

vv256: Bessel's equation. Series solutions

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UM-SJTU Joint Institute

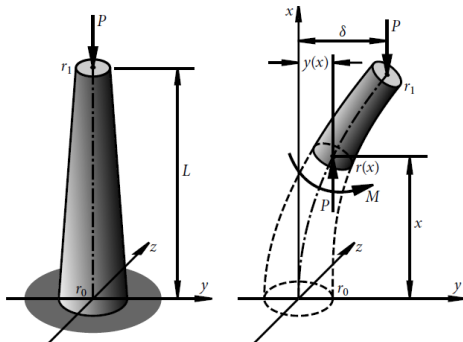
October 29, 2019

Motivating Examples: Buckling of a Tapered Column

Consider the stability of a tapered column of length L fixed at the base $x = 0$ and free at the top $x = L$. The column is subjected to an axial compressive load P at the top.

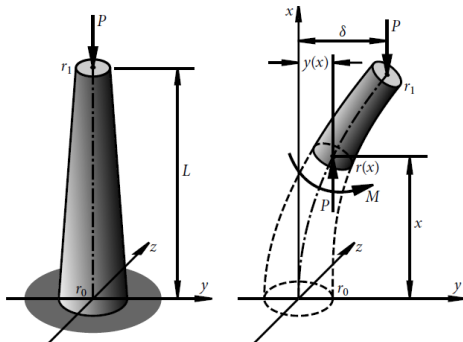
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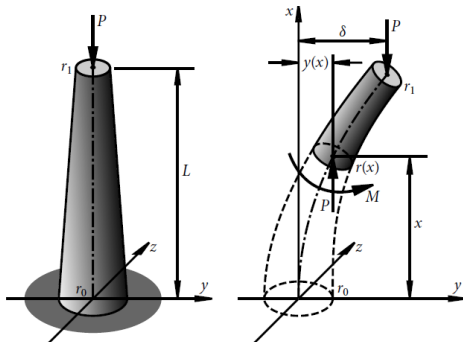
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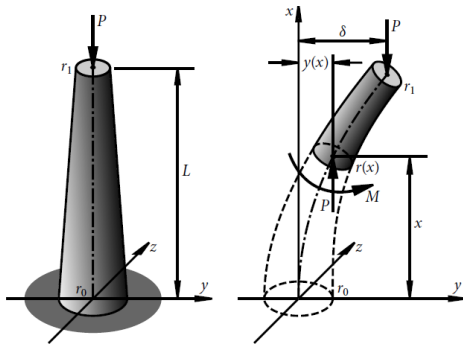
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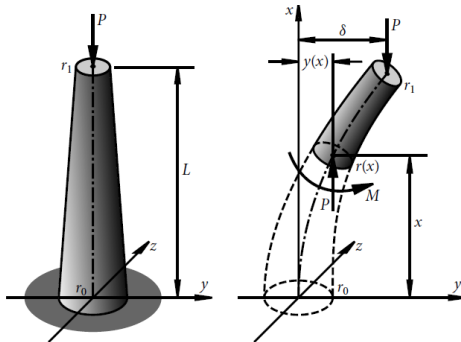
The cross-section of the column is of circular shape, with radii r_0 at the base and $r_1 < r_0$ at the top, respectively, varying linearly along the length x . The modulus of elasticity for the column material is E . Determine the **buckling load** P_{cr} when the column loses its stability.

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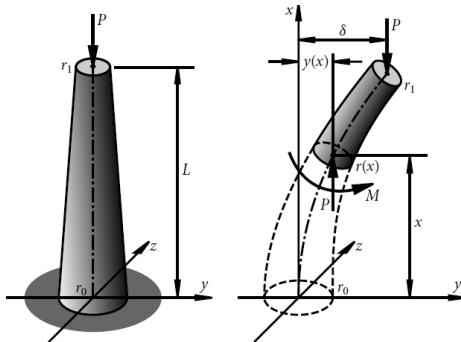
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- ▶ The bending moment at x is

$$M(x) = P(\delta - y(x)),$$

where δ is the deflection at the free end of the column.

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where $I(x)$ is the **moment of inertia** of the circular cross-section at x given by

$$I(x) = \frac{\pi}{4}r^4(x) = \frac{\pi}{4} \left(r_0 \left(1 - \frac{r_0 - r_1}{r_0} \cdot \frac{x}{L} \right) \right)^4 = I_0(1 - k_1\bar{x})^4$$

with

$$I_0 = \frac{\pi r_0^4}{4}, \quad k_1 = \frac{r_0 - r_1}{r_0}, \quad \bar{x} = \frac{x}{L}$$

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- ▶ We obtain a second-order differential equation

$$El_0(1 - k_1\bar{x})^4 \cdot \frac{1}{L} \frac{d^2\eta}{d\bar{x}^2} + P[L\eta(\bar{x})] = P(L\bar{\delta})$$

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- ▶ The general solution is

$$\eta(\bar{x}) = \eta_C(\bar{x}) + \eta_P(\bar{x}),$$

where $\eta_P(\bar{x})$ is a particular solution and $\eta_C(\bar{x})$ is the complementary solution, which is the solution of the homogeneous equation

$$(1 - k_1\bar{x})^4 \frac{d^2\eta}{d\bar{x}^2} + k^2\eta(\bar{x}) = 0.$$

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with $\alpha = 1/2$, $\beta = K$, $\rho = -1$, $\nu = 1/2$.

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- ▶ Transform the equation into the [Bessel equation](#) (see the hint in the end).

Solutions of Bessel's Equation

Bessel's equation

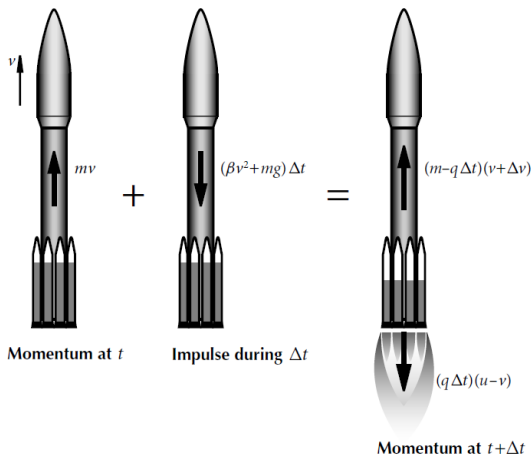
The equation

$$x^2 y'' + xy' + (x^2 - \nu^2)y = 0, x > 0,$$

$\nu = \text{const} \geq 0$, is called Bessel's equation

Motivating Examples: Ascending Motion of a Rocket

Consider the **ascending motion** of a rocket of initial mass m_0 (including shell and fuel). The fuel is consumed at a constant rate $q = -dm/dt$ and is expelled at a constant speed u relative to the rocket. At time t , the mass of the rocket is $m(t) = m_0 - qt$. If the velocity of the rocket is $v = v_0$ at $t = t_0$, determine the velocity $v(t)$.



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- ▶ Note that the drag force F_d depends not on the velocity but on the velocity squared.
 - ▶ If the fluid properties are considered constant, the drag force can be written as

$$F_d = \beta v^2,$$

where β is the damping coefficient.

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a first-order nonlinear differential equation with variable coefficients

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to transform the equation of the motion into the [Bessel equation](#)

$$\tau^2 \frac{d^2 V(\tau)}{d\tau^2} + \tau \frac{dV(\tau)}{d\tau} + (\tau^2 - \nu^2) V(\tau) = 0, \quad \nu = 2 \sqrt{\frac{\beta u}{q}}$$

Series Solution about an Ordinary Point

So, we need to know how to solve Bessel's equation!

- ▶ We introduce a concept of a **series solution** for a general n th-order equation with variable coefficients.
- ▶ We shall distinguish between the cases of ordinary, singular, regular singular and irregular points in series solutions.

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We shall consider two motivating examples to demonstrate that power series can be used to solve ODEs.

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where the constants a_n to be determined.

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$$y'(x) = \sum_{n=1}^{\infty} a_n n x^{n-1} = \sum_{m=0}^{\infty} a_{m+1} (m+1) x^m$$

change the index of summation

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- For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero:

$$x^0 : \quad a_1 - a_0 = 0 \Rightarrow$$

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$$\begin{aligned} a_1 &= a_0, \\ a_2 &= \frac{1}{2}a_1 = \frac{1}{2!}a_0, \\ a_3 &= \frac{1}{3}a_2 = \frac{1}{3} \cdot \frac{1}{2!}a_0, \end{aligned}$$

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Series Solution about an Ordinary Point

- The solution of the equation is

$$y(x) = a_0 + a_0x + \frac{1}{2!}a_0x^2 + \frac{1}{3!}a_0x^3 + \dots + \frac{1}{n!}a_0x^n + \dots =$$

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- Substitute the derivative into the equation to obtain

$$\sum_{n=0}^{\infty} a_{n+2}(n+2)(n+1)x^n + \sum_{n=0}^{\infty} a_n x^n =$$

Series Solution about an Ordinary Point

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$$x^5 : \quad 7 \cdot 6 a_7 + a_5 = 0 \Rightarrow \quad a_7 = -\frac{1}{7 \cdot 6} a_5 = -\frac{1}{7!} a_1,$$

Series Solution about an Ordinary Point

► In general,

$$a_{2k} = (-1)^k \frac{1}{(2k)!} a_0, \quad a_{2k+1} = (-1)^k \frac{1}{(2k+1)!} a_1$$

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- In general,

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- The solution is

$$\begin{aligned} y(x) &= \sum_{n=0}^{\infty} a_n x^n = a_0 \underbrace{\sum_{k=0}^{\infty} (-1)^k \frac{x^{2k}}{(2k)!}}_{\cos x} + a_1 \underbrace{\sum_{k=0}^{\infty} (-1)^k \frac{x^{2k+1}}{(2k+1)!}}_{\sin x} = \\ &= a_0 \cos x + a_1 \sin x, \quad a_0, a_1 \text{ are constants} \end{aligned}$$

Series Solution about an Ordinary Point

Definition

Consider the n th-order linear ordinary differential equation

$$y^{(n)}(x) + p_{n-1}(x)y^{(n-1)}(x) + \dots + p_0(x)y(x) = f(x).$$

A point x_0 is called an **ordinary point** of the given differential equation if each of the coefficients $p_0(x), p_1(x), \dots, p_{n-1}(x)$ and $f(x)$ are analytic at $x = x_0$, that is

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$$p_i(x) = \sum_{n=0}^{\infty} p_{i,n}(x - x_0)^n, \quad f(x) = \sum_{n=0}^{\infty} f_n(x - x_0)^n.$$

Theorem: Series Solution about an Ordinary Point

Theorem

If x_0 is an ordinary point of n th-order linear ordinary differential equation

$$y^{(n)}(x) + p_{n-1}(x)y^{(n-1)}(x) + \dots + p_0(x)y(x) = f(x),$$

then any solution of the equation can be expressed as a power series in $x - x_0$

$$y(x) = \sum_{n=0}^{\infty} a_n(x - x_0)^n, \quad |x - x_0| < R$$

and this representation is unique. Here $R \geq r$ is the radius of convergence.

Example: Legendre Equation

Find the power series solution in x of the **Legendre equation**

$$(1 - x^2)y'' - 2xy' + p(p + 1)y = 0, \quad p > 0.$$

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Legendre Equation

Therefore, $x = 0$ is an ordinary point and unique power series solution exists

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To determine the coefficients a_n , $n = 0, 1, \dots$, substitute the series solution into the equation to obtain

$$(1-x^2) \sum_{n=2}^{\infty} n(n-1)a_n x^{n-2} - 2x \sum_{n=1}^{\infty} n a_n x^{n-1} + p(p+1) \sum_{n=0}^{\infty} a_n x^n = 0.$$

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Changing the index of summation in the one part of the first term

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Changing the index of summation in the one part of the first term

$$\sum_{n=2}^{\infty} n(n-1)a_n x^{n-2} = \sum_{m=0}^{\infty} (m+2)(m+1)a_{m+2} x^m,$$

we obtain

Legendre Equation

$$\sum_{n=0}^{\infty} (n+2)(n+1)a_{n+2}x^n - \sum_{n=2}^{\infty} n(n-1)a_nx^n -$$
$$-2 \sum_{n=1}^{\infty} na_nx^n + p(p+1) \sum_{n=0}^{\infty} a_nx^n = 0.$$

Legendre Equation

$$\sum_{n=0}^{\infty} (n+2)(n+1)a_{n+2}x^n - \sum_{n=2}^{\infty} n(n-1)a_nx^n - \\ -2 \sum_{n=1}^{\infty} na_nx^n + p(p+1) \sum_{n=0}^{\infty} a_nx^n = 0.$$

As in the examples above, this equation is true if the coefficients of x^n , $n = 0, 1, \dots$, equal zero

x^0 :

Legendre Equation

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Let us calculate two more of them

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In general,

$$a_{2k} = (-1)^k \frac{p(p+1)(p-2)(p+3)\dots(p-2k+2)(p+2k-1)}{(2k)!} a_0$$

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Therefore, the power series solution of Legendre equation is

$$\begin{aligned}y(x) &= a_0 \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k)!} \prod_{i=1}^k ((p-2i+2)(p+2i-1)) x^{2k} \\&+ a_1 \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k+1)!} \prod_{i=1}^k ((p-2i+1)(p+2i)) x^{2k+1}, \quad |x| < 1.\end{aligned}$$

Series Solution about an Ordinary Point

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$x = 0$ is an ordinary point!

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Series Solution about an Ordinary Point

$$n = 3 : a_5 = \frac{1}{5 \cdot 4} \left(a_3 + \frac{a_2}{2} + \frac{a_1}{3} + \frac{a_0}{4} \right) =$$

Series Solution about an Ordinary Point

$$n = 3 : a_5 = \frac{1}{5 \cdot 4} \left(a_3 + \frac{a_2}{2} + \frac{a_1}{3} + \frac{a_0}{4} \right) = \frac{7a_0}{240} + \frac{a_1}{40}$$
$$n = 4 : a_6 = \frac{1}{6 \cdot 5} \left(a_4 + \frac{a_3}{2} + \frac{a_2}{3} + \frac{a_1}{4} + \frac{a_0}{5} \right) =$$

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$$\begin{aligned}n = 3 : \quad a_5 &= \frac{1}{5 \cdot 4} \left(a_3 + \frac{a_2}{2} + \frac{a_1}{3} + \frac{a_0}{4} \right) = \frac{7a_0}{240} + \frac{a_1}{40} \\n = 4 : \quad a_6 &= \frac{1}{6 \cdot 5} \left(a_4 + \frac{a_3}{2} + \frac{a_2}{3} + \frac{a_1}{4} + \frac{a_0}{5} \right) = \frac{43a_0}{2700} + \frac{a_1}{80}\end{aligned}$$

Series Solution about an Ordinary Point

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► The solution of the equation is

$$\begin{aligned}y(x) &= \sum_{n=0}^{\infty} a_n x^n = a_0 + a_1 x + \frac{a_0}{2} x^2 + \left(\frac{a_0}{12} + \frac{a_1}{6} \right) x^3 + \\&+ \left(\frac{5a_0}{72} + \frac{a_1}{24} \right) x^4 + \left(\frac{7a_0}{240} + \frac{a_1}{40} \right) x^5 + \left(\frac{43a_0}{2700} + \frac{a_1}{80} \right) x^6 + \dots\end{aligned}$$

Series Solution about a Regular Singular Point

Definition

Consider the n th-order linear ordinary differential equation

$$y^{(n)}(x) + p_{n-1}(x)y^{(n-1)}(x) + \dots + p_0(x)y(x) = f(x).$$

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A point x_0 is called a **singular point** of the given differential equation if it is not an ordinary point, that is, not all of the coefficients $p_0(x), p_1(x), \dots, p_{n-1}(x)$ are analytic at $x = x_0$.

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A point x_0 is called a **regular singular point** of the given differential equation if it is not an ordinary point, BUT all of $(x - x_0)^{n-k}p_k(x)$ are analytic for $k = 0, 1, \dots, n - 1$.

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A point x_0 is **irregular point** of the given differential equation if it is neither an ordinary point nor a regular singular point.

Fuchs' Th.: Series Solution about a Regular Singular Point

Fuch's Theorem

For the second-order linear homogeneous ordinary differential equation

$$y''(x) + P(x)y'(x) + Q(x)y(x) = 0,$$

if $x = 0$ is a regular singular point then

$$xP(x) = \sum_{n=0}^{\infty} P_n x^n, \quad x^2Q(x) = \sum_{n=0}^{\infty} Q_n x^n, \quad |x| < r.$$

Let the **indicial equation**

$$\alpha(\alpha - 1) + \alpha P_0 + Q_0 = 0$$

has two real roots $\alpha_1 \geq \alpha_2$. Then the DE has at least one Frobenius series solution given by

Fuchs' Th.: Series Solution about a Regular Singular Point

Fuch's Theorem

$$y_1(x) = x^{\alpha_1} \sum_{n=0}^{\infty} a_n x^n, \quad a_0 \neq 0, \quad 0 < x < r,$$

where the coefficients a_n can be determined by substituting $y_1(x)$ into the differential equation. A second linearly independent solution is obtained as follows:

1. If $\alpha_1 - \alpha_2$ is not equal to an integer, then a second **Frobenius series solution** is given by

$$y_2(x) = x^{\alpha_2} \sum_{n=0}^{\infty} b_n x^n, \quad 0 < x < r,$$

in which the coefficients b_n can be determined by substituting $y_2(x)$ into the differential equation.

Fuchs' Th.: Series Solution about a Regular Singular Point

Fuch's Theorem

2. If $\alpha_1 = \alpha_2 = \alpha$, then

$$y_2(x) = y_1(x) \ln x + x^\alpha \sum_{n=0}^{\infty} b_n x^n, \quad 0 < x < r,$$

where b_n can be determined by substituting $y_2(x)$ into the differential equation, once $y_1(x)$ is known. In this case, the second solution $y_2(x)$ is not a Frobenius series solution.

3. If $\alpha_1 - \alpha_2$ is a positive integer, then

$$y_2(x) = a y_1(x) \ln x + x^{\alpha_2} \sum_{n=0}^{\infty} b_n x^n, \quad 0 < x < r,$$

where b_n and a can be determined by substituting y_2 into the differential equation. The parameter a may be zero, in which case the second solution $y_2(x)$ is also a Frobenius series solution.

Solutions of Bessel's Equation

Bessel's equation

The equation

$$x^2 y'' + xy' + (x^2 - \nu^2)y = 0, x > 0,$$

$\nu = \text{const} \geq 0$, is called **Bessel's equation**

Solutions of Bessel's Equation

Bessel's equation is of the form

$$y'' + P(x)y' + Q(x)y = 0, \quad P(x) = \frac{1}{x}, \quad Q(x) = \frac{x^2 - \nu^2}{x^2}.$$

$x = 0$ is not an ordinary point (Why?)

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Both $xP(x)$, $x^2Q(x)$ are analytic at $x = 0$ and can be expanded as power series convergent for $|x| < \infty$.

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$$\alpha(\alpha - 1) + \alpha \cdot 1 - \nu^2 = 0 \Rightarrow$$

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Change the index of summation in the first part of the last term

$$\sum_{n=0}^{\infty} a_n x^{n+\nu+2} \Rightarrow \sum_{m=2}^{\infty} a_{m-2} x^{m+\nu} = \sum_{n=2}^{\infty} a_{n-2} x^{n+\nu}$$

Solutions of Bessel's Equation

$$x^\nu \left(\sum_{n=0}^{\infty} [(n+\nu)(n+\nu-1) + (n+\nu) - \nu^2] a_n x^n + \sum_{n=2}^{\infty} a_{n-2} x^n \right) = 0 \quad x^\nu \neq 0$$

$$\sum_{n=0}^{\infty} n(n+2\nu) a_n x^n + \sum_{n=2}^{\infty} a_{n-2} x^n = 0$$

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$, must be zero:

$$\begin{aligned} x^0 : \quad & 0 \cdot (0 + 2\nu) a_0 = 0 \Rightarrow & a_0 \neq 0 \text{ is arbitrary,} \\ x^1 : \quad & 1 \cdot (1 + 2\nu) a_1 = 0 \Rightarrow & a_1 = 0, \\ x^n : \quad & n \cdot (n + 2\nu) a_n + a_{n-2} = 0 \Rightarrow & a_n = -\frac{a_{n-2}}{n(n+2\nu)}, \quad n \geq 2 \end{aligned}$$

Solutions of Bessel's Equation

Therefore, $a_{2n+1} = 0$, $n = 0, 1, \dots$ and

$$a_2 = -\frac{a_0}{2(2+2\nu)} = -\frac{a_0}{2^2 \cdot 1(1+\nu)}$$

$$a_4 = -\frac{a_2}{4(4+2\nu)} = -\frac{a_2}{2^2 \cdot 2(1+\nu)} = (-1)^2 \frac{a_0}{2^4 \cdot 2!(1+\nu)(2+\nu)}$$

...

$$a_{2n} = (-1)^n \frac{a_0}{2^{2n} \cdot n!(1+\nu)(2+\nu) \dots (n+\nu)}$$

and

$$y_1(x) = a_0 x^\nu \sum_{n=0}^{\infty} (-1)^n \frac{1}{n!(1+\nu)(2+\nu) \dots (n+\nu)} \left(\frac{x}{2}\right)^{2n},$$

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Properties of the Gamma function

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Define the Gamma function by

$$\Gamma(\nu + 1) = \int_0^{\infty} t^{\nu} e^{-t} dt, \nu > 0.$$

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Therefore,

$$\begin{aligned}\Gamma(n + \nu + 1) &= (n + \nu) \Gamma(n + \nu) = (n + \nu)(n + \nu - 1) \Gamma(n + \nu - 1) \dots \\ &= (n + \nu)(n + \nu - 1) \dots (1 + \nu) \Gamma(1 + \nu)\end{aligned}$$

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The function $J_\nu(x)$ is called the **Bessel function of the first kind of order ν** .

What about the second linearly independent solution? Fuch's theorem!

Case 1. $\alpha_1 - \alpha_2 = 2\nu$ is not an integer

A second Frobenius solution is

$$y_2(x) = x^{-\nu} \sum_{n=0}^{\infty} b_n x^n, \quad 0 < x < \infty.$$

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We can follow the same procedure to show that

$$\begin{aligned} a_{2n-1} &= 0, \quad a_{2n} = (-1)^n \frac{b_0}{2^{2n} n! (1-\nu)(2-\nu) \dots (n-\nu)}, \\ y_2(x) &= b_0 x^{-\nu} \sum_{n=0}^{\infty} (-1)^n \frac{1}{n! (1-\nu)(2-\nu) \dots (n-\nu)} \left(\frac{x}{2}\right)^{2n}, \\ &\quad 0 < x < \infty. \end{aligned}$$

Letting $b_0 = [2^{-\nu} \Gamma(1-\nu)]^{-1}$,

$$y_2(x) = \sum_{n=0}^{\infty} (-1)^n \frac{1}{n! \Gamma(n-\nu+1)} \left(\frac{x}{2}\right)^{2n-\nu} = J_{-\nu}(x)$$

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Note, that the general solution can be also written in the form

$$y(x) = D_1 J_\nu(x) + D_2 Y_\nu(x),$$

where

$$Y_\nu(x) = \frac{J_\nu \cos \nu\pi - J_{-\nu}(x)}{\sin \nu\pi}$$

is the Bessel function of the second kind of order ν .

Case 2. $\alpha_1 = \alpha_2 \Rightarrow \nu = 0$

The first Frobenius series solution is simplified as

$$y_1(x) = J_0(x) = \sum_{n=0}^{\infty} (-1)^n \frac{1}{(n!)^2} \left(\frac{x}{2}\right)^{2n}, \quad 0 < x < \infty$$

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Substituting y_2, y_2', y_2'' into Bessel's equation results in

$$\begin{aligned} (x^2 y_1'' + x y_1' + x^2 y_1) \ln x + 2x y_1' + \sum_{n=2}^{\infty} n(n-1) b_n x^n + \\ + \sum_{n=1}^{\infty} n b_n x^n + \sum_{n=0}^{\infty} b_n x^{n+2} = 0. \end{aligned}$$

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Since $x^2 y_1'' + x y_1' + x^2 y_1 = 0$ and

Case 2. $\alpha_1 = \alpha_2 \Rightarrow \nu = 0$

$$2xy'_1 = 2x \sum_{n=1}^{\infty} (-1)^n \frac{1}{(n!)^2} \frac{2n \cdot x^{2n-1}}{2^{2n}} = \sum_{n=1}^{\infty} (-1)^n \frac{4n}{(n!)^2} \left(\frac{x}{2}\right)^{2n},$$

we obtain

$$\sum_{n=1}^{\infty} (-1)^n \frac{4n}{(n!)^2} \left(\frac{x}{2}\right)^{2n} + \sum_{n=2}^{\infty} n(n-1)b_n x^n + \sum_{n=1}^{\infty} n b_n x^n + \sum_{n=0}^{\infty} b_n x^{n+2} = 0$$

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$\begin{aligned} x^1 : \quad & 1 \cdot b_1 = 0 \Rightarrow b_1 = 0, \\ x^n, n \geq 1 : \quad & b_{2n+1} = 0 \text{ and} \end{aligned}$$

$$(-1)^n \frac{4n}{(n!)^2} \left(\frac{1}{2}\right)^{2n} + [2n(2n-1) + 2n]b_{2n} + b_{2n-2} = 0$$

and

$$b_{2n} = (-1)^{n+1} \frac{1}{n(n!)^2} \left(\frac{1}{2}\right)^{2n} - \frac{b_{2n-2}}{(2n)^2}$$

Case 2. $\alpha_1 = \alpha_2 \Rightarrow \nu = 0$

For simplicity, take $b_0 = 0$. Using mathematical induction, it can be shown that

$$b_{2n} = (-1)^{n+1} \frac{1 + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{n}}{(n!)^2} \left(\frac{1}{2}\right)^{2n}$$

Hence, a second linearly independent solution is

$$y_2(x) = J_0(x) \ln x + \sum_{n=1}^{\infty} (-1)^{n+1} \frac{1 + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{n}}{(n!)^2} \left(\frac{x}{2}\right)^{2n}$$

or, in terms of the [Bessel function of the second kind of order 0](#), $Y_0(x)$,

$$y_2(x) = \frac{\pi}{2} Y_0(x) + (\ln 2 - \gamma) J_0(x), \quad 0 < x, \infty$$

Case 2. $\alpha_1 = \alpha_2 \Rightarrow \nu = 0$

The Bessel function of the second kind of order 0 is defined as

$$Y_0(x) = \frac{2}{\pi} \left(\left(\ln \frac{x}{2} + \gamma \right) J_0(x) + \sum_{n=1}^{\infty} (-1)^{n+1} \frac{1 + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{n}}{(n!)^2} \left(\frac{x}{2} \right)^{2n} \right)$$

in which

$$\gamma = 0.57721566490153\dots = \lim_{n \rightarrow \infty} \left(\sum_{k=1}^n \frac{1}{k} - \ln n \right)$$

is the [Euler constant](#).

The general solution is

$$y(x) = C_1 J_0(x) + C_2 Y_0(x).$$

Case 3. ν is a positive integer

The first Frobenius series solution is simplified as

$$y_1(x) = J_\nu(x) = \sum_{n=0}^{\infty} (-1)^n \frac{1}{n!(n+\nu)!} \left(\frac{x}{2}\right)^{2n+\nu}, \quad 0 < x < \infty$$

Case 3. ν is a positive integer

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A second linearly independent solution is

$$y_2(x) = ay_1(x) \ln x + x^{-\nu} \sum_{n=0}^{\infty} b_n x^n, \quad 0 < x < \infty$$

Case 3. ν is a positive integer

The first Frobenius series solution is simplified as

$$y_1(x) = J_\nu(x) = \sum_{n=0}^{\infty} (-1)^n \frac{1}{n!(n+\nu)!} \left(\frac{x}{2}\right)^{2n+\nu}, \quad 0 < x < \infty$$

A second linearly independent solution is

$$y_2(x) = ay_1(x) \ln x + x^{-\nu} \sum_{n=0}^{\infty} b_n x^n, \quad 0 < x < \infty$$

$$y_2'(x) = a(y_1'(x) \ln x + \frac{y_1}{x}) + \sum_{n=0}^{\infty} (n - \nu) b_n x^{n-\nu-1},$$

Case 3. ν is a positive integer

The first Frobenius series solution is simplified as

$$y_1(x) = J_\nu(x) = \sum_{n=0}^{\infty} (-1)^n \frac{1}{n!(n+\nu)!} \left(\frac{x}{2}\right)^{2n+\nu}, \quad 0 < x < \infty$$

A second linearly independent solution is

$$y_2(x) = ay_1(x) \ln x + x^{-\nu} \sum_{n=0}^{\infty} b_n x^n, \quad 0 < x < \infty$$

$$y_2'(x) = a(y_1'(x) \ln x + \frac{y_1}{x}) + \sum_{n=0}^{\infty} (n-\nu) b_n x^{n-\nu-1},$$

$$y_2''(x) = a \left(y_1''(x) \ln x + \frac{2y_1'}{x} - \frac{y_1}{x^2} \right) + \sum_{n=0}^{\infty} (n-\nu)(n-\nu-1) b_n x^{n-\nu-2}.$$

Case 3. ν is a positive integer

Substituting into Bessel's equation results in

$$\begin{aligned} a[x^2 y_1'' + x y_1' + (x^2 - \nu^2) y_1] \ln x + 2a x y_1' + \sum_{n=0}^{\infty} (n - \nu)(n - \nu - 1) b_n x^{n - \nu} + \\ + \sum_{n=0}^{\infty} (n - \nu) b_n x^{n - \nu} + \sum_{n=0}^{\infty} b_n x^{n - \nu + 2} - \sum_{n=0}^{\infty} \nu^2 b_n x^{n - \nu} = 0 \end{aligned}$$

Case 3. ν is a positive integer

Substituting into Bessel's equation results in

$$\begin{aligned} a[x^2 y_1'' + x y_1' + (x^2 - \nu^2) y_1] \ln x + 2a x y_1' + \sum_{n=0}^{\infty} (n - \nu)(n - \nu - 1) b_n x^{n - \nu} + \\ + \sum_{n=0}^{\infty} (n - \nu) b_n x^{n - \nu} + \sum_{n=0}^{\infty} b_n x^{n - \nu + 2} - \sum_{n=0}^{\infty} \nu^2 b_n x^{n - \nu} = 0 \end{aligned}$$

y_1 is a solution of the Bessel's equation $\Rightarrow x^2 y_1'' + x y_1' + x^2 y_1 = 0$
and

Case 3. ν is a positive integer

Substituting into Bessel's equation results in

$$\begin{aligned} a[x^2 y_1'' + x y_1' + (x^2 - \nu^2) y_1] \ln x + 2axy_1' + \sum_{n=0}^{\infty} (n-\nu)(n-\nu-1)b_n x^{n-\nu} + \\ + \sum_{n=0}^{\infty} (n-\nu)b_n x^{n-\nu} + \sum_{n=0}^{\infty} b_n x^{n-\nu+2} - \sum_{n=0}^{\infty} \nu^2 b_n x^{n-\nu} = 0 \end{aligned}$$

y_1 is a solution of the Bessel's equation $\Rightarrow x^2 y_1'' + x y_1' + x^2 y_1 = 0$
and

$$2axy_1' =$$

Case 3. ν is a positive integer

Substituting into Bessel's equation results in

$$\begin{aligned} a[x^2 y_1'' + x y_1' + (x^2 - \nu^2) y_1] \ln x + 2ax y_1' + \sum_{n=0}^{\infty} (n-\nu)(n-\nu-1) b_n x^{n-\nu} + \\ + \sum_{n=0}^{\infty} (n-\nu) b_n x^{n-\nu} + \sum_{n=0}^{\infty} b_n x^{n-\nu+2} - \sum_{n=0}^{\infty} \nu^2 b_n x^{n-\nu} = 0 \end{aligned}$$

y_1 is a solution of the Bessel's equation $\Rightarrow x^2 y_1'' + x y_1' + x^2 y_1 = 0$
and

$$2ax y_1' = 2ax \sum_{n=1}^{\infty} (-1)^n \frac{(2n+\nu) \cdot x^{2n+\nu-1}}{n!(n+\nu)! 2^{2n+\nu}} =$$

Case 3. ν is a positive integer

Substituting into Bessel's equation results in

$$\begin{aligned} a[x^2 y_1'' + x y_1' + (x^2 - \nu^2) y_1] \ln x + 2a x y_1' + \sum_{n=0}^{\infty} (n - \nu)(n - \nu - 1) b_n x^{n - \nu} + \\ + \sum_{n=0}^{\infty} (n - \nu) b_n x^{n - \nu} + \sum_{n=0}^{\infty} b_n x^{n - \nu + 2} - \sum_{n=0}^{\infty} \nu^2 b_n x^{n - \nu} = 0 \end{aligned}$$

y_1 is a solution of the Bessel's equation $\Rightarrow x^2 y_1'' + x y_1' + x^2 y_1 = 0$
and

$$2a x y_1' = 2a x \sum_{n=1}^{\infty} (-1)^n \frac{(2n + \nu) \cdot x^{2n + \nu - 1}}{n!(n + \nu)! 2^{2n + \nu}} = \sum_{n=1}^{\infty} (-1)^n \frac{2a(2n + \nu)}{n!(n + \nu)!} \left(\frac{x}{2}\right)^{2n}$$

Case 3. ν is a positive integer

Substituting into Bessel's equation results in

$$\begin{aligned} a[x^2 y_1'' + xy_1' + (x^2 - \nu^2)y_1] \ln x + 2axy_1' + \sum_{n=0}^{\infty} (n-\nu)(n-\nu-1)b_n x^{n-\nu} + \\ + \sum_{n=0}^{\infty} (n-\nu)b_n x^{n-\nu} + \sum_{n=0}^{\infty} b_n x^{n-\nu+2} - \sum_{n=0}^{\infty} \nu^2 b_n x^{n-\nu} = 0 \end{aligned}$$

y_1 is a solution of the Bessel's equation $\Rightarrow x^2 y_1'' + xy_1' + x^2 y_1 = 0$
and

$$2axy_1' = 2ax \sum_{n=1}^{\infty} (-1)^n \frac{(2n+\nu) \cdot x^{2n+\nu-1}}{n!(n+\nu)! 2^{2n+\nu}} = \sum_{n=1}^{\infty} (-1)^n \frac{2a(2n+\nu)}{n!(n+\nu)!} \left(\frac{x}{2}\right)^{2n}$$

and multiplying the equation by x^ν

$$\sum_{n=1}^{\infty} (-1)^n \frac{2^{\nu+1} a(2n+\nu)}{n!(n+\nu)!} \left(\frac{x}{2}\right)^{2(n+\nu)} + \sum_{n=0}^{\infty} n(n-2\nu)b_n x^n + \sum_{n=2}^{\infty} b_{n-2} x^n = 0$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 :$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 :$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow b_1 = 0,$$

$$x^n, 2 \leq n < 2\nu :$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow b_1 = 0,$$

$$x^n, 2 \leq n < 2\nu : \quad n(n - 2\nu)b_n + b_{n-2} = 0 \Rightarrow$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow b_1 = 0,$$

$$x^n, 2 \leq n < 2\nu : \quad n(n - 2\nu)b_n + b_{n-2} = 0 \Rightarrow b_n = \frac{b_{n-2}}{n(2\nu - n)}$$

$$n = 2 :$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow b_1 = 0,$$

$$x^n, 2 \leq n < 2\nu : \quad n(n - 2\nu)b_n + b_{n-2} = 0 \Rightarrow b_n = \frac{b_{n-2}}{n(2\nu - n)}$$

$$n = 2 : \quad b_2 = \frac{b_0}{2(2\nu - 2)} =$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow b_1 = 0,$$

$$x^n, 2 \leq n < 2\nu : \quad n(n - 2\nu)b_n + b_{n-2} = 0 \Rightarrow b_n = \frac{b_{n-2}}{n(2\nu - n)}$$

$$n = 2 : \quad b_2 = \frac{b_0}{2(2\nu - 2)} = \frac{1}{2^2 \cdot 1(\nu - 1)}$$

$$n = 4 :$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow b_1 = 0,$$

$$x^n, 2 \leq n < 2\nu : \quad n(n - 2\nu)b_n + b_{n-2} = 0 \Rightarrow b_n = \frac{b_{n-2}}{n(2\nu - n)}$$

$$n = 2 : \quad b_2 = \frac{b_0}{2(2\nu-2)} = \frac{1}{2^2 \cdot 1(\nu-1)}$$

$$n = 4 : \quad b_4 = \frac{b_2}{4(2\nu-4)} =$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow b_1 = 0,$$

$$x^n, 2 \leq n < 2\nu : \quad n(n - 2\nu)b_n + b_{n-2} = 0 \Rightarrow b_n = \frac{b_{n-2}}{n(2\nu - n)}$$

$$n = 2 : \quad b_2 = \frac{b_0}{2(2\nu-2)} = \frac{1}{2^2 \cdot 1(\nu-1)}$$

$$n = 4 : \quad b_4 = \frac{b_2}{4(2\nu-4)} = \frac{1}{2^4 \cdot 2!(\nu-1)(\nu-2)}$$

...

$$n = 2k :$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow b_1 = 0,$$

$$x^n, 2 \leq n < 2\nu : \quad n(n - 2\nu)b_n + b_{n-2} = 0 \Rightarrow b_n = \frac{b_{n-2}}{n(2\nu - n)}$$

$$n = 2 : \quad b_2 = \frac{b_0}{2(2\nu-2)} = \frac{1}{2^2 \cdot 1(\nu-1)}$$

$$n = 4 : \quad b_4 = \frac{b_2}{4(2\nu-4)} = \frac{1}{2^4 \cdot 2!(\nu-1)(\nu-2)}$$

\dots

$$n = 2k : \quad b_{2k} = \frac{b_{2k-2}}{2^{2k}(2\nu-2k)} =$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow b_1 = 0,$$

$$x^n, 2 \leq n < 2\nu : \quad n(n - 2\nu)b_n + b_{n-2} = 0 \Rightarrow b_n = \frac{b_{n-2}}{n(2\nu - n)}$$

$$n = 2 : \quad b_2 = \frac{b_0}{2(2\nu-2)} = \frac{1}{2^2 \cdot 1(\nu-1)}$$

$$n = 4 : \quad b_4 = \frac{b_2}{4(2\nu-4)} = \frac{1}{2^4 \cdot 2!(\nu-1)(\nu-2)}$$

...

$$n = 2k : \quad b_{2k} = \frac{b_{2k-2}}{2^{2k}(2\nu-2k)} = \frac{1}{2^{2k}k!(\nu-1)(\nu-2)\dots(\nu-k)} = \frac{(\nu-k-1)!}{2^{2k}k!(\nu-1)}$$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow b_1 = 0,$$

$$x^n, 2 \leq n < 2\nu : \quad n(n - 2\nu)b_n + b_{n-2} = 0 \Rightarrow b_n = \frac{b_{n-2}}{n(2\nu - n)}$$

$$n = 2 : \quad b_2 = \frac{b_0}{2(2\nu-2)} = \frac{1}{2^2 \cdot 1(\nu-1)}$$

$$n = 4 : \quad b_4 = \frac{b_2}{4(2\nu-4)} = \frac{1}{2^4 \cdot 2!(\nu-1)(\nu-2)}$$

...

$$n = 2k : \quad b_{2k} = \frac{b_{2k-2}}{2^{2k}(2\nu-2k)} = \frac{1}{2^{2k}k!(\nu-1)(\nu-2)\dots(\nu-k)} = \frac{(\nu-k-1)!}{2^{2k}k!(\nu-1)}$$

and $b_{2k+1} = 0$, $n = 0, 1, 2, \dots$

Case 3. ν is a positive integer

For this equation to be true, the coefficient of x_n , $n = 0, 1, \dots$ must be zero.

$$x^0 : \quad 0 \cdot (0 - 2\nu)b_0 = 0 \Rightarrow b_0 = \text{is arbitrary, take } b_0 = 1,$$

$$x^1 : \quad 1 \cdot (1 - 2\nu)b_1 = 0 \Rightarrow b_1 = 0,$$

$$x^n, 2 \leq n < 2\nu : \quad n(n - 2\nu)b_n + b_{n-2} = 0 \Rightarrow b_n = \frac{b_{n-2}}{n(2\nu - n)}$$

$$n = 2 : \quad b_2 = \frac{b_0}{2(2\nu-2)} = \frac{1}{2^2 \cdot 1(\nu-1)}$$

$$n = 4 : \quad b_4 = \frac{b_2}{4(2\nu-4)} = \frac{1}{2^4 \cdot 2!(\nu-1)(\nu-2)}$$

...

$$n = 2k : \quad b_{2k} = \frac{b_{2k-2}}{2^{2k}(2\nu-2k)} = \frac{1}{2^{2k}k!(\nu-1)(\nu-2)\dots(\nu-k)} = \frac{(\nu-k-1)!}{2^{2k}k!(\nu-1)!}$$

and $b_{2k+1} = 0$, $n = 0, 1, 2, \dots$

From the coefficient of $x^{2\nu}$,

$$\frac{2^{\nu+1}a\nu}{\nu!} \left(\frac{x}{2}\right)^{2\nu} + b_{2\nu-2} = 0 \Rightarrow a = -\frac{1}{2^{\nu-1}(\nu-1)!}$$

The value of $b_{2\nu}$ is arbitrary, for simplicity, take $b_{2\nu} = 0$.

Case 3. ν is a positive integer

From the coefficient of $x^{2(\nu+n)}$, $n \geq 1$,

$$(-1)^{n+1} \frac{2^{\nu+1} a(2n+\nu)}{n!(n+\nu)!} \left(\frac{1}{2}\right)^{2(n+\nu)} + (2n+2\nu)(2n)b_{2\nu+2n} + b_{2(n-1+\nu)} = 0$$

Using mathematical induction, it can be shown that

$$b_{2(n+\nu)} = (-1)^{n+1} \frac{2^{\nu-1} a A_n}{n!(n+\nu)!} \left(\frac{1}{2}\right)^{2(n+\nu)},$$

where

$$A_n = \left(\frac{1}{1} + \frac{1}{2} + \dots + \frac{1}{n}\right) + \left(\frac{1}{1+\nu} + \frac{1}{2+\nu} + \dots + \frac{1}{n+\nu}\right)$$

Case 3. ν is a positive integer

Hence, a second linearly independent solution is

$$y_2(x) = aJ_\nu(x) \ln x + x^{-\nu} \left\{ \sum_{n=0}^{\nu-1} \frac{(\nu - n - 1)!}{n!(\nu - 1)!} \left(\frac{x}{2}\right)^{2n} + \right. \\ \left. + a \sum_{n=1}^{\infty} (-1)^{n+1} \frac{2^{\nu-1} A_n}{n!(n + \nu)!} \left(\frac{x}{2}\right)^{2(n+\nu)} \right\}$$

Case 3. ν is a positive integer

Hence, a second linearly independent solution is

$$y_2(x) = aJ_\nu(x) \ln x + x^{-\nu} \left\{ \sum_{n=0}^{\nu-1} \frac{(\nu - n - 1)!}{n!(\nu - 1)!} \left(\frac{x}{2}\right)^{2n} + \right. \\ \left. + a \sum_{n=1}^{\infty} (-1)^{n+1} \frac{2^{\nu-1} A_n}{n!(n + \nu)!} \left(\frac{x}{2}\right)^{2(n+\nu)} \right\}$$

Using the notation

$$1 + \frac{1}{2} + \dots + \frac{1}{n} = \psi(n+1) + \gamma, \quad \psi(1) = -\gamma,$$

where $\psi(n) = \Gamma'(n)/\Gamma(n)$ is the [psi function](#), $y_2(x)$ can be expressed in terms of the Bessel function of the second kind of order ν , $Y_\nu(x)$

Case 3. ν is a positive integer

$$y_2(x) = a \left\{ \frac{\pi}{2} Y_\nu(x) - \frac{1}{2} [\gamma - \psi(\nu + 1) - 2 \ln 2] J_\nu(x) \right\},$$

where, for $0 < x < \infty$

$$Y_\nu(x) = \frac{2}{\pi} J_\nu \ln \frac{x}{2} - \frac{1}{\pi} \sum_{n=0}^{\nu-1} \frac{(\nu - n - 1)!}{n!} \left(\frac{x}{2}\right)^{2n-\nu} - \\ - \frac{1}{\pi} \sum_{n=0}^{\infty} (-1)^n \frac{\psi(n+1) + \psi(n+\nu+1)}{n!(n+\nu)!} \left(\frac{x}{2}\right)^{2n+\nu}$$

The general solution is

$$y(x) = C_1 J_\nu(x) + C_2 Y_\nu(x).$$

This is the case that, when $\alpha_1 - \alpha_2$ is a positive integer, the second solution CONTAINS the logarithmic term $\ln x$.

Case 4. $\nu = k + 1/2$, $k = 0, 1, \dots$, and $\alpha_1 - \alpha_2 = 2k + 1$ is a positive integer

The first Frobenius series solution becomes

$$y_1(x) = J_{k+1/2}(x) = \sum_{n=0}^{\infty} (-1)^n \frac{1}{n! \Gamma(n + k + \frac{3}{2})} \left(\frac{x}{2}\right)^{2n+k-\frac{1}{2}}$$

and a second linearly independent solution is

$$y_2(x) = ay_1(x) \ln x + x^{-(k+\frac{1}{2})} \sum_{n=0}^{\infty} b_n x^n, \quad 0 < x < \infty$$

Show that

$$y_2(x) = \sum_{n=0}^{\infty} \frac{(-1)^n}{n! \Gamma(n - \nu + 1)} \left(\frac{x}{2}\right)^{2n-\nu} = J_{-\nu}(x),$$

that is the same as in the Case 1 when 2ν is not an integer.

This is the case that, when $\alpha_1 - \alpha_2$ is a positive integer, the second solution DOES NOT CONTAIN the logarithmic term $\ln x$.

Remark

As we have already seen, Bessel's equation may not appear in the standard form in practice. **How to transform the equation**

$$\frac{d^2y}{dx^2} + \frac{1-2\alpha}{x} \frac{dy}{dx} + \left[(\beta\rho x^{\rho-1})^2 + \frac{\alpha^2 - \nu^2 \rho^2}{x^2} \right] y = 0, x > 0$$

to the Bessel equation

$$\xi^2 \frac{d^2\eta}{d\xi^2} + \xi \frac{d\eta}{d\xi} + (\xi^2 - \nu^2)\eta = 0, \xi > 0?$$

Remark

As we have already seen, Bessel's equation may not appear in the standard form in practice. **How to transform the equation**

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And the differential equation becomes

$$x^{2-\alpha} \frac{d^2 y}{dx^2} + (1 - 2\alpha) x^{1-\alpha} \frac{dy}{dx} + \alpha^2 x^{-\alpha} y + (\beta^2 x^{2\rho} - \nu^2) \rho^2 x^{-\alpha} y = 0$$

and it is the considered equation.

Summary: Solutions of Bessel's Equation

Denote the solution of Bessel's equation

$$x^2 y'' + xy' + (x^2 - \nu^2)y = 0, \quad x > 0, \quad \nu = \text{const} \geq 0$$

as $y(x) = \mathcal{B}_\nu(x)$ where

$$\mathcal{B}_\nu(x) = C_1(x)J_\nu(x) + C_2Y_\nu(x),$$

$J_\nu(x)$, $Y_\nu(x)$ are the Bessel functions of the first and second kinds of order ν .

If $\nu \neq 0, 1, 2, \dots$, the solution can be written as

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is given by

$$y(x) = x^\alpha \mathcal{B}_\nu(\beta x^\rho)$$

Properties of the Bessel function of the first kind



$$J_{\nu-1}(x) + J_{\nu+1}(x) = \frac{2\nu}{x} J_{\nu}(x),$$



$$\begin{aligned} J'_{\nu}(x) &= J_{\nu-1}(x) - \frac{\nu}{x} J_{\nu}(x) = -J_{\nu+1}(x) + \frac{\nu}{x} J_{\nu}(x) = \\ &= \frac{1}{2}[J_{\nu-1}(x) - J_{\nu+1}(x)], \end{aligned}$$



$$\left(\frac{d}{x dx}\right)^m [x^{\nu} J_{\nu}(x)] = x^{\nu-m} J_{\nu-m}(x),$$



$$\left(\frac{d}{x dx}\right)^m [x^{-\nu} J_{\nu}(x)] = (-1)^m x^{-\nu-m} J_{\nu+m}(x),$$

Example: Buckling of a Tapered Column

Recall, that buckling of a tapered column is described by the equation with the general solution

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$$\alpha = 1/2, \beta = K, \rho = -1, \nu = 1/2.$$

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And the deflection of the column is

$$\eta(\xi) = \xi^{1/2} [C_1 J_{\frac{1}{2}}(K\xi^{-1}) + C_2 J_{-\frac{1}{2}}(K\xi^{-1})] + \bar{\delta}.$$

Example: Buckling of a Tapered Column

Determine $C_1, C_2, \bar{\delta}$ using boundary conditions

$$\text{at } x = 0 \text{ or } \xi = 1 : \quad \eta(\xi) = 0, \eta'(\xi) = 0,$$

$$\text{at } x = L \text{ or } \xi = 1 - k_1 : \quad \eta(\xi) = \bar{\delta}.$$

Note that

$$J'_{\frac{1}{2}}(x) = J_{\frac{1}{2}-1}(x) - \frac{\frac{1}{2}}{x} J_{\frac{1}{2}}(x) = J_{-\frac{1}{2}}(x) - \frac{1}{2x} J_{\frac{1}{2}}(x),$$

$$J'_{-\frac{1}{2}}(x) = -J_{-\frac{1}{2}+1}(x) + \frac{-\frac{1}{2}}{x} J_{-\frac{1}{2}}(x) = -J_{\frac{1}{2}}(x) - \frac{1}{2x} J_{-\frac{1}{2}}(x),$$

Example: Buckling of a Tapered Column

Differentiate the solution $\eta(\xi)$:

$$\eta'(\xi) = \frac{1}{2\sqrt{\xi}} \left[C_1 J_{\frac{1}{2}} \left(\frac{K}{\xi} \right) + C_2 J_{-\frac{1}{2}} \left(\frac{K}{\xi} \right) \right] + \sqrt{\xi} \left\{ C_1 \left[J_{-\frac{1}{2}} \left(\frac{K}{\xi} \right) - \frac{\xi}{2K} J_{\frac{1}{2}} \left(\frac{K}{\xi} \right) \right] + C_2 \left[-J_{\frac{1}{2}} \left(\frac{K}{\xi} \right) - \frac{\xi}{2K} J_{-\frac{1}{2}} \left(\frac{K}{\xi} \right) \right] \right\}$$

Therefore,

$$\eta'(1) = C_1 \left[J_{\frac{1}{2}}(K) - K J_{-\frac{1}{2}}(K) \right] + C_2 \left[J_{-\frac{1}{2}}(K) + K J_{\frac{1}{2}}(K) \right] = 0$$

At $x = L$ or $\xi = 1 - k_1$

$$\eta(1-k_1) = \sqrt{1-k_1} \left[C_1 J_{\frac{1}{2}}(K_L) + C_2 J_{-\frac{1}{2}}(K_L) \right] + \bar{\delta} = \bar{\delta}, \quad K_L = \frac{K}{1-k_1}.$$

Thus, the second equation to determine C_1, C_2 is

$$C_1 J_{\frac{1}{2}}(K_L) + C_2 J_{-\frac{1}{2}}(K_L) = 0$$

Example: Buckling of a Tapered Column

The determinant of the linear homogeneous system for C_1 , C_2 should be zero to have nontrivial solutions

$$\begin{vmatrix} J_{\frac{1}{2}}(K) - KJ_{-\frac{1}{2}}(K) & J_{-\frac{1}{2}}(K) + KJ_{\frac{1}{2}}(K) \\ J_{\frac{1}{2}}(K_L) & J_{-\frac{1}{2}}(K_L) \end{vmatrix} = 0$$

The equation

$$J_{-\frac{1}{2}}(K_L)[J_{\frac{1}{2}}(K) - KJ_{-\frac{1}{2}}(K)] - J_{\frac{1}{2}}(K_L)[J_{-\frac{1}{2}}(K) + KJ_{\frac{1}{2}}(K)] = 0$$

is called the **buckling equation**.

For a given value of $k_1 = (r_0 - r_1)/r_0$, the roots of this algebraic equation K_n , $n = 1, 2, \dots$, can be determined, from which the n th buckling load can be found.

$$K = \frac{k}{k_1}, \quad k^2 = \frac{PL^2}{El_0} \Rightarrow P_n = (p_n\pi)^2 \frac{El_0}{L^2}, \quad p_n = \frac{k_1 K_n}{\pi}, \quad n = 1, 2, \dots$$

Motivating Examples: Ascending Motion of a Rocket

Recall that the equation of ascending motion of a rocket is Bessel's equation

$$\tau^2 \frac{d^2 V(\tau)}{d\tau^2} + \tau \frac{dV(\tau)}{d\tau} + (\tau^2 - \nu^2) V(\tau) = 0, \quad \nu = 2\sqrt{\frac{\beta u}{q}}.$$

The solution of the Bessel's equation is

$$V(\tau) = C_1 J_\nu(\tau) + C_2 Y_\nu(\tau),$$

$$\frac{dV(\tau)}{d\tau} = C_1 \left(\frac{\nu}{\tau} J_\nu(\tau) - J_{\nu+1}(\tau) \right) + C_2 \left(\frac{\nu}{\tau} Y_\nu(\tau) - Y_{\nu+1}(\tau) \right)$$

The velocity of the rocket is

$$\begin{aligned} \nu(\tau) &= \frac{m(\tau) \dot{V}(\tau)}{\beta V(\tau)} = \frac{m(\tau) \left(-\frac{2\beta g}{q\tau} \frac{dV(\tau)}{d\tau} \right)}{\beta V(\tau)} = \\ &= \frac{\frac{q^2 \tau^2}{4\beta g} \left(-\frac{2\beta g}{q\tau} \right) \left(\frac{\nu}{\tau} [C_1 J_\nu(\tau) + C_2 Y_\nu(\tau)] - [C_1 J_{\nu+1}(\tau) + C_2 Y_{\nu+1}(\tau)] \right)}{\beta (C_1 J_\nu(\tau) + C_2 Y_\nu(\tau))} \end{aligned}$$

Exercises

A. Show that the general solution of the **Airy equation**

$$y'' - xy = 0$$

is, for $|x| < \infty$, $y(x) =$

$$= a_0 \left(1 + \sum_{n=1}^{\infty} \frac{\prod_{k=1}^n (3k-2)}{(3n)!} x^{3n} \right) + a_1 \left(x + \sum_{n=1}^{\infty} \frac{\prod_{k=1}^n (3k-1)}{(3n+1)!} x^{3n+1} \right)$$

B. Determine the general solution of the following differential equations in terms of power series about $x = 0$.

$$1. y''' + xy = 0, \quad 2. (1 - x^2)y'' + y = 0, \quad 3. y'' - 2x^2y = 0$$

$$4. y'' - 2x^2y' + xy = 0, \quad 5. (x^2 - 1)y'' + (4x - 1)y' + 2y = 0$$

C. Determine two linearly independent solutions of the following equations using the series solution approach

$$1. x^2y'' + (x - 2x^2)y' - xy = 0 \quad 2. x^2y'' - (2x + x^2)y' + 2y = 0$$

$$3. x^2y'' + \left(\frac{1}{2}x + x^2\right)y' + xy = 0$$