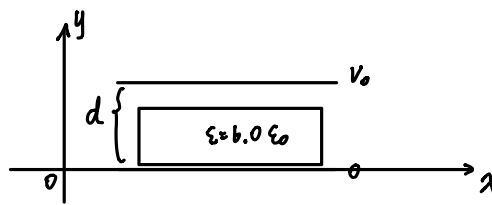


P4-1 a). Since $\nabla^2 V = 0$, we can assume that the V in air is $V_{\text{air}} = ay + b$, and the V in dielectric is $V_{\text{die}} = my + n$. Then $E = -\nabla V$, we have $\vec{E}_{\text{air}} = -a\hat{y}$ and $\vec{E}_{\text{die}} = -m\hat{y}$.



Moreover, we know that

$$\begin{cases} V_{\text{die}}(0) = 0 \Rightarrow n = 0 \\ V_{\text{air}}(d) = ad + b = V_0 \\ V_{\text{die}}(0.8d) = V_{\text{air}}(0.8d) \Rightarrow a \times 0.8d + b = m \times 0.8d \\ 1.0 \epsilon_0 E_{\text{die}}(0.8d) = \epsilon_0 E_{\text{air}}(0.8d) \Rightarrow b \times 0.8m = 0.8a \end{cases}$$

We can get that

$$\begin{cases} a = \frac{3V_0}{d} \\ b = -2V_0 \\ m = \frac{V_0}{2d} \\ n = 0 \end{cases}$$

Then, in dielectric slab: $V = \frac{V_0}{2d} y$ $\vec{E} = -\frac{V_0}{2d} \hat{y}$

b). $V = \frac{3V_0}{d} y - 2V_0$, $\vec{E} = -\frac{3V_0}{d} \hat{y}$

c). $P_{\text{upper}} = \frac{3\epsilon_0 V_0}{d}$

$P_{\text{lower}} = -\frac{V_0}{2d} \times 1.0 \epsilon_0 = -\frac{3\epsilon_0 V_0}{2d}$

d). Without slab: $V = \frac{V_0}{d} y$

$\vec{E} = -\frac{V_0}{d} \hat{y}$

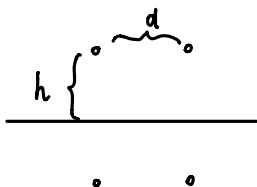
P4-5 a). $V(x, y, z) = \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{\sqrt{x^2 + (y-d)^2 + z^2}} - \frac{1}{\sqrt{x^2 + (y+d)^2 + z^2}} \right)$

$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0$.

b). $V(x, y, z) = \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{R_+} - \frac{1}{R_-} \right) + V_0$

c). $\vec{F} = -\hat{y} \cdot \frac{Q}{4\pi\epsilon_0 (2d)^2} = \frac{-Q^2}{16\pi\epsilon_0 d^2} \hat{y}$

P4-11



$\Delta V = \frac{P}{2\pi\epsilon_0} \ln\left(\frac{d}{a}\right) + \frac{-P}{2\pi\epsilon_0} \ln\left(\frac{d}{a}\right) + \frac{-P}{2\pi\epsilon_0} \ln\left(\frac{\sqrt{4h^2 + d^2}}{2h}\right) + \frac{P}{2\pi\epsilon_0} \ln\left(\frac{2h}{\sqrt{4h^2 + d^2}}\right)$

$C = \frac{P}{\Delta V} = \frac{\pi\epsilon_0}{\ln\left(\frac{2dh}{a\sqrt{4h^2 + d^2}}\right)}$

P4-14 a). Let l_1 denotes the distance to $+p_i$
and l_2 denotes the distance to $-p_i$

$$V = \frac{\rho_l}{2\pi\epsilon_0} \ln \frac{l_2}{l_1}$$

When on the surface of the long wire,

$$l_2 = b + (c_1 - a_1) \quad l_1 = b + (c_1 - a_1)$$

When on the surface of the circular tunnel,

$$l_2 = b + (c_2 - a_2) \quad l_1 = b - (c_2 - a_2)$$

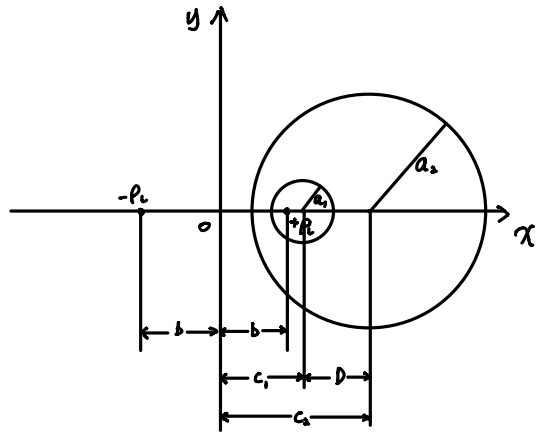
$$\text{Then } V_1 - V_2 = \frac{\rho_1}{2\pi\epsilon_0} \ln \left[\frac{b+(c_1-a_1)}{b-(c_1-a_1)} \cdot \frac{b-(c_2-a_2)}{b+(c_2-a_2)} \right]$$

We know that from the text book P. 168-169.

$$\begin{cases} b^2 = c_1^2 - a_1^2 = c_2^2 - a_2^2 \\ c_2 - c_1 = D \end{cases} \Rightarrow \begin{cases} c_1 = \frac{1}{2D} (a_2^2 - a_1^2 - D) \\ c_2 = \frac{1}{2D} (a_2^2 - a_1^2 + D) \end{cases}$$

$$b = \sqrt{c_1^2 - a_1^2} = \frac{1}{2D} \sqrt{(a_2^2 - a_1^2 - D)^2 - 4D^2 a_1^2}$$

$$C = \frac{P_c}{V_1 - V_2} = \frac{2\pi\epsilon_0}{\ln \left[\frac{b + (c_1 - a_1)}{b - (c_1 - a_1)} \cdot \frac{b - (c_2 - a_2)}{b + (c_2 - a_2)} \right]}$$



$$b). F = \frac{\rho_i^2}{2\pi\epsilon_0 (2b)^2} = \frac{\rho_i^2 D^2}{2\pi\epsilon_0 [(a_2^2 - a_1^2 - D)^2 - 4D^2 a_1^2]}$$

P4-17 Suppose the point is (x, y, z) .

$$V_1 = \frac{1}{4\pi\epsilon_0} \left(\frac{Q}{\sqrt{(x-d)^2 + y^2 + z^2}} - \frac{Q_1}{\sqrt{(x+d)^2 + y^2 + z^2}} \right)$$

$$V_2 = \frac{1}{4\pi\epsilon_0\epsilon_0} \left(\frac{Q + Q_1}{\sqrt{(x-d)^2 + y^2 + z^2}} \right)$$

$$\nabla^2 V_1 = \nabla^2 V_2 = 0.$$

$$\epsilon_0 \epsilon_1 \left. \frac{\partial V_1}{\partial x} \right|_{x=0} = \epsilon_0 \epsilon_2 \left. \frac{\partial V_2}{\partial x} \right|_{x=0}.$$

$$\frac{Q - Q_1}{\epsilon_1} = \frac{Q + Q_2}{\epsilon_2}$$

$$Q_1 = Q_2 = \frac{\epsilon_2 - \epsilon_1}{\epsilon_2 + \epsilon_1} Q$$

P4-23 a). $V = a\phi + b$

$$\left. \begin{aligned} V(0) &= b = 0 \\ V(\alpha) &= a\alpha + b = V_0 \end{aligned} \right\} \Rightarrow \begin{cases} a = \frac{V_0}{\alpha} \\ b = 0 \end{cases}$$

Therefore, $V = \frac{V_0}{\alpha} \cdot \phi$ for $0 < \phi < \alpha$

b) $\left. \begin{aligned} V(\alpha) &= a\alpha + b = V_0 \\ V(2\pi) &= a \cdot 2\pi + b = 0 \end{aligned} \right\} \Rightarrow \begin{cases} a = -\frac{V_0}{2\pi - \alpha} \\ b = \frac{2\pi V_0}{2\pi - \alpha} \end{cases}$

Therefore, $V = -\frac{V_0}{2\pi - \alpha} \phi + \frac{2\pi V_0}{2\pi - \alpha}$

P4-28 a). $V(b, \theta) = V_0$

$V(R, \theta) = -E_0 R \cos \theta + V_0$ for $R \geq b$.

$$V(R, \theta) = \frac{B_0}{R} + \left(\frac{B_1}{R^2} - E_0 \right) \cos \theta - \sum_{n=2}^{\infty} B_n R^{-(n+1)} P_n(\cos \theta)$$

$B_0 = bV_0$

$B_1 = E_0 b^3$

$B_n = 0$ when $n \geq 2$

$V(R, \theta) = \frac{bV_0}{R} + \left(\frac{E_0 b^3}{R^2} - E_0 \right) R \cos \theta$.

b). $\vec{E}_R = -\frac{\partial V}{\partial R} = \frac{bV_0}{R^2} + \left(\frac{E_0 b^3}{R^3} + E_0 \right) \cos \theta$

$\vec{E}_\theta = -\frac{\partial V}{\partial \theta} = \left(\frac{E_0 b^3}{R^2} - E_0 \right) R \sin \theta$

$\vec{E}(R, \theta) = \hat{a}_R \left(\frac{bV_0}{R^2} + \left(\frac{E_0 b^3}{R^3} + E_0 \right) \cos \theta \right) + \hat{a}_\theta \left(\frac{E_0 b^3}{R^2} - E_0 \right) R \sin \theta$