碰撞

Thermal Physics Sheet

Log Creative

1 压强

状态方程

$$pV = \nu RT \tag{1.1}$$

微观公式

$$p = nkT = \frac{2}{3}n\bar{\varepsilon} = \frac{2}{3}n\left(\frac{1}{2}m\bar{v^2}\right)$$
$$= \frac{2}{3}n\left(\frac{3}{2}m\bar{v_x^2}\right) \quad (1.2) \quad \mathbf{3}$$

范德瓦尔斯

$$\left(p + \nu^2 \frac{a}{V_m^2}\right) (V - \nu b) = \nu RT \qquad (1.3)$$

$$b = 4N_A \frac{4}{3}\pi \left(\frac{d}{2}\right)^2 \tag{1.4}$$

Boltzmann

$$p = p_0 e^{-\frac{Mgz}{RT}} (\text{Const. } T) \tag{1.5}$$

$$n = n_0 e^{-\frac{\varepsilon_0}{kT}} (\text{Const. } T) \tag{1.6}$$

2 分量

方均根(速率最大)

$$\sqrt{\bar{v^2}} = \sqrt{\frac{3RT}{M}} \tag{2.1}$$

$$\sqrt{\bar{v_x^2}} = \sqrt{\frac{RT}{M}} \tag{2.2}$$

均值(速率中等)

$$\bar{v} = \sqrt{\frac{8RT}{\pi M}} \tag{2.3}$$

$$\bar{v_x} = \sqrt{\frac{RT}{2\pi M}} \tag{2.4}$$

最可几速率(最概然速率,速率最小)

$$v_p = \sqrt{\frac{2RT}{M}} \tag{2.5}$$

(分量最概然速度大小为0.)

Maxwell

$$\frac{\mathrm{d}N_v}{N} = \left(\frac{m}{2\pi kT}\right)^{\frac{3}{2}} e^{-\frac{mv^2}{2kT}} 4\pi v^2 \mathrm{d}v$$
$$= f(v)\mathrm{d}v \tag{2.6}$$

$$\frac{\mathrm{d}N_{v_x}}{N} = \left(\frac{m}{2\pi kT}\right)^{\frac{1}{2}} e^{-\frac{mv_x^2}{2kT}} \mathrm{d}v$$
$$= f(v_x)\mathrm{d}v \tag{2.7}$$

$$N = \frac{1}{4}n\bar{v} \tag{2.8}$$

$$\Delta N_{0 \sim \beta v_{p(x)}} = \frac{N}{2} \operatorname{erf}(\beta)$$
 (2.9)

$$\Delta N_{0 \sim \beta v_p} = N \left[\operatorname{erf}(\beta) - \frac{2}{\sqrt{\pi}} \beta e^{-\beta^2} \right]$$
(2.10)

$$\operatorname{erf}(\beta) = \frac{2}{\sqrt{\pi}} \int_0^{\beta} e^{-x^2} dx$$
 (2.11)

能量按自由度均分定理

$$\frac{1}{2}m\bar{v^2} = \frac{3}{2}kT \tag{2.12}$$

$$\frac{1}{2}m\bar{v_x^2} = \frac{1}{2}kT \tag{2.13}$$

3 多项

摩尔平均总能量(平动、转动、振动)

$$\bar{\varepsilon}_{\rm m} = \frac{1}{2} (t + r + 2s) RT$$

$$= \frac{3}{2} RT (\text{single})$$

$$= \frac{5}{2} RT (\text{hard dual})$$

$$= \frac{7}{2} RT (\text{elastic dual})$$

$$= \frac{6}{2} RT (\text{multi})$$
 (3.1)

热容

$$C = mc (3.2)$$

$$C_{\rm m} = Mc \tag{3.3}$$

$$C_{V,m} = \frac{1}{2} (t + r + 2s) R$$
 (3.4)

扩散

$$\sigma = \pi d^2 \tag{3.5}$$

$$\bar{Z} = \sqrt{2}\sigma\bar{v}n\tag{3.6}$$

$$\bar{\lambda} = \frac{1}{\sqrt{2}\sigma n} \tag{3.7}$$

分子按自由程的分布: N_0 中自由程大于x的分子数,求微分得连续分布

$$N = N_0 e^{-\frac{x}{\lambda}} \tag{3.8}$$

粘滞、导热、扩散系数

$$dy = -c \left(\frac{dx}{dz}\right)_{z_0} dS dt \qquad (3.9)$$

$$\eta = \frac{1}{3}\rho\bar{v}\bar{\lambda} \tag{3.10}$$

$$\kappa = \eta c \tag{3.11}$$

$$D = \frac{\eta}{\rho} \tag{3.12}$$

$$c = \frac{C_{V,m}\nu}{m} \tag{3.13}$$

4 热力学定律

第一定律 加进一个的系统中的热量+对系统所做的功=系统内能的增加

$$Q + A = \Delta U \tag{4.1}$$

第二定律 不可能有这样一个过程,它的唯一结果**只是**从一个热库取出热量,并把它转化为功.

没有任何一台热机, 在从 T_1 取得热量 Q_1 ,而在 T_2 放出热量 Q_2 的过程中所做的功比可逆机更大.对于可逆机,

$$W = Q_1 - Q_2 = Q_1 \left(1 - \frac{T_2}{T_1} \right) \quad (4.2)$$

系统的熵 如果 ΔQ 是可逆地加在温度 为T的系统中的热量,那么这个系统的 熵增为

$$\Delta S = \frac{\Delta Q}{T} \tag{4.3}$$

熵为:

$$S(\boldsymbol{V}, \boldsymbol{T}) = R\left(\ln \boldsymbol{V} + \frac{1}{\gamma - 1}\ln \boldsymbol{T}\right) + a$$
(4.4)

当T = 0时,S = 0(**第三定律**).此时的熵 定义为:

在**可逆变化**中,系统所有部分(包括热库)的总熵**不变**.

在**不可逆变化**中,系统的总熵始终不断 **增加.无摩擦的准静**态过程是可逆的.

通用公式

$$Q_V = \Delta U = \nu C_{V,m} \Delta T \tag{4.5}$$

$$Q_p = \Delta H = \nu C_{p,m} \Delta T \tag{4.6}$$

$$A = -\int_{V_1}^{V_2} p dV$$
 (4.7)

$$H = U + pV \tag{4.8}$$

$$TdS = dU + pdV \tag{4.9}$$

理想气体

$$\Delta U = \nu C_{V \,\mathrm{m}} \Delta T \qquad (4.10)$$

$$C_{p,m} - C_{V,m} = R$$
 (4.11)

$$\Delta H = \nu C_{p,m} \Delta T \qquad (4.12)$$

范德瓦尔斯气体

$$\Delta U = \nu \left[C_{V,m} \Delta T - a \Delta \left(\frac{1}{V_{m}} \right) \right]$$
(4.13)

$$C_{p,m} - C_{V,m} = \frac{R}{1 - \frac{2a(V_m - b)^2}{RTV_m^3}}$$
 (4.14)

$$\Delta H_{\rm m} = (C_{V,m} + R)\Delta T - a\Delta \left(\frac{1}{V_{\rm m}}\right) + \frac{RT_2b}{V_2 - b} - \frac{RT_1b}{V_1 - b} \quad (4.15)$$

热力学过程 5

A:外界对系统所做的功 Q:系统从外界吸收的热量 C_m :摩尔热容

 ΔS :熵变 (理想气体)

等容过程 V = Const.

$$A = 0 (5.1)$$

(5.2)

(5.8)

$$Q = \nu C_{V,m} (T_2 - T_1)$$

$$C_{V,\mathrm{m}} = \frac{R}{\gamma - 1} \tag{5.3}$$

$$\Delta S = \nu C_{V,\text{m}} \ln \frac{T_2}{T_1} \tag{5.4}$$

等压过程 p = Const.

$$A = -p(V_2 - V_1)$$

= $-\nu R(T_2 - T_1)$ (5.5)

$$Q = \nu C_{p,m} (T_2 - T_1)$$
 (5.6)

$$C_{V,\mathrm{m}} = \frac{\gamma R}{\gamma - 1} \tag{5.7}$$

$$\Delta S = \nu C_{p,\text{m}} \ln \frac{T_2}{T_1}$$

等温过程 T = Const.

$$A = -p_1 V_1 \ln \frac{V_2}{V_1}$$

$$= -\nu R T_1 \ln \frac{V_2}{V_1}$$
(5.9)

$$Q = -A \tag{5.10}$$

$$C_{\rm m} = \infty \tag{5.11}$$

$$\Delta S = \nu R \ln \frac{V_2}{V_1} \tag{5.12}$$

绝热过程 Q=0

泊松方程

 $C_m = 0$

$$pV^{\gamma} = \text{Const.}$$
 (5.13)

$$TV^{\gamma-1} = \text{Const.}$$
 (5.14)

$$\frac{p^{\gamma - 1}}{T^{\gamma}} = \text{Const.} \tag{5.15}$$

$$A = \frac{p_1 V_1}{\gamma - 1} \left[\left(\frac{V_1}{V_2} \right)^{\gamma - 1} - 1 \right]$$

$$= \nu C_{V,m}(T_2 - T_1) \tag{5.16}$$

$$Q = 0 (5.17)$$

(5.18)

$$\Delta S = 0 \tag{5.19}$$

$$|$$
多方过程 $pV^n = \text{Const.}$

$$A = \frac{p_1 V_1}{n - 1} \left[\left(\frac{V_1}{V_2} \right)^{n - 1} - 1 \right]$$

$$= \frac{\nu R}{n-1}(T_2 - T_1) \qquad (5.20) \qquad dp = -\frac{Mgp}{RT}dz \text{(Boltzmann)}$$

$$Q = \nu \left(C_{V,m} - \frac{R}{n-1} \right) (T_2 - T_1) \qquad h = \frac{C_{p,m} T_0}{Mg} \left[1 - \left(\frac{p}{p_0} \right)^{\frac{\gamma - 1}{\gamma}} \right]$$
(5.21)

$$C_m = \frac{\gamma - n}{1 - n} C_{V, \mathbf{m}} \tag{5}$$

$$\Delta S = \nu C_{V,m}(\gamma - n) \ln \frac{V_2}{V_1}$$
 (5.23)

自由膨胀 A=0

这一过程不是准静态过程

绝热节流 H = Const.

焦汤系数

$$\alpha \equiv \lim_{\Delta p \to 0} \left(\frac{\Delta T}{\Delta p} \right)_H = \left(\frac{\partial T}{\partial p} \right)_H \quad (5.24)$$

α	类型(室温下)	效应
+	氮、氧、空气	制冷效应、正效应
_	氢气	制温效应、负效应
0	理想气体	零效应*

非理想气体的对应温度为**转换温度**.上转换 温度 $T^{\circ} = \frac{2a}{Rb}$

其他热力学方程

效率、冷却系数

$$\eta = \frac{A}{Q_1} = 1 - \frac{|Q_2|}{Q_1} \tag{6.1}$$

$$\varepsilon = \frac{Q_2}{Q_1 - Q_2} \tag{6.2}$$

卡诺、可逆

$$\eta = 1 - \frac{T_2}{T_2}$$
(6.3)

$$\varepsilon = \frac{T_2}{T_1 - T_2} \qquad (6.4)$$

$$\oint_{\text{invertible cycle}} \frac{\mathrm{d}Q}{T} = 0 \tag{6.5}$$

热力学定律常用(以及 4.9)

$$\left(\frac{\partial U}{\partial V}\right)_{V} = T \left(\frac{\partial p}{\partial T}\right)_{V} - p \tag{6.6}$$

$$dU = \left(\frac{\partial U}{\partial T}\right)_{V} dT + \left(\frac{\partial U}{\partial V}\right)_{T} dU$$

$$W = \frac{N!}{\prod_{i=1}^{n} N_i!} \tag{6.8}$$

$$S = k \ln W \tag{6.9}$$

大气温度梯度

$$\frac{\mathrm{d}T}{\mathrm{d}z} = -\frac{\gamma - 1}{\gamma} \frac{T}{p} \rho g \tag{6.10}$$

$$dp = -\frac{Mgp}{RT}dz(Boltzmann) \qquad (6.11)$$

$$h = \frac{C_{p,m}T_0}{Mg} \left[1 - \left(\frac{p}{p_0}\right)^{\frac{\gamma - 1}{\gamma}} \right] \quad (6.12)$$

(5.22) **7** 相变

克拉珀龙方程衍生式

$$\frac{\mathrm{d}p}{\mathrm{d}T} = \frac{l}{T(v_2 - v_1)}$$

$$= \frac{u_2 - u_1 + p(v_2 - v_1)}{T(v_2 - v_1)}$$

$$\approx \frac{l}{Tv_2}$$

$$\approx \frac{\Delta p}{LT}$$
(7.1)

_ 理想气体蒸气压方程

$$\ln p = -\frac{L}{RT} + \text{Const.} \qquad (7.2)$$

范德瓦尔斯相变临界点参量

$$\begin{cases}
T_k = \frac{8a}{27bR} \\
V_{mk} = 3b \\
p_k = \frac{a}{27b^2}
\end{cases}$$
(7.3)

$$\frac{RT_k}{p_k V_{\text{m}k}} = \frac{8}{3} \tag{7.4}$$

常数 8

$$R = \frac{k}{N_A} \tag{8.1}$$

	Value
R	8.314
k	1.381×10^{-23}
N_A	6.022×10^{23}

$$n = \frac{N_A}{V_{\rm m}} \tag{8.2}$$

$$\begin{array}{c|c} & \text{Value} \\ \hline V_{\text{m}}(\text{STP}) & 22.41 \end{array}$$

$$A = eU \tag{8.3}$$

	Value	
e	1.602×10^{-19} C	
eV	$1.602 \times 10^{-19} \text{J}$	