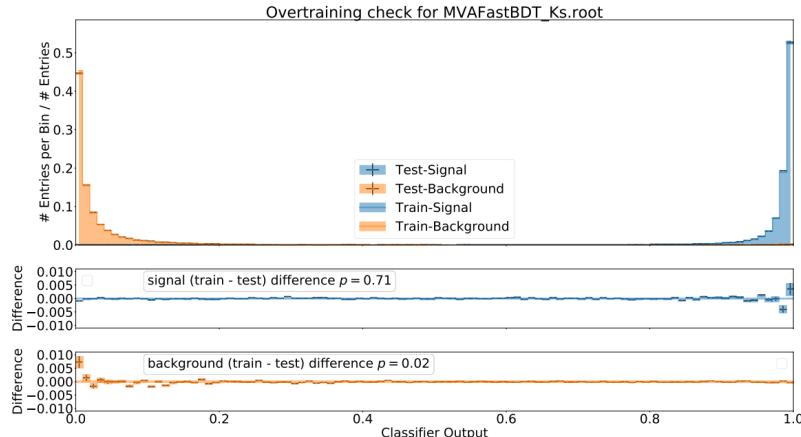
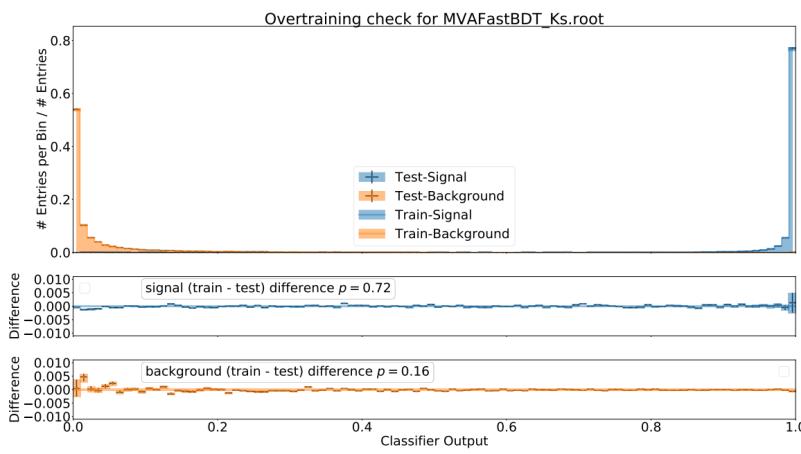


Figure 3-10: The left is ROC curve (blue for training and orange for testing) and the right is efficiency and purity (blue for efficiency and orange for purity) depending on cut of *KsFinderoutput*. Results are from *generic MC* sample.

With increasing the efficiency, the cut on the output of *KsFinder* is getting loose. The background rejection only starts to drop when the efficiency exceeds about 90% in both training and testing sample. To be noted, the curves are consistent in training and testing samples. While the ROC curve has shown the absence of noticeable overfitting in classification, the detailed check can be made by comparing the distributions of classifier output on true and fake K_S^0 in training and testing samples. Therefore, the distribution of signal and background in training and testing sample with respect to the *KsFinder* output is plotted, where a distinctive separation for both *signal MC* and *generic MC* is shown and no over-training is found, as shown in Figure 3-11.



a) Over-fitting check for *signal MC* sample.



b) Over-fitting check for *generic MC* sample.

Figure 3-11: The over-training check based on the comparison between training/testing data points in both signal and *generic MC*.

The cut value for $FBDT_Ks$ is determined by maximizing the “Figure of Merit” (FOM), as shown Equation 3.3, where S and B is the number of true and fake K_S^0 after the cut, respectively. The FOM distribution depending on the cut value of $FBDT_Ks$ is shown in Figure 3-12. The maximum FOM is achieved at $FBDT_Ks = 0.74$ in *signal MC*, which is going to be used as the cut value to further reject fake K_S^0 .

$$FOM = \frac{S}{\sqrt{S+B}} \quad (3.3)$$

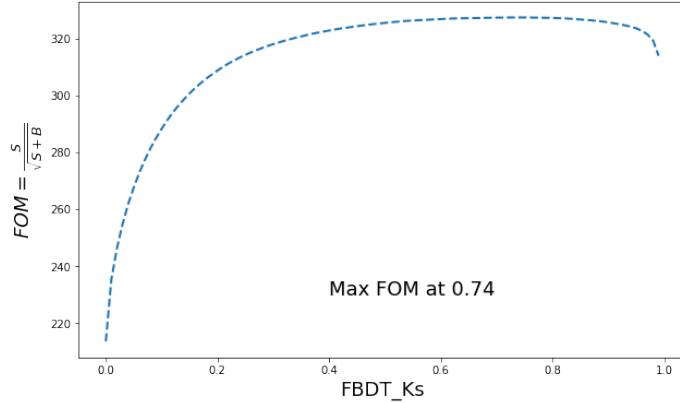


Figure 3-12: FOM of classifier output ($FBDT_Ks$) in *signal MC*, the maximum value is achieved at 0.74 and the curve is almost flat between $0.5 \sim 0.9$.

In the *signal MC* sample, the true K_S^0 fraction before applying *KsFinder* cut is 39%, and 95% of them are kept after the cut is applied. In the meantime, the fake K_S^0 fraction before applying the cut is 61%, and 98% of them are rejected after the cut is applied. The purity of the K_S^0 candidates is improved largely as shown in Figure 3-13.

3.2.6 Data Validation for *KsFinder*

The results from MC studies show an excellent performance of *KsFinder*. However, the validation of such a tool on the real experiment data is necessary. Since there’s no MC truth on target variable in real data, the FastBDT method is based on variables in MC samples. If these variables shows close distribution among MC and data, the classification performance is expected to be similar.

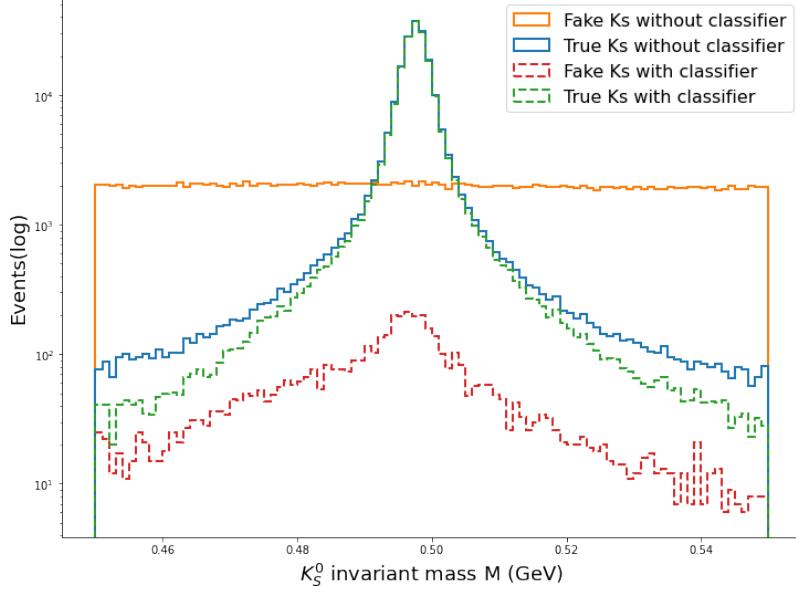


Figure 3-13: K_S^0 purity improvement with cut value of $FBDT_Ks$ at 0.74 applied. The blue solid line is true K_S^0 without $KsFinder$ and green dashed line is the true K_S^0 with the cut applied. The orange solid line is fake K_S^0 without the cut and the red dashed line is fake K_S^0 with the cut. About 95% of true K_S^0 are kept while 98% of the fake ones are rejected by applying the cut.

In addition, due to the fact that K_S^0 candidates are used for further reconstruction of B^0 , the mass and energy distributions may change after applying the cut, thus the validation that approves no clear bias on B^0 's variables that are used for signal extraction is also required. For comparison between MC and data, a small data sample from Belle II experiment 7 and 8 is used. The integral luminosity at $\Upsilon(4S)$ resonance for this data sample is about 5.17 fb^{-1} . The MC sample is extracted from *generic MC* with equivalent luminosity.

The Figure 3-14 shows the invariant mass and momentum distributions from data and MC samples, where the data and *generic MC* agree well. The uncertainties in data are calculated based on three times the Poisson standard deviation in each bin. The variable with the highest importance is the cosine angle between K_S^0 vertex and the direction of momentum, named as *cosVertexMomentum*, of which distribution is shown in Figure 3-15. In these comparison plots, the *generic MC* is shown in blue solid lines with no $KsFinder$ cut used. Similarly, data without using cut of $KsFinder$ are shown in yellow dots, which are closely distributed as the *generic MC*, indicating a

good data MC consistency. Since the FastBDT algorithm relies on the probability density functions to separate signal and backgrounds in each tree node, the similar distribution can lead to close classification power in data. The purple solid lines are presenting the true K_S^0 distribution in *generic MC*, while the red dots are the K_S^0 in data after using the cut value at 0.74. The reduced fraction of K_S^0 in data is close to the *signal MC*. All distributions comparison between data and *generic MC* by using *KsFindercut* are shown in Appendix B.

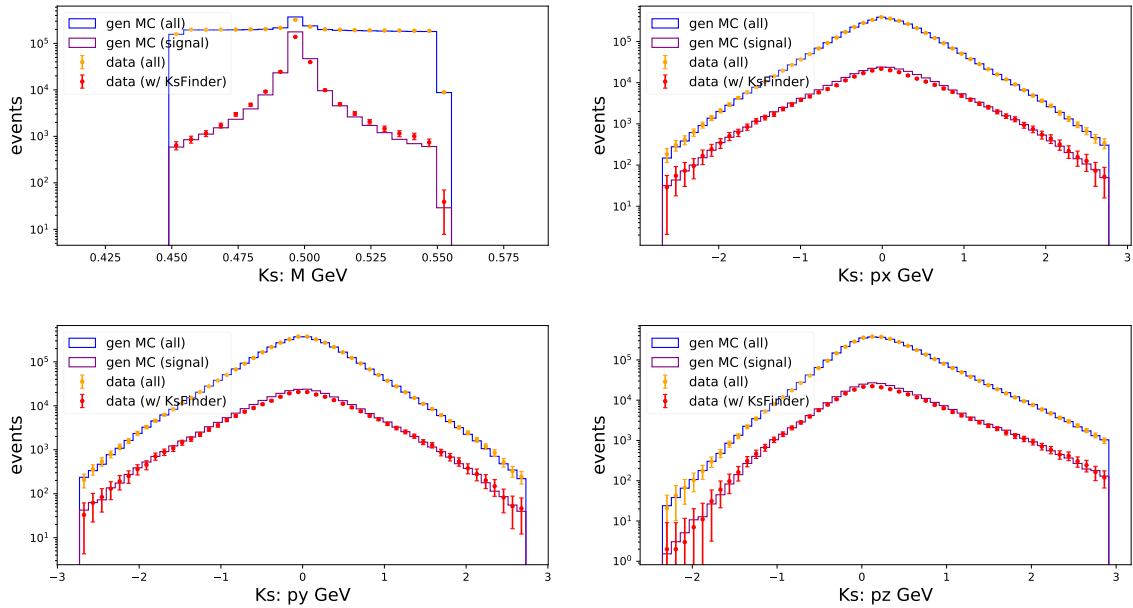


Figure 3-14: The distribution of invariant mass from charged pions and the momentum of K_S^0 in x, y, z directions. The blue line is from all *generic MC* where the purple line is the true K_S^0 in it. The yellow dots are data with no *KsFindercut* applied and the solid red dots are data after applying *KsFindercut* at 0.74.

3.2.7 Data and MC correction by *KsFinder*

Implementing *KsFindercut* on data may induce bias on the event numbers for K_S^0 because the training set of *KsFinder* is extracted from MC. To compensate such potential effect, a ratio as the data and MC correction is calculated based on the expected signal yield after using *KsFinder*. A maximum likelihood fit on invariant mass $M_{\pi^+\pi^-}$ for K_S^0 after applying cut value at 0.74 of *KsFinder* output is performed, where sig-

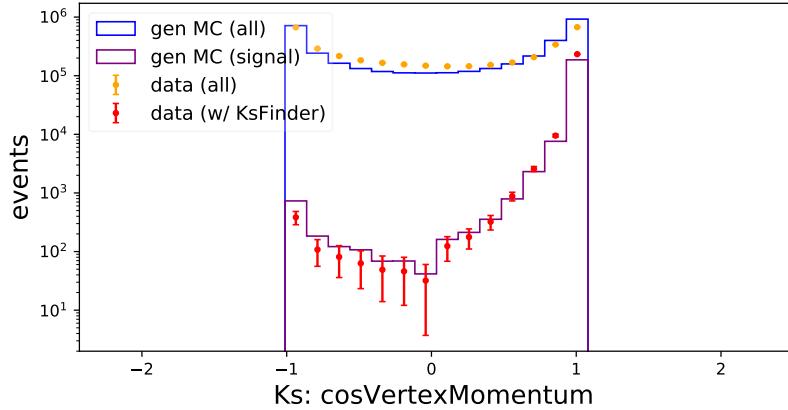


Figure 3-15: The distribution of $\cos\text{VertexMomentum}$ in data and MC with/without $KsFinder$ cut applied.

nal shape is modeled as a triple-Gaussian and background shape is modeled as a Chebyshev polynomial. The signal yield fraction is defined as Equation 3.4.

$$f_{K_S} = \frac{N_{sig}}{N_{tot}} \quad (3.4)$$

where N_{sig} is the signal number from the fit result and N_{tot} is the total events number. The fit is performed on both *generic MC* and data to obtain f_{K_S} , respectively. The \mathcal{R}_{K_S} is defined as the ratio of signal yield fraction f_{K_S} from MC and data as shown in Equation 3.5. The ratio at cut value of 0.74 is $\mathcal{R}_{K_S} = 1.009 \pm 0.011$ from the fit, as shown in Figure 3-16. Since the final state consists of three K_S^0 , the expected ratio between MC and data for B^0 is expected to be the cube of \mathcal{R}_{K_S} , with the uncertainty propagated from the uncertainty of \mathcal{R}_{K_S} . Therefore, the upper and lower limit for the correction ratio \mathcal{R}_{B^0} is expected to be 1.060 and 0.994. The result of \mathcal{R}_{B^0} is 1.027 ± 0.033 , which is close to 1 within its uncertainty. Hence, the correction \mathcal{R}_{B^0} is not applied in signal extraction of B^0 , but the uncertainty is taken into account as a possible systematic uncertainty term.

$$\mathcal{R}_{K_S} = \frac{f_{K_S}^{MC}}{f_{K_S}^{data}} \quad (3.5)$$

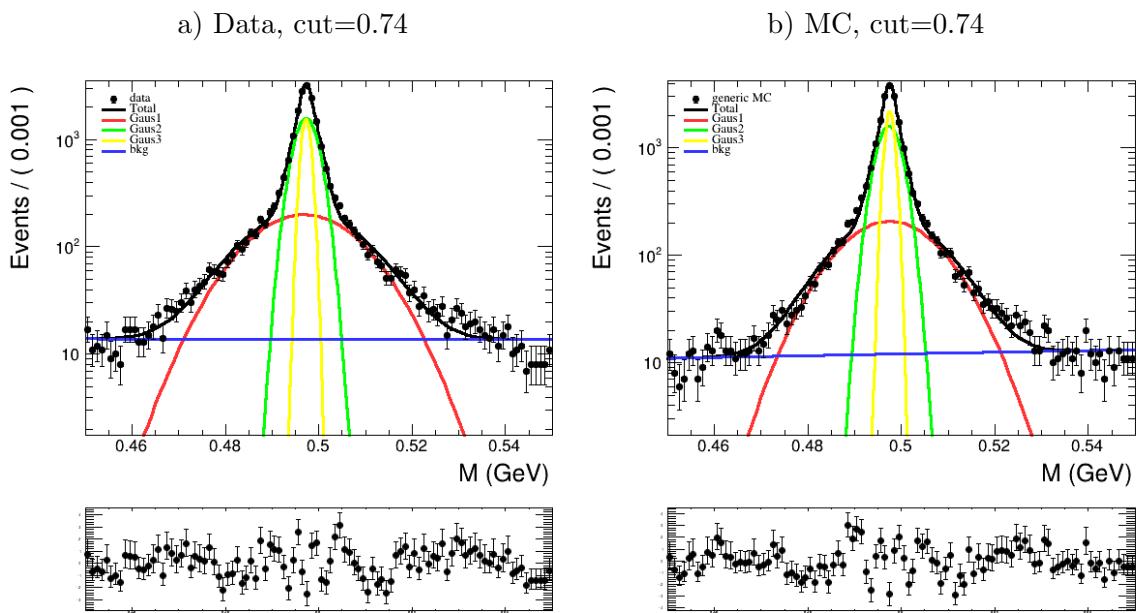


Figure 3-16: The fit on invariant mass $M_{\pi^+\pi^-}$ where signal component is modeled as a triple-Gaussian and background component is modeled as a Chebyshev polynomial. The signal fraction is slight higher in MC compared to that in data.

Chapter 4

Analysis strategies

4.1 Data Sample and Event Selection

As introduced in section 2.9, the branching fraction of $\mathcal{B}(B^0 \rightarrow K_S^0 K_S^0 K_S^0) = 6.0 \times 10^{-6}$. The simulation takes the $\Upsilon(4S)$ as the mother particle and generate its decay process to two scalar B^0 mesons with mixing. $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ is simulated based on only the possible phase-space of kinematics that final states could have, which no CP violation is assigned for $\mathcal{S}(\sin 2\phi_1)$ and \mathcal{A} in both *signal MC* and *generic MC*.

4.1.1 K_S^0 Selection

K_S^0 reconstruction is first reconstructed by the cut-based reconstruction which contains a large fraction of fake candidates, as discussed in chapter 3. In addition, a momentum cut on K_S^0 is used considering distribution in $B^0 \rightarrow K_S^0 K_S^0 K_S^0$. Only the K_S^0 candidates with momentum larger than 0.05 GeV are selected, as shown in Figure 4-1. Next, the vertex fit of K_S^0 is performed using *TreeFit* package[39]. Using fitted K_S^0 momentum and energy, fitted invariant mass is different from the one obtained directly using daughters' 4-vector. This quantity often receives impact of the measurement uncertainties. If we check out the distribution of fitted invariant mass based on daughters' SVD hits as section 3.1 discussed, *SVD00* K_S^0 shows a large dispersion from the central region of the distribution of invariant mass after vertex

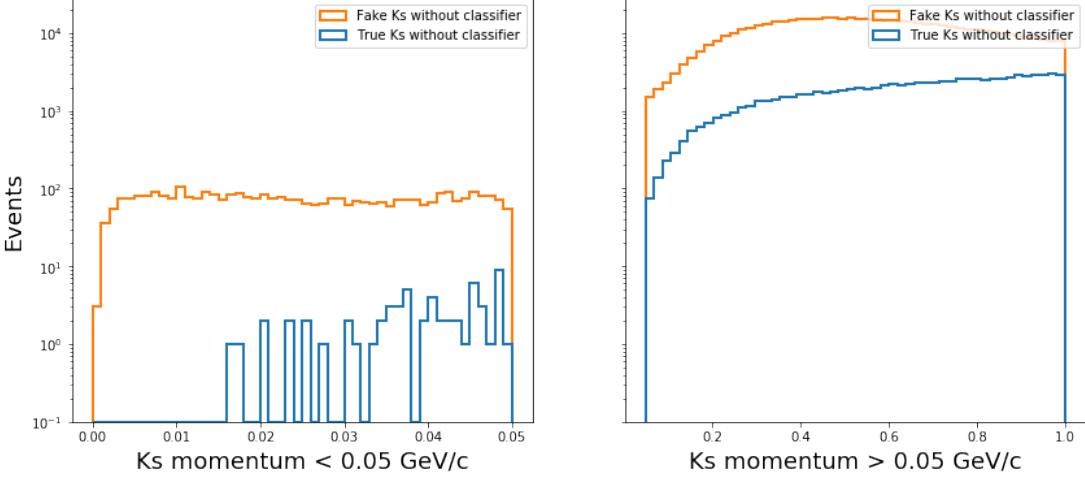


Figure 4-1: The distribution of K_S^0 momentum. Candidates smaller than 0.05 GeV/c are rejected. (Both plots shares the same y-axis scale on the left side.)

fit, while *SVD11* K_S^0 shows a small dispersion as shown in Figure 4-2. Therefore, considering the K_S^0 candidates with different SVD hits, see Figure 3-3 and Table 3.2, reflecting the different tracking quality of daughter π^\pm , the different cut on invariant mass $M_{\pi^+\pi^-}$ after vertex fit is applied depending on the SVD hits number of pion tracks. As shown in Figure 4-3, the sideband regions where fake K_S^0 is much higher than true K_S^0 are excluded. The cut windows are listed in Table 4.1. This improves the reconstruction purity.

In summary, K_S^0 are first reconstructed using BASF2 standard library with all converged vertex fit candidates are kept using *TreeFit*. Then we reject K_S^0 candidates with momentum smaller than 0.05 GeV/c. After this, tuned cuts on invariant mass after vertex fit for each K_S^0 based on their SVD hits of pions are applied according to Table 4.1. Eventually, only K_S^0 with *FBDT_Ks* larger than 0.74 are kept based on Figure 3-12.

K_S^0 type	SVD11	SVD10	SVD01	SVD00
$M_{\pi^+\pi^-}$ window (GeV)	(0.45,0.55)	(0.38,0.7)	(0.38,0.7)	(0.3,0.7)

Table 4.1: Invariant mass cut window after vertex fit for K_S^0 based on Figure 4-3.

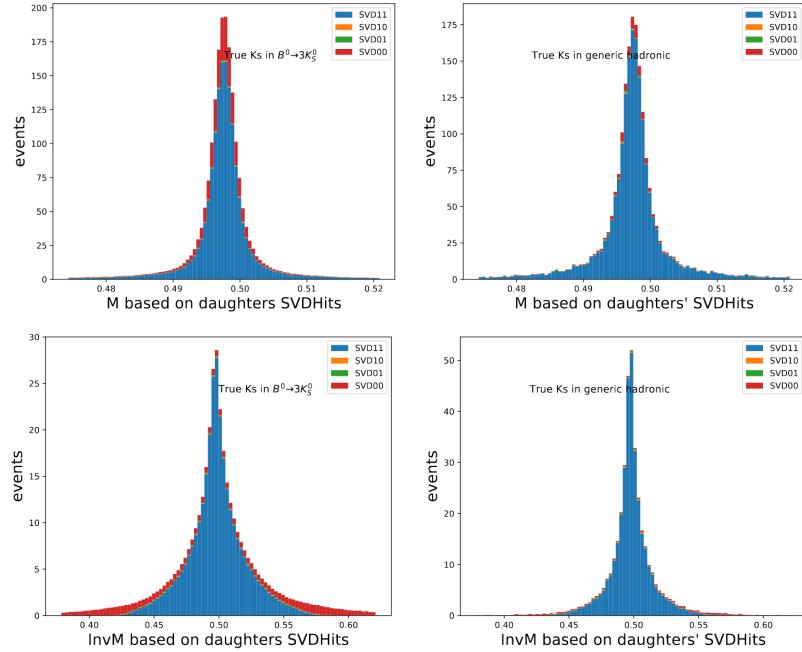


Figure 4-2: (Top two plots)The invariant mass from daughters' 4-vector, the left is $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ signal MC and the right is generic MC; (bottom two plots)The invariant mass after vertex fit, the left is $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ signal MC and the right is generic MC; In both cases, the red shows the clear dispersion on SVD00 type K_S^0 .

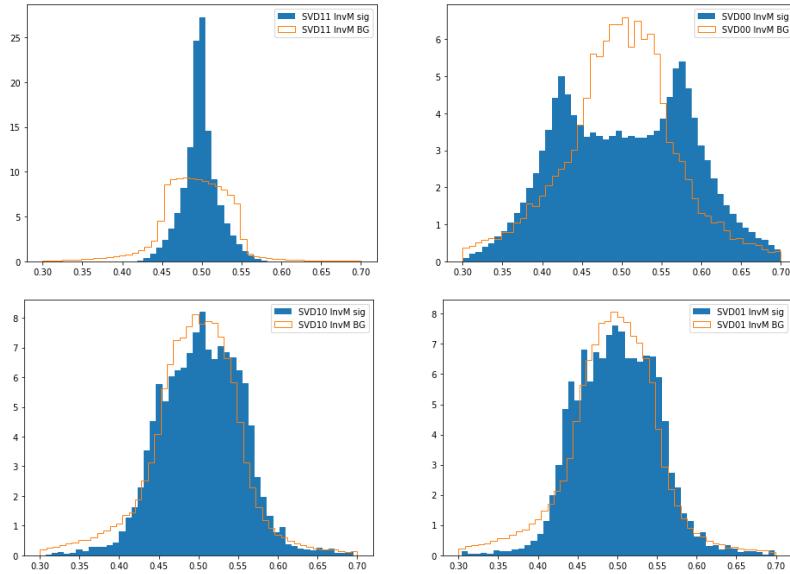


Figure 4-3: Invariant mass after vertex fit of K_S^0 , the sideband is excluded in each distribution to further reject fake K_S^0

4.1.2 B^0 Reconstruction

By combining three K_S^0 particles from selected dataset, we can reconstruct B^0 . The beam-constraint mass M_{bc} and energy difference ΔE are used to extract signal, as defined in Equation 4.1 and 4.2, respectively. For M_{bc} , s is defined as the invariant mass and p_B^* is the reconstructed B momentum, both in the center-of-mass frame of e^+e^- . For ΔE , E_B^* is the reconstructed energy in the center-of-mass frame of e^+e^- . These two variables are quite useful for discriminating signal and background events for hadronic B decay with fully reconstructed final states. In Belle II, the B^0 candidates with $M_{bc} > 5.2$ GeV and $|\Delta E| < 0.2$ GeV are requested to be reconstructed. The B^0 vertex fit using *TreeFit* is performed on each B^0 candidate and a quality index named *chiProb* is calculated. Only B^0 with converged vertex fit result is kept by a very loose cut of "chiProb > 0.001", which is essential to obtain the decay vertex positions of B mesons for further CP violation analysis. When multiple B^0 candidates are obtained in a single event, the best candidates selection (BCS) is performed by ranking their vertex fit quality. Since the BCS is based on the vertexing quality that might introduce bias in the vertex positions for CP fit, we check the distribution of the vertex χ^2 , as shown in Figure 4-4 top right where the data and *generic MC* present a good consistence within 1σ on average. The distribution of candidates number per event without BCS is shown in top left of Figure 4-4 as well, showing an agreement between data and *generic MC* within 1σ . The distribution of candidates per event is also in an agreement with that from *signal MC* (bottom left of Figure 4-4). The 2D distribution of M_{bc} and ΔE from $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ *signal MC* is shown in Figure 4-4 bottom right, where the correlation factor is about 15% between two observables.

$$M_{bc} = \sqrt{\frac{s}{4} - p_B^{*2}} \quad (4.1)$$

$$\Delta E = E_B^* - \frac{\sqrt{s}}{2} \quad (4.2)$$

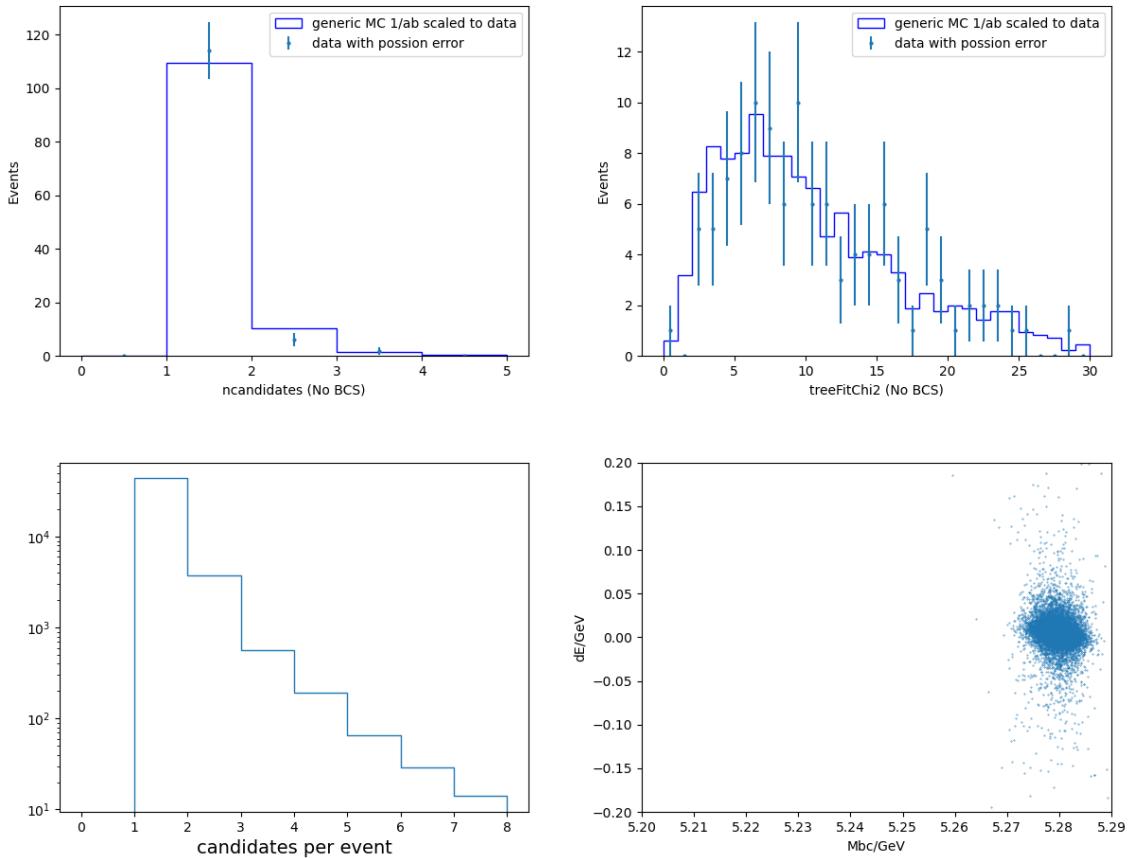


Figure 4-4: Top left is the candidates per event in data and *generic MC* before the BCS. Top Right is the χ^2 for data and *generic MC* before BCS. Bottom left is the number of candidates B^0 per event from *signal MC*. Bottom right is the 2D M_{bc} and ΔE distribution from *signal MC*.

4.1.3 Continuum Suppression

The production cross-section of $B\bar{B}$ from $\Upsilon(4S)$ receives a sizable contribution from other flavor of quarks other than b quark. This calls a demand to distinguish a specific $B\bar{B}$ decay events from combinatorial background from $e^+e^- \rightarrow q\bar{q}$, so called continuum suppression (CS). The rejection is essential because it's the dominated background. In the case of $b \rightarrow s$ charmless decay like $B^0 \rightarrow K_S^0 K_S^0 K_S^0$, the number of continuum background can exceed the signals by a few orders of magnitudes without suppression. In a $B\bar{B}$ event, two mesons are produced almost at rest in the CMS frame since the resonance state $\Upsilon(4S)$ is just slightly lighter than beam energy. As

a result , decay products are emitted more isotropically, compared to continuum background. The ARGUS and CLEO collaboration[40] developed a set of variables to suppress the continuum, which has also been implemented into BASF2 framework.

CLEO cones momentum can be presented as Equation 4.3, where p_i is momentum of i-th particle in Rest-Of-Event (ROE) particles in a event except for the ones used to reconstructed CP -side B^0 , θ_i is angle against momentum thrust of reconstructed CP -side of B meson.

$$L_n = \sum_{i \in ROE} p_i \times |\cos\theta_i| \quad (4.3)$$

Besides, the modified Super Fox-wolfram momentum named KSFW momentum, are calculated in each event. The KSFW momentum are defined as shown in Equation 4.4.

$$KSFW = \sum_{l=0}^4 (R_l^{so} + R_l^{oo}) + \gamma \sum_{n=1}^{N_t} |P(t)_n| \quad (4.4)$$

where the first term is shown in Equation 4.5.

$$R_l^{so} = \frac{\alpha_{cl} H_{cl}^{so} + \alpha_{nl} H_{nl}^{so} + \alpha_{ml} H_{ml}^{so}}{E_{beam}^* - \Delta E} \quad (4.5)$$

when l is odd in Equation 4.5:

$$H_{nl}^{so} = H_{ml}^{so} = 0 \quad (4.6)$$

and H_{cl}^{SO} is defined as shown in Equation 4.7:

$$H_{cl}^{so} = \sum_i \sum_{jx} Q_i Q_{jx} |p_{jx}| P_l(\cos\theta_{i,jx}) \quad (4.7)$$

i runs over B daughter particles and jx for other particles in ROE. Q is charge and p_{jx} is momentum for each particle. $P_l(\cos\theta_{i,jx})$ is the i -th order Legendre polynomial of cosine of i and jx -th particles. On the other hand, for l is even, H_{xl}^{SO} can be written

in Equation 4.8.

$$H_{xl}^{so} = \sum_i \sum_{jx} |p_{jx}| P_l(\cos\theta_{i,jx}) \quad (4.8)$$

The second term in Equation 4.4, when l is odd, can be defined as Equation 4.9.

$$R_l^{oo} = \sum_j \sum_k \beta_l Q_j Q_k |p_j| |p_k| P_l(\cos\theta_j, k) \quad (4.9)$$

j and k runs over ROE particles and others are same as Equation 4.7. For an even l :

$$R_l^{oo} = \sum_j \sum_k \beta_l |p_j| |p_k| P_l(\cos\theta_j, k) \quad (4.10)$$

β is Fisher coefficients to be determined. Using above definitions, we can form the possibility density functions for KSFW, cosine angle against B meson thrust $\cos\theta_B$ and ΔZ of two side vertices. Then based on each event's variables' value, we can calculate a ratio \mathcal{R} as Equation 4.11, where the likelihood L of signal(L_S) and background(L_B) are obtained from the possibility density functions defined in Equation 4.12. The \mathcal{R} is:

$$\mathcal{R} = \frac{L_S}{L_S + L_B} \quad (4.11)$$

$$L_{S/B} = P(KSFW)_{S/B} \times P(\cos\theta_B)_{S/B} \times P(\Delta Z)_{S/B} \quad (4.12)$$

where P is probability density function for signal and continuum, depending on the discriminating variables in the parentheses. For example, the distribution of a variable called R_2 shown in Figure 4-5 where the possibility density function is different for signal and continuum events in *generic MC*.

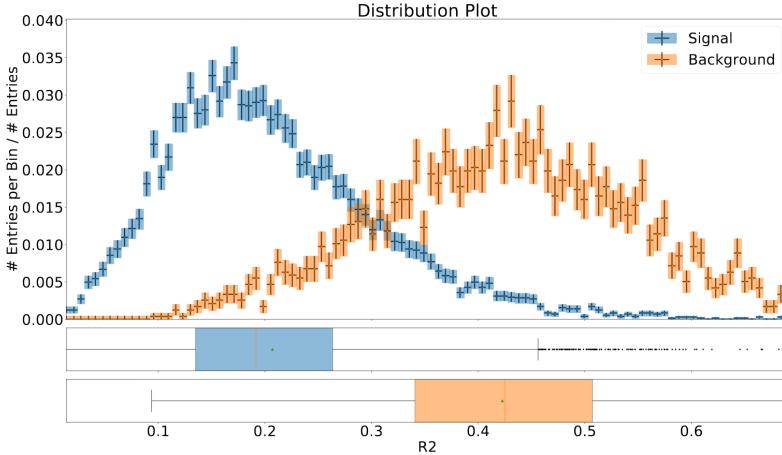


Figure 4-5: R_2 is the ratio of the second to the zeroth KSFW momentum in Equation 4.4 of which the distribution in *signal MC* sample which serves as the highest weight as a variable in discriminating the continuum events, having a quite different distribution between signal and background.

In order to maximize suppression power in this analysis, these variables (KSFW, CLEO cone momentum and angular distributions) are combined as an input for FastBDT classifier. The targeted variable is the continuum event truth. The MC samples using signal $B\bar{B}$ events from *signal MC* and continuum events from *generic MC* ($q\bar{q}$ components) are prepared in a ratio of their cross-section at $\Upsilon(4S)$ energy. The same events reconstruction procedures for B^0 is applied for both MC samples. Events passing the reconstruction for B^0 using M_{bc} and ΔE are used for training the continuum suppression classifier. The fraction of signal and background is set to 1:1 during the training. The output of continuum suppression classifier is renamed as *FBDT_CS*. Then we determine the cut value at 0.66 based on the maximum of *FOM* curve, as shown in Figure 4-9. The variables used in training are listed in Table 4.2 with their abbreviations and the rank of important variables is in Table 4.3.

The correlation between these training variables are shown in Figure 4-6. The correlation among the variables are varied in signal and background. The ROC curve and the efficiency/purity with respect to the classifier output are shown in Figure 4-7.

Overtraining check is made by comparing the distribution of signal and background depending on the classifier output in both training and testing samples. The

Table 4.2: Variables and the abbreviations for CS.

Observables	Abbreviations
CleoConeCS(9,)	CleoC1
KSFWVariables(hoo1,)	KSFWV1
CleoConeCS(7,)	CleoC2
CleoConeCS(5,)	CleoC3
KSFWVariables(hso22,)	KSFWV2
KSFWVariables(hoo3,)	KSFWV3
CleoConeCS(4,)	CleoC4
KSFWVariables(hoo4,)	KSFWV4
CleoConeCS(3,)	CleoC5
CleoConeCS(6,)	CleoC6
CleoConeCS(8,)	CleoC7
KSFWVariables(hso14,)	KSFWV5
KSFWVariables(hso00,)	KSFWV6
KSFWVariables(et,)	KSFWV7
KSFWVariables(hso24,)	KSFWV8
KSFWVariables(hso04,)	KSFWV9
KSFWVariables(hso20,)	KSFWV10
KSFWVariables(mm2,)	KSFWV11
KSFWVariables(hoo2,)	KSFWV12
thrustOm	thrus1
cosTBz	cosTB1
CleoConeCS(1,)	CleoC8
CleoConeCS(2,)	CleoC9
KSFWVariables(hso02,)	KSFWV13
KSFWVariables(hoo0,)	KSFWV14
KSFWVariables(hso12,)	KSFWV15
KSFWVariables(hso10,)	KSFWV16
cosTBTO	cosTB2
thrustBm	thrus2
R2	R2

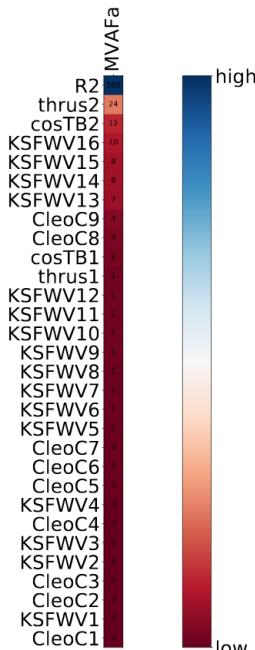


Table 4.3: The rank of important variables for CS.

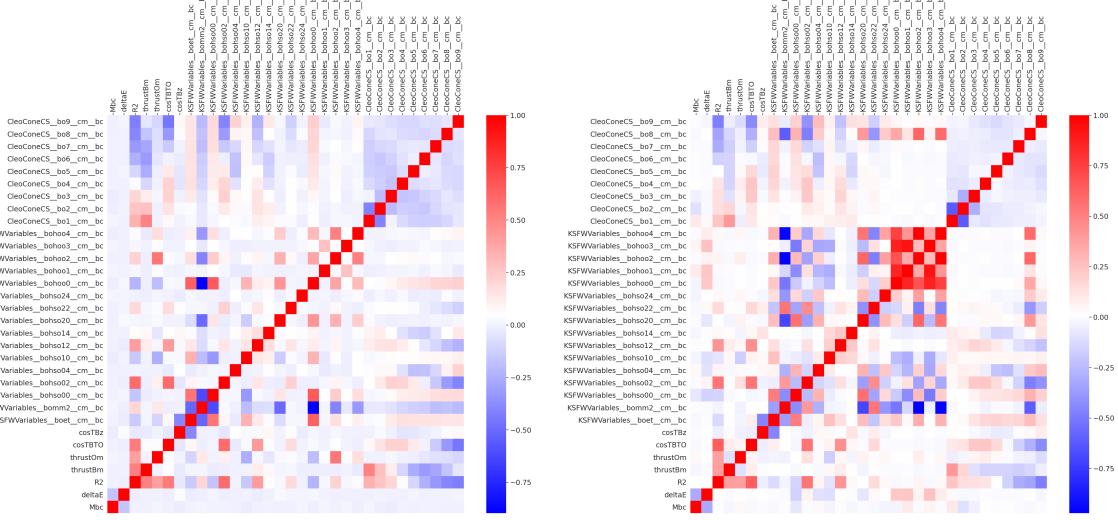


Figure 4-6: The correlation in variables for continuum suppression. The left is for signal and the right is for background.

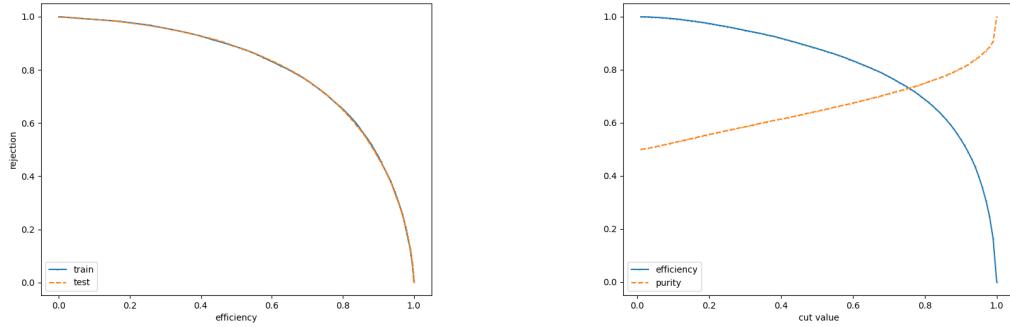


Figure 4-7: The left is the ROC curve (blue for training and orange for testing) and the right is the efficiency(blue) and purity(orange) regarding the classifier output *FBDT_CS*.

testing samples show about 1% lower in each bin for both signal and background events, which is within the acceptable range.

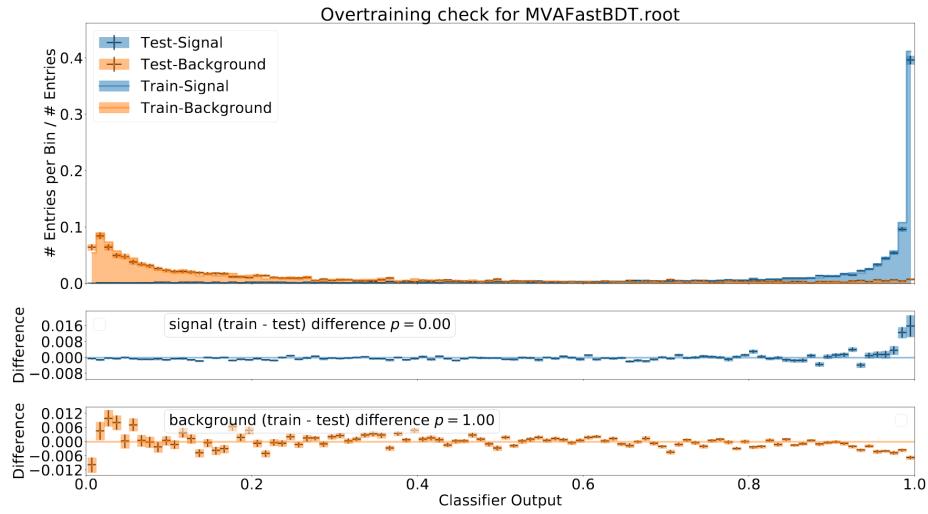


Figure 4-8: Over-training check of continuum classifier, where a very small difference in training and testing (1%) is shown.

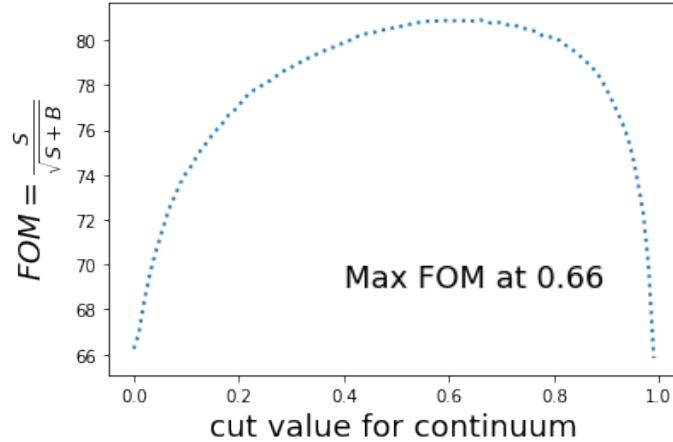


Figure 4-9: FOM depending on the cut value of continuum classifier output, cut value at 0.66 is used for continuum suppression.

The summary of B^0 selections is listed in Table 4.4, including the application of *KsFinder* (by *FBDT_Ks*) and continuum suppression (by *FBDT_CS*).

B^0	M_{bc}/GeV	$\Delta E/\text{GeV}$	<i>chiProb</i>	<i>Rank</i>	<i>FBDT_CS</i>	<i>FBDT_Ks</i>
Selection	$> 5.20 \& < 5.29$	$ \Delta E < 0.2$	> 0.001	= 1	> 0.66	> 0.74

Table 4.4: B^0 selection criteria, *chiProb* is from B^0 vertex fit and *Rank* is from best candidate selection.

Explicitly, the reconstruction performance of B^0 is summarized in Table 4.5, the efficiency, purity, fraction of multiplicity events and best candidates fraction of B^0 is slight improved in Belle II compared to Belle.

event selection	efficiency	purity	f_{MB}	BCS
Belle Standard	35%(33%)	96%(99%)	6%(6%)	83%(96%)
Belle II ($BG1$)	36%(34%)	96%(98%)	(4%)(4%)	95%(96%)
Belle II ($BG0$)	40%(36%)	96%(99%)	(3%)(3%)	97%(97%)

Table 4.5: The efficiency is defined by the fraction of best candidates among the MC input number. Purity is the fraction of true B^0 in best candidates. f_{MB} stands for multiple B^0 events fraction in true signal events. BCS is the fraction of best candidates being a true signal. All values in the parenthesis are calculated in $|M_{bc}| - 5.28 < 0.1$ and $|\Delta E| < 0.1$, called as “signal region” where efficiency is lower but purity is higher, compared to the full range of M_{bc} and ΔE in Table 4.4.

4.1.4 Resonance Background

Besides the major contribution from continuum background, charmonium resonance that mediates through $b \rightarrow c$ transition brings odd CP eigenvalue in the final states as same as $B^0 \rightarrow K_S^0 K_S^0 K_S^0$. Monitoring their contribution is also important. Basically, one needs to check the resonance states formed by two K_S^0 with corresponding invariant mass. In $B^0 \rightarrow X(K_S^0 K_S^0)K_S^0$, there are two types of resonant events that give out same final states, one is resonant signal and the other is resonant background. For $b \rightarrow s$ transitions as resonance signal because of the CP -even final states, X could be $f_2(1270)$, $f_0(1500)$, $f'_2(1525)$, $f_0(980)$, $f_0(1710)$ and $f_2(2010)$. For $b \rightarrow c$ transition as resonance background because of CP -odd final states, X could be D^0 , J/ψ , $\psi(2S)$, χ_{c0} , χ_{c1} , and χ_{c2} .

The number of these background in signal reconstruction could be further reduced by implementing veto on invariant mass of $2K_S^0$. However, such veto should be carefully validated with data. The distribution of invariant mass of X should agree well in MC and data, which is hard to check in the low luminosity. The distribution of $2K_S^0$ invariant mass in *generic MC* and data are shown in the Appendix

Some of these resonance have not been implemented inside *generic MC* production in the current Belle II simulation. Given the very limited statistics of data accumulation we used in this analysis, we only present the expected number of these resonances in 400 fb^{-1} luminosity (about 2.14×10^8 events) from generic $\Upsilon(4S)$ events. These numbers should be re-checked in the future when data accumulation increases, and veto must be based on the structure of $2K_S^0$ invariant mass from data as well. Details about the expected yields can be found in Table 4.6. Currently there is no veto applied for rejecting these resonant background considering the estimated background number is about 1 event in the current luminosity.

4.1.5 $B\bar{B}$ background and self-cross feed

Another possible contribution of backgrounds are from $B\bar{B}$ events including the charged and the neutral particles. The estimated contributions of these types can be checked with charged $B\bar{B}$ samples and the mixed samples. For this channel, the number of the events is very limited. Self-cross feed backgrounds stands for the events from the signal-like events but the tag-side particle(s) is associated as a fake signal. The combined contributions from $B\bar{B}$ background and self-cross feed is about 3% in the channel and therefore we don't perform special treatment on them.

4.1.6 Signal Extraction

The event selections defined in Table 4.4 is applied to *signal MC*, *generic MC* and experiment data for signal extraction. The integral luminosity in *generic MC* is 1 ab^{-1} and experiment data used in this analysis is about 62.8 fb^{-1} from the latest official processing.

The unbinned maximum likelihood fit using RooFit is performed to extract the signal. The 2D fit using both M_{bc} and ΔE are done by taking the probability density function as shown in Equation 4.13.

$$\mathcal{P}(M_{bc}, \Delta E) = f_{sig} \times \mathcal{P}_{sig}^{M_{bc}} \times \mathcal{P}_{sig}^{\Delta E} + (1 - f_{sig}) \mathcal{P}_{bkg}^{M_{bc}} \times \mathcal{P}_{bkg}^{\Delta E} \quad (4.13)$$

Table 4.6: Expected yield for signal and background resonances $2.14 \times 10^8 B\bar{B}$ in *generic MC*. The branching fraction of $B \rightarrow X K_S$ and $X \rightarrow 2K_S$ are listed for both PDG value and value in Belle II generic decay profile (see section 2.9). The events from CP -odd contamination is expected to be very low at current luminosity (62.8 fb^{-1}).

Resonances	$\text{Br}(B \rightarrow X K_S)\text{PDG}$	$\text{Br}(X \rightarrow 2K_S)$	$\text{Br}(B \rightarrow X K_S)\text{Dec.}$	$\text{Br}(X \rightarrow 2K_S)\text{Dec.}$	$B\bar{B}$ pairs	Expected yields
$D^0 K_S$	2.6×10^{-5}	1.7×10^{-4}	2.6×10^{-5}	1.8×10^{-4}	2.14×10^8	0.134
ηK_S	3.45×10^{-4}	$< 3.1 \times 10^{-4}$	4×10^{-4}	No Value	2.14×10^8	No Value
$J/\psi K_S$	4.35×10^{-4}	$< 1.4 \times 10^{-8}$	4.35×10^{-4}	0	2.14×10^8	0
$\psi(2S) K_S$	2.9×10^{-4}	$< 4.6 \times 10^{-6}$	2.9×10^{-4}	0	2.14×10^8	0
$\chi_{c0} K_S$	7.3×10^{-5}	3.16×10^{-3}	7.35×10^{-5}	3.1×10^{-3}	2.14×10^8	6.21
$\chi_{c1} K_S$	1.96×10^{-4}	6×10^{-5}	1.96×10^{-4}	1×10^{-5}	2.14×10^8	0.05
$\chi_{c2} K_S$	7.5×10^{-6}	2.6×10^{-4}	7.5×10^{-6}	5.5×10^{-4}	2.14×10^8	0.11
$f_2(1270) K_S$	1.35×10^{-6}	1.15×10^{-2}	1.35×10^{-6}	1.15×10^{-2}	2.14×10^8	0.42
$f'_2(1525) K_S$	1.5×10^{-7}	2.22×10^{-2}	No value	0.22	2.14×10^8	No Value
$f_2(2010) K_S$	5×10^{-7}	No Value	No Value	No Value	2.14×10^8	No Value
$f_0(980) K_S$	2.7×10^{-6}	No Value	2.75×10^{-6}	No Value	2.14×10^8	43.3
$f_0(1710) K_S$	5×10^{-7}	No Value	No Value	No Value	2.14×10^8	No Value
$f_0(1500) K_S$	6.5×10^{-5}	0.022	No Value	0.022	2.14×10^8	No Value
Total	-	-	-	-	-	$\simeq 50$

where $\mathcal{P}_{sig}^{M_{bc}}$ and $\mathcal{P}_{sig}^{\Delta E}$ are the single Gaussian and triple Gaussian functions. The f_{sig} is fraction of signal events based on M_{bc} and ΔE . The $\mathcal{P}_{bkg}^{M_{bc}}$ is primarily continuum events, and presented as Argus distribution as Equation 4.14 shows, with a preset mass threshold at $c = 5.29$ GeV.

$$f(x; \chi, c) = \frac{\chi^3}{\sqrt{2\pi}\Psi(\chi)} \cdot \frac{x}{c^2} \sqrt{1 - \frac{x^2}{c^2}} \cdot \exp\left\{-\frac{1}{2}\chi^2(1 - \frac{x^2}{c^2})\right\} \quad (4.14)$$

x is defined in $0 < x < c$. χ and c are parameters of the distribution, $\Psi(\chi) = \Phi(\chi) - \chi\phi(\chi) - \frac{1}{2}$ where $\Phi(\chi)$ and $\phi(\chi)$ cumulative distribution and probability density functions of the standard normal distribution, respectively.

The $\mathcal{P}_{bkg}^{\Delta E}$ is modeled by the first order Chebyshev polynomials. The shape parameters of signal events are determined by fitting to *signal MC*, and then fixed as constants in fitting of Equation 4.13 on M_{bc} and ΔE for *signal MC* and experiment data. Fitting results on *signal MC* are shown in Figure 4-10. The continuum back-

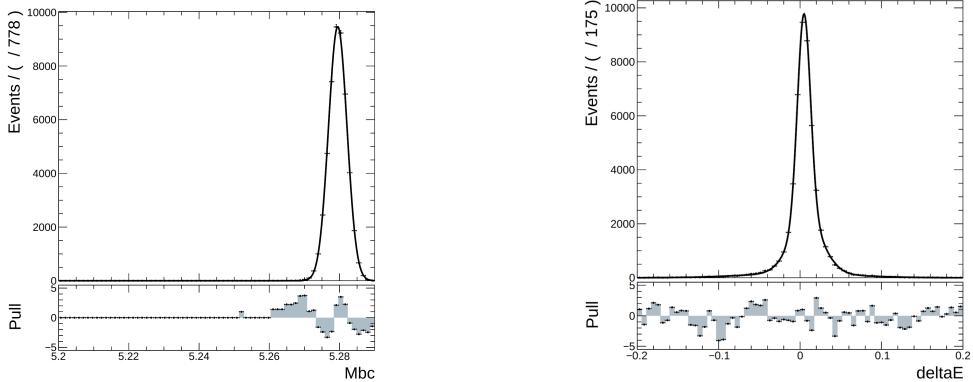


Figure 4-10: The distribution of M_{bc} and ΔE of *signal MC* of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ fitted with single and triple Gaussian functions respectively.

ground is fitted by using $q\bar{q}$ events from *generic MC* to determine the shapes then fix them as constants for 2D fit as shown in Figure 4-11.

Then we set the events number for signal and background as floating parameters and use Equation 4.13 as 2D fit model on 1 ab^{-1} *generic MC* and experiment data, which is also done by using unbinned maximum likelihood fit. For B^0 in *generic MC*, the stacked histogram of each contribution and the 2D fit result projected on M_{bc} and ΔE is shown in Figure 4-12.

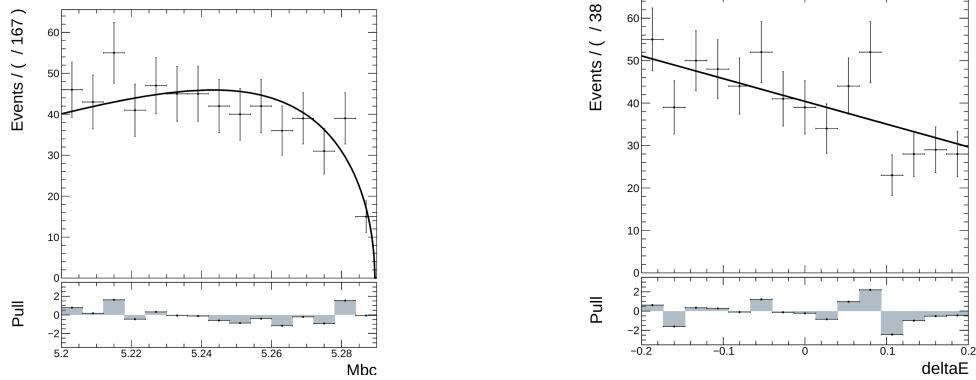


Figure 4-11: The distribution of M_{bc} and ΔE of continuum events in *signal MC* fitted with Argus and Chebyshev polynomial, respectively.

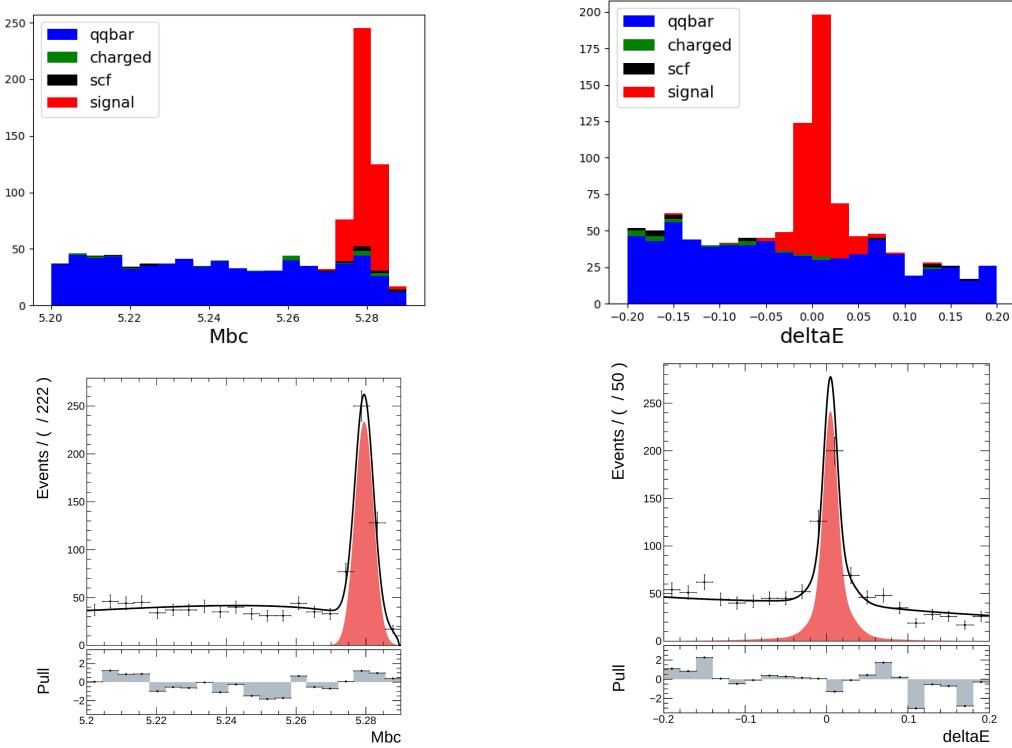


Figure 4-12: Top is the stacked plots for *generic MC* of M_{bc} and ΔE , where each background components are stacked with signal. The bottom is the 2D fit on 1 ab^{-1} *generic MC* projected on M_{bc} and ΔE , the red is signal component from the fit result in both plots.

Before perform fitting on experiment data, the distribution of K_S^0 invariant mass from the reconstructed B^0 candidates is compared between *generic MC* and experiment data. The distributions are shown in Figure 4-13, where the *generic MC* is

scaled to the luminosity of experiment data and an agreement within $\sim 1\sigma$ is observed on average.

The 2D fit of experiment data projected on M_{bc} and ΔE is in Figure 4-14.

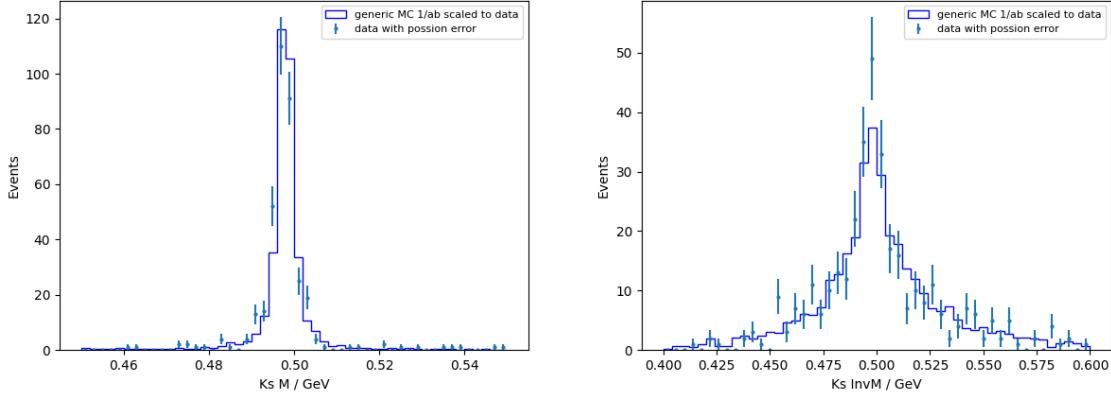


Figure 4-13: Invariant mass before(left) and after vertex fit(right) from *generic MC* and experiment data.

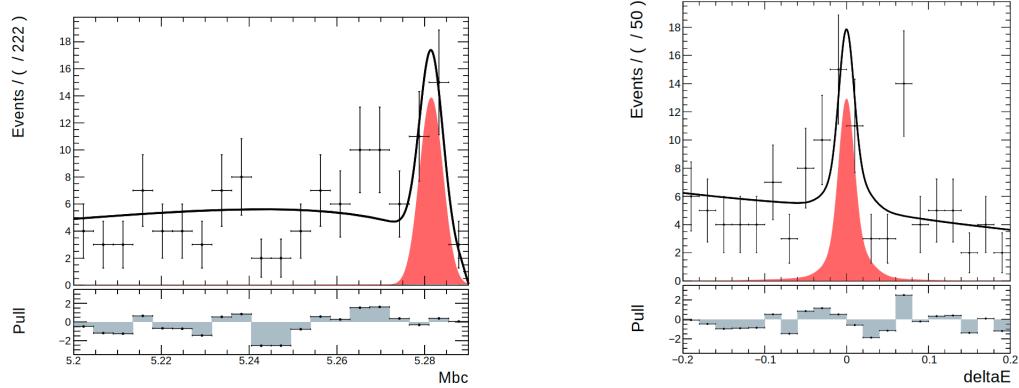


Figure 4-14: M_{bc} and ΔE 2D fit on 62.8 fb^{-1} data, the red is the signal component.

The number of signal events is extracted by the integral of fit model over the signal region which is defined as $5.27 < M_{bc} < 5.29 \text{ GeV}$ and $-0.1 < \Delta E < 0.1 \text{ GeV}$. The expected signal events with $\sim 35\%$ efficiency is calculated as Equation 4.15.

$$\mathcal{B}(B^0 \rightarrow K_S^0 K_S^0 K_S^0) = \frac{N_{sig}}{\mathcal{B}(K_S^0 \rightarrow \pi^+ \pi^-)^3 \times \epsilon_{rec} \times N_{B\bar{B}}} \quad (4.15)$$

In 1 ab^{-1} *generic MC*, the expected signal number is $7.7 \times 10^8 \times 6 \times 10^{-6} \times 21\% \times 35\% \simeq 339$. The 2D fit result from M_{bc} and ΔE yields 341 ± 20 events which agrees with

expected number within 1σ . The event number in sideband defined as $M_{bc} < 5.26$ GeV in *generic MC* is 507. Compared to Belle result with 772×10^6 ($\sim 1 \text{ ab}^{-1}$) $B\bar{B}$ pairs used, signal from data yields 327 ± 19 . In 62.8 fb^{-1} data fit in Belle II, we extract $N_{sig} = 17.4 \pm 4.2$ in signal region. The sideband region $M_{bc} < 5.26$ GeV contains 60 events in data.

To check linearity of the event number fitted from the M_{bc} and ΔE in this low statistics case, we extract the fraction of continuum backgrounds from *generic MC* sample rescaled to the experimental data luminosity, which includes about 46 continuum events. Then the number of signal events from 5 to 30 with 5 events per step are injected into the continuum events, to perform the M_{bc} and ΔE fit to check the output signal events number. The M_{bc} and ΔE distributions and fit in each injection test are shown in Figure 4-15. The fitted signal and background events depending on the injected numbers (linearity test) are presented in Figure 4-16, where the dependence on both signal and background events number are fitted with linear functions. The fit results show a good linearity on the input and output of signal numbers while the background numbers remain constant close to the input number.

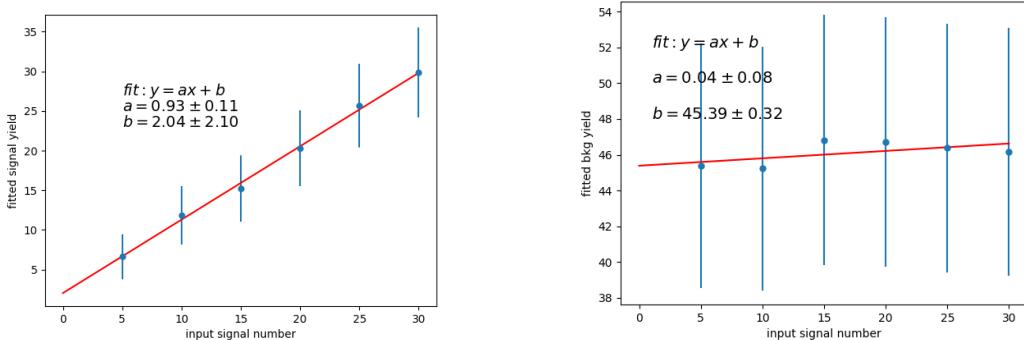
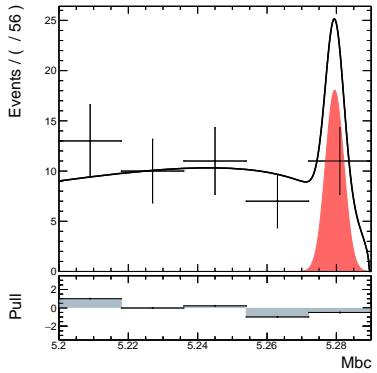


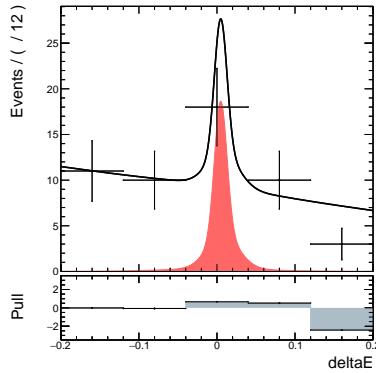
Figure 4-16: Injection test for signal extraction. The linearity is clear between input and output signal events number.

4.1.7 Kinematics and Vertexing Dependence on *KsFinder*

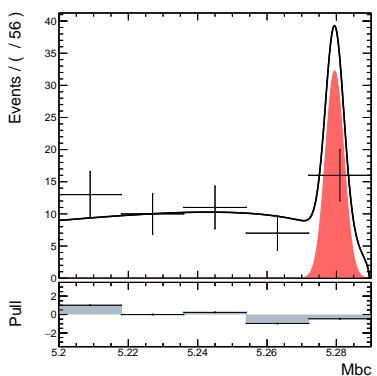
KsFinder largely reduce the combinatorial background of B^0 by improving K_S^0 purity. The previous section shows a good reconstruction performance at low statistics in



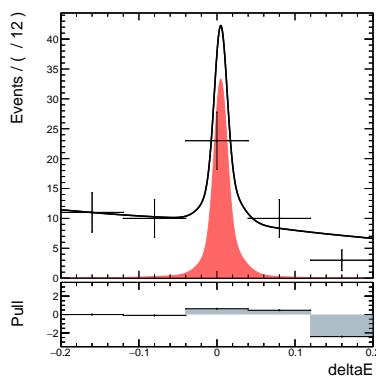
a) signal injected: 5



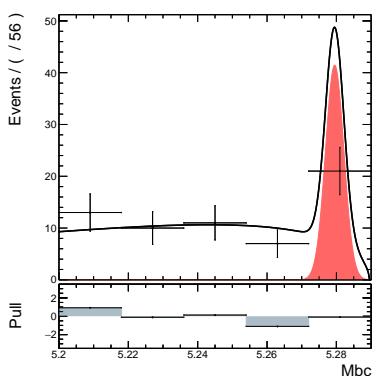
b) signal injected: 5



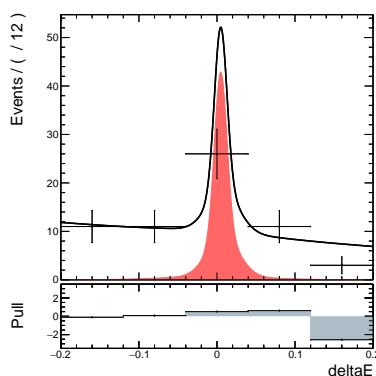
c) signal injected: 10



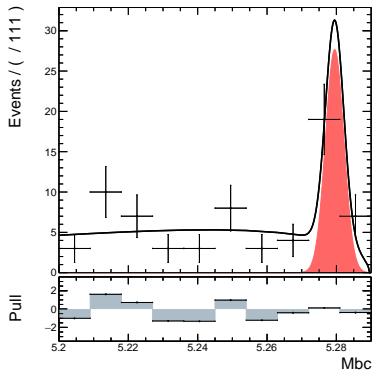
d) signal injected: 10



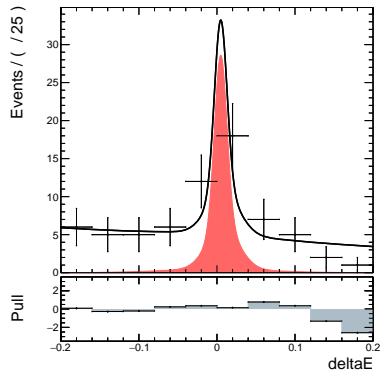
e) signal injected: 15



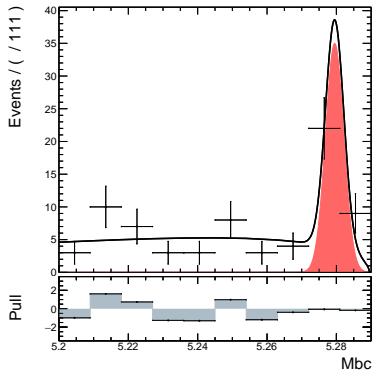
f) signal injected: 15



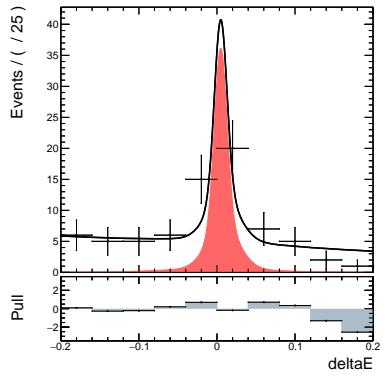
g) signal injected: 20



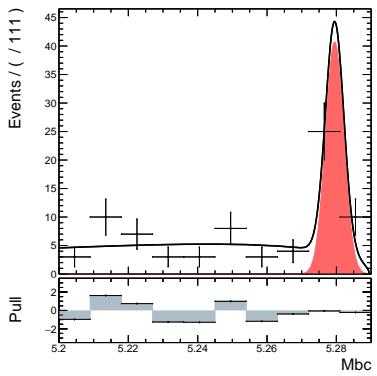
h) signal injected: 20



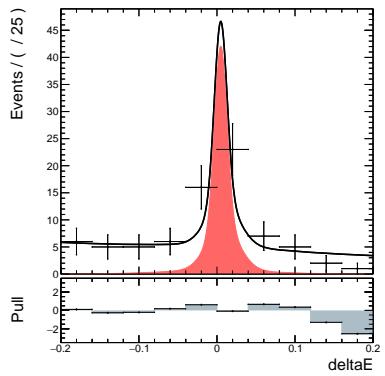
i) signal injected: 25



j) signal injected: 25



k) signal injected: 30



l) signal injected: 30

Figure 4-15: The fit results of M_{bc} and ΔE in signal injection test, where signal events from 5 to 30 with 5 per step are injected with 46 continuum events.

early phase 3 data. Without the power of rejection provided by K_S^0 finder, rediscovery of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ in early phase 3 of Belle II won't be feasible.

However, it's essential to check the potential impact on M_{bc} and ΔE , as well as vertex positions on z -axis of B^0 due the implementation of *KsFinder*. The K_S^0 classification uses information such as invariant mass and decay vertex positions which may propagate bias into B^0 signal extraction, eventually may affect the measurement of CP parameters. Given each type of B^0 based on how many CDC-only tracks it has in the final states, the comparison on M_{bc} and ΔE with or without *KsFinder* is performed by fitting the distribution in *signal MC*. M_{bc} and ΔE are modeled by signal and double Gaussian, respectively. Comparing corresponding fit results, no clear bias on M_{bc} and ΔE is found by using *KsFinder* where fit results are agreed well within one standard deviation. The fit results are shown in Figure 4-17 and 4-18. To be noted, the ΔE distributions show a small shift positively when more CDC-only tracks are used to reconstructed B^0 because the reconstructed energy from charged pions is slightly higher due to the less energy loss in VXD.

Similar to the comparison of M_{bc} and ΔE , the z direction vertex position and the vertex position difference between CP and tag sides are also checked, in which no clear bias are spotted either. The z and Δz are modeled using single Gaussian with the same mean but different standard deviation. The results are shown in Figure 4-19 and 4-20. It is obvious that in Figure 4-19, the CP -side resolution of vertex on z -axis is wider when the final states of B^0 have more CDC-only tracks, especially when all the tracks only contains CDC hits (6 CDC-only tracks).

Above all, no strong appearance of bias on M_{bc} and vertex positions from using *KsFinder* has been found, *KsFinder* may implement a small shift on the vertex position which is negligible compared to the large statistical uncertainty due to the current low luminosity. Hence, there's no correction on these observables are applied in this analysis, which the systematic uncertainty from *KsFinder* is evaluated.

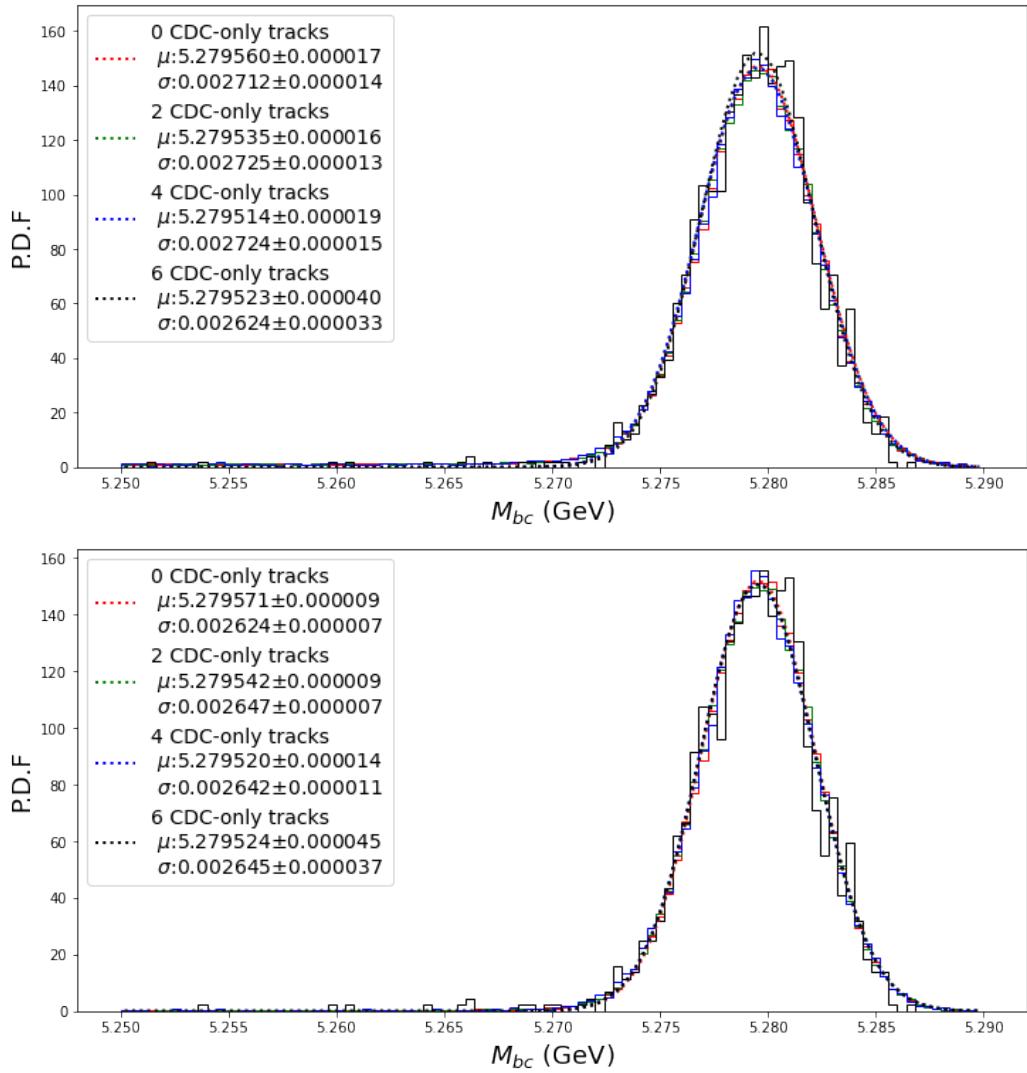


Figure 4-17: M_{bc} distribution based on the number of CDC-only tracks in final states.
Top: no *KsFinder* used; Bottom: *KsFinder* used.

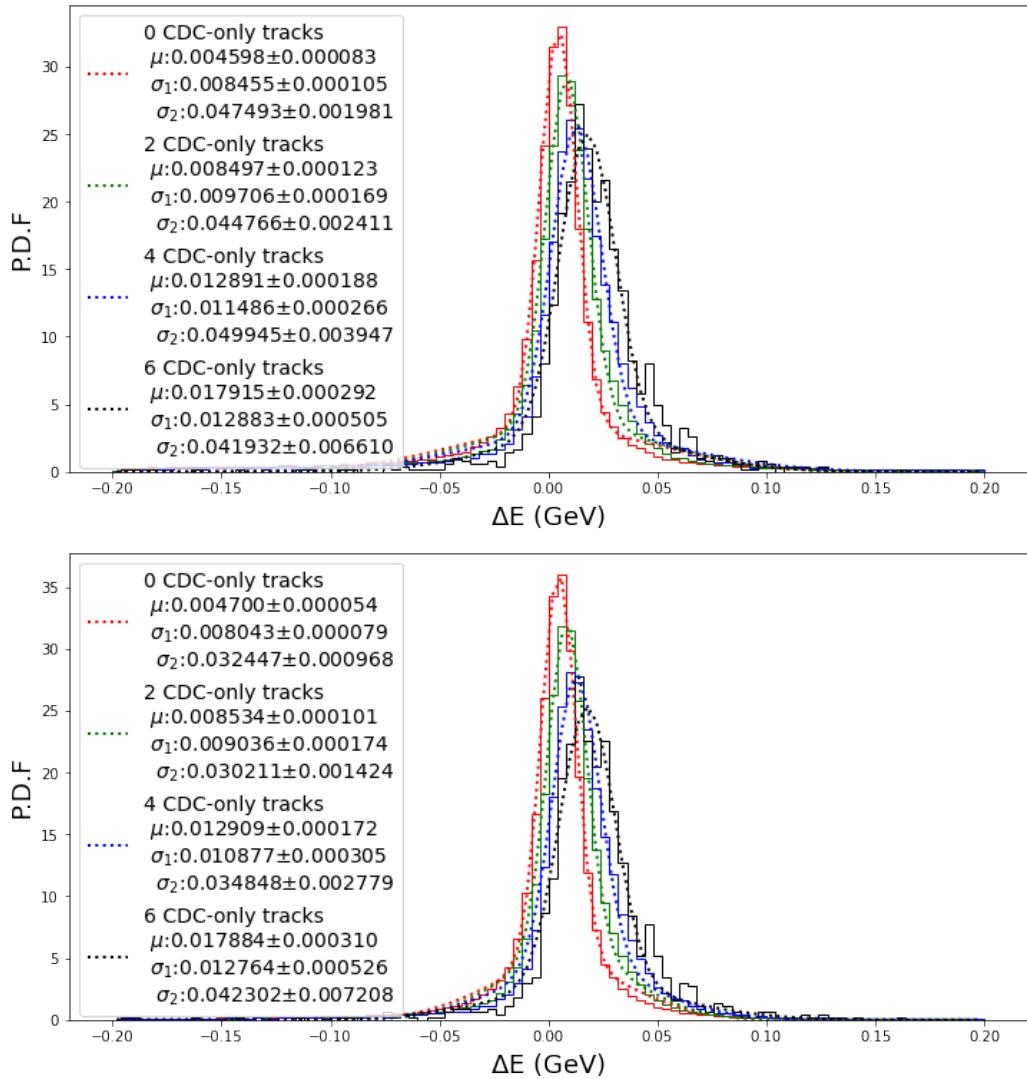


Figure 4-18: ΔE distribution based on the number of CDC-only tracks in final states.
Top: no *KsFinder*; Bottom: *KsFinder* used.

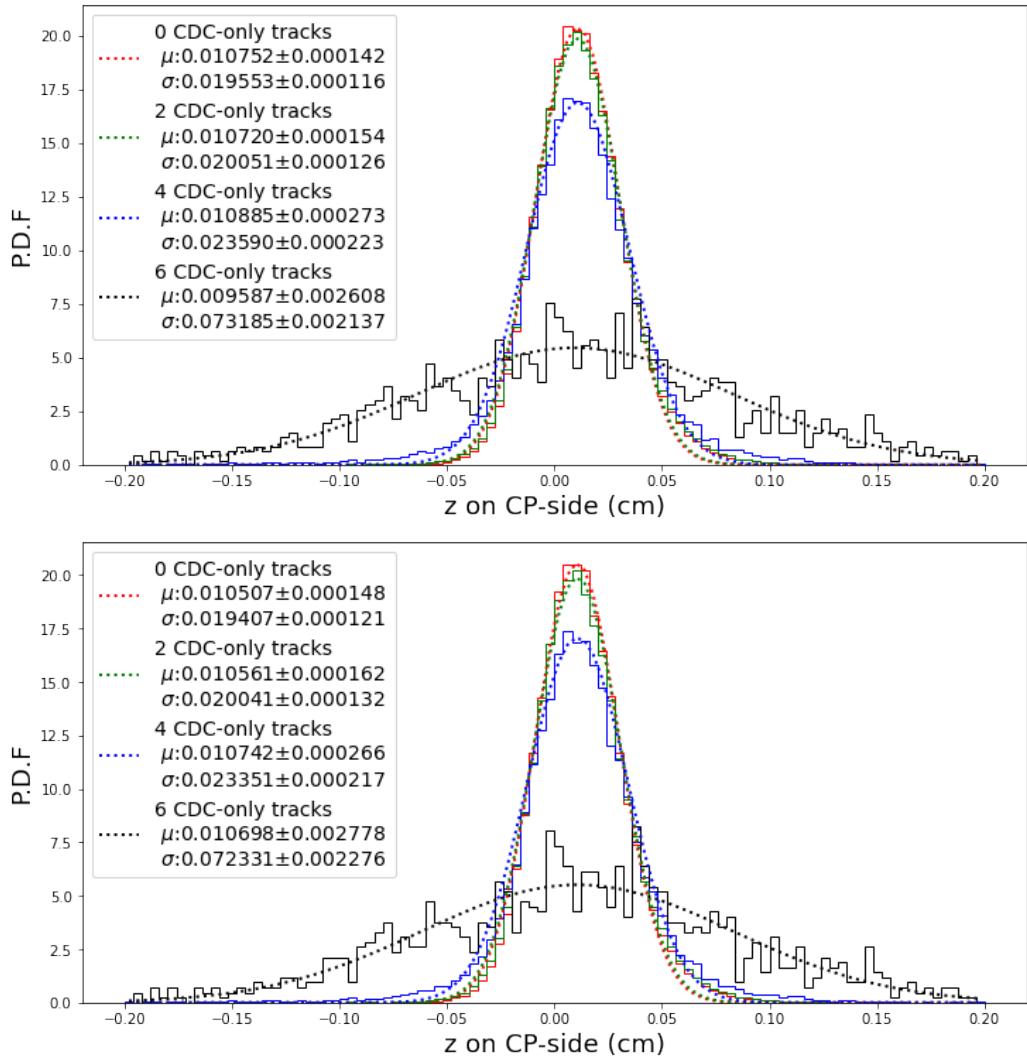


Figure 4-19: Δz distribution based on the number of CDC-only tracks in final states.
Top: no *KsFinder*; Bottom: *KsFinder* used.

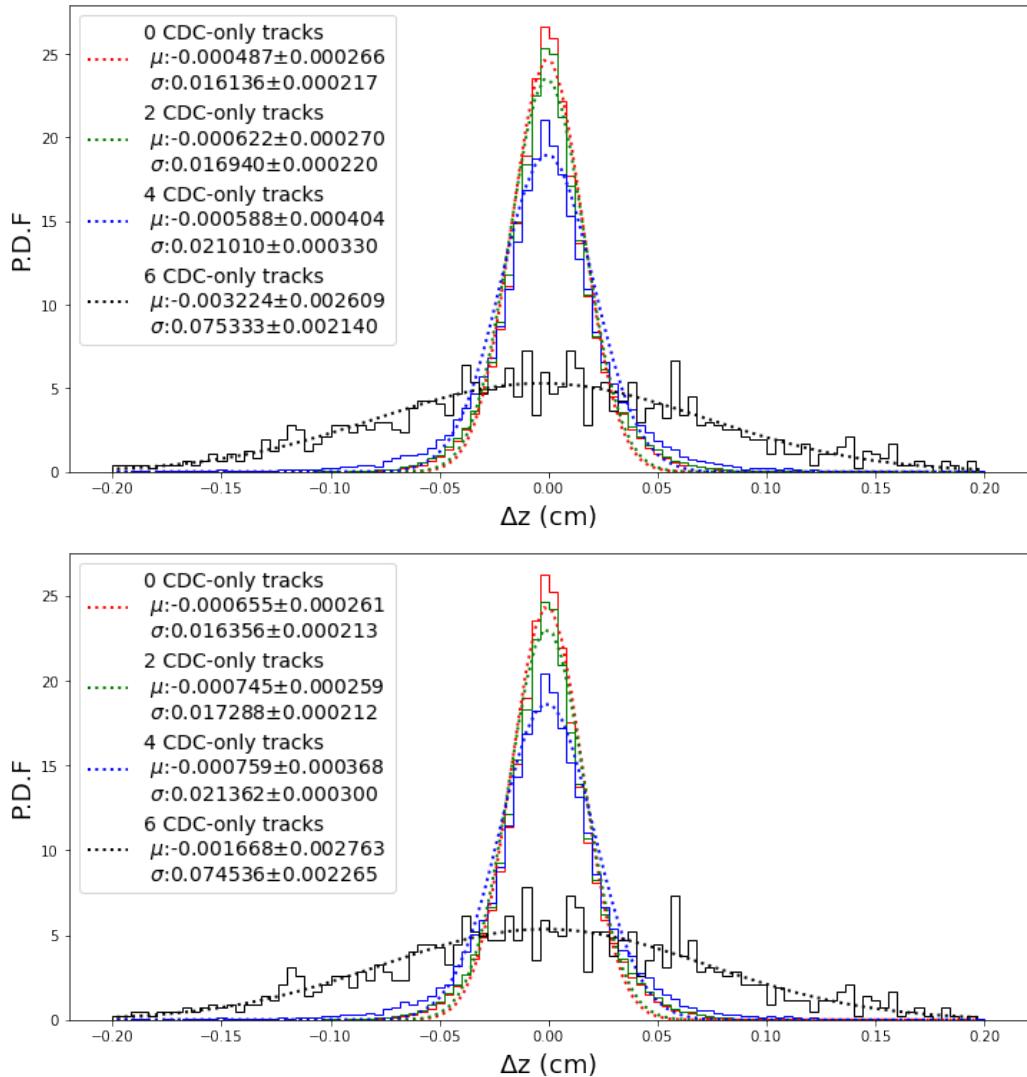


Figure 4-20: Δz distribution based on the number of CDC-only tracks in final states.
Top: no *KsFinder*; Bottom: *KsFinder* used.

Chapter 5

CP parameters measurement

The measurement of CP parameters \mathcal{S} and \mathcal{A} are performed by fitting Equation 5.1 to the distribution of events with respect to the decay time difference Δt and flavor q , where $\Delta t = t_{CP} - t_{tag}$ and $q = +1(-1)$ when the tag-side B meson is $B^0(\bar{B}^0)$.

$$\mathcal{P}_{sig}(\Delta t, q) = \frac{e^{-|\Delta t|/\tau_{B^0}}}{4\tau_{B^0}} \left\{ 1 + q \cdot [\mathcal{S} \sin(\Delta M_d \Delta t) + \mathcal{A} \cos(\Delta M_d \Delta t)] \right\} \quad (5.1)$$

The Equation 5.1 describes the physics distribution of signal events only. To perform the unbinned maximum likelihood fit on data, a complete model for i -th event that includes the overlay of background components and outlier bands can be defined as Equation 5.2.

$$\begin{aligned} \mathcal{P}(\Delta t_i, q_i, f_i^{sig}, \mathcal{S}, \mathcal{A}) &= (1 - f_{ol}) \left[f_{sig} \mathcal{P}_{sig}(\Delta t_i, q_i, \mathcal{S}, \mathcal{A}) + (1 - f_{sig}) \mathcal{P}_{bkg}(\Delta t_i) \right] \\ &\quad + f_{ol} \mathcal{P}_{ol}(\Delta t_i) \end{aligned} \quad (5.2)$$

where f_{sig} and f_{ol} are the fraction of signal and outlier components, respectively. The \mathcal{P}_{bkg} and \mathcal{P}_{ol} are defined by Equation 5.3 and 5.4.

$$\mathcal{P}_{bkg}(\Delta t_i) = f_{bkg}^\delta \delta(\Delta t_i - \mu_{bkg}^\delta) + (1 - f_{bkg}^\delta) \frac{1}{2\tau_{bkg}} e^{-|\Delta t_i - \mu_{bkg}^{bkg}|/\tau_{bkg}} \quad (5.3)$$

$$\mathcal{P}_{ol}(\Delta t_i) = G(\Delta t_i, \sigma_{ol}) \quad (5.4)$$

where $\delta(\Delta t_i - \mu_{bkg}^\delta)$ is Dirac δ function and G is single Gaussian. The outlier component is to improve the fit quality with large Δt events.

5.1 Vertex Resolution Model

The Equation 5.2 presents an ideal distribution of Δt_i for each event without considering the difference between measured and the true position of the vertex. The difference can be described by introducing resolution functions, turning Equation 5.2 into Equation 5.5.

$$\begin{aligned} \mathcal{P}(\Delta t_i, q_i, f_i^{sig}, \mathcal{S}, \mathcal{A}) = & (1 - f_{ol})[f_{sig} \mathcal{P}_{sig}(\Delta t_i) \otimes R_{sig}(\Delta t_i) \\ & + (1 - f_{sig}) \mathcal{P}_{bkg}(\Delta t_i) \otimes R_{bkg}(\Delta t_i)] \\ & + f_{ol} \mathcal{P}_{ol}(\Delta t_i) \otimes R_{ol}(\Delta t_i) \end{aligned} \quad (5.5)$$

The R_{sig} stands for the resolution function for signal events, which receives smearing effect from CP and tag side separately, namely R_{cp} and R_{tag} . The treatment of CP side and tag side is different because of vertexing strategies. For CP side, vertex of B^0 is reconstructed by fully fitting all the daughter particles. Instead, in tag side, there's no full reconstruction of B^0 so vertex fit is applied for the selected charged tracks in the rest-of-event. The background events have its own resolution model which is independent from CP violation parameters. The outlier is used to smooth fit for large Δt events. In the low statistics case as the current luminosity is, the outlier is not included in the fit to have a more realistic model for data.

For signal events, the resolution functions are studied for CP -side and tag-side based on each possible degradation such as detector resolutions, effect of tracks from non-primary B vertex and so on. Such a method is used in Belle analysis and named as artificial model. Details are summarized in [41]. Considered the vertex position difference Δz for signal events as shown in Equation 5.6.

$$\Delta z = \Delta z' + (z_{cp} - z'_{cp}) - (z_{tag} - z'_{tag}) \quad (5.6)$$

where the primed ones stands for physics truth of the position and the non-primed is the measured value, the resolution function receives contribution from both CP and tag-side effects. In the meantime, the resolution functions on both sides also depend on the applied constraint. Considering that the fine structure of IP profile is not yet fully understood and small discrepancies have been observed between data and simulation[42], there's no IP constraint applied for both sides in vertex fit, which avoids potential bias from IP profile under this low statistical situation. The combined contributions can be presented as Equation 5.7.

$$R_{sig} = R_{cp} \otimes R_{tag} \quad (5.7)$$

5.1.1 CP -side resolution function

CP -side vertex is fitted with all tracks from a reconstructed B^0 , thus the resolution models only depend on detectors' effect. For each event, the resolution effect can be different based on event-by-event reconstruction quality, primarily presented by the reduced χ^2 called χ^2/N from *TreeFit*, which N is the degree of freedom of the fit. The distribution of χ^2/N in data are shown in Figure 5-1.

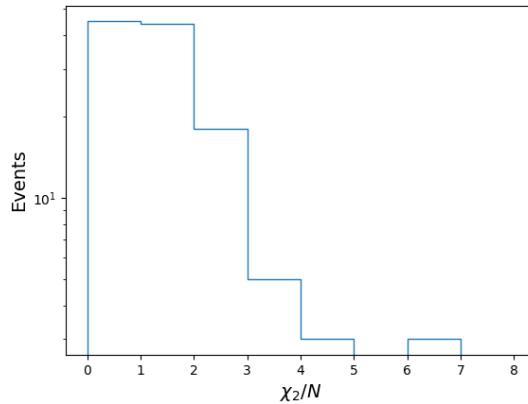


Figure 5-1: χ^2/N of selected events from data.

Therefore, we model the resolution functions on *CP*-side by using a double Gaussian function, where the mean is fixed to zero and the standard deviation is scaled by χ^2/N and the error of reconstructed vertex $\sigma_{z_{cp}}$, as shown in Equation 5.8.

$$R_{cp}(\delta z_{cp}) = (1 - f_{cp}^{tail})G(0, s_{cp}^{main}) + f_{cp}^{tail}G(0, s_{cp}^{tail}) \quad (5.8)$$

where s_{cp}^{main} and s_{cp}^{tail} are defined in Equation 5.9.

$$\begin{aligned} s_{cp}^{main} &= (s_0^{main} + s_1^{main} \cdot \chi_{cp}^2/N) \cdot \sigma_{z_{cp}} \\ s_{cp}^{tail} &= (s_0^{tail} + s_1^{tail} \cdot \chi_{cp}^2/N) \cdot \sigma_{z_{cp}} \end{aligned} \quad (5.9)$$

The dependence of resolution models on χ^2/N is shown in Figure 5-2. Restrictively speaking, the *CP*-side resolution for $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ is slight different from $B^0 \rightarrow J/\psi K_S^0$, due to the absence of the direct charged tracks from the B^0 vertex. The modification of the resolution function on *CP*-side will be further studied when more data becomes available in future. Given the current low statistics, the Equation 5.8 works well as an approximation. By fitting the resolution function using *signal MC* on *CP*-side, the parameters are fitted which are listed in Table 5.1.

Table 5.1: Parameters in R_{cp} .

f_{cp}^{tail}	0.07424 ± 0.0008
s_0^{main}	0.9151 ± 0.0077
s_1^{main}	0.2142 ± 0.0064
s_0^{tail}	2.0477 ± 0.0779
s_1^{tail}	1.3470 ± 0.0720

5.1.2 Tag-side resolution function

For the tag-side, the vertexing is done by using *KFit* and no IP constraint used. Due to the charged tracks from non-primary B vertex, the resolution functions on tag-side not only receives contribution from detectors' effect R_{det}^{tag} but also the resolution degradation from secondary vertex, called R_{np}^{tag} . To the contrary, if all tracks that

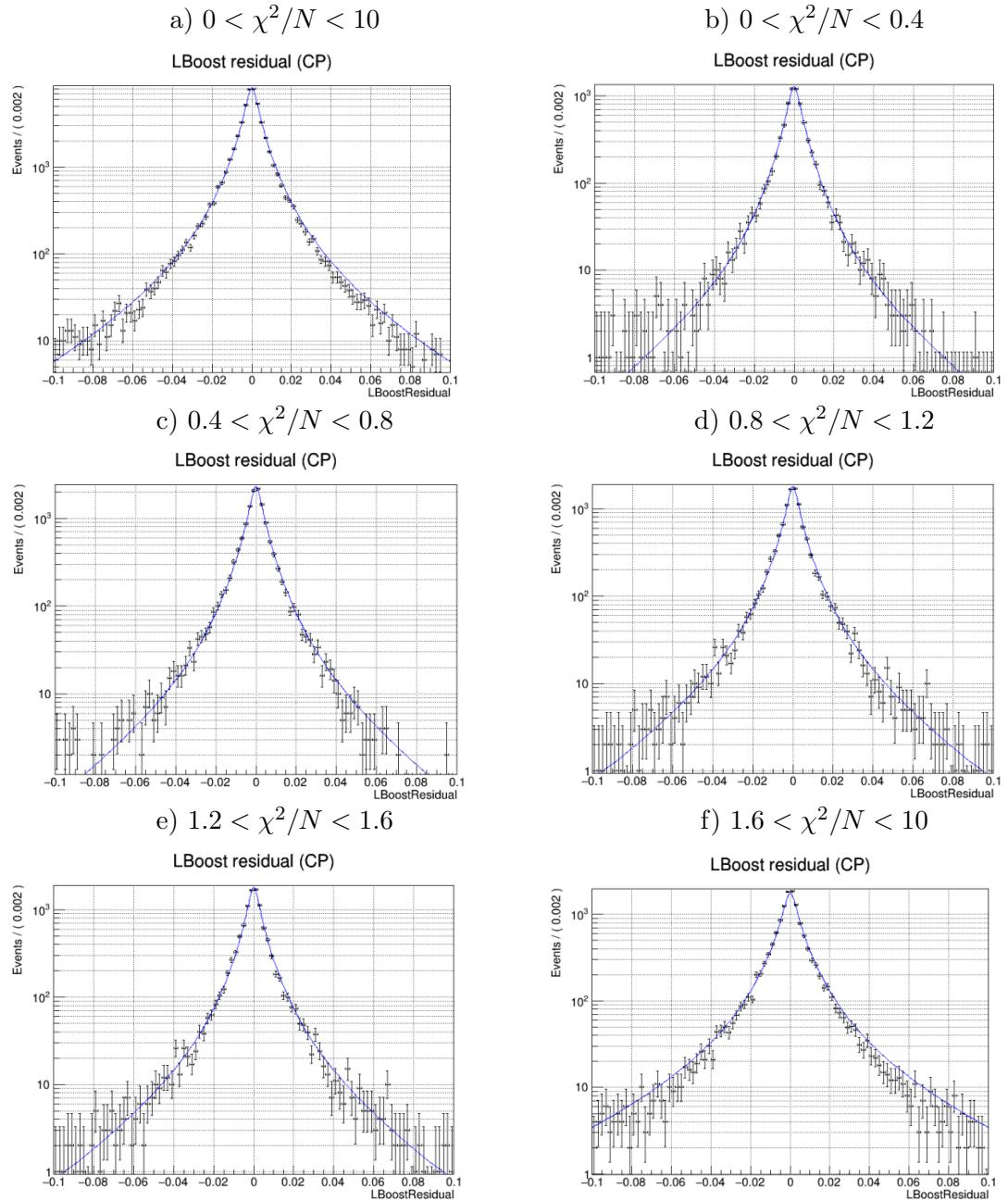


Figure 5-2: Resolution functions on *CP*-side, which shows dependence on the χ^2/N

are used for tag-side vertexing are primary tracks, the resolution will only be affected by the detectors' effect. The vertex position difference is defined as Equation 5.10. Therefore, the effects from both detectors and non-primary tracks contributes to the total resolution on tag-side as Equation 5.11 shows.

$$\begin{aligned} z_{tag} - z'_{tag} &= (z'_{tag} + \delta z_{tag}^{det} + \delta z_{tag}^{np}) - z'_{tag} \\ &= \delta z_{tag}^{det} + \delta z_{tag}^{np} \end{aligned} \quad (5.10)$$

$$R_{tag}(z_{tag} - z'_{tag}) = R_{det}^{tag}(\delta z_{tag}^{det}) \otimes R_{np}^{tag}(\delta z_{tag}^{np}) \quad (5.11)$$

Similarly to CP -side resolution function, detectors' effect is presented in Equation 5.12

$$R_{det}^{tag}(\delta z_{tag}^{det}) = (1 - f_{tag}^{tail})G(0, s_{tag}^{main} \cdot \sigma_{z_{tag}}) + f_{tag}^{tail}G(0, s_{tag}^{tail} \cdot \sigma_{z_{tag}}) \quad (5.12)$$

where main and tail Gaussian functions have the same central value at zero, but the standard deviation is scaled by χ_{tag}^2/N on the tag-side as shown in Equation 5.13.

$$s_{tag}^{main/tail} = s_0^{main/tail} + s_1^{main/tail} \cdot \chi_{tag}^2/N \quad (5.13)$$

Technically R_{det}^{tag} can be fitted with MC samples of which tag-side tracks are all from primary vertex. After obtaining the fitted parameters of R_{det}^{tag} , R^{tag} will only be dependent on R_{np}^{tag} . The fit model of R_{np}^{tag} is shown in Equation 5.14. It consists of three functions, including one Dirac δ function and two single-side exponential functions E_p and E_n . The $E_p(x, \tau_p) = (1/\tau_p)e^{-x/\tau_p}$ when $x > 0$ and the $E_n(x, \tau_n) = (1/\tau_n)e^{x/\tau_n}$ when $x < 0$. The exponential factors in both positive and negative components are scaled by the tag-side vertex uncertainty $\sigma_{z_{tag}}$.

$$R_{np}^{tag}(\delta z_{tag}^{np}) = f_\delta \delta(\delta z_{tag}^{np}) + (1 - f_\delta)[f_p E_p(\delta z_{tag}^{np}, \tau_p \cdot \sigma_{z_{tag}}) + (1 - f_p)E_n(\delta z_{tag}^{np}, \tau_n \cdot \sigma_{z_{tag}})] \quad (5.14)$$

Also, since tag-side has no dependence on how CP -side is reconstructed, the res-

olution functions on tag-side are almost mode-independent. Thus these parameters are obtained by fitting to the control sample data. The details about control sample study are here[42]. The fit plots for tag-side resolution functions are shown in Figure 5-3 and 5-4. The parameters obtained from the fit are listed in Table 5.2 and Table 5.3.

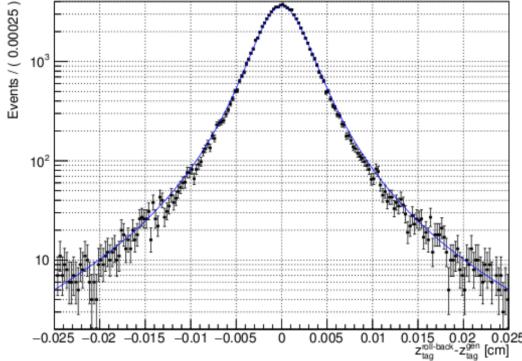


Figure 5-3: R_{det}^{tag} fit

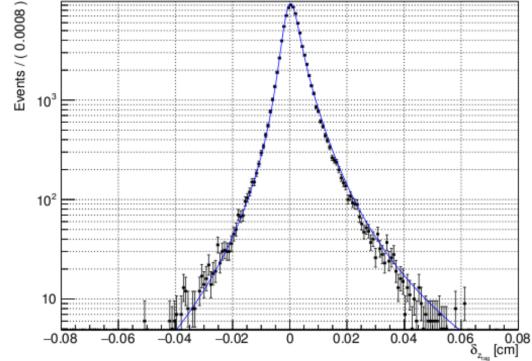


Figure 5-4: R_{np}^{tag} fit

Table 5.2: Parameters in R_{det}^{tag}

f_{tag}^{tail}	0.0523 ± 0.0025
s_0^{main}	1.1446 ± 0.0061
s_1^{main}	0.0443 ± 0.0022
s_0^{tail}	3.4480 ± 0.0897
s_1^{tail}	0.2666 ± 0.0276

Table 5.3: Parameters in R_{np}^{tag}

f_δ	0.6256 ± 0.0049
f_p	0.8316 ± 0.0051
τ_n	2.9141 ± 0.0758
τ_p	2.4846 ± 0.0269

The boost direction of each event is not constant event-by-event, so the position of vertex may not be optimized by calculating $\Delta t_i = \Delta z / \beta \gamma c$. This effect can be reduced by replacing vertex position difference on z-axis with the relative distance along the boosting direction, or introducing another resolution function called R_k [43]. The R_k has not been implemented in Belle II resolution model. Therefore, Δz projection on the boosted direction of each event is used for reducing this kinematics effect on resolution function.

5.1.3 Background events Δt distribution

The R_{bkg} is uncorrelated to vertex reconstruction method approximately. Because the background mainly comes from continuum events passing the selection, it's reasonable to model its resolution by a Gaussian-like function. A double-Gaussian with its standard deviation scaled by the measured uncertainties from both sides is used as Equation 5.15. To be noted, unlike resolution functions on CP or tag-side, the standard deviations of the double Gaussian are scaled by both the vertex position uncertainties $\sigma_{z_{cp}}$ and $\sigma_{z_{tag}}$.

$$R_{bkg} = (1 - f_{tail}^{bkg})G(\Delta t_i, \sigma_{main}^{bkg} \sqrt{\sigma_{z_{cp}}^2 + \sigma_{z_{tag}}^2}) + f_{tail}^{bkg}G(\Delta t_i, \sigma_{tail}^{bkg} \sqrt{\sigma_{z_{cp}}^2 + \sigma_{z_{tag}}^2}) \quad (5.15)$$

The background events Δt shapes $\mathcal{P}_{bkg} \otimes R_{bkg}$ can be determined by fitting to side-band data. There are totally seven floating parameters which are listed in Table 5.4 with fitted values using 60 sideband events at $M_{bc} < 5.26$ GeV, shown in Figure 5-5

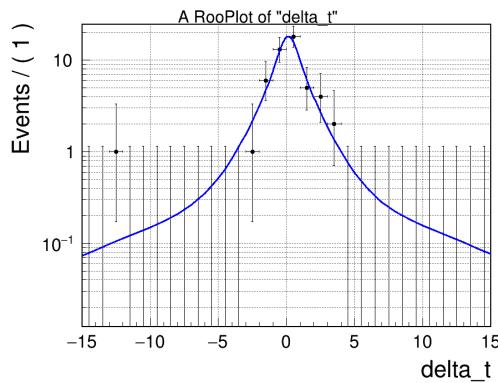


Figure 5-5: $\mathcal{P}_{bkg} \otimes R_{bkg}$ fit using 60 sideband events at $M_{bc} < 5.26$ GeV.

μ_{δ}^{bkg}	0.1310 ± 0.1902
μ_l^{bkg}	0.1638 ± 0.5030
τ_{bkg}	1.0541 ± 0.4370
f_{δ}^{bkg}	0.5861 ± 0.2570
f_{tail}^{bkg}	0.0417 ± 0.0408
σ_{main}^{bkg}	1.4348 ± 0.3940
σ_{tail}^{bkg}	28.0930 ± 8.8221

Table 5.4: Parameters in Background Δt distribution.

5.2 Flavor Tagging

In order to determine the flavor of tag side B^0 , flavor tagging algorithm is being developed. The flavor tagging uses information from μ^\pm, π^\pm, K^\pm and Λ which are categorized into 13 different types as illustrated in Figure 5-6.

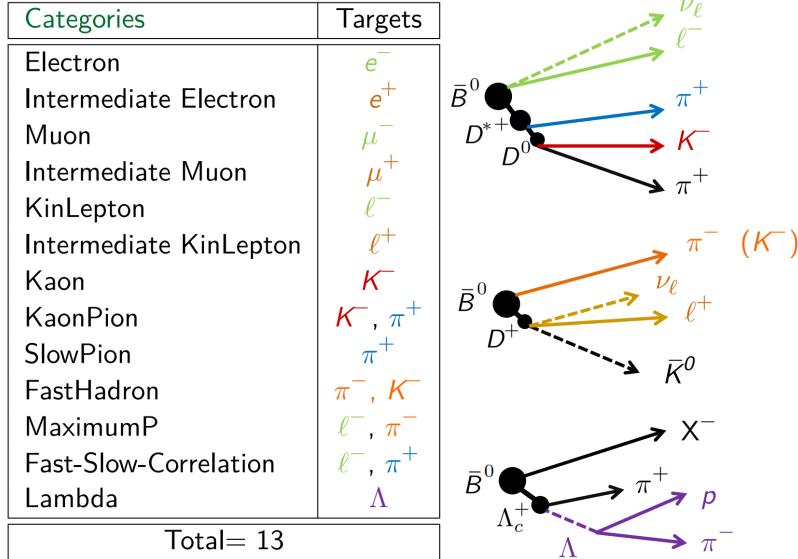


Figure 5-6: Particles and their categories used in flavor tagging algorithm[44].

For each particle that has been used from above categories, PID and kinematics information are extracted and feed to the combiner as training variables, to obtained a classifier response corresponding to each category. Then for all responses from

these categories, a total classifier is trained to present the likelihood of flavor q . This algorithm is called category-based method and used in this thesis. After the reconstruction on the CP -side B^0 is done, the rest-of-events tracks used to form the particle lists are selected⁵⁻⁶. The FastBDT as the back-end algorithm is chosen for performing training on the classifier of flavor tagging. Targeted variable is true q of tag-side neutral B in MC. To minimize impact of the reconstruction performance on CP -side, MC sample of $B^0 \rightarrow \nu\nu$ is used as the training sample where the final state in CP -side are completely invisible.

Considering the limited power of flavor tagging accuracy, there is a certain fraction of events that are wrongly tagged, among all events that can be flavor tagged based on the final states. Thus, the flavor tagging efficiency ϵ and wrong tag fraction w are defined, respectively. Taking into account of the performance of flavor tagging, the observed distribution of Equation 5.2 becomes Equation 5.16.

$$\mathcal{P}_{sig}^{obs}(\Delta t, q, \epsilon, w) = \frac{e^{-|\Delta t|/\tau_{B^0}}}{4\tau_{B^0}} \epsilon \left\{ 1 - q \cdot \Delta w + q(1 - 2w) \cdot [\mathcal{S} \sin(\Delta M_d \Delta t) + \mathcal{A} \cos(\Delta M_d \Delta t)] \right\} \quad (5.16)$$

Compared to the original, the term with \mathcal{S} and \mathcal{A} is scaled by factor $r \equiv |1 - 2w|$, defined as the dilution factor. The statistical uncertainty of \mathcal{S} now receives contribution from ϵ and w : $\delta(\mathcal{S}) \propto \frac{1}{\epsilon(1-2w)N_{rec}}$, which means it's quite important to correctly measure ϵ and w . The uncertainty of w is much larger than ϵ which makes w an important source of systematic uncertainty. The validation of flavor tagger using flavor specific control sample is summarized here[44]. The w for each single event is defined as a probability which can be presented by the average wrong tag fraction within a close range. The binned values of dilution factor r is defined for the calculation of w as $[0, 0.1, 0.25, 0.5, 0.625, 0.75, 0.875, 1]$, which is called r -bin. For all events that have been successful tagged, they are projected into histogram of r -bin, and w is calculated in each bin by the fraction of events with $q \cdot r$ opposite to its MC flavor. The distribution of $q \cdot r$ is shown in Figure 5-7 using *signal MC* of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$.

Besides, w can be different between B^0 and $\bar{B^0}$, where $\bar{w} = (w_{B^0} + w_{\bar{B^0}})/2$ and

$\Delta w = w_{B^0} - w_{\bar{B}^0}$. Due to the small value of Δw , the contribution from Δw can be treated as zero in Equation 5.16. Similarly, for ϵ , the values calculated based on each r -bin are used in Equation 5.16 as well. The difference $\mu = \epsilon_{B^0} - \epsilon_{\bar{B}^0}$ is about 1 % to 2 % in each r -bin, thus treated as zero, too. The distributions of w , Δw , ϵ , and μ in each r -bin are shown in Figure 5-8.

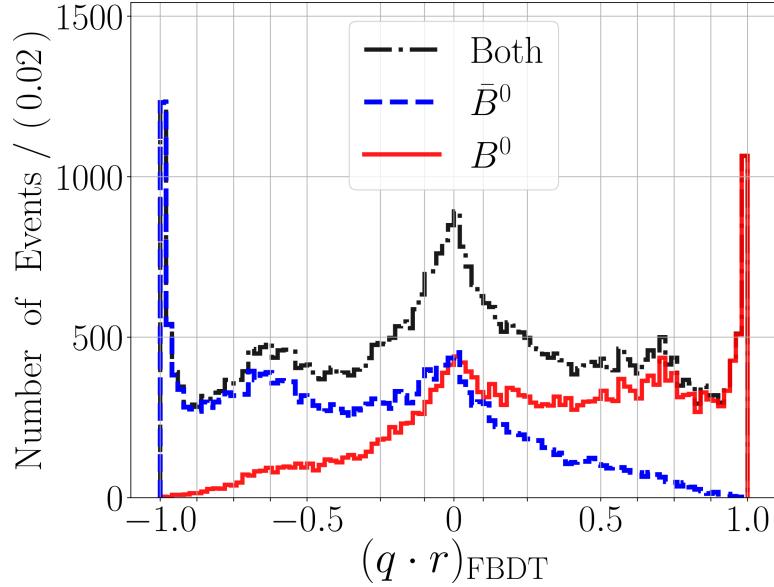


Figure 5-7: The distribution of flavor tagger output ($q \cdot r$) for both tag-side of B^0 and \bar{B}^0

5.3 CP Fitter

The parameters that are needed for measuring \mathcal{S} and \mathcal{A} are studied and obtainable. Using observed Δt distribution from selected events, Equation 5.5 can be fitted using unbinned maximum likelihood fit which takes Δt , signal fraction f_{sig} , the flavor charge q as observables. In the meantime the vertexing error $\sigma_{z_{cp}}$, $\sigma_{z_{tag}}$ and χ^2/N are used as event-by-event conditional variables that are accessed during the fitting. For Belle II, a new CP fitter is developed based on Python and RooFit, which is naturally easy to use and maintained with BASF2. The fitter requires a configuration files which contains all the parameters' definitions including their ranges, initial values, floating states and uncertainties.

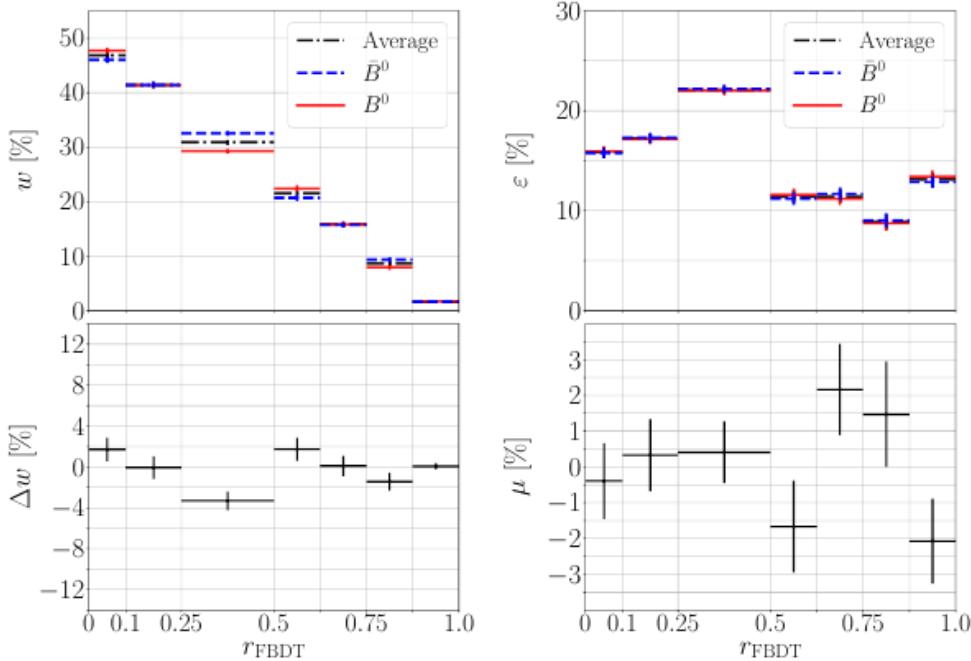


Figure 5-8: The flavor tagging efficiency, wrong tagging fraction, and their difference between different flavors in each of the r -bin.

5.4 Blind analysis and fit

As a required procedure to make sure the CP parameters are measured without bias due to the preconceived results, a blind analysis procedure is conducted before the fit is actually performed using the experimental data. The blind fit procedure includes the CP fit on *signal MC* and *generic MC*, with different number of events used. To check the reliability of fit result from CP fitter, a linearity test and toy MC study are also performed.

5.4.1 CP fit on MC samples

Using CP fitter, we first perform the CP fit on events in *signal MC* and *generic MC*. The *signal MC* and *generic MC* are generated with phase-space model which contains zero CP violation ($\mathcal{S} = \mathcal{A} = 0$). The events that pass the selections in Table 5.5 are used for CP parameters fit.

We have 10000 (8873 passing selections) events from signal sample and 415 (373

Observables	Selections
Δt	$-70 < \Delta t < 70$ ps
$CP\text{-side } \chi^2/N$	$0 < (\chi^2/N)_{cp} < 8$
tag-side χ^2/N	$0 < (\chi^2/N)_{tag} < 50$
$\sigma_{z_{tag}}$	$\sigma_{z_{tag}} < 0.1$ cm
signal region	$5.27 < M_{bc} < 5.29$ GeV and $ \Delta E < 0.1$ GeV

Table 5.5: The selection criteria for events that are used for CP parameters fit.

passing selections) events from 1 ab^{-1} *generic MC* to fit CP parameters. To mimic the events number expected in data sample, 30 events randomly taken from *generic MC* are used to perform the fit as well. The plots are shown in Figure 5-9, 5-10 and 5-11. The fit results of \mathcal{S} and \mathcal{A} are summarized in Table 5.6.

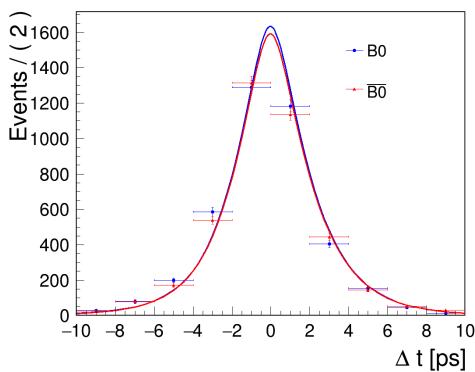


Figure 5-9: CP fit on 8873 *signal MC*.

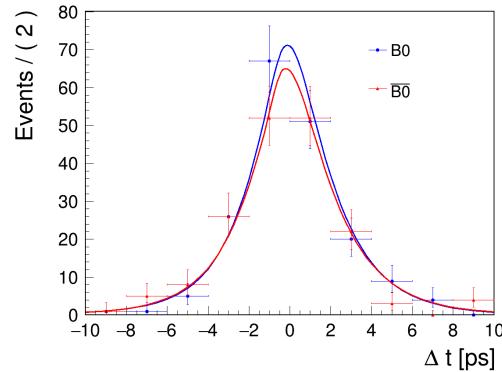


Figure 5-10: CP fit on 373 *generic MC*.

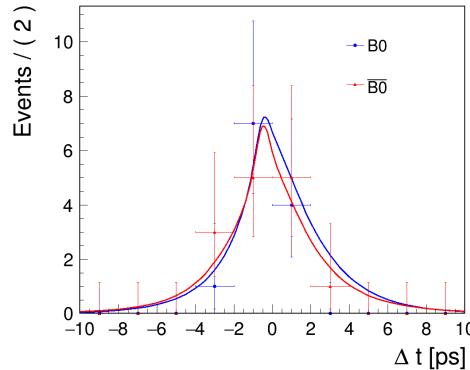


Figure 5-11: CP fit on 30 *generic MC*.

The fit results are consistent with expectation in non- CP violation from MC input, and the statistical uncertainties has the tendency $\delta \propto 1/\sqrt{N}$ as poission distribution, where N is events number used for CP fit. To test fit on non-zero CP violating MC,

Table 5.6: The CP fit results using *signal MC* and *generic MC* with only statistical uncertainties.

Sample (events)	\mathcal{S}	\mathcal{A}
<i>signal MC</i> (8873)	$\sin(2\phi_1) = 0.00 \pm 0.04$	$\mathcal{A} = -0.01 \pm 0.02$
<i>generic MC</i> (373)	$\sin(2\phi_1) = 0.00 \pm 0.21$	$\mathcal{A} = -0.05 \pm 0.07$
<i>generic MC</i> (30)	$\sin(2\phi_1) = 0.20 \pm 0.85$	$\mathcal{A} = -0.06 \pm 0.30$

the fit on $B^0 \rightarrow J/\psi K_S^0$ *signal MC* is also done, the details of events selection as well as fit model determination can be found[42]. The fit result over 10000 events is shown in Figure 5-12, which results in $\sin(2\phi_1) = 0.70 \pm 0.05$ and $\mathcal{A} = -0.01 \pm 0.02$. The results agree with the input.

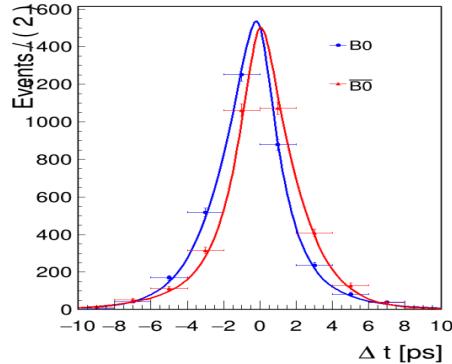


Figure 5-12: CP fit over 10000 $B^0 \rightarrow J/\psi K_S^0$ *signal MC*.

5.4.2 Linearity Test

To validate the CP fit linearity, a series of toy MC samples is generated, which the χ^2 from vertex fit, events number N and vertex errors on CP and tag-side are sampled from the distribution of *signal MC*. The resolution functions parameters are kept as same as CP fit on *generic MC*. The input \mathcal{A} is set to zero while the input value of $\sin(2\phi_1)$ is running from 0.1 to 0.9. Each dataset contains 10000 events. The dependence between input and output are shown in Figure 5-13. The linearity fit shows a good agreement between input and output.

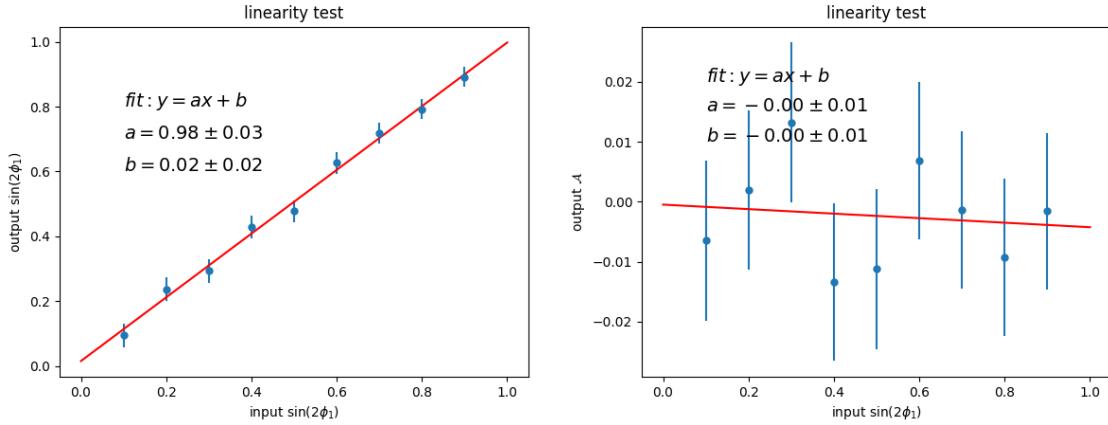


Figure 5-13: Linearity test of CP fit.

Also, we fix $\sin(2\phi_1)$ at zero while floating A from 0.1 to 0.9, the dependence between input and output are as Figure 5-14 shows. The linearity fit shows a good agreement as well.

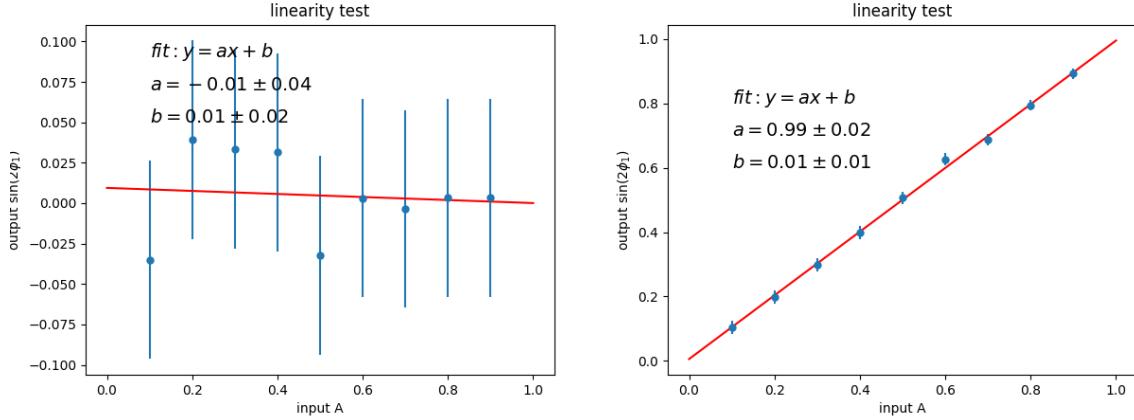


Figure 5-14: Linearity test of CP fit.

5.4.3 Toy MC Fit Pull

In order to check the fit bias with input-output method, a series of 1000 dataset of toy MC has been created containing about 26 events in each. The event number is set based on the expected number from signal region in data after the selection. The χ^2 from vertex fit, events number N and vertex errors on CP and tag-side are sampled from the distribution of data. The fit to dataset is performed with zero input $\sin(2\phi_1)$

and \mathcal{A} as floating parameters. We expect to use the normal distribution to fit the pull of $\sin(2\phi_1)$ and \mathcal{A} .

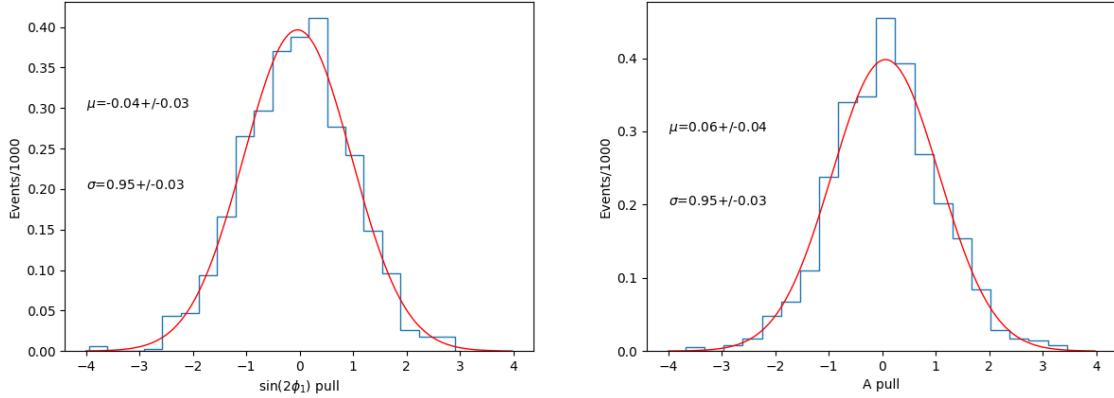


Figure 5-15: Pull of $\sin(2\phi_1)$ and \mathcal{A} fitted with the standard normal distribution.

The fit results shows a good recovery of input $\sin(2\phi_1)$ and \mathcal{A} with no clear bias is spotted.

5.4.4 Lifetime and Δm_d Fit

Before looking at CP parameters in data, we need to check if the physics parameters are consistent when setting the CP fitter to fit them in float. To test lifetime fit, first we use 10000 *signal MC* events which is generated by $\tau_{B^0} = 1.520$ from PDG value. The $\sin(2\phi_1)$ and \mathcal{A} are fixed at zero during the fit, for which the generator level CP violation is zero. This is equivalent fit to Equation 5.17.

$$\mathcal{P}(\Delta t, \tau_{B^0}) = \frac{e^{-|\Delta t|/\tau_{B^0}}}{4\tau_{B^0}} \quad (5.17)$$

The fit result on *signal MC* is 1.537 ± 0.024 ps which is consistent with the input. We perform the lifetime fit on data in signal region, and the CP parameters are fixed based on PDG values to: $\sin(2\phi_1) = 0.69$ and $\mathcal{A} = 0$. The fitted lifetime from $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ is 1.431 ± 0.382 ps. The result is consistent with PDG value. The distribution of Δt in lifetime fit is shown as Figure 5-16. The B^0 and B^+ lifetime fit

using control sample is also performed and summarized in here[42]. The results are consistent with PDG values as input in MC generator.

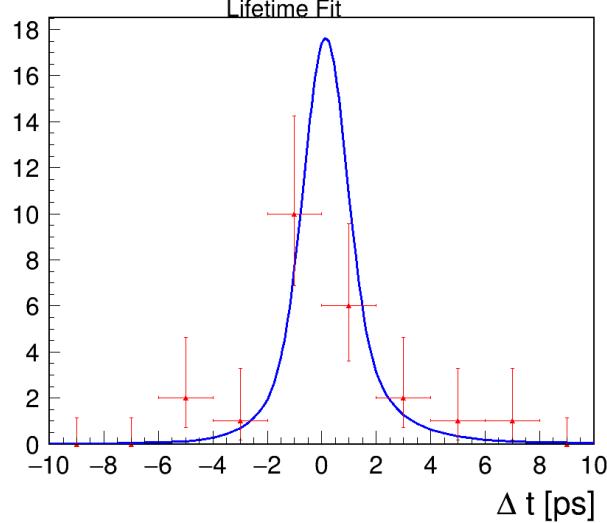


Figure 5-16: Lifetime fit on data

To test the fit on physics parameter Δm_d , we generate 200 toy MC sets of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ with input $\Delta m_d = 0.507 \text{ GeV}/c^2$ where each set contains 26 events as same as data. The fit result is close to normal distribution and the pull of Δm_d is shown in Figure 5-17.

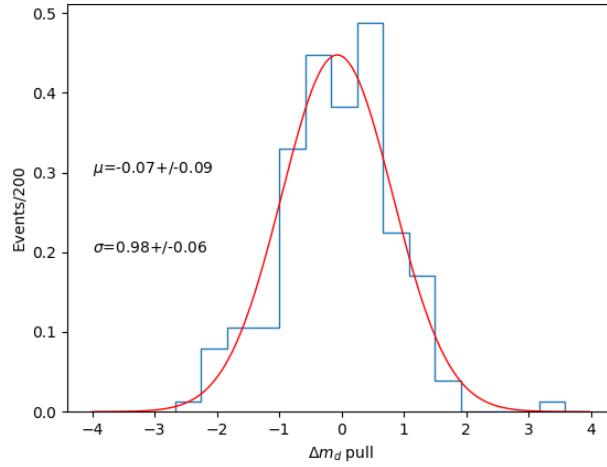


Figure 5-17: Pull of Δm_d

5.5 CP fit on data

After the CP fit procedures are reviewed by Belle II collaboration, the permission of measuring CP parameters using 62.8 fb^{-1} Belle II data is granted. The events number used for the CP fit is 26, and the fit result is shown Figure 5-18.

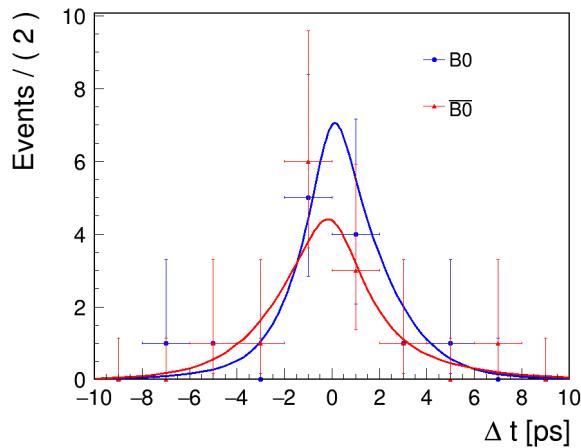


Figure 5-18: The CP fit from data.

The results of CP parameters are:

$$\begin{aligned} \sin(2\phi_1) &= 0.82 \pm 0.85(\text{stat}) \\ \mathcal{A} &= -0.21 \pm 0.28(\text{stat}) \end{aligned} \tag{5.18}$$

5.6 Systematic Uncertainty

The systematic uncertainty that affects the fit results may come from many aspects of the measurement setup. The contributions at the current stage of Belle II are summarized as Table 5.7.

For the listed sources of systematic uncertainty, if the parameters are defined with MC study, we float the value by $\pm 2\sigma$ of their uncertainty, and if the parameters are defined by data, we float the value by $\pm 1\sigma$, where σ is the uncertainty of the parameter. The impact of CP parameters are separately estimated from each sources with positive and negative differentials. By adding in quadrature of each term, the

Sources	$\delta\mathcal{S}$	$\delta\mathcal{A}$
signal fraction	0.04	0.04
background Δt shapes	0.04	0.04
signal Δt shapes	0.03	0.01
fit bias	0.01	0.01
flavor tagging	0.004	0.004
physics parameters	0.007	0.001
<i>KsFinder</i> impact on data	0.004	0.001
vertex reconstruction	0.019	0.021
Total	0.07	0.06

Table 5.7: The contributions of each source of systematic uncertainty.

overall systematic uncertainty is obtained.

The signal resolution functions' parameters are determined from MC study for signal component. The impact on fit results is summarized as follows Table 5.8.

Table 5.8: systematic uncertainty from signal Δt shapes

source	$+\delta\mathcal{S}$	$+\delta\mathcal{A}$	$-\delta\mathcal{S}$	$-\delta\mathcal{A}$
f_{cp}^{tail}	-0.000096	-0.000057	0.000014	0.000056
s_0^{main}	0.005443	0.001299	-0.005675	-0.001404
s_1^{main}	0.019934	-0.000903	-0.020204	0.000633
s_0^{tail}	-0.003233	-0.001623	0.00327	0.001596
f_{tag}^{tail}	0.00314	-0.001257	-0.003117	0.001266
s_0^{main}	0.002011	-0.001395	-0.001956	0.001398
s_1^{main}	0.005059	-0.00084	-0.004969	0.000825
s_0^{tail}	-0.000135	-0.000393	0.00010	0.000435
s_1^{tail}	0.000101	0.000027	-0.000472	0.000129
f_δ	-0.007248	-0.000552	0.007231	0.000591
f_p	0.003037	0.004347	-0.003069	-0.004314
τ_n	-0.00101	-0.002841	0.000937	0.00294
τ_p	0.004497	0.002502	-0.004648	-0.002478

The background Δt shapes' parameters are determined from data sideband $M_{bc} < 5.26$ GeV. The impact on fit results is summarized as follows:

Table 5.9: systematic uncertainty from background Δt shapes

source	$+\delta\mathcal{S}$	$+\delta\mathcal{A}$	$-\delta\mathcal{S}$	$-\delta\mathcal{A}$
μ_δ^{bkg}	-0.014294	-0.016581	0.006758	0.006537
μ_l^{bkg}	-0.002798	-0.012567	0.003789	0.012783
τ_{bkg}	0.001377	0.001689	-0.004159	0.000085
f_δ^{bkg}	-0.011315	0.001365	0.011187	-0.001395
f_{tail}^{bkg}	-0.002661	0.00153	0.00248	-0.001368
σ_{main}^{bkg}	0.0207015	0.022041	-0.0236175	-0.01569
σ_{tail}^{bkg}	-0.000275	-0.000159	0.000179	0.000141

The flavor tagging parameters wrong tagging fraction w in each rbin is determined from *signal MC*. The impact in each rbin on fit results is summarized as follows:

Table 5.10: systematic uncertainty from wrong tagging fraction

source	$+\delta\mathcal{S}$	$+\delta\mathcal{A}$	$-\delta\mathcal{S}$	$-\delta\mathcal{A}$
w_1	-0.0018919	0.001911	0.0018549	-0.002004
w_2	-0.0016448	0.001104	0.0016085	-0.001155
w_3	-0.0004899	0.001344	0.0004726	-0.001341
w_4	0.0006556	0.000264	-0.0006542	-0.000255
w_5	-0.0001228	0.000204	0.0001225	-0.000195
w_6	0.0000948	0.000054	0.0000957	-0.000045
w_7	0.0001911	-0.000396	-0.0001907	0.000402

The physics parameters Δm_d and τ_{B^0} uncertainties are included using the PDG average value. The impact on fit results is summarized as follows:

Table 5.11: systematic uncertainty from physics parameters

source	$+\delta\mathcal{S}$	$+\delta\mathcal{A}$	$-\delta\mathcal{S}$	$-\delta\mathcal{A}$
Δm_d	-0.001767	-0.000687	0.001778	0.000696
τ_{B^0}	-0.004561	-0.000546	0.004565	0.000555

The signal fraction is determined using 2D fit results of M_{bc} and ΔE from data. The impact on fit results is summarized as follows:

Table 5.12: systematic uncertainty from signal fraction

source	$+\delta\mathcal{S}$	$+\delta\mathcal{A}$	$-\delta\mathcal{S}$	$-\delta\mathcal{A}$
mu1_mbc	0.000822	-0.003888	-0.0007965	0.003849
sigma1_mbc	0.0004755	0.008442	-0.000628	-0.008733
m0_argus	-0.000707	0.00414	0.001448	-0.005781
c_argus	-0.005544	0.001449	0.000922	-0.000078
f1_de	0.0278255	0.020589	-0.0192365	-0.008409
f2_de	0.020809	0.017649	-0.0161285	-0.007005
mu1_de	-0.000443	-0.000153	0.0004955	0.000088
mu2_de	-0.000563	0.001446	0.0005905	-0.001446
mu3_de	-0.0031635	-0.000834	0.003354	0.000981
sigma1_de	-0.0001715	-0.000966	0.000206	0.000906
sigma2_de	-0.0031495	0.002958	0.0026345	-0.002475
sigma3_de	-0.001926	-0.00255	0.0024695	0.002985
a0_cheb	0.0009515	0.000057	-0.0008925	-0.000102
N_sig_f	-0.0046395	0.003987	0.004922	-0.003504

The fit bias uncertainties is determined by the fit error of 100000 *signal MC* events, which is $\delta\mathcal{S} = 0.009817$ and $\delta\mathcal{A} = 0.005702$.

Table 5.13: systematic uncertainty from fit bias

source	$\delta\mathcal{S}$	$\delta\mathcal{A}$
fit bias	0.009817	0.005702

Applying KsFinder cut at 0.74 based on MC study may introduce small impact on data due to the different response on the classifier between data and MC. Therefore the contribution of systematic uncertainty from KsFinder is considered. At cut value 0.74, the $\mathcal{R}_{B'}$ presenting MC and data signal yield ratio is $\mathcal{R}_{B^0} = 1.027 \pm 0.033$,

where the upper and lower limit is 1.060 and 0.994, respectively. These two ratios are applied on the signal fraction obtained by data to repeat the fit, and the difference of fit results compared to the original values are used as systematic uncertainty, see Table 5.14.

Table 5.14: systematic uncertainty from KsFinder.

source	$\delta\mathcal{S}$	$\delta\mathcal{A}$
$\mathcal{R}_{B^0} = 1.06$	0.004826	-0.000606
$\mathcal{R}_{B^0} = 0.994$	-0.000508	0.000007

For the contributions from vertex reconstruction, the impacts from the selections in Table 5.5 are considered. Given the fact that cut values in Table 5.5 are very loose and the statistics from data is very limited, the changing of the these values doesn't affect events collected from data so that systematic uncertainty can not be reflected correctly. Therefore, 1 ab^{-1} *generic MC* is used with the modified ranges to estimate the potential systematic uncertainty from vertex reconstruction. Besides, due to the absence of IP constraint in vertex fit, the impact from the IP constraint options as well as the potential bias are not considered. The summarized systematic uncertainties are listed in Table 5.15, where the zero values appear due to the unchanged events input under the modified ranges.

Table 5.15: systematic uncertainty from vertex reconstruction

source	$\delta\mathcal{S}$	$\delta\mathcal{A}$
$\sigma_{z_{tag}} < 0.05 \text{ cm}$	0.004369	-0.003599
$\sigma_{z_{tag}} < 0.15 \text{ cm}$	0.000000	0.000000
$\chi^2/N(CP) < 3$	0.018197	-0.020242
$\chi^2/N(CP) < 13$	0.000000	0.000000
$\chi^2/N(\text{tag}) < 40$	0.000000	0.000000
$\chi^2/N(\text{tag}) < 60$	0.000000	0.000000
$ \Delta t < 50 \text{ ps}$	0.003325	-0.000396
$ \Delta t < 90 \text{ ps}$	0.000000	0.000000
IP constraint	0.000000	0.000000

Chapter 6

Conclusion, discussion and prospect

The CP parameters measurement is performed based on the validation of analysis strategies by blind analysis. The blind fit is performed which shows a consistent fit result for CP parameters compared to the simulation input. The linearity and pull of the CP fit are checked to validate the reliability of the fit procedures. The fit result on B^0 lifetime using experiment data is also agreed with the current value in PDG with a relatively large statistical uncertainty due to the low statistics from data.

After the CP fit procedures are validated, the permission of CP fit using data is given by the Belle II collaboration, the CP parameters \mathcal{S} and \mathcal{A} using Belle II early data in 2019 and 2020 spring and summer is performed. The result is shown in Equation 6.1.

$$\begin{aligned}\mathcal{S} &= -\sin(2\phi_1) = -0.82 \pm 0.85(\text{stat}) \pm 0.07(\text{syst}) \\ \mathcal{A} &= -0.21 \pm 0.28(\text{stat}) \pm 0.06(\text{syst})\end{aligned}\tag{6.1}$$

The result agrees with the prediction of the Standard Model and the previous results from Belle and BaBar. The systematics study is performed considering the main contributing sources at the moment. The CP parameters precision from this study

is majorly limited due to the large statistical uncertainties from very low statistics, which leads to no clear evidence or hint on the NP effects.

What is worth of noticing is that main analysis tools required by performing the CP measurement on this channel are in a good stage of the development. The newly developed *KsFinder* contributes much in improving the signal significance by effectively rejecting fake K_S^0 , which is also useful in background rejection for other channels with K_S^0 in the final state. The model of the resolution of vertex positions has been studied using MC sample and sideband data, which is benefiting from the understanding of vertex reconstruction performance in the current Belle II detectors. To make a proper use of the reconstructed vertex information and perform CP fit compactly, a new CP fitter is built and being validated, which will serve as a multi-functional analysis tool for Belle II CP violation study in future.

6.0.1 Improvements in statistical uncertainty

This study has shown a good potential of performing CP measurement in Belle II for the incoming years with more and more data recorded. The precision on \mathcal{S} requires the large luminosity as shown in Figure 1-5 and the statistical uncertainty in this thesis fits in the scale. With 50 ab^{-1} luminosity from the full Belle II data sample in future, the statistical uncertainty is expected to be reduced. The CP fit on the MC sample with different mount of events used reflects that the statistical uncertainty is reduced proportionally around factor of $\frac{1}{\sqrt{N}}$, where N is the events used in CP fit. Therefore, the expected reduction of statistical uncertainty due to increased data sample with current reconstruction efficiency is predicted as Figure 6-1. The projected statistical uncertainty is slight reduced compared with the Belle result at 1 ab^{-1} . In future Belle II with full data sample, the estimated statistical uncertainty is about 0.0302 as shown in Figure 6-1, which is also improved compared to the original estimation at 0.037 from Belle II technical design report in Table 1.1. From Table 4.5, the current B^0 reconstruction efficiency is about 35% which could be further optimized mainly by improving K_S^0 efficiency from the current CDC-only tracking. Due to the enlarged VXD volume of Belle II, the expected K_S^0 acceptance region is larger than Belle,

which could contribute more on B^0 signal yield in future to help reduce the statistical uncertainty.

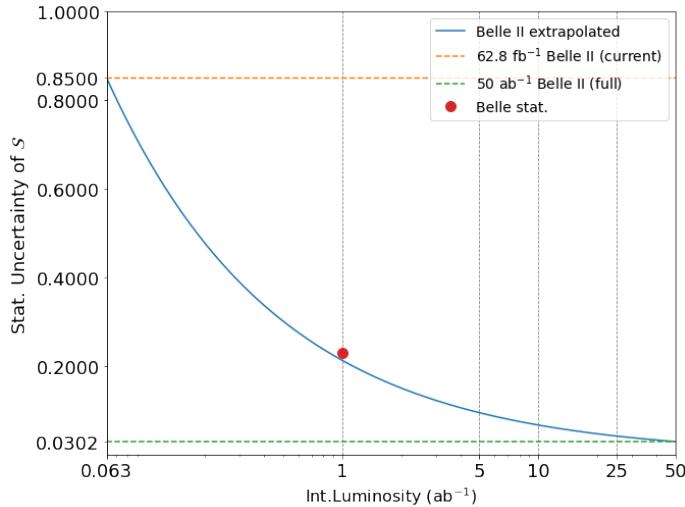


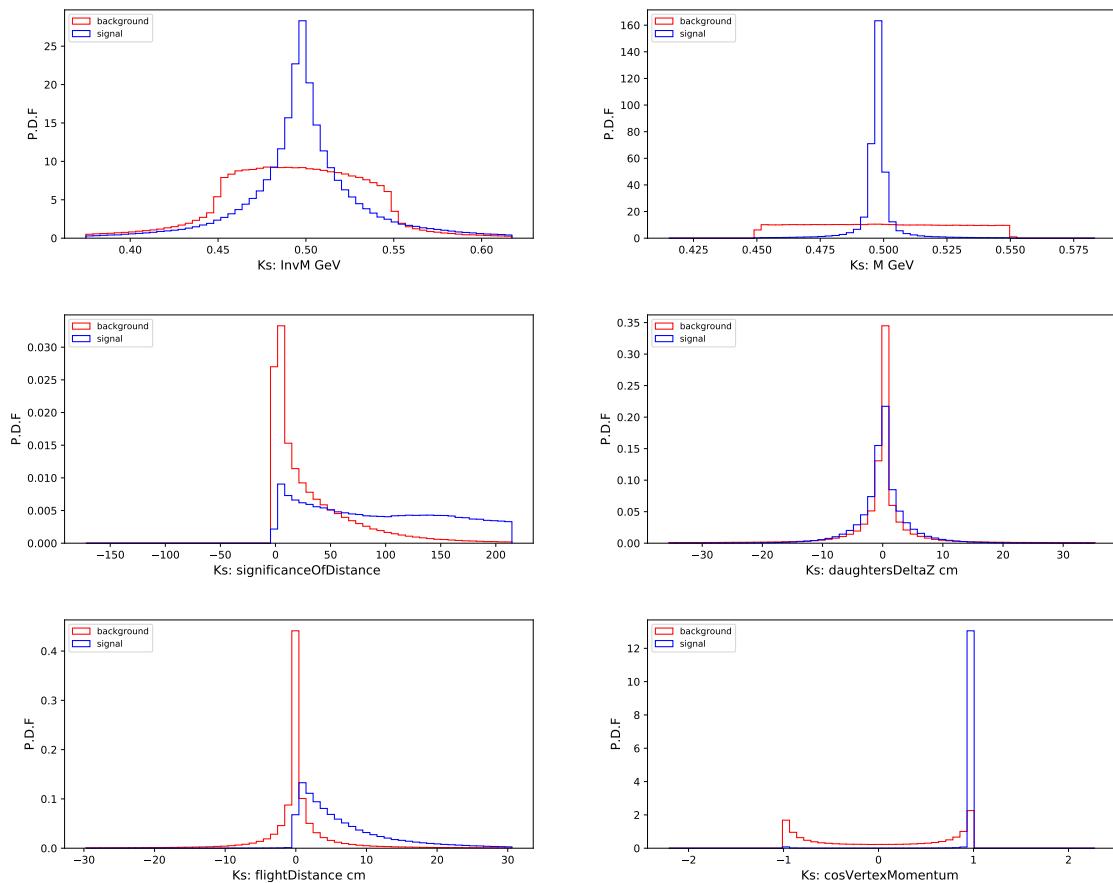
Figure 6-1: Statistical uncertainty of \mathcal{S} extrapolation based on the current result of $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ in Belle II.

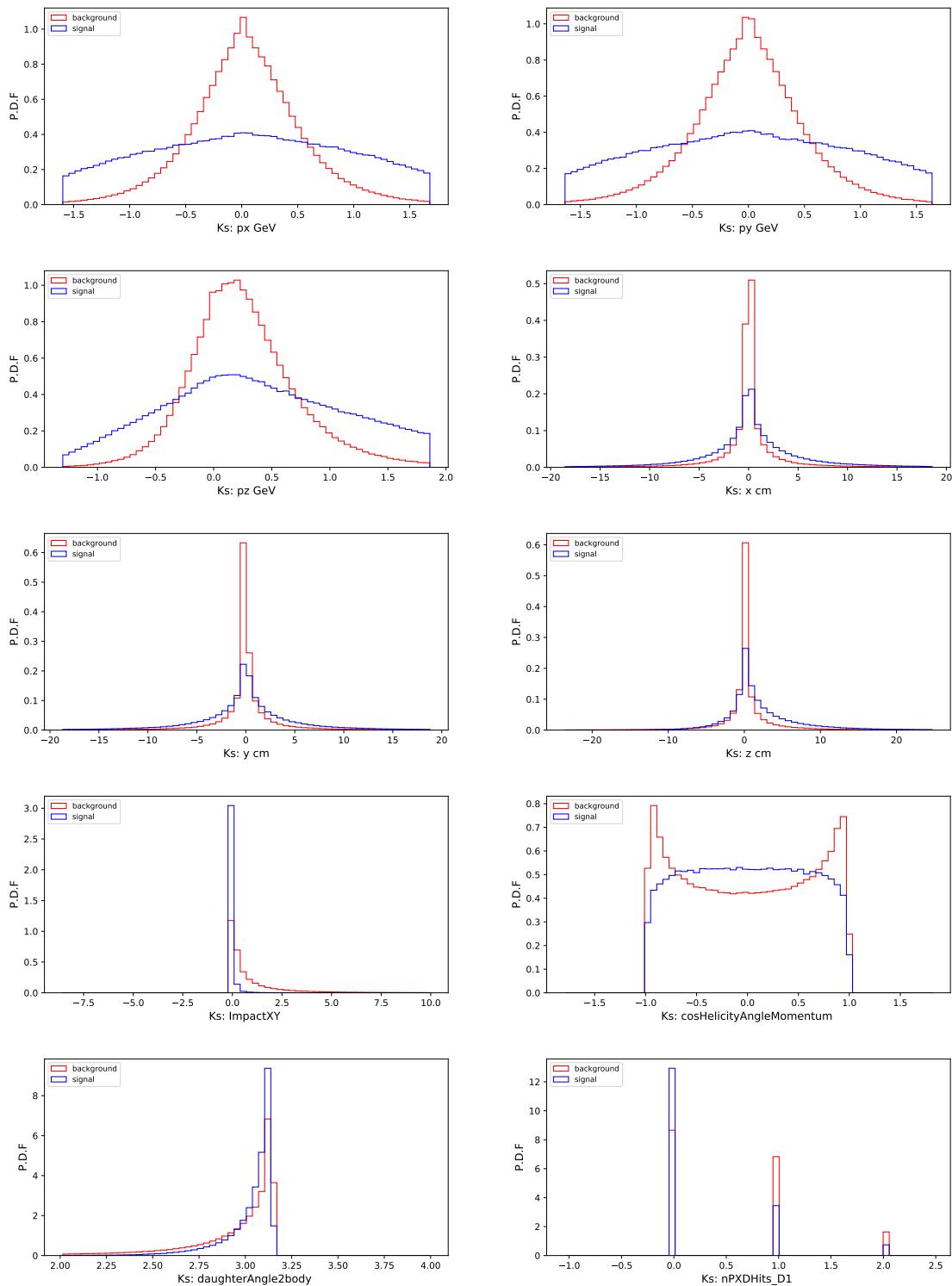
Along with the data collection continues, we will be able to finely test and improve the analysis strategies and tools to a better stage using data, such as improving the reconstruction efficiency and purity of B^0 when the luminosity ramps up to much higher level with much higher backgrounds. At integral luminosity at 50 ab^{-1} level, the statistical uncertainty of this decay on $\Delta\mathcal{S}$ would be trimmed down to a comparable value around 0.03 which is close the Standard Model correction, offering a much better probe on whether New Physics is influential at this level of precision. The progress that has been made so far in this thesis paves a well-constructed and solid path towards future results. From the current result, the chance of having a much precised measurement in the next a few years on this channel is very promising and searching New Physics effect in penguin-mode $b \rightarrow s$ transition from Belle II is proven to be an exciting and important topic.

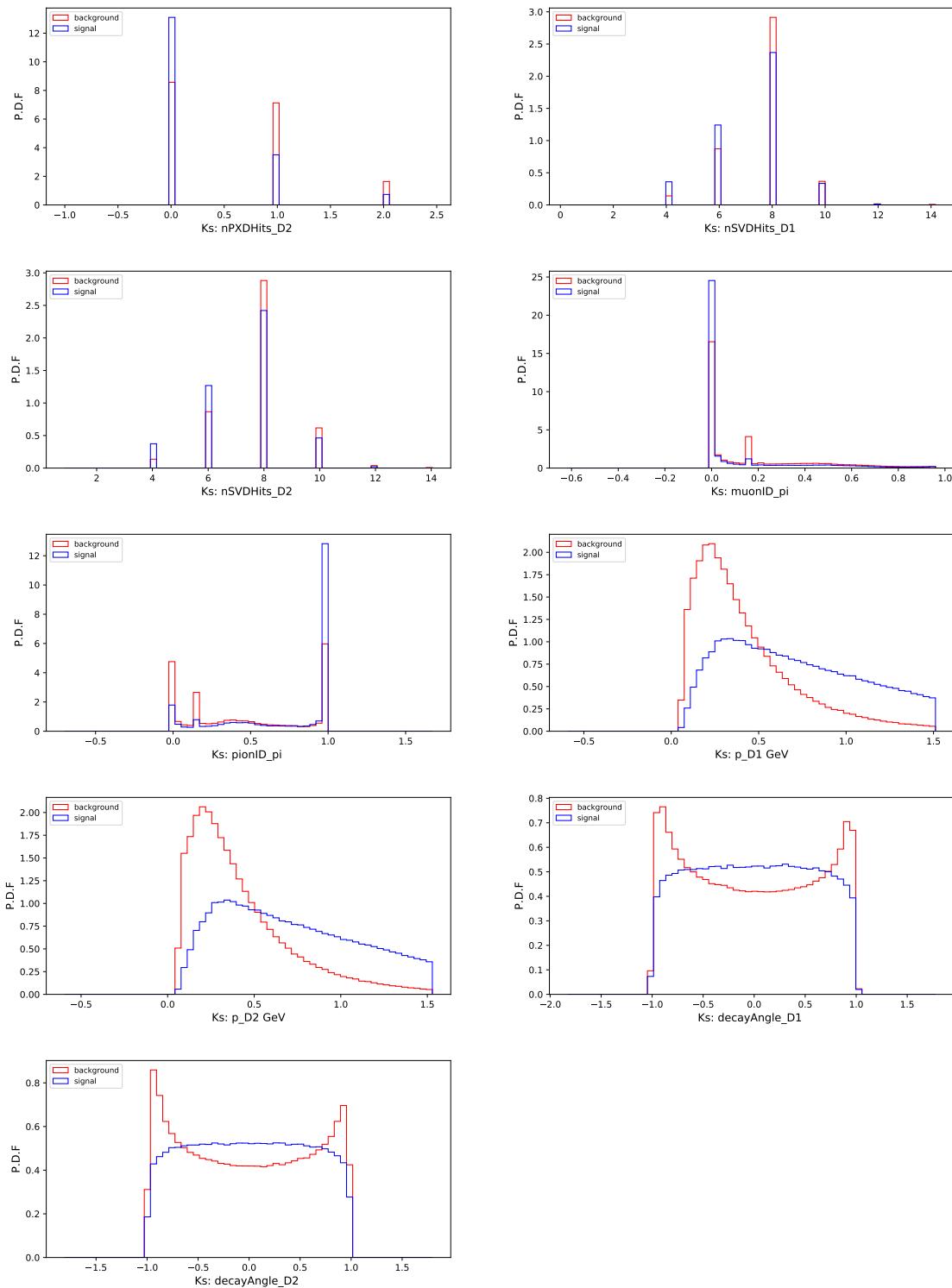
Appendix A

Training Observables for K_S^0 Classifier

Figure A-1: The distribution of input variables in signal MC for KsFinder. The red is the from fake K_S^0 and the blue is from true K_S^0



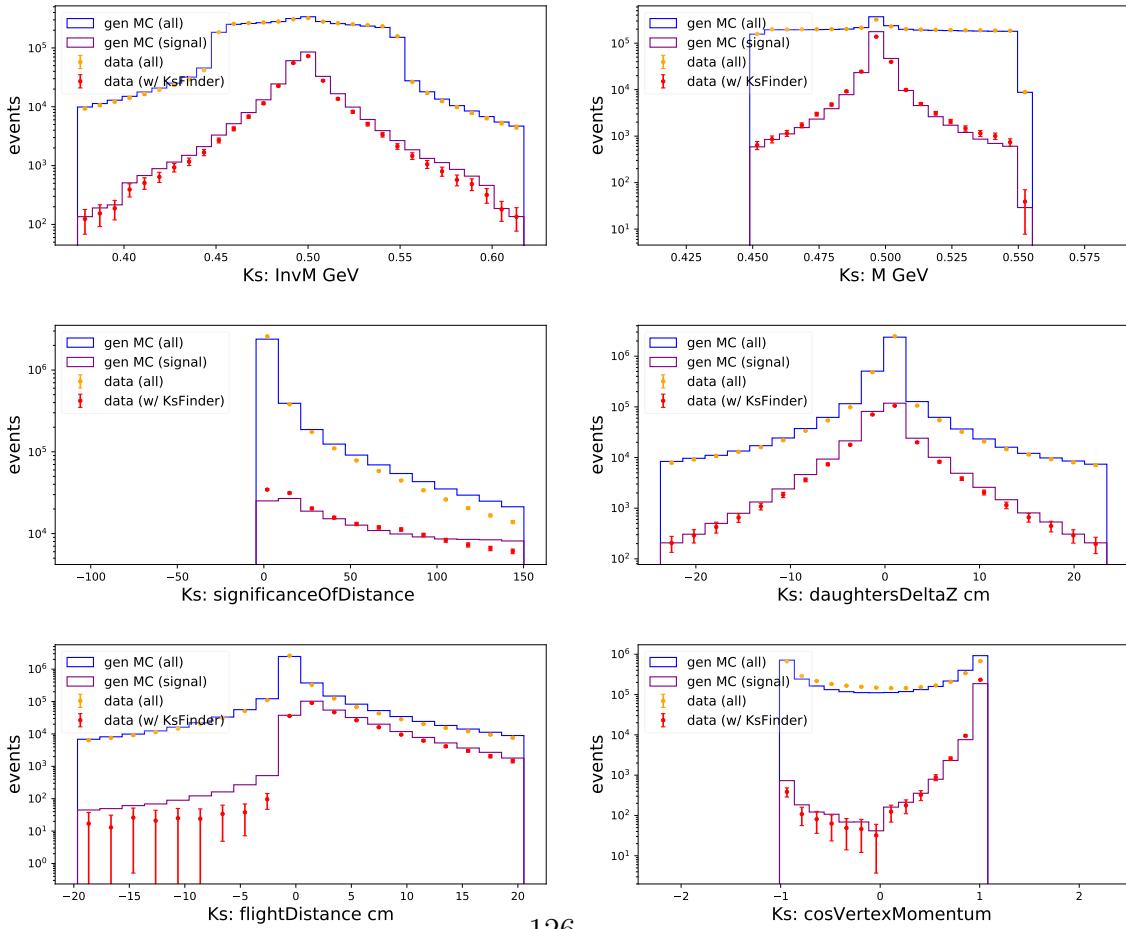


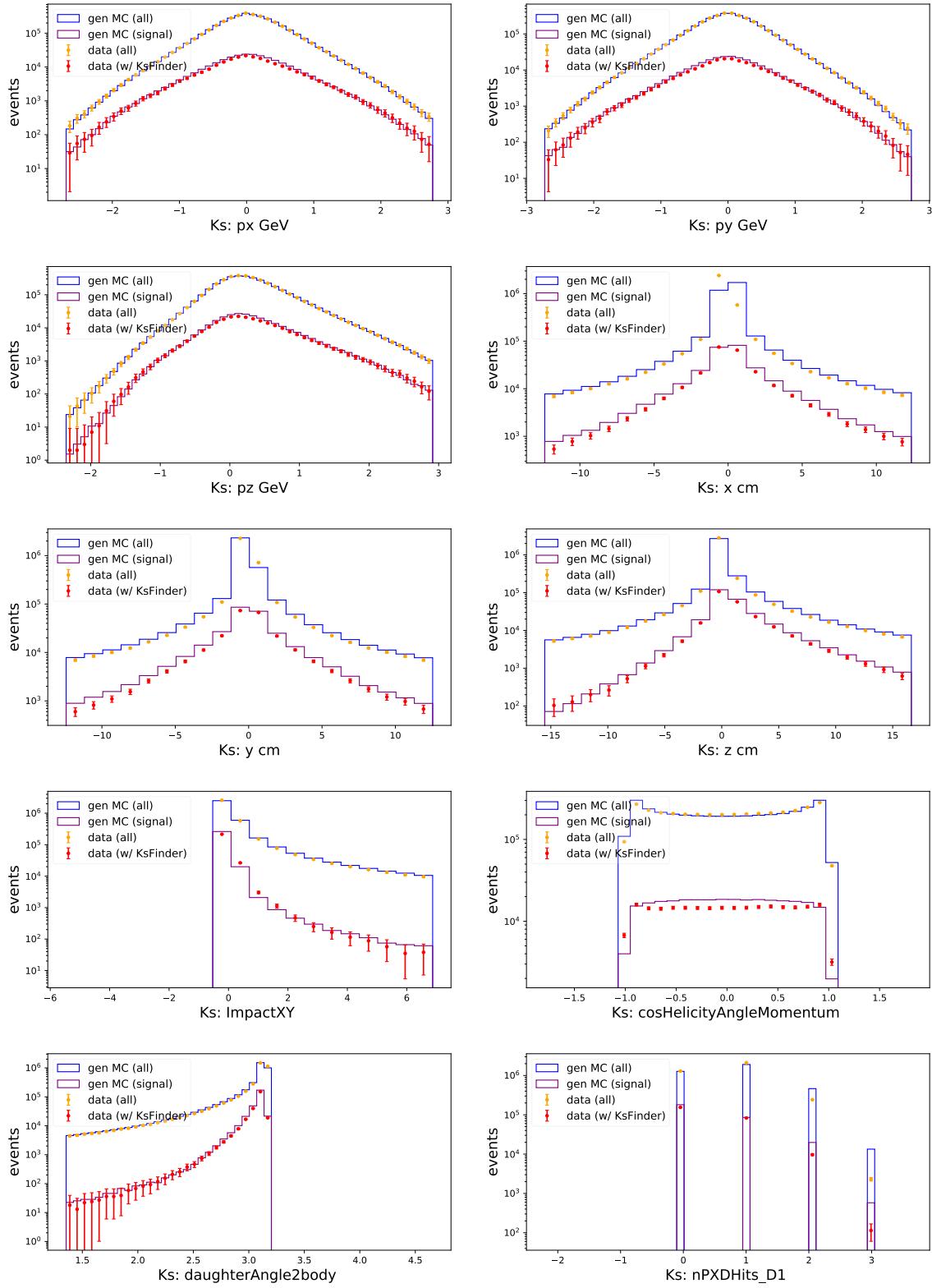


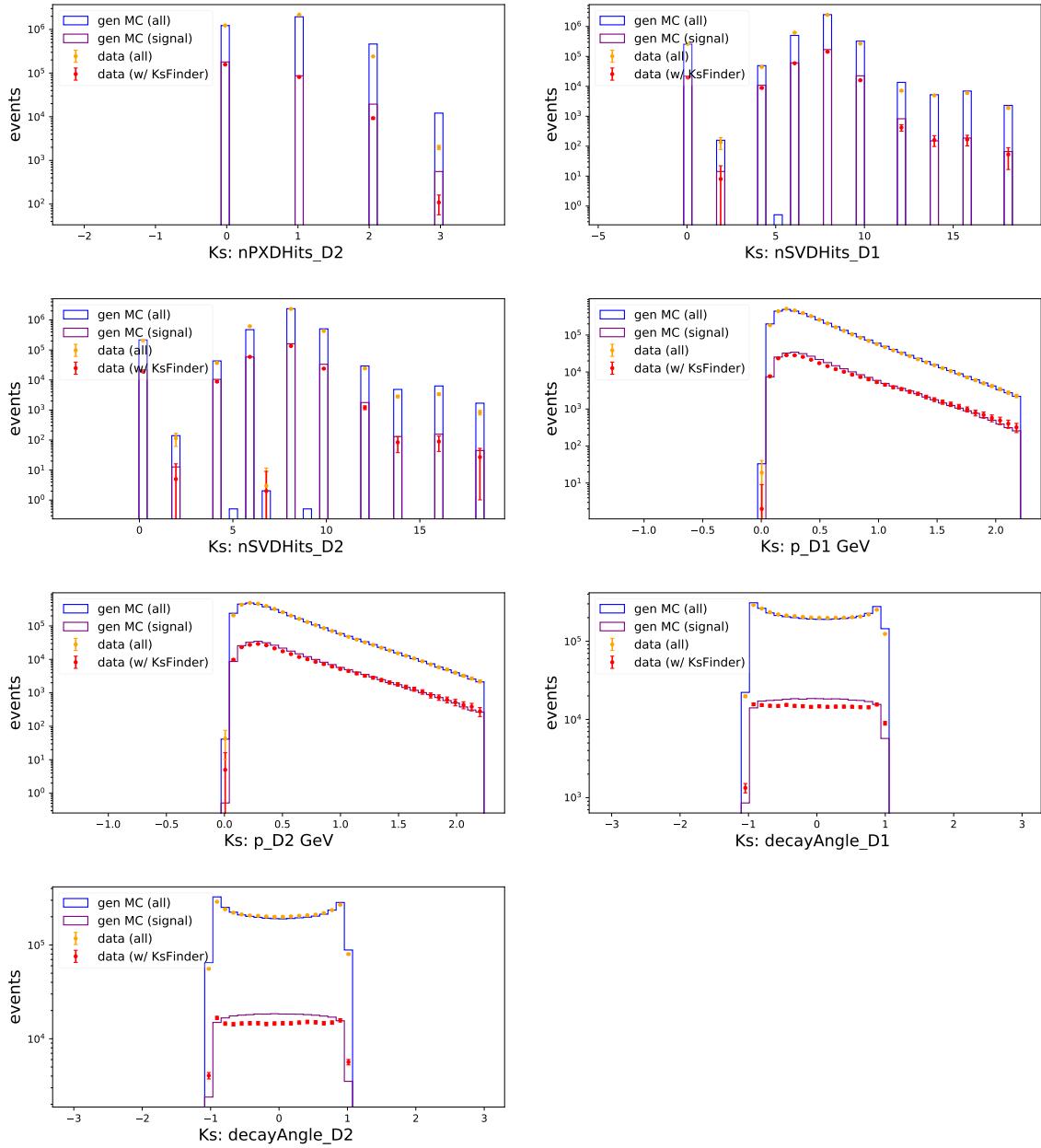
Appendix B

Data Validation Plots for K_S^0

Figure B-1: The distribution of the training variables in KsFinder. The blue and purple solid lines are the total and true K_S^0 distributions from generic MC, respectively. The yellow and red dots are the data distribution before and after applying KsFinder cut. The uncertainties in data are taken as three times the Poisson standard deviation.







Appendix C

$2K_S^0$ invariant mass distribution

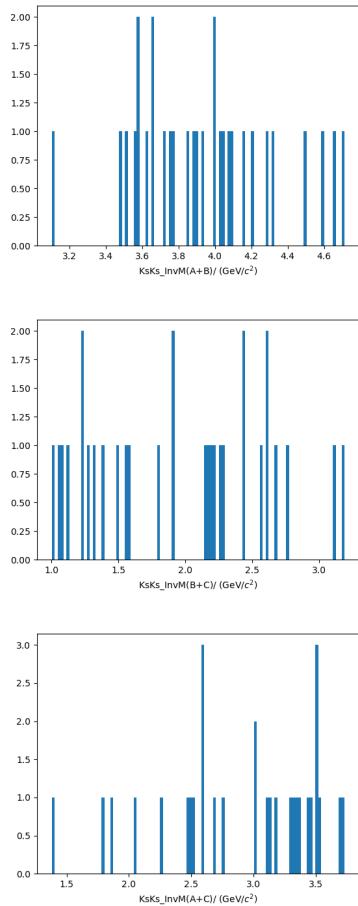


Figure C-1: Experiment data in signal region

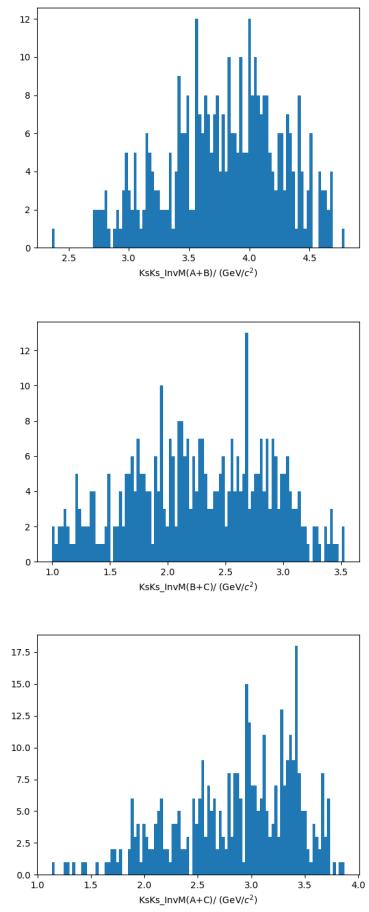


Figure C-2: Generic MC in signal region.

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