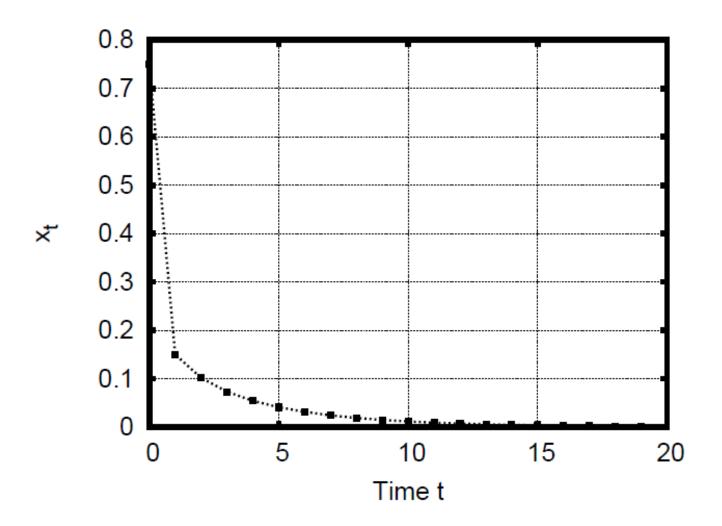


# **Dynamical Systems Lecture 5.02**

EEU45C09 / EEP55C09 Self Organising Technological Networks

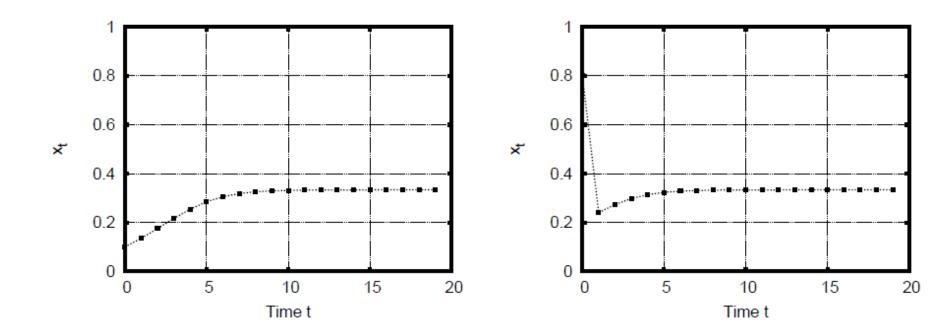
> Nicola Marchetti nicola.marchetti@tcd.ie

- Logistic equation: f(x) = rx(1-x).
- A simple model of resource-limited population growth.
- The population x is expressed as a fraction of the carrying capacity.
   0 \le x \le 1.
- r is a parameter—the growth rate—that we will vary.
- Let's first see what happens if r = 0.8.

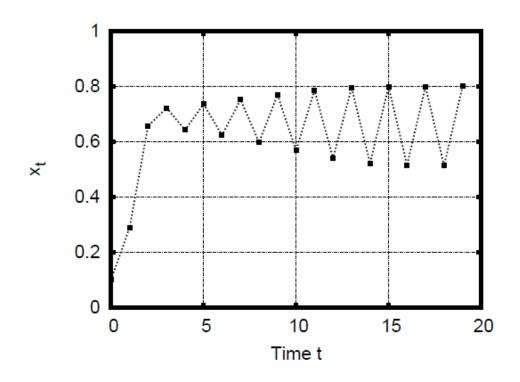


0 is an attracting fixed point.

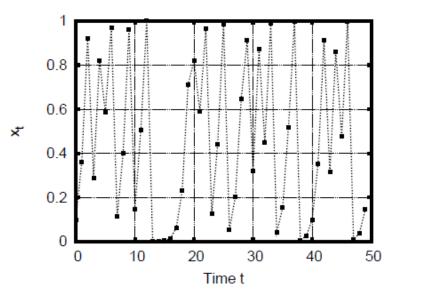
• Logistic equation, r = 1.5.

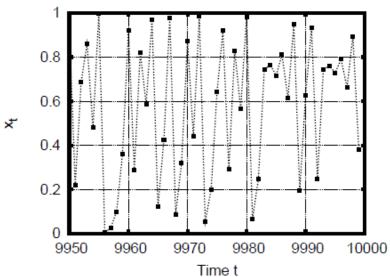


- All initial conditions are pulled toward 0.33.
- 0.33 is an attracting fixed point.



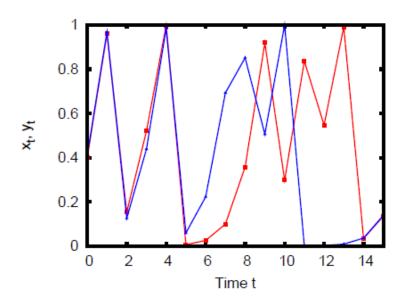
- Logistic equation, r = 3.2.
- Initial conditions are pulled toward a cycle of period 2.
- The orbit oscillates between 0.513045 and 0.799455.
- This cycle is an attractor. Many different initial conditions get pulled to it.

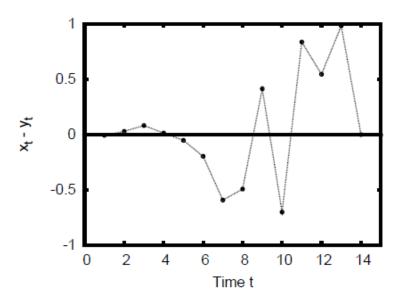




- Logistic equation, r = 4.0.
- What's going on here?!
- The orbit is not periodic. In fact, it never repeats.
- This is a rigorous result; it doesn't rely on computers.
- What happens if we try different initial conditions?

• Two slightly different initial conditions,  $x_0 = 0.4$  and  $x_0 = 0.41$ .

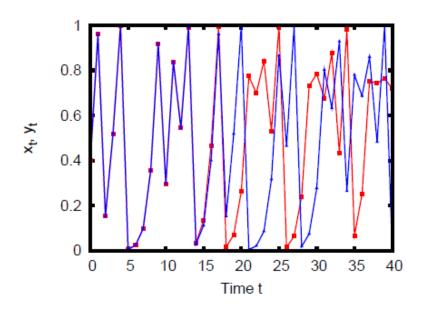


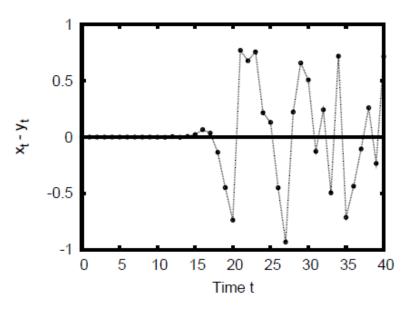


- The right graph plots the difference between the two orbits on the left
- Note that the difference between the two orbits grows.
- Can think of the blue as the true values, and the red as the predicted values.
- The plot on the right can be viewed as prediction error over time.
- How can we improve our predictions?

#### Sensitive Dependence on Initial Conditions

• Two different initial conditions,  $x_0 = 0.4$  and  $x_0 = 0.4000001$ .

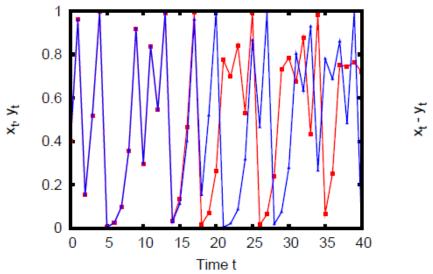


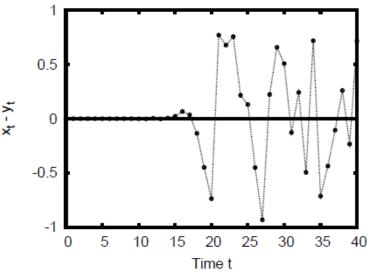


- The two initial conditions differ by one part in one million
- The orbits differ significantly after around 20 iterations, whereas before they differed after around 4 iterations.
- Increasing the accuracy of the initial condition by a factor of  $10^5$  allow us to predict the outcome 5 times further.

## Sensitive Dependence on Initial Conditions

• Two different initial conditions,  $x_0 = 0.4$  and  $x_0 = 0.4000001$ .





- Thus, for all practical purposes, this system is unpredictable, even though it is deterministic.
- This phenomena is known as Sensitive Dependence on Initial Conditions, or, more colloquially, The Butterfly Effect.
- Arbitrarily small differences in initial conditions grow to become arbitrarily large.

#### **Formal Definition of SDIC**

 A dynamical system has sensitive dependence on initial conditions (SDIC) if arbitrarily small differences in initial conditions eventually lead to arbitrarily large differences in the orbits.

#### More formally

- Let X be a metric space, and let f be a function that maps X to itself:  $f: X \mapsto X$ .
- The function f has SDIC if there exists a  $\delta>0$  such that  $\forall x_1\in X$  and  $\forall \epsilon>0$ , there is an  $x_2\in X$  and a natural number  $n\in N$  such that  $d[x_1,x_2]<\epsilon$  and  $d[f^n(x_1),f^n(x_2)]>\delta$ .
- In other words, two initial conditions that start  $\epsilon$  apart will, after n iterations, be separated by a distance  $\delta$ .

#### **Formal Definition of Chaos**

There is not a 100% standard definition of chaos. But here is one of the most commonly used ones:

An iterated function is **chaotic** if:

- The function is deterministic.
- The system's orbits are bounded.
- 3. The system's orbits are **aperiodic**; i.e., they never repeat.
- The system has sensitive dependence on initial conditions.

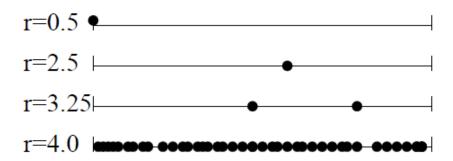
Other properties of a chaotic dynamical system  $(f: X \mapsto X)$  that are sometimes taken as defining features:

- 1. **Dense periodic points:** The periodic points of f are dense in X.
- 2. **Topological transitivity:** For all open sets  $U, V \in X$ , there exists an  $x \in U$  such that, for some  $n < \infty$ ,  $f_n(x) \in V$ . I.e., in any set there exists a point that will get arbitrarily close to any other set of points.

We have seen several possible long-term behaviors for the logistic equation:

- 1. r = 0.5: attracting fixed point at 0.
- 2. r = 2.5: attracting fixed point at 0.6.
- 3. r = 3.25: attracting cycle of period 2.
- 4. r = 4.0: chaos.

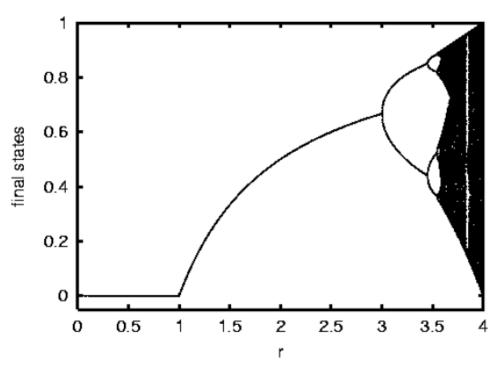
Graphically, we can illustrate this as follows:



- I.e., for each r, iterate and plot the final x values as dots on the number line.
- What else can the logistic equation do??

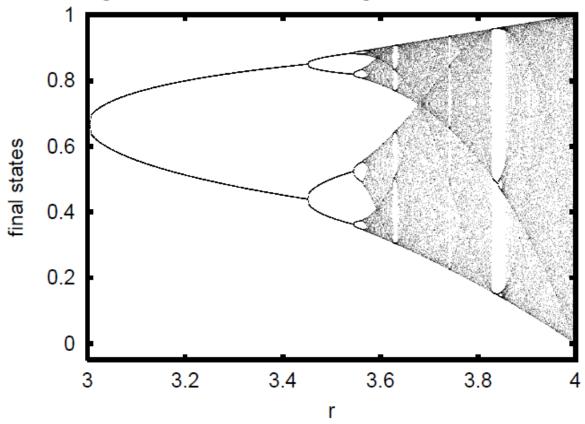
$$r=4.0$$

- ullet Do this for more and more r values and "glue" the lines together.
- Turn sideways and ...



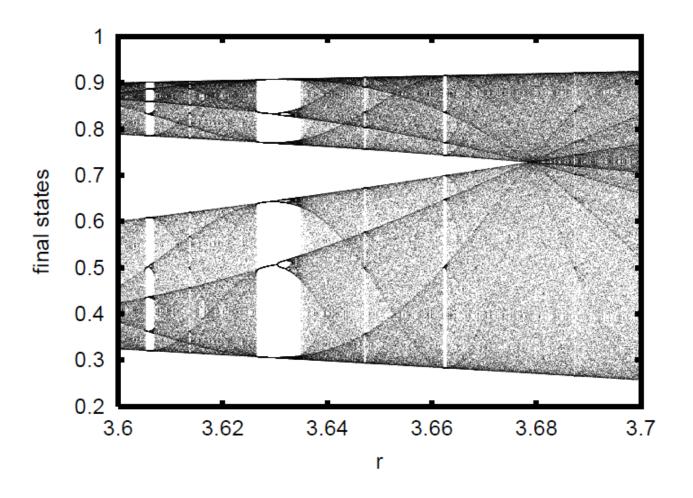
- The bifurcation diagram shows all the possible long-term behaviors for the logistic map.
- 0 < r < 1, the orbits are attracted to zero.
- 1 < r < 3, the orbits are attracted to a non-zero fixed point.
- 3 < r < 3.45, orbits are attracted to a cycle of period 2.
- Chaotic regions appear as dark vertical lines.

Let's zoom in on a region of the bifurcation diagram:



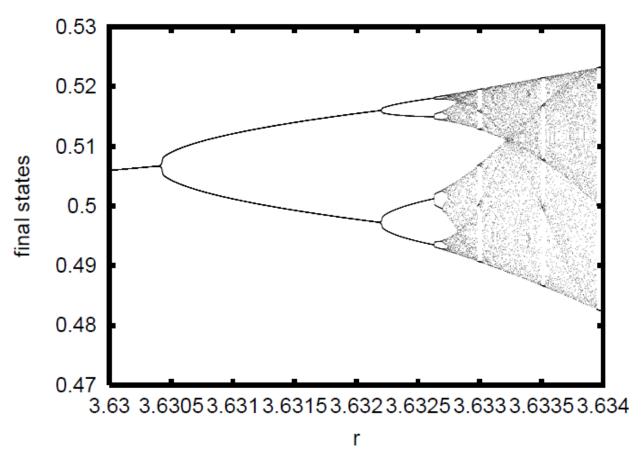
- The sudden qualitative changes are known as bifurcations.
- There are **period-doubling bifurcations** at  $r \approx 3.45$ ,  $r \approx 3.544$ , etc.
- Note the window of period 3 near r = 3.83.

Let's zoom in again:



Note the sudden changes from chaotic to periodic behavior.

Let's zoom in once more:



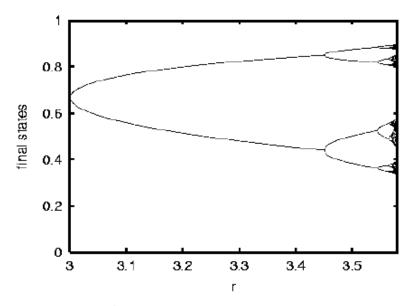
- Note the small scales on the vertical axis, and the tiny scale on the y axis.
- Note the self-similar structure. As we zoom in we keep seeing pitchforks.

# **Bifurcation Diagram - Summary**

- As we vary r, the logistic equation shuffles suddenly between chaotic and periodic behaviors, but the bifurcation diagram reveals that these transitions appear in a structured, or regular, way.
- This is an example of a sort of "order within chaos."
- Bifurcations—a sudden, qualitative change in behavior as a parameter is continuously varied—is a generic feature of non-linear systems.
- In the next few slides we'll examine one of the regularities in the bifurcation diagram: The period-doubling route to chaos.

# **Period-Doubling Route to Chaos**

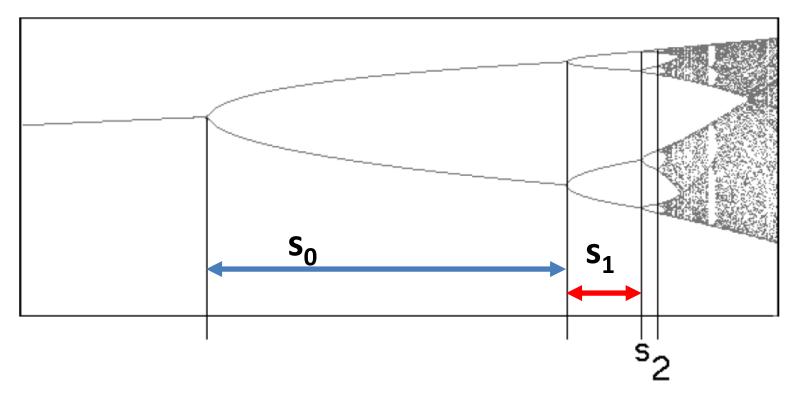
As r is increased from 3, a sequence of period doubling bifurcations occur.



- At  $r=r_{\infty}\approx 3.569945672$  the periods "accumulate" and the map becomes chaotic.
- ullet For  $r > r_{\infty}$  it has SDIC. For  $r < r_{\infty}$  it does not.
- This is a type of phase transition: a sudden qualitative change in a system's behavior as a parameter is varied continuously.

# Period-Doubling Route to Chaos - Geometric Scaling

 Let's examine the ratio of the lengths of the pitchfork tines in the bifurcation diagram.

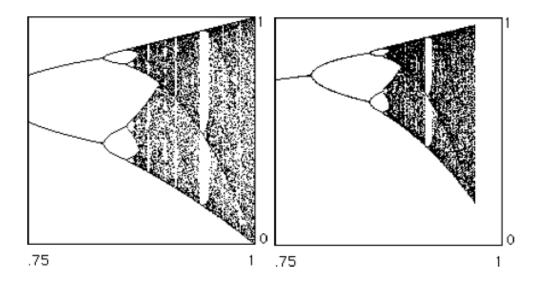


- The first ratio is:  $\delta_1 = \frac{s_1 s_0}{s_2 s_1}$ .
- The n<sup>th</sup> ratio is:  $\delta_n = \frac{s_n s_{n-1}}{s_{n+1} s_n}$ .

# Feigenbaum's Constant

- This ratio approaches a limit:  $\lim_{n\to\infty} \delta_n = 4.669201609...$  This is known as **Feigenbaum's constant**  $\delta$ .
- This means that the bifurcations occur in a regular way.
- ullet Amazingly, the value of  $\delta$  is **universal**: it is the same for any period-doubling route to chaos!

### Universality



- The figure on the left is the bifurcation diagram for  $f(x) = r \sin(\pi x)$ .
- ullet The figure on the right is the bifurcation diagram for  $f(x)=rac{27}{4}rx^2(1-x)$ .
- ullet The bifurcation diagrams are very similar: **both have**  $\deltapprox 4.6692$ .
- Mathematically, things are constrained so that there is, in some sense, only
  one possible way for a system to undergo a period-doubling to chaos.

# **Experimental Verification of Universality**

- ullet Universality isn't just a mathematical curiosity. Physical systems undergo period-doubling order-chaos transitions. Almost miraculously, these systems also appear to have a universal  $\delta$ .
- Experiments have been done on fluids, circuits, acoustics:
  - Water:  $4.3 \pm .8$
  - Mercury:  $4.4 \pm .1$
  - Diode:  $4.5 \pm .6$
  - Transistor:  $4.5 \pm .3$
  - Helium:  $4.8 \pm .6$

Data from Cvitanović, Universality in Chaos, World Scientific, 1989.

- A very simple equation, the logistic equation, has produced a quantitative prediction about complicated systems (e.g., fluid turbulence) that has been verified experimentally.
- Nature is somehow constrained.

#### **Chaos - Deterministic Source of Randomness**

- A chaotic system behaves as if it is random, not governed by a deterministic rule.
- For r=4, the symbolic dynamics of the logistic equation produce a sequence of 0's and 1's that is indistinguishable from a fair coin toss.
- Symbolic dynamics: 0 if  $x < \frac{1}{2}$ , 1 if  $x > \frac{1}{2}$ .
- The apparent randomness arises because the system is so deterministic.
   Determinism gives rise to SDIC.

### **Strange Attractors**

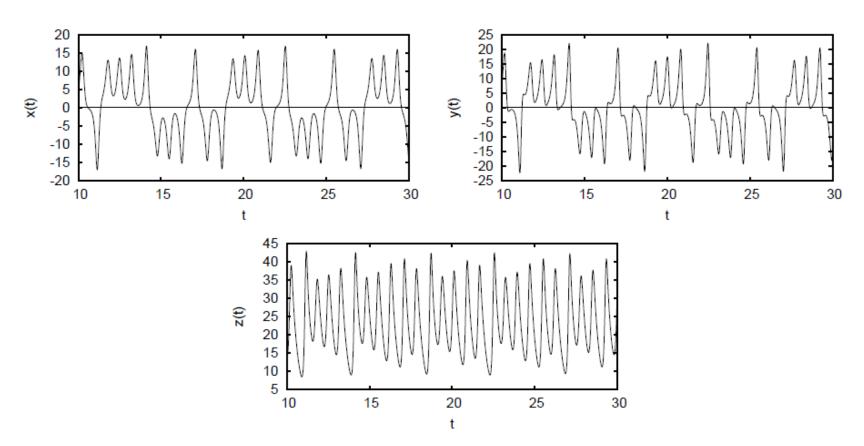
$$\frac{dx}{dt} = \sigma(y - x)$$

$$\frac{dy}{dt} = x(\rho - z) - y$$

$$\frac{dz}{dt} = xy - \beta z .$$
(1)

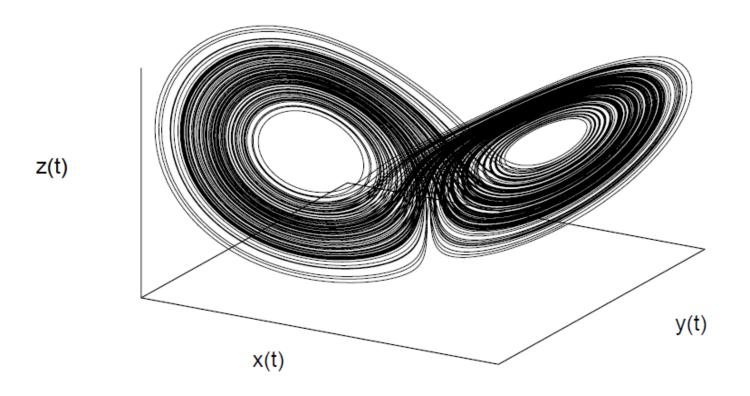
- The Lorenz Equations: introduced by Edward Lorenz in the early 1960's as a very simple model of a weather system.
- Here there are three variables, x, y, and z that change in time.
- The variables are continuous: defined for every time t, not just discrete times.
- The Lorenz equations are differential equations, a type of dynamical system where the rate of change at every instant is specified.
- From this rate of change, one can figure out how the variables themselves change.

## **Lorenz Equations - Aperiodic Trajectories**



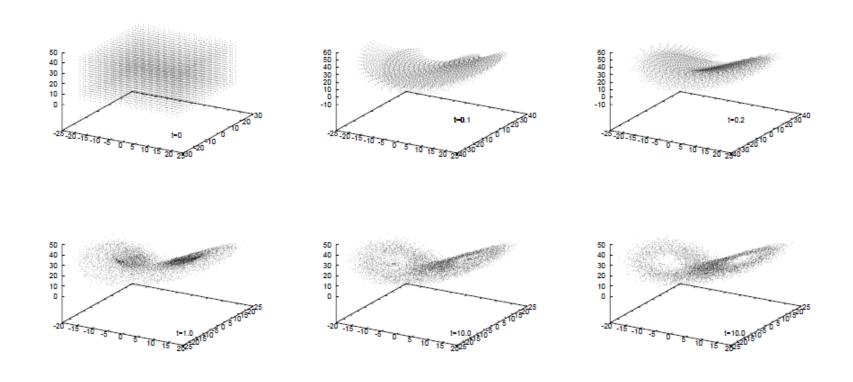
- x, y, and z are all aperiodic. They do not repeat.
- How are x, y, and z related? To see this, let's plot the three variables on the same graph.

# **Lorenz Equations - Strange Attractor**



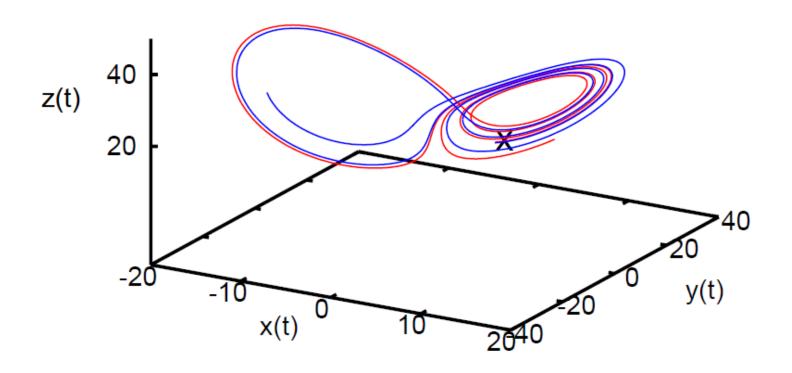
- Although individually x, y, and z move seemingly at random, when plotted together one can see a complicated relationship between them.
- The trajectory weaves through space but never repeats.

# **Lorenz Equations - Strange Attractor**



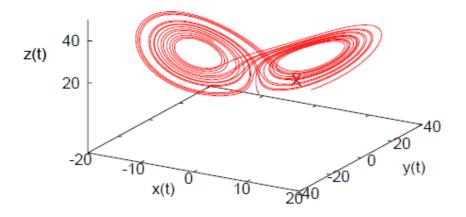
- This shape is an attractor. Orbits get pulled to it.
- Plot of 8000 different initial conditions uniformly distributed in a cube.
- The orbits are pulled to the attractor.

## **Lorenz Equations - Butterfly Effect**



- The Lorenz Equations show the butterfly effect.
- The blue and red orbit start at almost exactly the same point, indicated by X.
- Very quickly the two orbits become quite different.

#### **A Chaotic Attractor**



- The attractor is stable; it attracts all orbits.
- But the dynamics on the attractor itself are chaotic.
- The system is a mix of order and unpredictability.
- Roughly speaking, unpredictability ≈ weather.
- Global structure, the shape of the attractor ≈ climate.
- Strange attractors are a sort of order hidden within chaos.

#### **Chaos - Conclusions**

- ullet Deterministic systems can produce random, unpredictable behavior. E.g., logistic equation with r=4.
- Simple systems can produce complicated behavior. E.g., long periodic behavior in logistic equation.
- Some features of dynamical systems are universal—the same for many different systems.
- Chaos and other structures can be stable.
- Aubin and Dahan Dalmedico: [C]haos has definitely blurred a number of old epistemological boundaries and conceptual oppositions hitherto seemingly irreducible such as order/disorder, random/nonrandom, simple/complex, local/global, stable/unstable, ....

Aubin and Dahan Dalmedico, Historia Mathematica 29 (2002), 167. doi:10.1006/hmat.2002.2351

## **Chaos & Complex Systems**

- Many researchers who did groundbreaking work in chaos in the 1970s and 1980s are now doing work in complex systems.
- Appreciation that complex behavior can have simple origins.
- Universality gives us some reason to believe that we can understand complicated and complex systems with simple models.
- More generally, order and disorder, simplicity and complexity, are seen to not be opposites or mutually exclusive categories.
- There is a surprising and delightful creativity to simple, iterated systems.
- Chaos and dynamical systems hint at how randomness, complexity, and structure may emerge out of a simple and deterministic(?) world.
- But it is just one thread in the complex systems tapestry.

# Ackowledgement

• David Feldman, Santa Fe Institute