

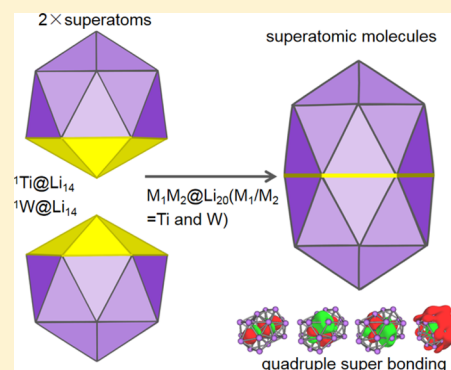
Face-Sharing Homo- and Hetero-Bitetrahexahedral Superatomic Molecules $M_1M_2@Li_{20}$ ($M_1/M_2 = Ti$ and W)

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Supporting Information

ABSTRACT: Superatoms, being a class of clusters with chemical behavior similar to atoms, can be used as building blocks for constructing novel functional materials. Here, superatomic molecules $M_1M_2@Li_{20}$ ($M_1/M_2 = Ti$ and W) are built with our recently discovered tetrahexahedral superatoms $Ti@Li_{14}$ and $W@Li_{14}$. A low-energy face-sharing bi-tetrahexahedral structure with a high symmetry of D_{6h} is identified by global minimum structure search and frequency calculations. The number of shared Li atoms amounts to 6, which is rarely seen before. Molecular orbital and chemical bonding analyses reveal that although isolated superatoms $Ti@Li_{14}$ and $W@Li_{14}$ are nonmagnetic, $Ti_2@Li_{20}$ is an open-shell superatomic molecule with a magnetic moment of $2 \mu_B$, whereas for $W_2@Li_{20}$, the electronic shell remains closed. In both $Ti_2@Li_{20}$ and $W_2@Li_{20}$, a quadruple superbonding between superatoms is found. Interestingly, an assembly of two hetero-tetrahexahedral superatoms $Ti@Li_{14}$ and $W@Li_{14}$ also gives a face-sharing bi-tetrahexahedral structure but with a notable dipole moment. This study provides a basic understanding for the superatomic bonding of $Ti@Li_{14}$ and $W@Li_{14}$, which may aid their application in developing multi-superatom molecules or even bulk crystals.



1. INTRODUCTION

By tuning size, geometry, composition, and charge state, the properties of atomic clusters can be easily tailored, offering abundant building blocks for cluster-assembled materials with targeted functions.^{1–3} As a special type of atomic clusters, superatoms are proposed to mimic the chemical behavior of simple elements in the periodic table. The valence electrons in superatoms fill in a series of atom-like orbitals, with the orbital shapes resemble to those of atoms.⁴ However, different from atoms, the electronic filling order of superatoms generally does not follow the Hund's rule because of the effects of atomic deformations. For metal clusters, their electronic sequences can be understood by the jellium model,^{5,6} where the valence electrons are subject to a uniform positive potential derived from the atomic nuclei and inner electrons. The resulting electronic levels of metal clusters are $|1S^2|1P^6|1D^{10}|2S^2|1F^{14}|$, and so forth, where 2, 8, 18, 20, 34, and 40, and so forth, electrons correspond to magic clusters or inert element superatoms. The superatom concept has been applied successfully to explain the stability of many metal clusters, such as the large peaks at the size of 8, 20, and 40 in mass spectra of sodium clusters,⁷ the resistant of Al_n^- ($n = 13, 23, 37$, etc.) toward oxygen etching,⁸ and the superhalogen character of Al_{13} .⁴

The high stability of superatoms makes them promising for designing cluster-assembled materials. However, the structures of pure superatoms are often distorted to low symmetry upon cluster assembly. A potential route to solve this problem is by ligand protection, such as thiolates (RS)⁹ and phosphines

(PR₃).¹⁰ A typical ligand-protected superatom is the icosahedral Au_{13}^{5+} (8e),^{11–13} which possesses high stability and has been isolated in the experiment. Via different bonding modes, many dimeric structures of Au_{13}^{5+} are constructed, for example, a vertex-sharing Au_{25}^{9+} (16e) core found in the $[Au_{25}(SR)_5(PPh_3)_{10}Cl_2]^{2+}$ cluster,^{14,15} a face-sharing Au_{23}^{9+} (14e) core formed in the $Au_{38}(SR)_{24}$ cluster,^{16,17} and the nonsharing Au_{26}^{10+} (16e) core identified in $Au_{40}(SR)_{24}$.^{18,19} With properly ligand protected, the superatom-assembled materials are accessible.

In a superatom context, binary superatoms, composed of two different elements, are another intriguing field, such as the endohedral $M@Si_n$ ^{20–25} and $X@Al_{12}$ clusters,^{26–31} where properties can be precisely controlled by changing the doped elements M/X while keeping the whole geometric structure. For the Ti-doped silicon $Ti@Si_{16}$ cluster,^{20,21,24} there are 68 valence electrons with a filling order of $1S^2|1P^6|1D^{10}|1F^{14}|2S^2|1G^{18}|2P^6|2D^{10}$, analogous to the noble gas atoms. When the central atom Ti is replaced by its neighboring element Sc or Ta, the corresponding properties of $Sc@Si_{16}$ cluster²⁰ and $Ta@Si_{16}$ cluster^{22,24,25} change to be halogen-like or alkali-metal-like, respectively. Similarly, the Cu- or Si-doped Al_{12} cluster can be tuned to mimic a phosphorus or noble gas atom, respectively.^{27,29} Above all, on the basis of the same geometric structures, the properties of binary superatoms can

Received: February 26, 2019

Revised: May 7, 2019

Published: May 29, 2019

be easily tailored by the variation of endohedral atoms. Moreover, compared with pure superatoms, the characteristics of binary superatoms are often augmented, being more suited to real-life applications.³² However, there are still few studies on the assembled materials from binary superatoms.

Noteworthy, using Li clusters as a test case, Cheng and Yang proposed the super valence bond (SVB) model in 2013, in which superatoms could form delocalized superbonding by sharing both valence pairs and nuclei with other superatoms or atoms.³³ On the basis of the SVB model, the electronic stability of some prolate metal clusters and some ligand-protected clusters has been successfully understood,^{34–37} where the shell closure of superatoms is obtained by superatoms–superatoms or superatoms–atom bonding. Moreover, taking double-icosahedral motif of 26e $\text{Li}_{20}\text{Mg}_3$ and 30e $\text{Li}_{18}\text{Mg}_3\text{Al}_2$ as examples,³⁸ Cheng et al. reveal the existence of quintuple superbonding ($\sigma, 2\pi, 2\delta$) between two icosahedral superatoms, further expanding the community of chemical bonds. As a way of assembling materials from binary superatoms, the SVB model gives a referable direction.

Recently, we have found a series of new binary superatoms, namely, $\text{TM}@\text{Li}_{14}$ (TM = Sc, Ti, V, Y, Zr, Nb, Hf, Ta, and W),³⁹ which present a D_{6d} symmetry tetrahexahedral structure with the cage composed of 14 lithium atoms and the center occupied by a transition-metal atom. Among them, $\text{Ti}@\text{Li}_{14}$ and $\text{W}@\text{Li}_{14}$, of which all of the effective valence electrons are paired, are chosen as building blocks and a new type of stable face-sharing bitetrahexahedral superatomic molecules $\text{M}_1\text{M}_2@\text{Li}_{20}$ ($\text{M}_1/\text{M}_2 = \text{Ti}$ and W) are constructed in this paper. The structural, electronic, and magnetic properties of $\text{M}_1\text{M}_2@\text{Li}_{20}$ ($\text{M}_1/\text{M}_2 = \text{Ti}$ and W) are systematically studied, which reveal three important findings: first, when the two superatoms bond, up to six atoms are shared, whereas in common clusters, the shared number of atoms is usually 0–3. Second, despite the nonmagnetic feature of the two superatoms, the superatomic molecules could be magnetic. Last, a quadruple superbonding between the two superatoms is discovered.

2. COMPUTATIONAL DETAILS

The low-energy structures of $\text{M}_1\text{M}_2@\text{Li}_{20}$ ($\text{M}_1/\text{M}_2 = \text{Ti}$ and W) clusters are located by Crystal structure AnaLYsis by Particle Swarm Optimization (CALYPSO),⁴⁰ which is an efficient structure prediction method. Then, the obtained low-energy isomers are fully optimized at the density functional theory level implemented in the Gaussian16 package.⁴¹ All of these low-energy isomers are confirmed as true local minima by frequency calculations. To choose valid gradient-corrected exchange–correlation functionals for transition-metal compounds,^{42,43} three functionals including pure functional PW91,⁴⁴ meta-GGA functional M06-L,⁴⁵ and hybrid functional TPSSH⁴⁶ are adopted with SDD as the basis set.⁴⁷ In addition, different spin multiplicities (SMs) are considered for each isomer, where the SM of ground state is identified with the lowest energy. All of the pictures presented in the paper are based on the PW91/SDD calculations, and molecular visualization is performed using Multiwfn_3.6.⁴⁸

3. RESULTS AND DISCUSSION

3.1. Geometric Structure. The geometric structures of the face-sharing bitetrahexahedral clusters $\text{Ti}_2@\text{Li}_{20}$ and $\text{W}_2@\text{Li}_{20}$ have a D_{6h} symmetry, whereas their corresponding building superatoms $\text{Ti}@\text{Li}_{14}$ and $\text{W}@\text{Li}_{14}$ possess the D_{6d} symmetry.

As shown in Figure 1, the formation of superatomic molecules can be viewed as the fusion of two $\text{Ti}@\text{Li}_{14}$ or $\text{W}@\text{Li}_{14}$

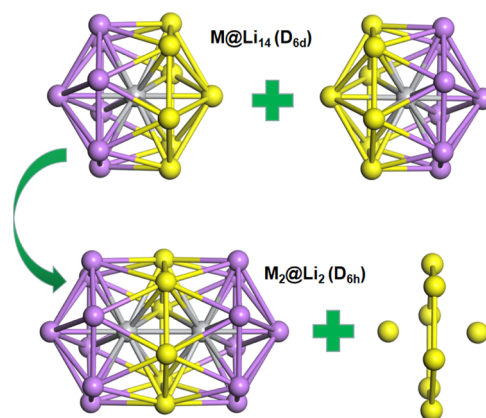


Figure 1. Diagrams of the formation of $\text{M}_1\text{M}_2@\text{Li}_{20}$ ($\text{M}_1, \text{M}_2 = \text{Ti}$ and W) from superatoms $\text{Ti}@\text{Li}_{14}$ or $\text{W}@\text{Li}_{14}$.

superatoms along the C_6 axis, where the approaching top atom of one superatom is replaced by the endohedral transition-metal atom of another superatom, whereas eight redundant lithium atoms are removed with a six-membered Li ring shared between the two superatoms. Frequency calculations check that they are all true local minima. Through structure prediction by CALYPSO code, the face-sharing bitetrahexahedral $\text{Ti}_2@\text{Li}_{20}$ is confirmed as the lowest energy isomer (A1 in Figure 2), where the nearest neighbor low-energy structure (A2) is 0.04 eV higher in energy at the calculations of PW91/SDD level. In fact, the A2 structure can

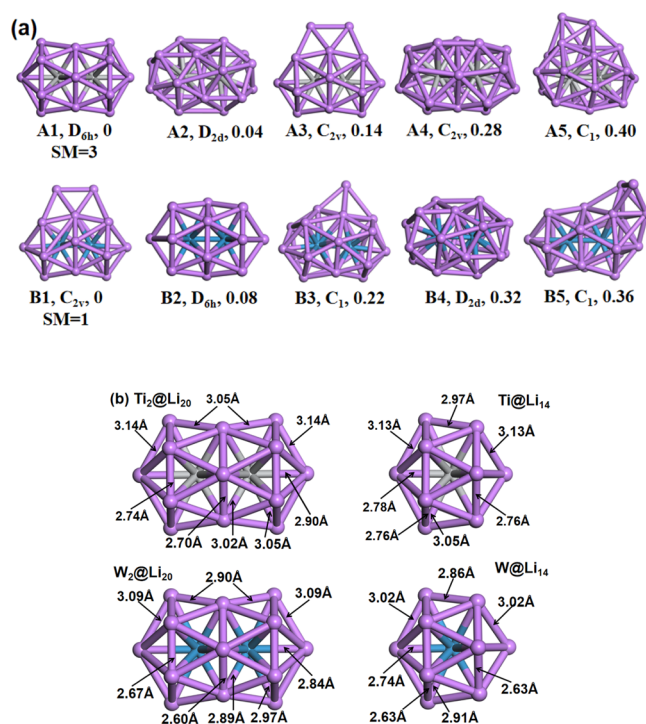


Figure 2. (a) Low-energy structures with the symmetry and the relative energies in eV for $\text{Ti}_2@\text{Li}_{20}$ and $\text{W}_2@\text{Li}_{20}$ at PW91/SDD levels, where the SMs of all of the isomers are 3 for $\text{Ti}_2@\text{Li}_{20}$ and 1 for $\text{W}_2@\text{Li}_{20}$. (b) Bond lengths of the D_{6h} symmetry $\text{Ti}_2@\text{Li}_{20}$, $\text{W}_2@\text{Li}_{20}$, and their respective building units.

Table 1. Results of the Calculated Symmetry, the Fragmentation Energy (E_f), the Binding Energy (E_b) per Atom, the Embedding Energy (D_e), HOMO–LUMO Gaps (E_{gap}), the Dipole Moments (d), and the Total Magnetic Moment (μ_s) of $\text{Ti}_2@$ Li_{20} and $\text{W}_2@$ Li_{20} with D_{6h} symmetry

M ₁ M ₂	methods	symm.	E _f (eV)	E _b (eV)	D _e (eV)	E _{gap} (eV)	d (D)	μ _s (μ _B)
Ti ₂	PW91	D _{6h}	1.25	1.58	11.6	0.35	0	2
	M06L		1.40	1.67	11.2	0.30		
	TPSSH		1.23	1.52	10.1	0.76		
TiW	PW91	C _{6v}	1.13	1.77	15.6	0.31	1.58	0
	M06L		1.31	1.81	14.3	0.40	2.15	
	TPSSH		1.00	1.67	13.6	0.76	1.71	
W ₂	PW91	D _{6h}	1.16	1.95	19.7	0.48	0	0
	M06L		1.28	1.95	17.5	0.65		
	TPSSH		1.07	1.84	17.3	0.85		
Li ₂₂ ^a	PW91	C ₁	1.27	1.17		0.38	0	1
	M06L		1.33	1.27		0.36		
	TPSSH		1.30	1.17		0.73		

^aThe lowest energy structures of Li_{22} , Li_{21} , and Li_{20} are obtained by CALYPSO code, Li_{20} refs.^{40,56}

be viewed as an isosupermolecule of the face-sharing bitetrahexahedral $\text{Ti}_2@$ Li_{20} , where isosupermolecules are proposed by Yang and co-workers⁴⁹ and defined as those clusters with difference structures but resembling each other in electronic structures. As for $\text{W}_2@$ Li_{20} clusters, the face-sharing bitetrahexahedral structure is metastable and lies above the C_{2v} symmetry lowest-energy isomer by 0.08 eV. The corresponding isosupermolecule (B4) is also found.

Compared to the superatom monomer, the Li–Li bond lengths of the face-sharing bitetrahexahedral $\text{Ti}_2@$ Li_{20} along the axis are elongated by 0.08 Å in the middle (2.97 Å for $\text{Ti}@$ Li_{14} and 3.05 Å for $\text{Ti}_2@$ Li_{20}), whereas they are almost unchanged at both ends (3.13 Å for $\text{Ti}@$ Li_{14} and 3.14 Å for $\text{Ti}_2@$ Li_{20}). Along the vertical direction of the axis, all of the Li–Li bond lengths are shortened by 0.06 and 0.02 Å at the sharing face and the ends (2.76 Å for $\text{Ti}@$ Li_{14} , 2.70 and 2.74 Å for $\text{Ti}_2@$ Li_{20}), respectively. Differently, for the face-sharing bitetrahexahedral $\text{W}_2@$ Li_{20} , all of the Li–Li bond lengths along the axis are elongated by 0.04 and 0.07 Å at the middle and ends (2.86 and 3.02 Å for $\text{W}@$ Li_{14} , and 2.90 and 3.09 Å for $\text{W}_2@$ Li_{20}), respectively. Along the vertical axis direction, the Li–Li bond lengths are shortened by 0.03 Å near the sharing face (2.63 Å for $\text{W}@$ Li_{14} and 2.60 Å for $\text{W}_2@$ Li_{20}), whereas at both ends, they are elongated by 0.04 Å (2.63 Å for $\text{W}@$ Li_{14} and 2.67 Å for $\text{W}_2@$ Li_{20}). The changes of Li–Li bond lengths imply that $\text{W}_2@$ Li_{20} is elongated larger along the axis than $\text{Ti}_2@$ Li_{20} , which affects their stability.

3.2. Relative Stability. To investigate the stability of the face-sharing bitetrahexahedral clusters $\text{M}_1\text{M}_2@$ Li_{20} ($\text{M}_1/\text{M}_2 = \text{Ti}$ and W), their average binding energy (E_b) and fragmentation energy (E_f) are calculated, which are respectively defined by the following formulas

$$E_b(\text{M}_1\text{M}_2@ \text{Li}_{20}) = \frac{E(\text{M}_1) + E(\text{M}_2) + 20E(\text{Li}) - E(\text{M}_1\text{M}_2@ \text{Li}_{20})}{22}$$

$$E_f(\text{M}_1\text{M}_2@ \text{Li}_{20}) = E(\text{M}_1\text{M}_2@ \text{Li}_{19}) + E(\text{Li}) - E(\text{M}_1\text{M}_2@ \text{Li}_{20})$$

In the above formulas, M_1 and M_2 denote the dopant Ti or W atom. $E(\text{M}_1\text{M}_2@ \text{Li}_{20})$ represents the total energy of the face-sharing bitetrahexahedral clusters $\text{M}_1\text{M}_2@ \text{Li}_{20}$ ($\text{M}_1/\text{M}_2 = \text{Ti}$ and W), and $E(\text{M}_1\text{M}_2@ \text{Li}_{19})$ is the total energy of their

corresponding fragments $\text{M}_1\text{M}_2@ \text{Li}_{19}$ by taking off one lithium atom from $\text{M}_1\text{M}_2@ \text{Li}_{20}$ ($\text{M}_1/\text{M}_2 = \text{Ti}$ and W). $E(\text{Li})$ and $E(\text{M}_1/\text{M}_2)$ represent the total energy of the free-state Li atom and dopant atom, respectively.

In addition, the embedding energy indicating the feasibility of fusing transition-metal atoms into a cage is also calculated, which is defined as

$$D_e(\text{M}_1\text{M}_2@ \text{Li}_{20}) = E(\text{shell}) + E(\text{core}) - E(\text{M}_1\text{M}_2@ \text{Li}_{20})$$

where the total energies of the outer shell, center dopants, and complexes are denoted as $E(\text{shell})$, $E(\text{core})$, and $E(\text{M}_1\text{M}_2@ \text{Li}_{20})$, respectively. Using such definition, a positive D_e value means that the encapsulation of dopants into the cage is favored thermodynamically.

The calculated results are summarized in Table 1. Compared to the pure Li_{22} cluster, the E_b of $\text{M}_1\text{M}_2@ \text{Li}_{20}$ ($\text{M}_1/\text{M}_2 = \text{Ti}$ and W) is substantially increased, which indicates that doping with transition-metal atoms can significantly enhance their stability. The E_f of $\text{M}_1\text{M}_2@ \text{Li}_{20}$ ($\text{M}_1/\text{M}_2 = \text{Ti}$ and W) is decreasing with the increasement of electron counts, so the clusters with more electrons dissociate easier. All of the values of D_e are positive, demonstrating the feasibility of doping the transition-metal atom into the Li_{20} cage.

The energy gaps (E_{gap}) between the highest occupied molecular orbital (HOMO) and lowest unoccupied molecular orbital (LUMO) are also calculated to examine their electronic stability. The E_{gap} values of $\text{Ti}_2@ \text{Li}_{20}$, $\text{TiW}@ \text{Li}_{20}$, and $\text{W}_2@ \text{Li}_{20}$ are 0.35, 0.31, and 0.48 eV, respectively, at the calculation level of PW91/SDD. Comparing to the gaps of typical magnetic superatoms, these gaps are moderate, for example, $\text{V}@ \text{Na}_8$ (0.69 eV),⁵⁰ VNa_8^- (0.42 eV),⁵¹ $\text{Pb}@ \text{Mn}_{12}@ \text{Pb}_{20}$ (0.38 eV),⁵² MnSr_9 (0.35 eV),⁵³ and so forth. Besides, the standard GGA functionals often underestimate the gaps, and thus the hybrid functional TPSSH is further applied to get more accurate gaps. At the calculation level of TPSSH/SDD, the E_{gap} values of $\text{Ti}_2@ \text{Li}_{20}$, $\text{TiW}@ \text{Li}_{20}$, and $\text{W}_2@ \text{Li}_{20}$ are found to be 0.76, 0.76, and 0.85 eV, respectively, as shown in Table 1.

3.3. Electronic Properties. On the basis of above discussions, the D_{6h} symmetry face-sharing bitetrahexahedral complexes $\text{Ti}_2@ \text{Li}_{20}$ and $\text{W}_2@ \text{Li}_{20}$ have high average binding energies, faired fragmentation energies, and moderate HOMO–LUMO gaps. All of these results suggest that $\text{Ti}_2@ \text{Li}_{20}$ and $\text{W}_2@ \text{Li}_{20}$ can be good candidates for the superatomic

molecule, which is further confirmed by the analyses of MOs and chemical bonding patterns.

For $\text{Ti}_2@Li_{20}$, the number of total effective valence electrons is 28, where one ($2s^1$) and four ($3d^2 4s^2$) effective valence electrons are offered by a lithium atom and a titanium atom, respectively. According to the SVB model,³³ the face-sharing bi-tetrahedral $\text{Ti}_2@Li_{20}$ is assembled by two $\text{Ti}@Li_{14}^{4+}$ (14e) superatoms with an open electronic configuration.

$$\text{Ti}_2@Li_{20}(28e) = 2 \times \text{Ti}@Li_{14}^{4+}(14e) - 8Li^+$$

The equation means that eight Li^+ nuclei (six on the sharing face and two at the top) are removed for the reduction of the electron volume. However, $\text{Ti}@Li_{14}^{4+}$ (14e) is not stable. For simplicity, the building unit $\text{Ti}@Li_{14}$ superatom is adopted, of which the electronic configuration is $1S^2 1P^6 1D^{10}$ with all the effective valence electrons paired following the 18-electron rule. The analysis of the face-sharing bi-tetrahedral $W_2@Li_{20}$ is similar, where the electronic configuration of the building unit $W@Li_{14}$ is $1S^2 1P^6 1D^{10} 2S^2$, fulfilling the 20-electron rule.

To discover the chemical bonding patterns of the face-sharing bitetrahedral complexes $\text{Ti}_2@Li_{20}$ and $W_2@Li_{20}$, their canonical Kohn–Sham MOs of the valence shell are first analyzed. The MO diagrams of $\text{Ti}_2@Li_{20}$ are given in Figure 3, and its schematic MO energy-level diagrams are plotted in Figure 4 for clarity. The first two sets MOs of $\text{Ti}_2@Li_{20}$ (σ_{1s} and σ_{1s}^*) are formed by two 1s superatomic orbitals (SAOs). The next two doubly degenerate sets MOs ($\pi_{1p_x, 1p_y}$ and $\pi_{1p_x, 1p_y}^*$) and two sets (σ_{1p_z} and $\sigma_{1p_z}^*$) are composed of four $1p_x$ and $1p_y$ SAOs and two $1p_z$ SAOs, respectively. The MOs constructed from 1d SAOs are indicated by a number for clear. Then, the two doubly degenerate sets MOs ($\delta_{1d_{x^2-y^2}, xy}$ and $\delta_{1d_{x^2-y^2}, xy}^*$) and two doubly degenerate sets MOs [$\pi_{1d_{xz}, yz}$ and $\pi_{1d_{xz}, yz}^*$ (unlabeled)] consist of four $1d_{x^2-y^2}, xy$ SAOs and four $1d_{xz}, yz$ SAOs, respectively, where the antibonding orbitals $\delta_{1d_{x^2-y^2}, xy}^*$ of β and $\pi_{1d_{xz}, yz}^*$ of $\alpha\beta$ are unoccupied. The 38th occupied MO of α is $\sigma_{1d_{z^2}}$, mainly constructed from $1d_{z^2}$ electrons and a slight 2s electrons, which corresponds to the 47th unoccupied MO for β , and their antibonding MOs $\sigma_{1d_{z^2}}^*$ are not marked because of locating at a high-energy level. The two sets MOs (σ_{2s} and σ_{2s}^*) are composed of two 2s SAOs, where σ_{2s} is the 38th occupied MO for β , corresponding to the first unoccupied MO of α . Therefore, the electronic configuration of $\text{Ti}_2@Li_{20}$ is $(\sigma_{1s})^2 (\sigma_{1s}^*)^2 (\pi_{1p_x, 1p_y})^4 (\pi_{1p_x, 1p_y}^*)^4 (\sigma_{1p_z})^2 (\delta_{1d_{x^2-y^2}, xy})^4 (\sigma_{1d_{z^2}})_\alpha^1 (\sigma_{2s})_\beta^1 (\pi_{1d_{xz}, yz})^4 (\sigma_{1p_z}^*)^2 (\delta_{1d_{x^2-y^2}, xy}^*)^2 (\delta_{1d_{x^2-y^2}, xy})_\beta^0 (\sigma_{2s})_\alpha^0 (\sigma_{2s}^*)_\alpha^0 (\sigma_{2s}^*)_\beta^0 (\sigma_{1d_{z^2}})_\beta^0$, where the energy difference between $(\delta_{1d_{x^2-y^2}, xy})_\alpha^*$ and $(\delta_{1d_{x^2-y^2}, xy})_\beta^*$ forms its HOMO–LUMO gap (0.76 eV/TPSSH). The face-sharing bitetrahedral $\text{Ti}_2@Li_{20}$ is with a magnetic moment of $2 \mu_B$.

For complex $W_2@Li_{20}$, the number of total effective valence electrons is 32, where each tungsten atom offers six ($5d^4 6s^2$) effective valence electrons. From the canonical MOs of $W_2@Li_{20}$ given in Figure 3 and its schematic representation of Figure 4, the eight MOs, including two sets (σ_{1s} and σ_{1s}^*), two doubly degenerate sets ($\pi_{1p_x, 1p_y}$ and $\pi_{1p_x, 1p_y}^*$), and two sets (σ_{1p_z}

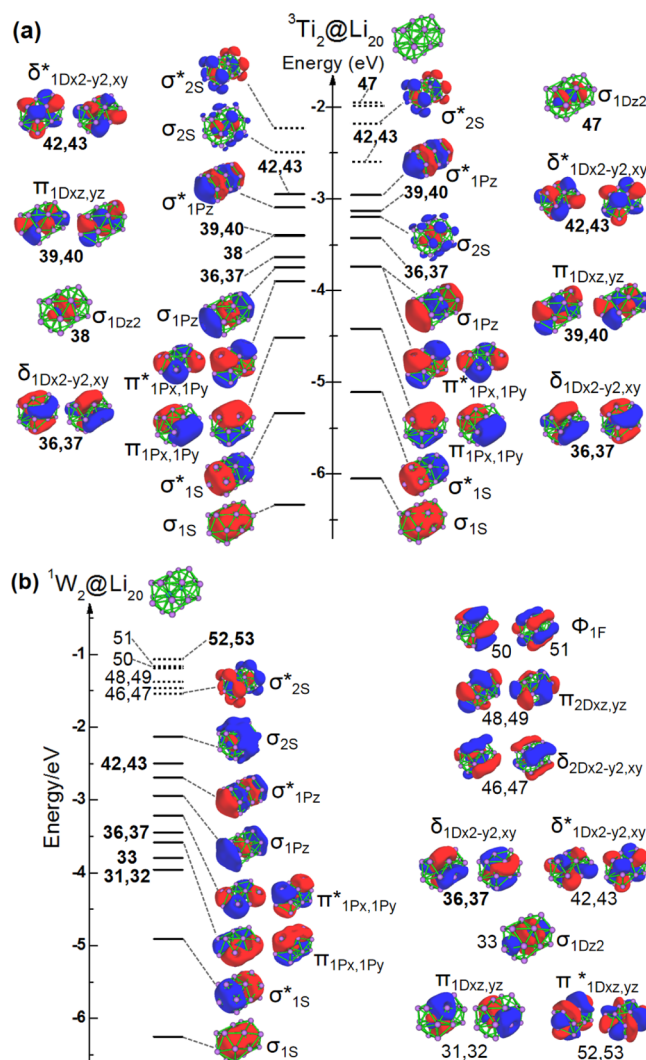


Figure 3. Kohn–Sham MOs diagrams of the face-sharing bitetrahedral superatomic molecules (a) $\text{Ti}_2@Li_{20}$ and (b) $W_2@Li_{20}$.

and $\sigma_{1p_z}^*$), are composed of two 1s SAOs, four $1p_x, y$ SAOs, and two $1p_z$ SAOs, respectively, which are corresponding to four super lone pair (LP) orbitals belonging to each superatom ($1S^2 1P^6$). The MOs marked with the numbers 31, 32, and 33 are bonding orbitals $\pi_{1d_{xz}, yz}$ and $\sigma_{1d_{z^2}}$, constructed from $1d_{xz}, yz$ and $1d_{z^2}$ SAOs, respectively, and their corresponding antibonding orbitals $\pi_{1d_{xz}, yz}^*$ and $\sigma_{1d_{z^2}}^*$ are 2.89 and 3.80 eV higher in energy, respectively. Thus, only the antibonding orbitals $\pi_{1d_{xz}, yz}^*$ (52 and 53) are shown in Figure 3. The two doubly degenerate sets MOs ($\delta_{1d_{x^2-y^2}, xy}$ and $\delta_{1d_{x^2-y^2}, xy}^*$) labeled with the number 36, 37, and 42, 43, respectively, are constituted by four $1d_{x^2-y^2}, xy$ SAOs. The next two sets MOs (σ_{2s} and σ_{2s}^*) are from two 2s SAOs, where the antibonding orbitals σ_{2s}^* are unoccupied. Hence, the electronic configuration of $W_2@Li_{20}$ is $(\sigma_{1s})^2 (\sigma_{1s}^*)^2 (\pi_{1d_{xz}, yz})^4 (\sigma_{1d_{z^2}})^2 (\pi_{1p_x, 1p_y})^4 (\delta_{1d_{x^2-y^2}, xy})^4 (\pi_{1p_x, 1p_y}^*)^4 (\sigma_{1p_z})^2 (\sigma_{1p_z}^*)^2 (\delta_{1d_{x^2-y^2}, xy}^*)^4 (\sigma_{2s})^2 (\sigma_{2s}^*)^0$. With all of the valence electrons paired, the face-sharing bitetrahedral $W_2@Li_{20}$ is nonmagnetic. The energy difference between σ_{2s} and σ_{2s}^* orbitals is its HOMO–LUMO gap (0.85 eV/TPSSH).

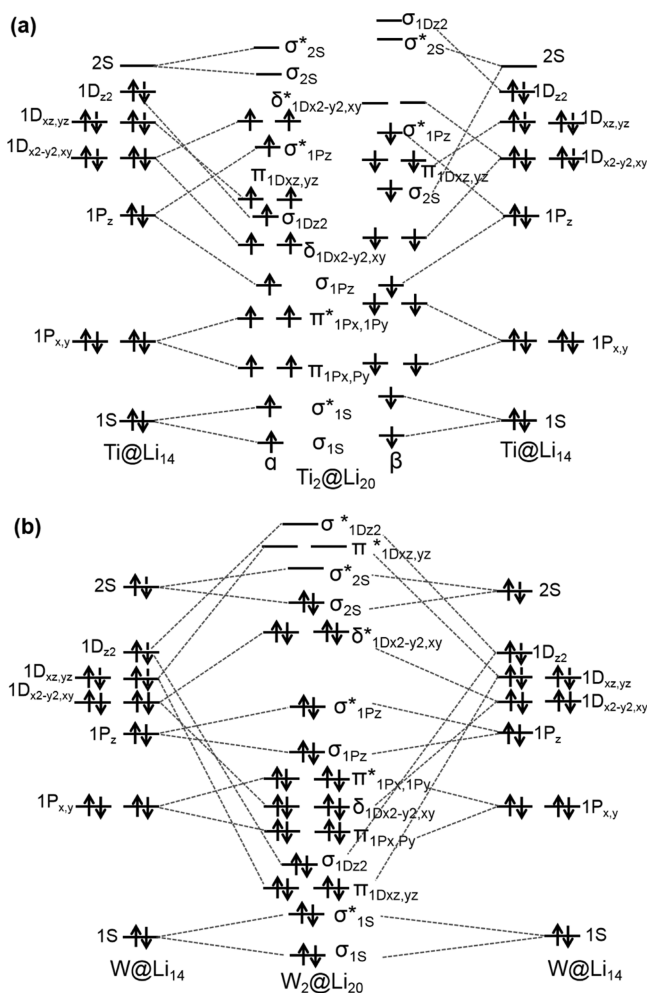


Figure 4. Schematic representation of the MO energy-level diagrams for the face-sharing bi-tetrahexahedral superatomic molecules (a) $\text{Ti}_2@Li_{20}$ and (b) $\text{W}_2@Li_{20}$.

From the canonical MOs, both $\text{Ti}_2@Li_{20}$ and $\text{W}_2@Li_{20}$ can be viewed as superatomic molecules.

To further confirm the existence of superatom–superatom superbonding in $\text{Ti}_2@Li_{20}$ and $\text{W}_2@Li_{20}$, a vivid view on their chemical bonding patterns is given by the adaptive natural density partitioning (AdNDP) method,⁵⁴ which is a chemical bonding analysis tool developed by Zubarev and Boldyrev. As shown in Figure 5, the $\text{Ti}_2@Li_{20}$ cluster has four 15c–2e super LPs (super s, p_x , p_y , and p_z) located in each superatom, and five 22c–2e superatom–superatom bonds (σ , 2π , 2δ) in the whole cluster. The remaining two electrons are two single electrons filled in the 42nd and 43rd MOs $\delta_{1d_{x^2-y^2,xy}}^*$ of α . Thus, the super bond order of $\text{Ti}_2@Li_{20}$ is four because of the occupied antibonding orbitals $\delta_{1d_{x^2-y^2,xy}}^*$ of α . For $\text{W}_2@Li_{20}$, the AdNDP analysis gives six 15c–2e super LPs (super s, p_x , p_y , p_z , $d_{x^2-y^2}$, and d_{xy}) in each superatom, four 22c–2e superatom–superatom bonds (2σ , 2π) in the whole cluster. Correspondingly, the super bond order of $\text{W}_2@Li_{20}$ is also four. The super LPs, super $\sigma 2\pi 2\delta$ -bonds, and super $2\sigma 2\pi$ -bonds for $\text{Ti}_2@Li_{20}$ and $\text{W}_2@Li_{20}$ are all with high occupancy numbers. Therefore, the face-sharing bitetrahexahedral complexes $\text{Ti}_2@Li_{20}$ and $\text{W}_2@Li_{20}$ can be viewed as superatomic molecules. Such a bonding pattern is similar to those of 26e $\text{Li}_{20}\text{Mg}_3$ and 30e $\text{Li}_{18}\text{Mg}_3\text{Al}_2$ clusters, which have a high bond

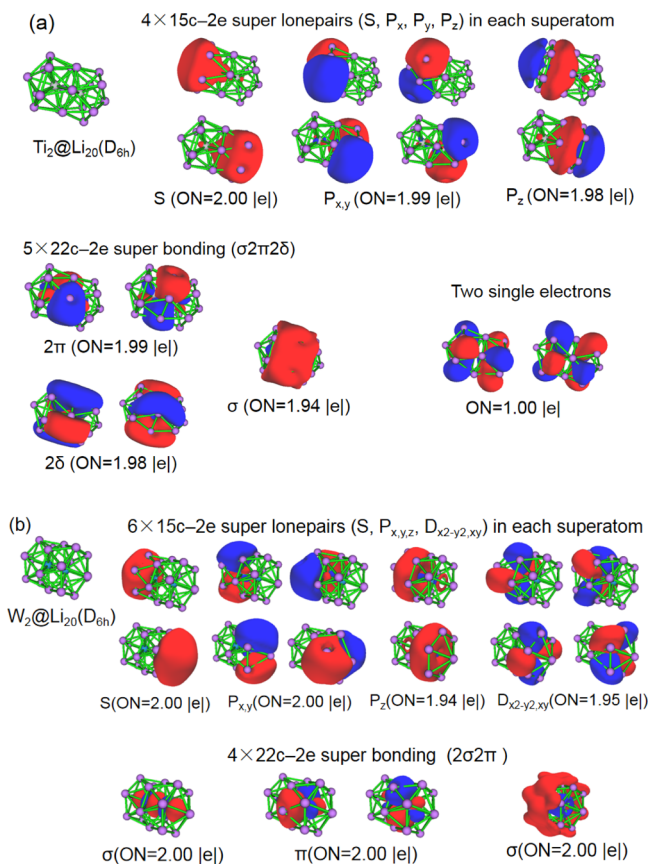


Figure 5. AdNDP localized natural bonding orbitals of the face-sharing bi-tetrahexahedral superatomic molecules (a) $\text{Ti}_2@Li_{20}$ and (b) $\text{W}_2@Li_{20}$.

order with quintuple superbonding ($\sigma 2\pi 2\delta$).³⁸ As comparison, the MO diagrams and bonding patterns of their corresponding isosuperatoms are given in Figures S1 and S2, which can be viewed as the assembly of two 14c superatoms $\text{Ti}@Li_{13}$ and $\text{W}@Li_{13}$ shared by six nuclei.

On the basis of the above analysis, the face-sharing bitetrahexahedral $\text{Ti}_2@Li_{20}$ and $\text{W}_2@Li_{20}$ clusters are identified as superatomic molecules in geometric and electronic shells. From SAOs to MOs, the orders of 1D energy level for $\text{W}_2@Li_{20}$ change a lot as the schematic representations of MO energy-level diagrams shown in Figure 4. The positions of $\pi_{1d_{xz,yz}}$ and $\sigma_{1d_{z^2}}$ MOs energy levels for $\text{W}_2@Li_{20}$ are lower than that of $\text{Ti}_2@Li_{20}$. The differences can be derived from their structures. Compared to $\text{Ti}_2@Li_{20}$, where the Li–Li bond lengths at both ends along the axis remain the lengths of its building unit $\text{Ti}@Li_{14}$, the corresponding Li–Li bond lengths of $\text{W}_2@Li_{20}$ are elongated by 0.07 Å. Meanwhile, the Li–W bond lengths of $\text{W}_2@Li_{20}$ in the inner away from the sharing face are elongated by 0.06 Å, whereas the corresponding Li–Ti bond lengths of $\text{Ti}_2@Li_{20}$ is almost unchanged. The more elongated geometric structure of $\text{W}_2@Li_{20}$ leads to lower positions of $\pi_{1d_{xz,yz}}$ and $\sigma_{1d_{z^2}}$ MOs, also suggesting that the interactions of its two superatoms are long range. Further, the stability of building units is needed to consider. For the tetrahexahedral superatoms, the 18-electron systems are easier to exist than the 20-electron systems, where all of the IVB transition-metal elements, namely, Ti, Zr, and Hf atoms embedded in the center of the Li_{14} cage, the formed

superatoms are with the lowest energy, whereas among the VIB transition-metal elements, only W atom.³⁹ This phenomenon predicts that it is more strenuous to maintain the superatomic structures if the systems have too many electrons. Thus, it is reasonable that the lowest energy isomer is not the face-sharing bitetrahexahedral motif, but the C_{2v} symmetry B1 for $W_2@Li_{20}$. Meanwhile, the existence of isosupermolecules gives another way to assemble materials from superatoms, and the binary superatoms are excellent choices as the building blocks, for example, VLi_n ($n = 1-13$) clusters, the value of n from 8 to 13 corresponding to superatom,⁵⁵ promoting the possibility of assembled materials.

Interestingly, when the endohedral transition-metal elements are a hybrid of Ti and W atoms, the face-sharing bitetrahexahedral structure for $TiW@Li_{20}$ is still retained, but with a C_{6v} symmetry because the hetero-Ti and W atoms decrease the symmetry. The corresponding electronic configuration of the $TiW@Li_{20}$ cluster is $(\sigma_{1s})^2(\sigma_{1s}^*)^2(\pi_{1p_x,1p_y})^4(\sigma_{1d_z})^2(\sigma_{1d_{xz,yz}})^4(\delta_{1d_{x^2-y^2,xy}})^4(\pi_{1p_x,1p_y}^*)^4(\sigma_{1p_z})^2(\sigma_{1p_z}^*)^2(\delta_{1d_{x^2-y^2,xy}}^*)^4(\sigma_{2s})^0(\sigma_{2s}^*)^0$, similar to the electron filling orders of superatomic molecule $W_2@Li_{20}$ except the orders of $\sigma_{1d_{xz,yz}}$ and σ_{1d_z} . The MO diagrams and chemical bonding patterns of $TiW@Li_{20}$ are shown in Figures S3 and S4, respectively, where the energy difference of $\delta_{1d_{x^2-y^2,xy}}^*$ and σ_{2s} forms its HOMO–LUMO gap (0.76 eV/TPSSH), equal to the gap of $Ti_2@Li_{20}$ at the TPSSH/SDD level. Different from them, the hetero-bitetrahexahedral superatomic molecule $TiW@Li_{20}$ produces a dipole moment (1.71 D/TPSSH), as shown in Table 1. The present results can be enhanced by proper selection of the dopants or the building blocks. Consequently, by varying the central transition-metal elements of $M_1M_2@Li_{20}$ ($M_1M_2 = Ti$ and W), their properties can be precisely tuned while with the geometric structures fixed.

4. CONCLUSIONS

A D_{6h} symmetry face-sharing bi-tetrahexahedral motif of superatomic molecules is found via the assembly of two $Ti@Li_{14}$ or $W@Li_{14}$ superatoms along the axis, where the middle six lithium atoms are shared by the two superatoms. Combining CALYSPO structure prediction method and density functional theory calculations, the face-sharing bitetrahexahedral $Ti_2@Li_{20}$ and $W_2@Li_{20}$ are proved to be low-energy isomers. According to MO diagrams and chemical bonding patterns, the superatomic molecules $Ti_2@Li_{20}$ or $W_2@Li_{20}$ are formed by quadruple superbonding between two tetrahexahedral superatoms. Interestingly, $Ti_2@Li_{20}$ is identified as a magnetic superatomic molecule with a magnetic moment of $2 \mu_B$, whereas $W_2@Li_{20}$ remains nonmagnetic. Moreover, when two hetero-tetrahexahedral superatoms $Ti@Li_{14}$ and $W@Li_{14}$ are chosen as the building blocks, the resulting superatomic molecule $TiW@Li_{20}$ still features in a face-sharing bi-tetrahexahedral structure, but with a reduced symmetry of C_{6v} and a net dipole moment (1.71 D/TPSSH). Our findings are encouraging because the superatomic molecules are stable despite plenty of atoms shared between two superatoms as well as no ligand protection. Meanwhile, the properties, such as the HOMO–LUMO gap, magnetic moment, and dipole moment, can be tailored by changing the endohedral transition-metal elements.

■ ASSOCIATED CONTENT

§ Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpca.9b01855.

Kohn–Sham MO diagrams and AdNDP localized natural bonding orbitals of the three superatomic molecules, namely, two isosupermolecules $Ti_2@Li_{20}$ (D_{2d} , A2) and $W_2@Li_{20}$ (D_{2d} , B4), and one hetero-superatomic molecule $TiW@Li_{20}$ (PDF)

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Notes

The author declares no competing financial interest.

■ ACKNOWLEDGMENTS

This work is supported by the Special Foundation for Theoretical Physics Research Program of China (no. 11847119) and the Innovation and Strong School Research Project of Guangdong Ocean University (no. GDOU2017052621).

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