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Steady laminar incompressible flow between adjacent discs:

- Continuity:

$$\frac{1}{r} \frac{\partial(rv_r)}{\partial r} + \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{\partial v_z}{\partial z} = 0$$

- Radial direction momentum:

$$v_r \frac{\partial v_r}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_r}{\partial \theta} + v_z \frac{\partial v_r}{\partial z} - \frac{v_\theta^2}{r} = -\frac{1}{\rho} \left(\frac{\partial P}{\partial r} \right) + \nu \left\{ \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_r}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 v_r}{\partial \theta^2} + \frac{\partial^2 v_r}{\partial z^2} - \frac{v_r}{r^2} - \frac{2}{r^2} \frac{\partial v_\theta}{\partial \theta} \right\} + f_r$$

- Theta direction momentum:

$$v_r \frac{\partial v_\theta}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_\theta}{\partial \theta} + v_z \frac{\partial v_\theta}{\partial z} + \frac{v_\theta v_r}{r} = -\frac{1}{\rho r} \left(\frac{\partial P}{\partial \theta} \right) + \nu \left\{ \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_\theta}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 v_\theta}{\partial \theta^2} + \frac{\partial^2 v_\theta}{\partial z^2} - \frac{v_\theta}{r^2} - \frac{2}{r^2} \frac{\partial v_r}{\partial \theta} \right\} + f_\theta$$

- Z direction momentum:

$$v_r \frac{\partial v_z}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_z}{\partial \theta} + v_z \frac{\partial v_z}{\partial z} = -\frac{1}{\rho} \left(\frac{\partial P}{\partial z} \right) + \nu \left\{ \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_z}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 v_z}{\partial \theta^2} + \frac{\partial^2 v_z}{\partial z^2} \right\} + f_z$$

Assumptions:

- Flow is taken to be two-dimensional: $v_z = 0$, v_r and v_θ are taken to be constant across the channel between rotor plates, and hence, are mean flow velocities at each r and θ .
- The flow field with v_r and v_θ velocity components is **treated as being inviscid**, with a body-force representation of the wall shear effects. The viscous drag exerted on the flow by the side walls of the channel between the rotors is modelled as a body force acting on the flow at each (r, θ) location.
- The flow field is radially symmetric. The inlet flow at the rotor outer edge is assumed uniform, resulting in a flow field that is the same at any angle θ . All θ derivatives of the flow quantities are therefore zero.
- Radial velocity derivatives and divisions in radial and theta directions are negligible as compared to Axial case.**
- Constant friction factor characteristics within the flow passage.
- Axial direction effects must not be ignored.

Simplified version:

- Continuity:

$$\frac{1}{r} \frac{\partial(rv_r)}{\partial r} = 0$$

- Radial direction momentum:

$$v_r \frac{\partial v_r}{\partial r} - \frac{v_\theta^2}{r} = -\frac{1}{\rho} \left(\frac{\partial P}{\partial r} \right) + \nu \left\{ \frac{\partial^2 v_r}{\partial z^2} + \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_r}{\partial r} \right) - \frac{v_r}{r^2} \right\}$$

- Theta direction momentum:

$$v_r \frac{\partial v_\theta}{\partial r} + \frac{v_\theta v_r}{r} = f_\theta$$

- Z direction momentum:

$$-\frac{1}{\rho} \left(\frac{\partial P}{\partial z} \right) = 0$$

Velocity definition:

$$v_r(r, z) = \bar{v}_r(r)\phi(z)$$

$$v_\theta(r, z) = \bar{v}_\theta(r)\phi(z) + U(r)$$

Where:

$$\phi(z) = \left(\frac{n+1}{n}\right) \left[1 - \left(\frac{2z}{b}\right)^n\right]$$

$$\bar{v}_r(r) = \frac{1}{b} \int_{-\frac{b}{2}}^{\frac{b}{2}} v_r(r, z) dz$$

$$\bar{v}_\theta(r) = \frac{1}{b} \int_{-\frac{b}{2}}^{\frac{b}{2}} v_\theta(r, z) dz$$

$$U(r) = \frac{r}{r_0} U_0$$

Integral simplifications for phi:

$$\int_{-\frac{b}{2}}^{\frac{b}{2}} \phi(z) dz = 2 \int_0^{\frac{b}{2}} \phi(z) dz = b$$

$$\int_{-\frac{b}{2}}^{\frac{b}{2}} \phi(z)^2 dz = 2 \int_0^{\frac{b}{2}} \phi(z)^2 dz = \frac{2(n+1)}{2n+1} b$$

$$\int_{-\frac{b}{2}}^{\frac{b}{2}} \frac{d^2 \phi(z)}{dz^2} dz = 2 \int_0^{\frac{b}{2}} \frac{d^2 \phi(z)}{dz^2} dz = -\frac{4(n+1)}{b}$$

Equation 5 yields that $rv_r = C_r$, and integrating it over the channel:

$$\int_{-\frac{b}{2}}^{\frac{b}{2}} rv_r dz = \int_{-\frac{b}{2}}^{\frac{b}{2}} r \bar{v}_r(r) \phi(z) dz = r \bar{v}_r(r) b = b C_r$$

mass conservation requires that $-2\pi r_0 \rho \int_{-b/2}^{b/2} v_r dz = \dot{m}_c$, where \dot{m}_c is the mass flow rate per channel:

$$-2\pi r_0 \rho \int_{-b/2}^{b/2} v_r dz = -2\pi r_0 \rho \bar{v}_r(r_0) \int_{-b/2}^{b/2} \phi(z) dz = -2\pi r_0 \rho \bar{v}_r(r_0) b = \dot{m}_c$$

Note: $\bar{v}_r(r_0) = \bar{v}_{r0}$

$$\bar{v}_{r0} = \frac{\dot{m}_c}{2\pi r_0 b \rho}$$

Hence:

$$\bar{v}_r = -\frac{r_0}{r} \bar{v}_{r0}$$

Theta direction wall friction force for fluid element in channel between rotors with volume V_e is given by:

$$F_\theta = \tau_w A_w = \frac{4\tau_w V_e}{D_H}$$

where D_H is the hydraulic diameter of the channel ($D_H = 4V_e/A_w$). For parallel plates, $D_H = 2b$. For Newtonian fluid, it follows that:

$$\tau_w = f \frac{\rho \bar{v}_\theta^2}{2}$$

$$f = \frac{\tau_w}{\rho \bar{v}_\theta^2 / 2} = \frac{2\mu \left[\frac{\partial(v_\theta(r, z) - U)}{\partial z} \right]_{z=b/2}}{\rho \bar{v}_\theta^2}$$

For the purpose of this analysis, tangential shear interaction of the flow within the disk surface is postulated

$$f = \frac{Po}{Re_c}$$

Where:

$$Re_c = \frac{\rho \bar{v}_\theta D_H}{\mu}$$

For laminar flow, $Po = 24$. For flow with roughened surface of turbulence action, an enhancement number F_{Po} is defined as such:

$$F_{Po} = \frac{Po}{24}$$

Relationship between n , F_{po} and Po

$$\begin{aligned} f &= \frac{2\mu \left[\frac{\partial(v_\theta(r, z) - U)}{\partial z} \right]_{z=\frac{b}{2}}}{\rho \bar{v}_\theta^2} \\ &= \frac{2\mu \bar{v}_\theta \left[\frac{\partial \Phi(z)}{\partial z} \right]_{z=\frac{b}{2}}}{\rho \bar{v}_\theta^2} \\ &= \frac{2\mu}{\rho \bar{v}_\theta} \left[-\frac{n+1}{n} \left(\frac{2}{b} \right) \left(n \left(\frac{2z}{b} \right)^{n-1} \right) \right]_{z=\frac{b}{2}} \\ &= \frac{2\mu}{\rho \bar{v}_\theta} \left[-\frac{2(n+1)}{b} \right] \end{aligned}$$

Combining with the remaining two equations (absolute terms):

$$\begin{aligned} \frac{Po}{Re_c} &= \frac{2\mu}{\rho \bar{v}_\theta} \left[\frac{2(n+1)}{b} \right] \\ \frac{\mu Po}{\rho \bar{v}_\theta D_H} &= \frac{2\mu}{\rho \bar{v}_\theta} \left[\frac{2(n+1)}{b} \right] \\ Po &= D_H \left[\frac{4(n+1)}{b} \right] \\ Po &= 8(n+1) \\ Po/8 &= (n+1) = 3F_{Po} \end{aligned}$$

Wall Shear Stress

$$\begin{aligned} \tau_w &= f \frac{\rho \bar{v}_\theta^2}{2} \\ \tau_w &= \frac{2\mu}{\rho \bar{v}_\theta} \left[-\frac{2(n+1)}{b} \right] \frac{\rho \bar{v}_\theta^2}{2} \\ \tau_w &= -\frac{2\mu(n+1)\bar{v}_\theta}{b} \end{aligned}$$

Wall Friction Term

$$f_\theta = \frac{F_\theta}{\rho V_e} = \frac{4\tau_w}{\rho D_H} = -\frac{4}{\rho D_H} \frac{2\mu(n+1)\bar{v}_\theta}{b} = -\frac{16\mu(n+1)\bar{v}_\theta}{\rho D_H^2}$$

Dimensionless Variables:

$$\xi = \frac{r}{r_0}$$

$$\widehat{W} = \frac{\bar{v}_\theta}{U_0}$$

$$\hat{P} = \frac{P - P_0}{\rho U_0^2 / 2}$$

$$V_{r0} = \frac{\bar{v}_{r0}}{U_0}$$

$$\varepsilon = 2b/r_0$$

$$\text{Re}_m^* = \frac{D_H}{r_0} \frac{\dot{m}_c D_H}{2\pi r_0 b \mu} = \frac{\dot{m}_c D_H}{r_0^2 \mu} = 4V_{r0} \frac{b^2 U_0}{r_0 \nu}$$

Note

From the above definitions:

$$v_r = -\frac{\nu \phi(z) \text{Re}_m^*}{D_H \xi \varepsilon} = -\frac{r_0}{r} \phi(z) \bar{v}_{r0} = -\frac{r_0}{r} \phi(z) V_{r0} U_0$$

$$\frac{\partial}{\partial r} = \frac{\partial}{\partial \xi} \frac{\partial \xi}{\partial r} = \frac{1}{r_0} \frac{\partial}{\partial \xi}$$

Tangential ODE

$$v_r \frac{\partial v_\theta}{\partial r} + \frac{v_\theta v_r}{r} = -\frac{16\mu(n+1)\bar{v}_\theta}{\rho D_H^2}$$

$$\frac{\partial v_\theta}{\partial r} + \frac{v_\theta}{r} = -\frac{1}{v_r} \frac{16\mu(n+1)\bar{v}_\theta}{\rho D_H^2}$$

$$\frac{\partial v_\theta}{\partial r} + \frac{v_\theta}{r} = \frac{D_H \xi \varepsilon}{\nu \phi(z) \text{Re}_m^*} \frac{16\mu(n+1)\bar{v}_\theta}{\rho D_H^2}$$

$$\frac{\partial v_\theta}{\partial r} + \frac{v_\theta}{r} = \frac{\xi \varepsilon}{\phi(z) \text{Re}_m^*} \frac{16(n+1)\bar{v}_\theta}{D_H}$$

$$\phi(z) \left[\frac{\partial \bar{v}_\theta}{\partial r} + \frac{\bar{v}_\theta}{r} \right] + \frac{2U_0}{r_0} = \frac{\xi \varepsilon}{\phi(z) \text{Re}_m^*} \frac{16(n+1)\bar{v}_\theta}{D_H}$$

$$\phi(z)^2 \left[\frac{\partial \bar{v}_\theta}{\partial r} + \frac{\bar{v}_\theta}{r} \right] + \phi(z) \frac{2U_0}{r_0} = \frac{\xi \varepsilon}{\text{Re}_m^*} \frac{16(n+1)\bar{v}_\theta}{D_H}$$

Integrating both side across the channel $\left(\frac{1}{b} \int_{-b/2}^{b/2} dz \right)$:

$$\frac{2(n+1)}{2n+1} \left[\frac{\partial \bar{v}_\theta}{\partial r} + \frac{\bar{v}_\theta}{r} \right] + \frac{2U_0}{r_0} = \frac{\xi \varepsilon}{\text{Re}_m^*} \frac{16(n+1)\bar{v}_\theta}{D_H}$$

$$\frac{2(n+1)}{2n+1} U_0 \left[\frac{\partial \widehat{W}}{\partial r} + \frac{\widehat{W}}{r} \right] + \frac{2U_0}{r_0} = \frac{U_0 \varepsilon}{\text{Re}_m^*} \frac{16\xi(n+1)\widehat{W}}{D_H}$$

$$\frac{2(n+1)}{2n+1} \left(\frac{U_0}{r_0} \right) \left[\frac{\partial \widehat{W}}{\partial \xi} + \frac{\widehat{W}}{\xi} \right] + 2 \left(\frac{U_0}{r_0} \right) = \left(\frac{U_0}{r_0} \right) \frac{16\xi(n+1)\widehat{W}}{\text{Re}_m^*}$$

Divide both side by $\left(\frac{U_0}{r_0} \right)$:

$$\frac{2(n+1)}{2n+1} \left[\frac{\partial \widehat{W}}{\partial \xi} + \frac{\widehat{W}}{\xi} \right] = \frac{16\xi(n+1)\widehat{W}}{\text{Re}_m^*} - 2$$

$$\left[\frac{\partial \widehat{W}}{\partial \xi} + \frac{\widehat{W}}{\xi} \right] = \frac{2n+1}{n+1} \left(\frac{8\xi(n+1)\widehat{W}}{\text{Re}_m^*} \right) - \frac{2n+1}{n+1}$$

$$\left[\frac{\partial \widehat{W}}{\partial \xi} + \frac{\widehat{W}}{\xi} \right] = \frac{8(2n+1)\xi\widehat{W}}{\text{Re}_m^*} - \frac{2n+1}{n+1}$$

$$-\frac{2n+1}{n+1} = \left[\frac{1}{\xi} - \frac{8(2n+1)\xi}{\text{Re}_m^*} \right] \widehat{W} + \frac{\partial \widehat{W}}{\partial \xi}$$

Radial ODE

$$v_r \frac{\partial v_r}{\partial r} - \frac{v_\theta^2}{r} = -\frac{1}{\rho} \left(\frac{\partial P}{\partial r} \right) + \nu \left\{ \frac{\partial^2 v_r}{\partial z^2} + \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_r}{\partial r} \right) - \frac{v_r}{r^2} \right\}$$

$$\begin{aligned} & \frac{(r_0 \phi(z) V_{r0} U_0)^2}{r} \frac{\partial}{\partial r} \left(\frac{1}{r} \right) - \left[\frac{(\bar{v}_\theta(r) \phi(z) + U(r))^2}{r} \right] \\ &= -\frac{1}{\rho} \left(\frac{\partial P}{\partial r} \right) + \nu \left[\frac{\partial^2}{\partial z^2} \left(-\frac{r_0}{r} \phi(z) V_{r0} U_0 \right) - (r_0 \phi(z) V_{r0} U_0) \left(\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \left(\frac{1}{r} \right) \right) - \frac{1}{r^3} \right) \right] \\ & - \frac{(r_0 \phi(z) V_{r0} U_0)^2}{r^3} - \left[\frac{(\widehat{W} U_0 \phi(z) + U(r))^2}{r} \right] \\ &= -\frac{1}{\rho} \left(\frac{\partial P}{\partial r} \right) + \nu \left[-\left(\frac{r_0 V_{r0} U_0}{r} \right) \frac{\partial^2 \phi(z)}{\partial z^2} - (r_0 \phi(z) V_{r0} U_0) \left(\frac{1}{r} \frac{\partial}{\partial r} \left(-\frac{1}{r} \right) - \frac{1}{r^3} \right) \right] \\ & - \frac{(r_0 \phi(z) V_{r0} U_0)^2}{r^3} - \left[\frac{(\widehat{W} U_0 \phi(z))^2 + 2\widehat{W} U_0 U(r) \phi(z) + U(r)^2}{r} \right] \\ &= -\frac{1}{\rho} \left(\frac{\partial P}{\partial r} \right) + \nu \left[-\left(\frac{r_0 V_{r0} U_0}{r} \right) \frac{\partial^2 \phi(z)}{\partial z^2} + 0 \right] \end{aligned}$$

Integrating both side across the channel $\left(\frac{1}{b} \int_{-b/2}^{b/2} dz \right)$:

$$\begin{aligned} & -\frac{2(n+1)}{2n+1} \frac{(r_0 V_{r0} U_0)^2}{r^3} - \left[\frac{2(n+1)}{2n+1} \frac{(\widehat{W} U_0)^2}{r} + \frac{2\widehat{W} U_0 U(r)}{r} + \frac{U(r)^2}{r} \right] \\ &= -\frac{1}{\rho} \left(\frac{\partial P}{\partial r} \right) + \nu \left(\frac{r_0 V_{r0} U_0}{r} \right) \left(\frac{4(n+1)}{b^2} \right) \end{aligned}$$

$$-\frac{2(n+1)}{2n+1} \frac{(V_{r0}U_0)^2}{r_0} \left(\frac{1}{\xi^3}\right) - \left[\frac{2(n+1)}{2n+1} \frac{(\widehat{W}U_0)^2}{r_0} + \frac{2\widehat{W}U_0U(r)}{r_0} + \frac{U(r)^2}{r_0} \right] \frac{1}{\xi}$$

$$= -\frac{1}{r_0\rho} \left(\frac{\partial P}{\partial \xi}\right) + v \left(\frac{V_{r0}U_0}{\xi}\right) \left(\frac{4(n+1)}{b^2}\right)$$

$$\frac{1}{\rho} \left(\frac{\partial P}{\partial \xi}\right) = v \left(\frac{r_0V_{r0}U_0}{\xi}\right) \left(\frac{4(n+1)}{b^2}\right) + \frac{2(n+1)}{2n+1} \left[\frac{(V_{r0}U_0)^2}{\xi^3} \right]$$

$$+ \left[\frac{2(n+1)(\widehat{W}U_0)^2}{2n+1} + 2\widehat{W}U_0U(r) + U(r)^2 \right] \frac{1}{\xi}$$

$$\frac{1}{\rho} \left(\frac{\partial P}{\partial \xi}\right) = v \left(\frac{r_0V_{r0}U_0}{\xi}\right) \left(\frac{4(n+1)}{b^2}\right) + \frac{2(n+1)}{2n+1} \left[\frac{(V_{r0}U_0)^2}{\xi^3} \right]$$

$$+ \left[\frac{2(n+1)(\widehat{W}U_0)^2}{2n+1} + 2\widehat{W}U_0^2\xi + U_0^2\xi^2 \right] \frac{1}{\xi}$$

$$\frac{U_0^2}{2} \left(\frac{\partial \hat{P}}{\partial \xi}\right) = \frac{2(n+1)}{2n+1} \left[\frac{(V_{r0}U_0)^2}{\xi^3} + \frac{(\widehat{W}U_0)^2}{\xi} \right] + 2\widehat{W}U_0^2 + U_0^2\xi + v \left(\frac{r_0V_{r0}U_0}{\xi}\right) \left(\frac{4(n+1)}{b^2}\right)$$

$$\frac{\partial \hat{P}}{\partial \xi} = \frac{4(n+1)}{(2n+1)\xi^3} [V_{r0}^2 + \xi^2\widehat{W}^2] + 4\widehat{W} + 2\xi + v \left(\frac{r_0V_{r0}}{\xi U_0}\right) \left(\frac{8(n+1)}{b^2}\right)$$

$$\frac{\partial \hat{P}}{\partial \xi} = \frac{4(n+1)}{(2n+1)\xi^3} [V_{r0}^2 + \xi^2\widehat{W}^2] + 4\widehat{W} + 2\xi + \left(\frac{1}{4V_{r0}}\right) \left(\frac{vr_0}{b^2U_0}\right) \left(\frac{32(n+1)V_{r0}^2}{\xi}\right)$$

$$\frac{\partial \hat{P}}{\partial \xi} = \frac{4(n+1)}{(2n+1)\xi^3} [V_{r0}^2 + \xi^2\widehat{W}^2] + 4\widehat{W} + 2\xi + \left(\frac{32(n+1)V_{r0}^2}{\text{Re}_m^*\xi}\right)$$

Power Generated Equation

From the definition of torque per rotor disc surface:

$$T = 2\pi \int_{r_i}^{r_0} \tau_w r^2 dr$$

$$T = 2\pi \int_{r_i}^{r_0} \left(\frac{2\mu(n+1)\bar{v}_\theta}{b} \right) r^2 dr$$

$$T = 2\pi \int_{\xi_i}^{\xi_0=1} \left(\frac{2\mu(n+1)\bar{v}_\theta}{b} \right) (r_0\xi)^2 r_0 d\xi$$

$$T = \left(\frac{4\pi\mu r_0^3(n+1)}{b} \right) \int_{\xi_i}^1 \bar{v}_\theta \xi^2 d\xi$$

$$T = \left(\frac{4\pi U_0 \mu r_0^3(n+1)}{b} \right) \int_{\xi_i}^1 \widehat{W} \xi^2 d\xi$$

Total Torque:

$$T_{total} = 2(n_{disc} - 1)T$$

Total Power:

$$P_{total} = \omega T_{total}$$

Mechanical Efficiency

$$\eta_{rm} = \frac{v_{\theta 0} U_0 - v_{\theta i} U_i}{v_{\theta 0} U_0} = 1 - \frac{v_{\theta i} U_i}{v_{\theta 0} U_0} = 1 - \frac{(\bar{v}_{\theta i} + U_i) U_i}{(\bar{v}_{\theta 0} + U_0) U_0}$$

Rearranging:

$$\eta_{rm} = 1 - \frac{\left(\hat{W}_i + \frac{U_i}{U_0}\right) \frac{U_i}{U_0}}{(\hat{W}_0 + 1)} = 1 - \frac{\left(\hat{W}_i + \frac{r_i}{r_0}\right) \frac{r_i}{r_0}}{(\hat{W}_0 + 1)} = 1 - \frac{(\hat{W}_i + \xi_i) \xi_i}{(\hat{W}_0 + 1)}$$

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