# vv256: Laplace Transform

Dr.Olga Danilkina

UM-SJTU Joint Institute



November 29, 2019

The Laplace Transform is the method for solving IVPs for ODEs and PDEs by converting the differential equation into an algebraic equation.

The Laplace Transform is the method for solving IVPs for ODEs and PDEs by converting the differential equation into an algebraic equation.

A general integral transform of a function f(t) into another function  $\bar{f}(s)$  is defined by

$$\bar{f}(s) = \int_{0}^{\beta} K(s,t)f(t) dt.$$

The Laplace Transform is the method for solving IVPs for ODEs and PDEs by converting the differential equation into an algebraic equation.

A general integral transform of a function f(t) into another function  $\bar{f}(s)$  is defined by

$$ar{f}(s) = \int_{lpha}^{eta} K(s,t) f(t) dt.$$

The function K(s, t) is called the kernel of the transform.

The Laplace Transform is the method for solving IVPs for ODEs and PDEs by converting the differential equation into an algebraic equation.

A general integral transform of a function f(t) into another function  $\bar{f}(s)$  is defined by

$$\bar{f}(s) = \int_{0}^{\beta} K(s,t)f(t) dt.$$

The function K(s,t) is called the kernel of the transform.

The Laplace transform is a special case with

$$\alpha = 0, \, \beta = \infty, \, K(s, t) = e^{-st} \Rightarrow$$

The Laplace Transform is the method for solving IVPs for ODEs and PDEs by converting the differential equation into an algebraic equation.

A general integral transform of a function f(t) into another function  $\bar{f}(s)$  is defined by

$$ar{f}(s) = \int_{lpha}^{eta} K(s,t) f(t) dt.$$

The function K(s,t) is called the kernel of the transform.

The Laplace transform is a special case with  $\alpha=0,\ \beta=\infty,\ K(s,t)=e^{-st}\Rightarrow$  improper integral We assume that f(t) as a real-valued function defined for t>0.

The Laplace transform of f(t) is

$$ar{f}(s) = L[f(t)\colon t \to s] = \int_0^\infty e^{-st} f(t) dt.$$

The function K(s, t) is called the kernel of the transform.

The Laplace transform of f(t) is

$$ar{f}(s) = L[f(t)\colon t \to s] = \int_0^\infty \mathrm{e}^{-st} f(t) \, dt.$$

The function K(s,t) is called the kernel of the transform. Reverse the procedure and obtain the inverse Laplace transform: if  $L[f(t)\colon t\to s]$  then

$$L^{-1}\left[\bar{f}(s)\colon s\to t\right]=f(t).$$

The Laplace transform of f(t) is

$$ar{f}(s) = L[f(t)\colon t \to s] = \int_0^\infty \mathrm{e}^{-st} f(t) \, dt.$$

The function K(s,t) is called the kernel of the transform.

Reverse the procedure and obtain the inverse Laplace transform: if I[f(t): t] = cl then

$$L[f(t): t \rightarrow s]$$
 then

$$L^{-1}\left[\bar{f}(s)\colon s\to t\right]=f(t).$$

1. What functions have a Laplace transform?

The Laplace transform of f(t) is

$$ar{f}(s) = L[f(t)\colon t \to s] = \int_0^\infty \mathrm{e}^{-st} f(t) \, dt.$$

The function K(s, t) is called the kernel of the transform.

Reverse the procedure and obtain the inverse Laplace transform: if  $L[f(t): t \rightarrow s]$  then

$$L^{-1}\left[\bar{f}(s)\colon s\to t\right]=f(t).$$

- 1. What functions have a Laplace transform?
- 2. Can two functions f(t), g(t) have the same Laplace transform?

The Laplace transform of f(t) is

$$\bar{f}(s) = L[f(t): t \to s] = \int_0^\infty e^{-st} f(t) dt.$$

The function K(s, t) is called the kernel of the transform.

Reverse the procedure and obtain the inverse Laplace transform: if  $L[f(t): t \rightarrow s]$  then

$$L^{-1}\left[\overline{f}(s)\colon s\to t\right]=f(t).$$

- 1. What functions have a Laplace transform?
- 2. Can two functions f(t), g(t) have the same Laplace transform?

The existence and uniqueness of the Laplace transform of a function f(t) is guaranteed if there exist real K, M and a such that

- 1. f(t) is piecewise continuous for t > 0,
- 2.  $|f(t)| \leq Ke^{at}$  for  $t \geq M$ .

$$L[1: t \to s] = \int_0^\infty e^{-st} = \frac{e^{-st}}{-s} \bigg|_0^\infty = \frac{1}{s}$$

**Example 1.** Find the Laplace transform of the function f(t) = 1.

$$L[1: t \to s] = \int_0^\infty e^{-st} = \left. \frac{e^{-st}}{-s} \right|_0^\infty = \frac{1}{s}$$

**Remark:** The result is limited to s > 0 only.

**Example 1.** Find the Laplace transform of the function f(t) = 1.

$$L[1: t \to s] = \int_0^\infty e^{-st} = \left. \frac{e^{-st}}{-s} \right|_0^\infty = \frac{1}{s}$$

**Remark:** The result is limited to s > 0 only.

**Example 1.** Find the Laplace transform of the function f(t) = 1.

$$L[1: t \to s] = \int_0^\infty e^{-st} = \left. \frac{e^{-st}}{-s} \right|_0^\infty = \frac{1}{s}$$

**Remark:** The result is limited to s > 0 only.

$$L[t: t \to s] = \int_0^\infty t e^{-st} =$$

**Example 1.** Find the Laplace transform of the function f(t) = 1.

$$L[1: t \to s] = \int_0^\infty e^{-st} = \frac{e^{-st}}{-s} \Big|_0^\infty = \frac{1}{s}$$

**Remark:** The result is limited to s > 0 only.

$$L[t: t \to s] = \int_0^\infty t e^{-st} = \left. \frac{t e^{-st}}{-s} \right|_0^\infty - \int_0^\infty \frac{e^{-st}}{(-s)} dt =$$

**Example 1.** Find the Laplace transform of the function f(t) = 1.

$$L[1: t \to s] = \int_0^\infty e^{-st} = \left. \frac{e^{-st}}{-s} \right|_0^\infty = \frac{1}{s}$$

**Remark:** The result is limited to s > 0 only.

$$L[t:t\to s] = \int_0^\infty t e^{-st} = \left. \frac{te^{-st}}{-s} \right|_0^\infty - \int_0^\infty \frac{e^{-st}}{(-s)} dt = \frac{1}{s} \left. \frac{e^{-st}}{-s} \right|_0^\infty$$

**Example 1.** Find the Laplace transform of the function f(t) = 1.

$$L[1: t \to s] = \int_0^\infty e^{-st} = \left. \frac{e^{-st}}{-s} \right|_0^\infty = \frac{1}{s}$$

**Remark:** The result is limited to s > 0 only.

$$L[t: t \to s] = \int_0^\infty t e^{-st} = \left. \frac{te^{-st}}{-s} \right|_0^\infty - \int_0^\infty \frac{e^{-st}}{(-s)} dt = \frac{1}{s} \left. \frac{e^{-st}}{-s} \right|_0^\infty$$
$$= \frac{1}{s^2}, \quad s > 0$$

**Example 1.** Find the Laplace transform of the function f(t) = 1.

$$L[1: t \to s] = \int_0^\infty e^{-st} = \left. \frac{e^{-st}}{-s} \right|_0^\infty = \frac{1}{s}$$

**Remark:** The result is limited to s > 0 only.

**Example 2.** Find the Laplace transform of the function f(t) = t.

$$L[t: t \to s] = \int_0^\infty t e^{-st} = \left. \frac{t e^{-st}}{-s} \right|_0^\infty - \int_0^\infty \frac{e^{-st}}{(-s)} dt = \frac{1}{s} \left. \frac{e^{-st}}{-s} \right|_0^\infty$$
$$= \frac{1}{s^2}, \quad s > 0$$

**Example 1.** Find the Laplace transform of the function f(t) = 1.

$$L[1: t \to s] = \int_0^\infty e^{-st} = \left. \frac{e^{-st}}{-s} \right|_0^\infty = \frac{1}{s}$$

**Remark:** The result is limited to s > 0 only.

**Example 2.** Find the Laplace transform of the function f(t) = t.

$$L[t: t \to s] = \int_0^\infty t e^{-st} = \left. \frac{te^{-st}}{-s} \right|_0^\infty - \int_0^\infty \frac{e^{-st}}{(-s)} dt = \frac{1}{s} \left. \frac{e^{-st}}{-s} \right|_0^\infty$$
$$= \frac{1}{s^2}, \quad s > 0$$

$$L[e^{\alpha t}\colon t\to s]=\int_0^\infty e^{\alpha t}e^{-st}=$$

**Example 1.** Find the Laplace transform of the function f(t) = 1.

$$L[1: t \to s] = \int_0^\infty e^{-st} = \left. \frac{e^{-st}}{-s} \right|_0^\infty = \frac{1}{s}$$

**Remark:** The result is limited to s > 0 only.

**Example 2.** Find the Laplace transform of the function f(t) = t.

$$L[t: t \to s] = \int_0^\infty t e^{-st} = \left. \frac{te^{-st}}{-s} \right|_0^\infty - \int_0^\infty \frac{e^{-st}}{(-s)} dt = \frac{1}{s} \left. \frac{e^{-st}}{-s} \right|_0^\infty$$
$$= \frac{1}{s^2}, \quad s > 0$$

$$L[e^{\alpha t}\colon t\to s] = \int_0^\infty e^{\alpha t} e^{-st} = \left. \frac{e^{(\alpha-s)t}}{\alpha-s} \right|_0^\infty =$$

**Example 1.** Find the Laplace transform of the function f(t) = 1.

$$L[1: t \to s] = \int_0^\infty e^{-st} = \left. \frac{e^{-st}}{-s} \right|_0^\infty = \frac{1}{s}$$

**Remark:** The result is limited to s > 0 only.

**Example 2.** Find the Laplace transform of the function f(t) = t.

$$L[t: t \to s] = \int_0^\infty t e^{-st} = \left. \frac{te^{-st}}{-s} \right|_0^\infty - \int_0^\infty \frac{e^{-st}}{(-s)} dt = \frac{1}{s} \left. \frac{e^{-st}}{-s} \right|_0^\infty$$
$$= \frac{1}{s^2}, \quad s > 0$$

$$L[e^{\alpha t} \colon t \to s] = \int_0^\infty e^{\alpha t} e^{-st} = \left. \frac{e^{(\alpha - s)t}}{\alpha - s} \right|_0^\infty = \frac{1}{s - \alpha}, \quad s > \alpha$$

$$L[\cos at + i\sin at: t \rightarrow s] = L[e^{iat}: t \rightarrow s] =$$

$$L[\cos at + i\sin at: t \rightarrow s] = L[e^{iat}: t \rightarrow s] = \frac{1}{s - ia} = \frac{1}{s - ia}$$

$$L[\cos at + i\sin at: t \to s] = L[e^{iat}: t \to s] = \frac{1}{s - ia} =$$

$$= \frac{s + ia}{s^2 + a^2}, \quad s > 0$$

**Example 4.** Find the Laplace transform of the functions  $f(t) = \sin at$  and  $f(t) = \cos at$ .

$$L[\cos at + i \sin at: t \to s] = L[e^{iat}: t \to s] = \frac{1}{s - ia} =$$

$$= \frac{s + ia}{s^2 + a^2}, \quad s > 0$$

Separate real and imaginary parts:

**Example 4.** Find the Laplace transform of the functions  $f(t) = \sin at$  and  $f(t) = \cos at$ .

$$L[\cos at + i\sin at: t \to s] = L[e^{iat}: t \to s] = \frac{1}{s - ia} =$$

$$= \frac{s + ia}{s^2 + a^2}, \quad s > 0$$

Separate real and imaginary parts:

$$L[\cos at: t \to s] = \frac{s}{s^2 + a^2}, \quad L[\sin at: t \to s] = \frac{a}{s^2 + a^2}$$

# Some General Properties

Linearity.

If 
$$\bar{f}(s) = L[f(t): t \to s]$$
 and  $\bar{g}(s) = L[g(t): t \to s]$  exist then  $L[af(t) + bg(t): t \to s]$  exists for all constants  $a$  and  $b$  and

$$L[af(t) + bg(t): t \rightarrow s] = a\bar{f}(s) + b\bar{g}(s)$$

# Some General Properties

Linearity.

If  $\bar{f}(s) = L[f(t): t \to s]$  and  $\bar{g}(s) = L[g(t): t \to s]$  exist then  $L[af(t) + bg(t): t \to s]$  exists for all constants a and b and

$$L[af(t) + bg(t): t \rightarrow s] = a\overline{f}(s) + b\overline{g}(s)$$

First shifting property.

Suppose  $\bar{f}(s) = L[f(t): t \to s]$  exists and that a is a constant then  $L[e^{at}f(t): t \to s]$  exists and

$$L\left[e^{at}f(t)\colon t\to s\right]=\bar{f}(s-a)$$

or

$$L^{-1}\left[\bar{f}(s-a)\colon s\to t\right]=e^{at}f(t)$$

# Laplace Transforms of Selected Functions

$$f(t) = L^{-1}\left[ar{f}(s)\colon s o t
ight] \quad ar{f}(s) = L\left[f(t)\colon t o s
ight]$$
 $1 \qquad \qquad rac{1}{s} \qquad Re(s) > 0$ 
 $e^{at} \qquad \qquad rac{1}{s-a} \qquad Re(s) > 0$ 
 $t^n, \, n\in\mathbb{Z}_+ \qquad \qquad rac{n!}{s^{n+1}} \qquad Re(s) > 0$ 
 $t^p, \, p > -1 \qquad \qquad rac{\Gamma(p+1)}{s^{p+1}} \qquad Re(s) > 0$ 
 $\cos at \qquad \qquad rac{s}{s^2+a^2} \qquad Re(s) > 0$ 
 $\sin at \qquad \qquad rac{a}{2a-a} \qquad Re(s) > 0$ 

# Laplace Transforms of Selected Functions

$$f(t) = L^{-1}\left[ar{f}(s)\colon s o t
ight] \quad ar{f}(s) = L\left[f(t)\colon t o s
ight]$$
 $\operatorname{cosh} at \qquad rac{s}{s^2-a^2} \qquad \operatorname{Re}(s) > |a|$ 
 $\operatorname{sinh} at \qquad rac{a}{s^2-a^2} \qquad \operatorname{Re}(s) > |a|$ 
 $\operatorname{e}^{at}\cos bt \qquad rac{s-a}{(s-a)^2+b^2} \qquad \operatorname{Re}(s) > a$ 
 $\operatorname{e}^{at}\sin bt \qquad rac{b}{(s-a)^2+b^2} \qquad \operatorname{Re}(s) > a$ 
 $t^n \operatorname{e}^{at}, \ n \in \mathbb{Z}_+ \qquad rac{n!}{(s-a)^{n+1}} \qquad \operatorname{Re}(s) > a$ 

#### Convolution

Let  $\bar{f}(s) = L[f(t): t \to s], \bar{g}(s) = L[g(t): t \to s]$ 

be Laplace transforms of f(t) and g(t).

#### Convolution

Let 
$$\bar{f}(s) = L[f(t): t \to s], \bar{g}(s) = L[g(t): t \to s]$$

be Laplace transforms of f(t) and g(t).

▶ The linearity of the Laplace transform operator guarantees:

$$L^{-1}\left[ar{f}(s)+ar{g}(s)\colon s o t
ight]=$$

#### Convolution

Let 
$$\bar{f}(s) = L[f(t): t \to s], \bar{g}(s) = L[g(t): t \to s]$$

be Laplace transforms of f(t) and g(t).

▶ The linearity of the Laplace transform operator guarantees:

$$L^{-1}\left[\bar{f}(s)+\bar{g}(s)\colon s\to t\right]=$$

$$=L^{-1}\left[ar{f}(s)\colon s o t
ight]+L^{-1}\left[ar{g}(s)\colon s o t
ight]=$$

Let 
$$\bar{f}(s) = L[f(t): t \to s], \bar{g}(s) = L[g(t): t \to s]$$

be Laplace transforms of f(t) and g(t).

▶ The linearity of the Laplace transform operator guarantees:

$$L^{-1}\left[\bar{f}(s)+\bar{g}(s)\colon s\to t\right]=$$

$$= L^{-1} \left[ \bar{f}(s) \colon s \to t \right] + L^{-1} \left[ \bar{g}(s) \colon s \to t \right] = f(t) + g(t)$$

Let 
$$\bar{f}(s) = L[f(t): t \to s], \bar{g}(s) = L[g(t): t \to s]$$

be Laplace transforms of f(t) and g(t).

▶ The linearity of the Laplace transform operator guarantees:

$$L^{-1}\left[\bar{f}(s)+\bar{g}(s)\colon s\to t\right]=$$

$$= L^{-1} \left[ \bar{f}(s) \colon s \to t \right] + L^{-1} \left[ \bar{g}(s) \colon s \to t \right] = f(t) + g(t)$$

But the inverse Laplace transform of a product does not equal the product of the inverse Laplace transforms of the factors:

$$L^{-1}\left[\bar{f}(s)\bar{g}(s)\colon s\to t\right]\neq f(t)g(t)$$

Let 
$$\bar{f}(s) = L[f(t): t \to s], \ \bar{g}(s) = L[g(t): t \to s]$$

be Laplace transforms of f(t) and g(t).

▶ The linearity of the Laplace transform operator guarantees:

$$L^{-1}\left[\bar{f}(s)+\bar{g}(s)\colon s\to t\right]=$$

$$= L^{-1} \left[ \bar{f}(s) \colon s \to t \right] + L^{-1} \left[ \bar{g}(s) \colon s \to t \right] = f(t) + g(t)$$

But the inverse Laplace transform of a product does not equal the product of the inverse Laplace transforms of the factors:

$$L^{-1}\left[\bar{f}(s)\bar{g}(s)\colon s\to t\right]\neq f(t)g(t)$$

For example,

$$ar{f}(s)=rac{1}{s},\,ar{g}(s)=rac{1}{s^2}\Rightarrow$$

Let 
$$\bar{f}(s) = L[f(t): t \to s], \bar{g}(s) = L[g(t): t \to s]$$

be Laplace transforms of f(t) and g(t).

▶ The linearity of the Laplace transform operator guarantees:

$$L^{-1}\left[\bar{f}(s)+\bar{g}(s)\colon s\to t\right]=$$

$$= L^{-1} \left[ \bar{f}(s) \colon s \to t \right] + L^{-1} \left[ \bar{g}(s) \colon s \to t \right] = f(t) + g(t)$$

But the inverse Laplace transform of a product does not equal the product of the inverse Laplace transforms of the factors:

$$L^{-1}\left[\bar{f}(s)\bar{g}(s)\colon s\to t\right]\neq f(t)g(t)$$

For example,

$$\bar{f}(s) = \frac{1}{s}, \ \bar{g}(s) = \frac{1}{s^2} \Rightarrow f(t) = 1, \ g(t) = t$$

Let 
$$\bar{f}(s) = L[f(t): t \to s], \bar{g}(s) = L[g(t): t \to s]$$

be Laplace transforms of f(t) and g(t).

▶ The linearity of the Laplace transform operator guarantees:

$$L^{-1}\left[\bar{f}(s)+\bar{g}(s)\colon s\to t\right]=$$

$$=L^{-1}\left[\bar{f}(s)\colon s\to t\right]+L^{-1}\left[\bar{g}(s)\colon s\to t\right]=f(t)+g(t)$$

But the inverse Laplace transform of a product does not equal the product of the inverse Laplace transforms of the factors:

$$L^{-1}\left[\bar{f}(s)\bar{g}(s)\colon s\to t\right]\neq f(t)g(t)$$

For example,

$$\bar{f}(s) = \frac{1}{s}, \ \bar{g}(s) = \frac{1}{s^2} \Rightarrow f(t) = 1, \ g(t) = t$$

Let 
$$\bar{f}(s) = L[f(t): t \to s], \ \bar{g}(s) = L[g(t): t \to s]$$

be Laplace transforms of f(t) and g(t).

▶ The linearity of the Laplace transform operator guarantees:

$$L^{-1}\left[\bar{f}(s)+\bar{g}(s)\colon s\to t\right]=$$

$$=L^{-1}\left[\bar{f}(s)\colon s\to t\right]+L^{-1}\left[\bar{g}(s)\colon s\to t\right]=f(t)+g(t)$$

But the inverse Laplace transform of a product does not equal the product of the inverse Laplace transforms of the factors:

$$L^{-1}\left[\bar{f}(s)\bar{g}(s)\colon s\to t\right]\neq f(t)g(t)$$

For example,

$$\bar{f}(s) = \frac{1}{s}, \ \bar{g}(s) = \frac{1}{s^2} \Rightarrow f(t) = 1, \ g(t) = t$$

$$L^{-1}\left[\bar{f}(s)\bar{g}(s)\colon s\to t\right]=$$

Let 
$$\bar{f}(s) = L[f(t): t \to s], \bar{g}(s) = L[g(t): t \to s]$$

be Laplace transforms of f(t) and g(t).

▶ The linearity of the Laplace transform operator guarantees:

$$L^{-1}\left[\bar{f}(s)+\bar{g}(s)\colon s\to t\right]=$$

$$=L^{-1}\left[\bar{f}(s)\colon s\to t\right]+L^{-1}\left[\bar{g}(s)\colon s\to t\right]=f(t)+g(t)$$

But the inverse Laplace transform of a product does not equal the product of the inverse Laplace transforms of the factors:

$$L^{-1}\left[\bar{f}(s)\bar{g}(s)\colon s\to t\right]\neq f(t)g(t)$$

For example,

$$\bar{f}(s) = \frac{1}{s}, \ \bar{g}(s) = \frac{1}{s^2} \Rightarrow f(t) = 1, \ g(t) = t$$

$$L^{-1}\left[\bar{f}(s)\bar{g}(s)\colon s\to t\right]=L^{-1}\left[1/s^3\colon s\to t\right]=$$

Let 
$$ar{f}(s) = L[f(t)\colon t o s]\,,\, ar{g}(s) = L[g(t)\colon t o s]$$

be Laplace transforms of f(t) and g(t).

▶ The linearity of the Laplace transform operator guarantees:

$$L^{-1}\left[\bar{f}(s)+\bar{g}(s)\colon s\to t\right]=$$

$$=L^{-1}\left[\bar{f}(s)\colon s\to t\right]+L^{-1}\left[\bar{g}(s)\colon s\to t\right]=f(t)+g(t)$$

But the inverse Laplace transform of a product does not equal the product of the inverse Laplace transforms of the factors:

$$L^{-1}\left[\bar{f}(s)\bar{g}(s)\colon s\to t\right]\neq f(t)g(t)$$

For example,

$$\bar{f}(s) = \frac{1}{s}, \ \bar{g}(s) = \frac{1}{s^2} \Rightarrow f(t) = 1, \ g(t) = t$$

$$L^{-1}\left[\bar{f}(s)\bar{g}(s)\colon s\to t\right] = L^{-1}\left[1/s^3\colon s\to t\right] = t^2/2\neq f(t)g(t) = t$$

Given two functions f(t) and g(t), the convolution  $f \star g$  of f and g is defined by the integral

$$(f \star g)(t) = \int_0^t f(t - u)g(u) du$$

Given two functions f(t) and g(t), the convolution  $f \star g$  of f and g is defined by the integral

$$(f \star g)(t) = \int_0^t f(t - u)g(u) du$$

whenever this integral exists.

•  $(f \star g)(t) = (g \star f)(t)$  verify by simple change of variable

Given two functions f(t) and g(t), the convolution  $f \star g$  of f and g is defined by the integral

$$(f \star g)(t) = \int_0^t f(t - u)g(u) du$$

- $(f \star g)(t) = (g \star f)(t)$  verify by simple change of variable
- $L[(f \star g)(t) \colon t \to s] = \bar{f}(s)\bar{g}(s)$

Given two functions f(t) and g(t), the convolution  $f \star g$  of f and g is defined by the integral

$$(f \star g)(t) = \int_0^t f(t - u)g(u) du$$

- $(f \star g)(t) = (g \star f)(t)$  verify by simple change of variable
- $L[(f \star g)(t) \colon t \to s] = \bar{f}(s)\bar{g}(s)$

$$L\left[(f\star g)(t)\colon t\to s\right]=$$

Given two functions f(t) and g(t), the convolution  $f \star g$  of f and g is defined by the integral

$$(f \star g)(t) = \int_0^t f(t - u)g(u) du$$

- $(f \star g)(t) = (g \star f)(t)$  verify by simple change of variable
- $L[(f \star g)(t) \colon t \to s] = \bar{f}(s)\bar{g}(s)$

$$L[(f\star g)(t)\colon t\to s]=\int_0^\infty\left(\int_0^t f(t-u)g(u)\,du\right)e^{-st}\,dt$$

Given two functions f(t) and g(t), the convolution  $f \star g$  of f and g is defined by the integral

$$(f \star g)(t) = \int_0^t f(t - u)g(u) du$$

- $(f \star g)(t) = (g \star f)(t)$  verify by simple change of variable
- $L[(f \star g)(t) \colon t \to s] = \bar{f}(s)\bar{g}(s)$

$$L[(f \star g)(t) \colon t \to s] = \int_0^\infty \left( \int_0^t f(t - u)g(u) \, du \right) e^{-st} \, dt$$
$$= \int_0^\infty g(u) \left( \int_u^\infty f(t - u)e^{-st} \, dt \right) \, du$$

Given two functions f(t) and g(t), the convolution  $f \star g$  of f and g is defined by the integral

$$(f \star g)(t) = \int_0^t f(t - u)g(u) du$$

- $(f \star g)(t) = (g \star f)(t)$  verify by simple change of variable
- $L[(f \star g)(t) \colon t \to s] = \bar{f}(s)\bar{g}(s)$

$$L[(f \star g)(t) \colon t \to s] = \int_0^\infty \left( \int_0^t f(t - u)g(u) \, du \right) e^{-st} \, dt$$
$$= \int_0^\infty g(u) \left( \int_u^\infty f(t - u)e^{-st} \, dt \right) \, du$$
$$= \int_0^\infty g(u) \left( \int_0^\infty f(w)e^{-s(w+u)} \, dw \right) \, du =$$

Given two functions f(t) and g(t), the convolution  $f \star g$  of f and g is defined by the integral

$$(f \star g)(t) = \int_0^t f(t - u)g(u) du$$

- $(f \star g)(t) = (g \star f)(t)$  verify by simple change of variable
- $L[(f \star g)(t) \colon t \to s] = \bar{f}(s)\bar{g}(s)$

$$L[(f \star g)(t) \colon t \to s] = \int_0^\infty \left( \int_0^t f(t - u)g(u) \, du \right) e^{-st} \, dt$$
$$= \int_0^\infty g(u) \left( \int_u^\infty f(t - u)e^{-st} \, dt \right) \, du$$
$$= \int_0^\infty g(u) \left( \int_0^\infty f(w)e^{-s(w+u)} \, dw \right) \, du = \bar{f}(s)\bar{g}(s)$$

**Example.** Compute  $L^{-1}\left[s/(s^2+1)^2\colon s\to t\right]$ .

**Example.** Compute  $L^{-1}\left[s/(s^2+1)^2\colon s \to t\right]$ .

The convolution property

$$L^{-1}\left[(f(s)\bar{g}(s)\colon s\to t\right]=(f\star g)(t)$$

**Example.** Compute  $L^{-1}\left[s/(s^2+1)^2\colon s\to t\right]$ .

The convolution property

$$L^{-1}\left[(f(\bar{s})\bar{g}(s)\colon s\to t\right]=(f\star g)(t)$$

$$L^{-1}\left[\frac{s}{(s^2+1)^2}\colon s\to t\right]=$$

**Example.** Compute  $L^{-1}\left[s/(s^2+1)^2\colon s\to t\right]$ .

The convolution property

$$L^{-1}\left[(f(s)\bar{g}(s)\colon s\to t\right]=(f\star g)(t)$$

$$L^{-1}\left[\frac{s}{(s^2+1)^2}\colon s\to t\right] = L^{-1}\left[\frac{s}{s^2+1}\frac{1}{s^2+1}\colon s\to t\right]$$

**Example.** Compute  $L^{-1}\left[s/(s^2+1)^2\colon s\to t\right]$ .

The convolution property

$$L^{-1}\left[(f(s)\bar{g}(s)\colon s\to t\right]=(f\star g)(t)$$

$$L^{-1}\left[\frac{s}{(s^2+1)^2} \colon s \to t\right] = L^{-1}\left[\frac{s}{s^2+1} \frac{1}{s^2+1} \colon s \to t\right] = (\sin \star \cos)(t)$$

**Example.** Compute  $L^{-1}\left[s/(s^2+1)^2\colon s\to t\right]$ .

The convolution property

$$L^{-1}\left[(f(s)\bar{g}(s)\colon s\to t\right]=(f\star g)(t)$$

$$L^{-1}\left[\frac{s}{(s^2+1)^2} \colon s \to t\right] = L^{-1}\left[\frac{s}{s^2+1} \frac{1}{s^2+1} \colon s \to t\right] = (\sin \star \cos)(t)$$

$$(\sin \star \cos)(t) = \int_0^t \sin(t - u) \cos u \, du$$

**Example.** Compute  $L^{-1}[s/(s^2+1)^2: s \to t]$ .

The convolution property

$$L^{-1}\left[(f(s)\bar{g}(s)\colon s\to t\right]=(f\star g)(t)$$

$$L^{-1}\left[\frac{s}{(s^2+1)^2} \colon s \to t\right] = L^{-1}\left[\frac{s}{s^2+1} \frac{1}{s^2+1} \colon s \to t\right] = (\sin \star \cos)(t)$$

$$(\sin \star \cos)(t) = \int_0^t \sin(t - u) \cos u \, du$$

$$=\sin t \int_0^t \cos^2 u \, du - \cos t \int_0^t \sin u \cos u \, du =$$

**Example.** Compute  $L^{-1}\left[s/(s^2+1)^2\colon s\to t\right]$ .

The convolution property

$$L^{-1}\left[(f(s)\bar{g}(s)\colon s\to t\right]=(f\star g)(t)$$

asserts that

$$L^{-1}\left[\frac{s}{(s^2+1)^2} \colon s \to t\right] = L^{-1}\left[\frac{s}{s^2+1} \frac{1}{s^2+1} \colon s \to t\right] = (\sin \star \cos)(t)$$

$$(\sin \star \cos)(t) = \int_0^t \sin(t - u) \cos u \, du$$

$$= \sin t \int_0^t \cos^2 u \, du - \cos t \int_0^t \sin u \cos u \, du = \frac{t \sin t}{2}$$

You can check the result  $L[(t/2)\sin t\colon t\to s]$  by direct calculation as well

**Example.** Show that the solution of the integral equation

$$f(t) + \int_0^t e^u f(t-u) du = g(t)$$

can be written in the form

$$f(t) = g(t) - \int_0^t g(u) du.$$

**Example.** Show that the solution of the integral equation

$$f(t) + \int_0^t e^u f(t-u) du = g(t)$$

can be written in the form

$$f(t) = g(t) - \int_0^t g(u) du.$$

Let  $\bar{f}(s)$  be the Laplace transform of f(t). Take Laplace transforms of the integral equation

$$\bar{f} + L[(f \star e)(t) \colon t \to s] = \bar{g}$$

**Example.** Show that the solution of the integral equation

$$f(t) + \int_0^t e^u f(t-u) du = g(t)$$

can be written in the form

$$f(t) = g(t) - \int_0^t g(u) du.$$

Let  $\bar{f}(s)$  be the Laplace transform of f(t). Take Laplace transforms of the integral equation

$$\bar{f} + L[(f \star e)(t) \colon t \to s] = \bar{g}$$

and use the property of the Laplace transform of a convolution

$$\bar{f} + \bar{f}L\left[e^t \colon t \to s\right] = \bar{g} \Rightarrow$$

**Example.** Show that the solution of the integral equation

$$f(t) + \int_0^t e^u f(t-u) du = g(t)$$

can be written in the form

$$f(t) = g(t) - \int_0^t g(u) du.$$

Let  $\bar{f}(s)$  be the Laplace transform of f(t). Take Laplace transforms of the integral equation

$$\bar{f} + L[(f \star e)(t) \colon t \to s] = \bar{g}$$

▶ and use the property of the Laplace transform of a convolution

$$[\bar{f} + \bar{f}L[e^t: t \to s] = \bar{g} \Rightarrow \bar{f} = \bar{g} - \frac{\bar{g}}{s}$$

**Example.** Show that the solution of the integral equation

$$f(t) + \int_0^t e^u f(t-u) du = g(t)$$

can be written in the form

$$f(t) = g(t) - \int_0^t g(u) du.$$

Let  $\bar{f}(s)$  be the Laplace transform of f(t). Take Laplace transforms of the integral equation

$$\bar{f} + L[(f \star e)(t): t \rightarrow s] = \bar{g}$$

▶ and use the property of the Laplace transform of a convolution

$$\bar{f} + \bar{f}L\left[e^t \colon t \to s\right] = \bar{g} \Rightarrow \bar{f} = \bar{g} - \frac{\bar{g}}{s}$$

$$f(t) = g(t) - L^{-1} \left[ \frac{\overline{g}(s)}{s} \colon s \to t \right] = g(t) + \int_0^t g(x) dx$$



$$L\left[\dot{f}(t)\colon t o s
ight]=$$

$$L\left[\dot{f}(t)\colon t\to s\right]=\int_0^\infty \dot{f}(t)e^{-st}\,dt=$$

$$L\left[\dot{f}(t)\colon t\to s\right] = \int_0^\infty \dot{f}(t)e^{-st}\,dt = \left[f(t)e^{-st}\right]_0^\infty +$$

$$+s\int_0^\infty f(t)e^{-st}\,dt =$$

$$L\left[\dot{f}(t)\colon t\to s\right] = \int_0^\infty \dot{f}(t)e^{-st}\,dt = \left[f(t)e^{-st}\right]_0^\infty +$$

$$+s\int_0^\infty f(t)e^{-st}\,dt = s\bar{f}(s) - f(0)$$

We assume that a function f has a Laplace transform and its derivative  $\dot{f}$  is piecewise continuous in the interval  $(0,\infty)$ . Then

$$L\left[\dot{f}(t)\colon t\to s\right] = \int_0^\infty \dot{f}(t)e^{-st}\,dt = \left[f(t)e^{-st}\right]_0^\infty +$$

$$+s\int_0^\infty f(t)e^{-st}\,dt = s\bar{f}(s) - f(0)$$

Similarly,

$$L\left[\ddot{f}(t)\colon t\to s\right]=\int_0^\infty \ddot{f}(t)e^{-st}\,dt=$$

We assume that a function f has a Laplace transform and its derivative  $\dot{f}$  is piecewise continuous in the interval  $(0,\infty)$ . Then

$$L\left[\dot{f}(t)\colon t\to s\right] = \int_0^\infty \dot{f}(t)e^{-st}\,dt = \left[f(t)e^{-st}\right]_0^\infty +$$

$$+s\int_0^\infty f(t)e^{-st}\,dt = s\bar{f}(s) - f(0)$$

Similarly,

$$L\left[\ddot{f}(t)\colon t\to s\right] = \int_0^\infty \ddot{f}(t)e^{-st} dt = \left[\dot{f}(t)e^{-st}\right]_0^\infty +$$

$$+s\int_0^\infty \dot{f}(t)e^{-st} dt =$$

# Application to Initial-Value Problems

We assume that a function f has a Laplace transform and its derivative  $\dot{f}$  is piecewise continuous in the interval  $(0,\infty)$ . Then

$$L\left[\dot{f}(t)\colon t\to s\right] = \int_0^\infty \dot{f}(t)e^{-st}\,dt = \left[f(t)e^{-st}\right]_0^\infty +$$

$$+s\int_0^\infty f(t)e^{-st}\,dt = s\bar{f}(s) - f(0)$$

Similarly,

$$L\left[\ddot{f}(t)\colon t\to s\right] = \int_0^\infty \ddot{f}(t)e^{-st} dt = \left[\dot{f}(t)e^{-st}\right]_0^\infty +$$

$$+s\int_0^\infty \dot{f}(t)e^{-st} dt = -\dot{f}(0) + sL\left[\dot{f}(t)\colon t\to s\right]$$

# Application to Initial-Value Problems

We assume that a function f has a Laplace transform and its derivative  $\dot{f}$  is piecewise continuous in the interval  $(0,\infty)$ . Then

$$L\left[\dot{f}(t)\colon t\to s\right] = \int_0^\infty \dot{f}(t)e^{-st}\,dt = \left[f(t)e^{-st}\right]_0^\infty +$$

$$+s\int_0^\infty f(t)e^{-st}\,dt = s\bar{f}(s) - f(0)$$

Similarly,

$$L\left[\ddot{f}(t)\colon t\to s\right] = \int_0^\infty \ddot{f}(t)e^{-st} dt = \left[\dot{f}(t)e^{-st}\right]_0^\infty +$$

$$+s\int_0^\infty \dot{f}(t)e^{-st} dt = -\dot{f}(0) + sL\left[\dot{f}(t)\colon t\to s\right]$$

$$= s^2\bar{f}(s) - sf(0) - \dot{f}(0)$$

▶ Both results are valid for s > 0.

**Example 1.** Find the solution of the initial value problem

$$\ddot{y} - 4y = 0$$
,  $y(0) = 1$ ,  $\dot{y}(0) = 2$ .

**Example 1.** Find the solution of the initial value problem

$$\ddot{y} - 4y = 0$$
,  $y(0) = 1$ ,  $\dot{y}(0) = 2$ .

Denote  $\bar{y}(s) = L[y(t)\colon t o s]$  and

**Example 1.** Find the solution of the initial value problem

$$\ddot{y} - 4y = 0$$
,  $y(0) = 1$ ,  $\dot{y}(0) = 2$ .

Denote  $\bar{y}(s)=L[y(t)\colon t\to s]$  and apply the Laplace transform operation to the differential equation

$$s^2\bar{y}(s) - sy(0) - \dot{y}(0) - 4\bar{y}(s) = 0.$$

**Example 1.** Find the solution of the initial value problem

$$\ddot{y} - 4y = 0$$
,  $y(0) = 1$ ,  $\dot{y}(0) = 2$ .

Denote  $\bar{y}(s)=L[y(t)\colon t\to s]$  and apply the Laplace transform operation to the differential equation

$$s^2\bar{y}(s) - sy(0) - \dot{y}(0) - 4\bar{y}(s) = 0.$$

Solve the algebraic equation for  $\bar{y}(s)$  to obtain

$$\bar{y}(s) = \frac{sy(0) + \dot{y}(0)}{s^2 - 4} =$$

**Example 1.** Find the solution of the initial value problem

$$\ddot{y} - 4y = 0$$
,  $y(0) = 1$ ,  $\dot{y}(0) = 2$ .

Denote  $\bar{y}(s)=L[y(t)\colon t\to s]$  and apply the Laplace transform operation to the differential equation

$$s^2\bar{y}(s) - sy(0) - \dot{y}(0) - 4\bar{y}(s) = 0.$$

Solve the algebraic equation for  $\bar{y}(s)$  to obtain

$$\bar{y}(s) = \frac{sy(0) + \dot{y}(0)}{s^2 - 4} = \frac{s + 2}{s^2 - 4} = \frac{1}{s - 2}.$$

**Example 1.** Find the solution of the initial value problem

$$\ddot{y} - 4y = 0$$
,  $y(0) = 1$ ,  $\dot{y}(0) = 2$ .

Denote  $\bar{y}(s) = L[y(t)\colon t \to s]$  and apply the Laplace transform operation to the differential equation

$$s^2\bar{y}(s) - sy(0) - \dot{y}(0) - 4\bar{y}(s) = 0.$$

Solve the algebraic equation for  $\bar{y}(s)$  to obtain

$$\bar{y}(s) = \frac{sy(0) + \dot{y}(0)}{s^2 - 4} = \frac{s + 2}{s^2 - 4} = \frac{1}{s - 2}.$$

It immediately follows that

$$y(t) = L^{-1} \left[ \frac{1}{s-2} \colon s \to t \right] = e^{2t}.$$

**Example 2.** Find the solution of the initial value problem

$$\ddot{y} + 4y = 0$$
,  $y(0) = 1$ ,  $\dot{y}(0) = 2$ .

**Example 2.** Find the solution of the initial value problem

$$\ddot{y} + 4y = 0$$
,  $y(0) = 1$ ,  $\dot{y}(0) = 2$ .

Let 
$$\bar{y}(s) = L[y(t): t \to s]$$
,

**Example 2.** Find the solution of the initial value problem

$$\ddot{y} + 4y = 0$$
,  $y(0) = 1$ ,  $\dot{y}(0) = 2$ .

Let  $\bar{y}(s) = L[y(t) \colon t \to s]$ , then the usual procedure shows that  $\bar{y}(s)$  satisfies

$$s^2\bar{y}(s) - sy(0) - \dot{y}(0) + 4\bar{y}(s) = 0$$

and

**Example 2.** Find the solution of the initial value problem

$$\ddot{y} + 4y = 0$$
,  $y(0) = 1$ ,  $\dot{y}(0) = 2$ .

Let  $\bar{y}(s) = L[y(t) \colon t \to s]$ , then the usual procedure shows that  $\bar{y}(s)$  satisfies

$$s^2\bar{y}(s) - sy(0) - \dot{y}(0) + 4\bar{y}(s) = 0$$

and therefore,

$$\bar{y}(s) = \frac{s+2}{s^2+4} = \frac{s}{s^2+4} + \frac{2}{s^2+4}.$$

**Example 2.** Find the solution of the initial value problem

$$\ddot{y} + 4y = 0$$
,  $y(0) = 1$ ,  $\dot{y}(0) = 2$ .

Let  $\bar{y}(s) = L[y(t) \colon t \to s]$ , then the usual procedure shows that  $\bar{y}(s)$  satisfies

$$s^2\bar{y}(s) - sy(0) - \dot{y}(0) + 4\bar{y}(s) = 0$$

and therefore,

$$\bar{y}(s) = \frac{s+2}{s^2+4} = \frac{s}{s^2+4} + \frac{2}{s^2+4}.$$

Thus

$$y(t) = \cos 2t + \sin 2t$$

As you can see, the method as useful as far we can take inverse Laplace transforms easily.

**Example 3.** Consider an electrical circuit with

- ▶ an inductor of inductance  $L = 10^3$  Henry,
- ightharpoonup a resistor of resistance R = 6000 Ohm,
- ▶ an uncharged capacitor of capacitance  $C = (1/9)10^3$  Farad and
- ▶ a battery of electromotive force  $U(t) = 250 \cos t$  Volt which are connected in a serial arrangement. The circuit is activated at t = 0. Find the charge on the capacitor and the electrical current flowing in the circuit for  $t \ge 0$ .

**Example 3.** Consider an electrical circuit with

- ▶ an inductor of inductance  $L = 10^3$  Henry,
- ightharpoonup a resistor of resistance  $R = 6000 \, Ohm$ ,
- ▶ an uncharged capacitor of capacitance  $C = (1/9)10^3$  Farad and
- ▶ a battery of electromotive force  $U(t) = 250 \cos t$  Volt

which are connected in a serial arrangement. The circuit is activated at t=0. Find the charge on the capacitor and the electrical current flowing in the circuit for t>0.

Apply Kirchhoff's Second Law (the sum of all voltage drops around a closed circuit is zero) to obtain the differential equation w.r.t the electrical current I and the rate of change of charge Q

$$L\dot{I}(t) + RI(t) + \frac{Q}{C} = U(t)$$

With  $\dot{Q}(t) = I(t)$ , initial conditions Q(0) = 0 and I(0) = 0 and the given values for L, R, C, the IVP becomes

$$\ddot{Q}(t) + 6\dot{Q}(t) + 9Q(t) = \frac{1}{4}\cos t, \ Q(0) = 0, \ \dot{Q}(0) = 0.$$

With  $\dot{Q}(t)=I(t)$ , initial conditions Q(0)=0 and I(0)=0 and the given values for  $L,\,R,\,C$ , the IVP becomes

$$\ddot{Q}(t) + 6\dot{Q}(t) + 9Q(t) = \frac{1}{4}\cos t, \ Q(0) = 0, \ \dot{Q}(0) = 0.$$

Denote  $\bar{Q}(s) = L[Q(t): t \to s]$  and apply the Laplace transform operation to the equation:

With  $\dot{Q}(t) = I(t)$ , initial conditions Q(0) = 0 and I(0) = 0 and the given values for L, R, C, the IVP becomes

$$\ddot{Q}(t) + 6\dot{Q}(t) + 9Q(t) = \frac{1}{4}\cos t, \ Q(0) = 0, \ \dot{Q}(0) = 0.$$

Denote  $\bar{Q}(s) = L[Q(t): t \to s]$  and apply the Laplace transform operation to the equation:

$$s^2 \bar{Q}(s) - sQ(0) - \dot{Q}(0) + 6(s\bar{Q}(s) - Q(0)) + 9\bar{Q}(s) = \frac{1}{4} \frac{s}{s^2 + 1}$$

With  $\dot{Q}(t) = I(t)$ , initial conditions Q(0) = 0 and I(0) = 0 and the given values for L, R, C, the IVP becomes

$$\ddot{Q}(t) + 6\dot{Q}(t) + 9Q(t) = \frac{1}{4}\cos t, \ Q(0) = 0, \ \dot{Q}(0) = 0.$$

Denote  $\bar{Q}(s) = L[Q(t): t \to s]$  and apply the Laplace transform operation to the equation:

$$s^2\bar{Q}(s) - sQ(0) - \dot{Q}(0) + 6(s\bar{Q}(s) - Q(0)) + 9\bar{Q}(s) = \frac{1}{4}\frac{s}{s^2 + 1}$$

Express  $\bar{Q}(s)$  and apply partial fraction decomposition:

With  $\dot{Q}(t) = I(t)$ , initial conditions Q(0) = 0 and I(0) = 0 and the given values for L, R, C, the IVP becomes

$$\ddot{Q}(t) + 6\dot{Q}(t) + 9Q(t) = \frac{1}{4}\cos t, \ Q(0) = 0, \ \dot{Q}(0) = 0.$$

Denote  $\bar{Q}(s) = L[Q(t)\colon t \to s]$  and apply the Laplace transform operation to the equation:

$$s^2ar{Q}(s) - sQ(0) - \dot{Q}(0) + 6(sar{Q}(s) - Q(0)) + 9ar{Q}(s) = rac{1}{4}rac{s}{s^2+1}$$

Express  $\bar{Q}(s)$  and apply partial fraction decomposition:

$$\bar{Q}(s) = \frac{s}{4(s^2+1)(s^2+6s+9)} =$$

$$= \frac{1}{200} \left( \frac{4s}{s^2+1} + \frac{3}{s^2+1} - \frac{4}{s+3} - \underbrace{\frac{15}{(s+3)^2}}_{\text{apply the shift formula}} \right)$$

Extract the solution

$$Q(t) = \frac{1}{200} (4\cos t + 3\sin t - 4e^{-3t} - 15te^{-3t})$$

Extract the solution

$$Q(t) = \frac{1}{200} (4\cos t + 3\sin t - 4e^{-3t} - 15te^{-3t})$$

The current I(t) follows immediately from  $I(t) = \dot{Q}(t)$  and is given by

$$I(t) = \frac{1}{200}(-4\sin t + 3\cos t - 3e^{-3t} + 45te^{-3t}).$$

Extract the solution

$$Q(t) = \frac{1}{200} (4\cos t + 3\sin t - 4e^{-3t} - 15te^{-3t})$$

The current I(t) follows immediately from  $I(t) = \dot{Q}(t)$  and is given by

$$I(t) = \frac{1}{200}(-4\sin t + 3\cos t - 3e^{-3t} + 45te^{-3t}).$$

The long-term behavior of the solution is described by the result

$$\lim_{t\to\infty}\left[Q(t)-\frac{1}{200}(4\cos t+3\sin t)\right]=0.$$

Extract the solution

$$Q(t) = \frac{1}{200} (4\cos t + 3\sin t - 4e^{-3t} - 15te^{-3t})$$

The current I(t) follows immediately from  $I(t) = \dot{Q}(t)$  and is given by

$$I(t) = \frac{1}{200}(-4\sin t + 3\cos t - 3e^{-3t} + 45te^{-3t}).$$

The long-term behavior of the solution is described by the result

$$\lim_{t\to\infty}\left[Q(t)-\frac{1}{200}(4\cos t+3\sin t)\right]=0.$$

The solution for large time is dominated by the inhomogeneous term in the differential equation, while the contributions from the complementary function are damped out.

Extract the solution

$$Q(t) = \frac{1}{200} (4\cos t + 3\sin t - 4e^{-3t} - 15te^{-3t})$$

The current I(t) follows immediately from  $I(t) = \dot{Q}(t)$  and is given by

$$I(t) = \frac{1}{200}(-4\sin t + 3\cos t - 3e^{-3t} + 45te^{-3t}).$$

The long-term behavior of the solution is described by the result

$$\lim_{t\to\infty}\left|Q(t)-\frac{1}{200}(4\cos t+3\sin t)\right|=0.$$

The solution for large time is dominated by the inhomogeneous term in the differential equation, while the contributions from the complementary function are damped out.

In this example, the forcing function  $U(t) = 250 \cos t$  is a continuous function of time. In many applications, such a function is a step function or an impulse function.

The unit step function (or Heaviside function) is defined by

$$H(t-a) = \begin{cases} 0 & t < a \\ 1 & t \ge a. \end{cases}$$

The unit step function (or Heaviside function) is defined by

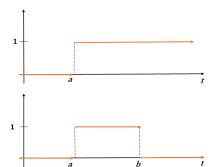
$$H(t-a) = \begin{cases} 0 & t < a \\ 1 & t \ge a. \end{cases}$$

The unit step function can be used to construct more complicated step functions, for example f(t) = H(t-a) - H(t-b).

The unit step function (or Heaviside function) is defined by

$$H(t-a) = \begin{cases} 0 & t < a \\ 1 & t \ge a. \end{cases}$$

The unit step function can be used to construct more complicated step functions, for example f(t) = H(t - a) - H(t - b).



The Laplace transform of the unit step function is

$$L[H(t-a)\colon t\to s]=\int_0^\infty e^{-st}H(t-a)\,dt=$$

The Laplace transform of the unit step function is

$$L[H(t-a): t \to s] = \int_0^\infty e^{-st} H(t-a) dt = \int_a^\infty e^{-st} dt =$$

The Laplace transform of the unit step function is

$$L[H(t-a): t \to s] = \int_0^\infty e^{-st} H(t-a) dt = \int_0^\infty e^{-st} dt = \frac{e^{-as}}{s}, s > 0$$

The Laplace transform of the unit step function is

$$L[H(t-a): t \to s] = \int_0^\infty e^{-st} H(t-a) dt = \int_a^\infty e^{-st} dt = \frac{e^{-as}}{s}, s > 0$$

The second shifting property

$$L[f(t-a)H(t-a)\colon t\to s]=e^{-as}\bar{f}(s)$$

The Laplace transform of the unit step function is

$$L[H(t-a): t \to s] = \int_0^\infty e^{-st} H(t-a) dt = \int_0^\infty e^{-st} dt = \frac{e^{-as}}{s}, s > 0$$

The second shifting property

$$L[f(t-a)H(t-a): t \to s] = e^{-as}\bar{f}(s)$$

**Example.** Compute 
$$L^{-1}\left[\frac{(1-e^{-\pi s/2})}{1+s^2}: s \to t\right]$$

The Laplace transform of the unit step function is

$$L[H(t-a): t \to s] = \int_{0}^{\infty} e^{-st} H(t-a) dt = \int_{0}^{\infty} e^{-st} dt = \frac{e^{-as}}{s}, s > 0$$

The second shifting property

$$L[f(t-a)H(t-a): t \rightarrow s] = e^{-as}\overline{f}(s)$$

**Example.** Compute  $L^{-1}\left[\frac{(1-e^{-\pi s/2})}{1+s^2}: s \to t\right]$ 

$$L^{-1}\left[\frac{1 - e^{-\pi s/2}}{1 + s^2} \colon s \to t\right] = L^{-1}\left[\frac{1}{1 + s^2} \colon s \to t\right] - L^{-1}\left[\frac{e^{-\pi s/2}}{1 + s^2} \colon s \to t\right]$$

The Laplace transform of the unit step function is

$$L[H(t-a): t \to s] = \int_0^\infty e^{-st} H(t-a) dt = \int_0^\infty e^{-st} dt = \frac{e^{-as}}{s}, s > 0$$

The second shifting property

$$L[f(t-a)H(t-a): t \rightarrow s] = e^{-as}\bar{f}(s)$$

**Example.** Compute  $L^{-1}\left[\frac{(1-e^{-\pi s/2})}{1+s^2}: s \to t\right]$ 

$$L^{-1}\left[\frac{1 - e^{-\pi s/2}}{1 + s^2} \colon s \to t\right] = L^{-1}\left[\frac{1}{1 + s^2} \colon s \to t\right] - L^{-1}\left[\frac{e^{-\pi s/2}}{1 + s^2} \colon s \to t\right]$$

$$= \sin t -$$

The Laplace transform of the unit step function is

$$L[H(t-a): t \to s] = \int_0^\infty e^{-st} H(t-a) dt = \int_0^\infty e^{-st} dt = \frac{e^{-as}}{s}, s > 0$$

The second shifting property

$$L[f(t-a)H(t-a): t \rightarrow s] = e^{-as}\overline{f}(s)$$

**Example.** Compute  $L^{-1}\left[\frac{(1-e^{-\pi s/2})}{1+s^2}: s \to t\right]$ 

$$L^{-1}\left[\frac{1 - e^{-\pi s/2}}{1 + s^2} : s \to t\right] = L^{-1}\left[\frac{1}{1 + s^2} : s \to t\right] - L^{-1}\left[\frac{e^{-\pi s/2}}{1 + s^2} : s \to t\right]$$

$$= \sin t - H(t - \pi/2)\sin(t - \pi/2) =$$

The Laplace transform of the unit step function is

$$L[H(t-a): t \to s] = \int_{0}^{\infty} e^{-st} H(t-a) dt = \int_{0}^{\infty} e^{-st} dt = \frac{e^{-as}}{s}, s > 0$$

The second shifting property

$$L[f(t-a)H(t-a): t \rightarrow s] = e^{-as}\overline{f}(s)$$

**Example.** Compute  $L^{-1}\left[\frac{(1-e^{-\pi s/2})}{1+s^2}: s \to t\right]$ 

$$L^{-1}\left[\frac{1 - e^{-\pi s/2}}{1 + s^2} \colon s \to t\right] = L^{-1}\left[\frac{1}{1 + s^2} \colon s \to t\right] - L^{-1}\left[\frac{e^{-\pi s/2}}{1 + s^2} \colon s \to t\right]$$

$$= \sin t - H(t - \pi/2) \sin(t - \pi/2) = \sin t + H(t - \pi/2) \cos t$$

**Example.** Solve the initial value problem

$$\ddot{y} + y = \begin{cases} 1 & 0 \le t < 1, \\ 0 & t \ge 1, \end{cases}, y(0) = 0, \dot{y}(0) = 0.$$

Such a problem may appear when an undamped oscillator is disturbed from rest in its equilibrium state by a constant external driving force which is active for a finite period only.

**Example.** Solve the initial value problem

$$\ddot{y} + y = \begin{cases} 1 & 0 \le t < 1, \\ 0 & t \ge 1, \end{cases}, y(0) = 0, \dot{y}(0) = 0.$$

Such a problem may appear when an undamped oscillator is disturbed from rest in its equilibrium state by a constant external driving force which is active for a finite period only.

Represent the driving force as

$$f(t) = H(t) - H(t-1)$$

and let  $\bar{y}(s) = L[y(t): t \to s]$ .

**Example.** Solve the initial value problem

$$\ddot{y} + y = \begin{cases} 1 & 0 \le t < 1, \\ 0 & t \ge 1, \end{cases}, y(0) = 0, \dot{y}(0) = 0.$$

Such a problem may appear when an undamped oscillator is disturbed from rest in its equilibrium state by a constant external driving force which is active for a finite period only. Represent the driving force as

$$f(t) = H(t) - H(t-1)$$

and let  $\bar{y}(s) = L[y(t): t \rightarrow s]$ .

Take the Laplace transform of the equation to obtain

$$s^2\bar{y}(s) - s\dot{y}(0) - y(0) + \bar{y}(s) = \frac{1}{s} - \frac{e^{-s}}{s}$$

Then

$$\bar{y}(s) = \frac{1 - e^{-s}}{s(s^2 + 1)} =$$

Then

$$\bar{y}(s) = \frac{1 - e^{-s}}{s(s^2 + 1)} = \frac{1}{s} - \frac{s}{s^2 + 1} - \underbrace{\frac{e^{-s}}{s} + \frac{se^{-s}}{s^2 + 1}}_{\text{apply the 2nd shifting pr.}}$$

Then

$$\bar{y}(s) = \frac{1 - e^{-s}}{s(s^2 + 1)} = \frac{1}{s} - \frac{s}{s^2 + 1} - \underbrace{\frac{e^{-s}}{s} + \frac{se^{-s}}{s^2 + 1}}_{\text{apply the 2nd shifting pr.}}$$

The solution of the IVP is

$$y(t) = 1 - \cos t - H(t-1) + H(t-1)\cos(t-1)$$

Then

$$\bar{y}(s) = \frac{1 - e^{-s}}{s(s^2 + 1)} = \frac{1}{s} - \frac{s}{s^2 + 1} - \underbrace{\frac{e^{-s}}{s} + \frac{se^{-s}}{s^2 + 1}}_{\text{apply the 2nd shifting pr.}}$$

The solution of the IVP is

$$y(t) = 1 - \cos t - H(t-1) + H(t-1)\cos(t-1)$$
 
$$y(t) = \begin{cases} 1 - \cos t & 0 \le t < 1, \\ \cos(t-1) - \cos t & t > 1, \end{cases}$$

Then

$$\bar{y}(s) = \frac{1 - e^{-s}}{s(s^2 + 1)} = \frac{1}{s} - \frac{s}{s^2 + 1} - \underbrace{\frac{e^{-s}}{s} + \frac{se^{-s}}{s^2 + 1}}_{\text{apply the 2nd shifting pr.}}$$

The solution of the IVP is

$$y(t) = 1 - \cos t - H(t-1) + H(t-1)\cos(t-1)$$

$$y(t) = \begin{cases} 1 - \cos t & 0 \le t < 1, \\ \cos(t-1) - \cos t & t \ge 1, \end{cases}$$

The solution is continuous at t = 1 even though the driving force f(t) is not.

Consider the unit impulse function  $\delta(t)$  (or delta or Dirac delta function) which is not an ordinary function but belongs to the class of so-called generalized functions or distributions.

Consider the unit impulse function  $\delta(t)$  (or delta or Dirac delta function) which is not an ordinary function but belongs to the class of so-called generalized functions or distributions.

The unit impulse function  $\delta(t)$  is defined through the integral relation

$$\int_{-\infty}^{+\infty} \delta(t-a)f(t)\,dt = f(a)$$

for any integrable function f(t).

Consider the unit impulse function  $\delta(t)$  (or delta or Dirac delta function) which is not an ordinary function but belongs to the class of so-called generalized functions or distributions.

The unit impulse function  $\delta(t)$  is defined through the integral relation

$$\int_{-\infty}^{+\infty} \delta(t-a)f(t)\,dt = f(a)$$

for any integrable function f(t).

Consider the unit impulse function  $\delta(t)$  (or delta or Dirac delta function) which is not an ordinary function but belongs to the class of so-called generalized functions or distributions.

The unit impulse function  $\delta(t)$  is defined through the integral relation

$$\int_{-\infty}^{+\infty} \delta(t-a)f(t)\,dt = f(a)$$

for any integrable function f(t).

Sequence 1: Let 
$$\delta_{\varepsilon}(t-a) = \left\{ \begin{array}{ll} \frac{1}{2\varepsilon} & |t-a| < \varepsilon, \\ 0 & |t-a| \geq \varepsilon. \end{array} \right.$$

Consider the unit impulse function  $\delta(t)$  (or delta or Dirac delta function) which is not an ordinary function but belongs to the class of so-called generalized functions or distributions.

The unit impulse function  $\delta(t)$  is defined through the integral relation

$$\int_{-\infty}^{+\infty} \delta(t-a)f(t)\,dt = f(a)$$

for any integrable function f(t).

Sequence 1: Let 
$$\delta_{\varepsilon}(t-a) = \begin{cases} \frac{1}{2\varepsilon} & |t-a| < \varepsilon, \\ 0 & |t-a| \geq \varepsilon. \end{cases}$$

$$\int_{-\infty}^{+\infty} \delta_{arepsilon}(t-a) f(t) \, dt = rac{1}{2arepsilon} \int_{a-arepsilon}^{a+arepsilon} f(t) \, dt =$$

Consider the unit impulse function  $\delta(t)$  (or delta or Dirac delta function) which is not an ordinary function but belongs to the class of so-called generalized functions or distributions.

The unit impulse function  $\delta(t)$  is defined through the integral relation

$$\int_{-\infty}^{+\infty} \delta(t-a)f(t)\,dt = f(a)$$

for any integrable function f(t).

**Sequence 1:** Let 
$$\delta_{\varepsilon}(t-a) = \begin{cases} \frac{1}{2\varepsilon} & |t-a| < \varepsilon, \\ 0 & |t-a| \geq \varepsilon. \end{cases}$$

$$\int_{-\infty}^{+\infty} \delta_{\varepsilon}(t-\mathsf{a}) f(t) \, dt = \frac{1}{2\varepsilon} \int_{\mathsf{a}-\varepsilon}^{\mathsf{a}+\varepsilon} f(t) \, dt = f(\mathsf{a}+\theta\varepsilon), \, |\theta| < 1$$

We used the fact that f is integrable and applied the integral mean value theorem above.

We used the fact that f is integrable and applied the integral mean value theorem above. Therefore,

$$\lim_{\varepsilon\to 0}\int_{-\infty}^{+\infty}\delta_{\varepsilon}(t-a)f(t)\,dt=f(a)$$

for any real integrable function f. Therefore,

$$\delta(t-a) = \lim_{\varepsilon \to 0} \delta_{\varepsilon}(t-a)$$

We used the fact that f is integrable and applied the integral mean value theorem above. Therefore,

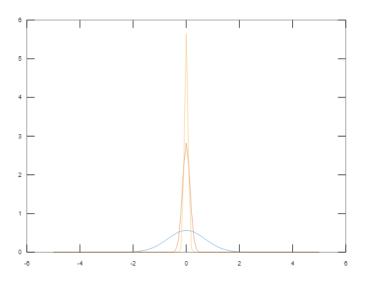
$$\lim_{\varepsilon\to 0}\int_{-\infty}^{+\infty}\delta_{\varepsilon}(t-a)f(t)\,dt=f(a)$$

for any real integrable function f. Therefore,

$$\delta(t-a) = \lim_{\varepsilon \to 0} \delta_{\varepsilon}(t-a)$$

**Sequence 2:** Another popular description of  $\delta(t)$  (a=0) is the limit of the sequence

$$\delta_n(t) = \lim_{n \to \infty} \frac{n}{\sqrt{\pi}} e^{-n^2 t^2}$$



The graph of the function  $(n/\sqrt{\pi})e^{-n^2t^2}$  for  $n=1,\,5,\,10.$ 

In the strict mathematical sense neither of these limits exists, but they provide useful conceptual ways towards an understanding of the delta function.

In the strict mathematical sense neither of these limits exists, but they provide useful conceptual ways towards an understanding of the delta function.

Letting f(t) = 1 in the definition of the delta function,

$$\int_{-\infty}^{+\infty} \delta(t-a) \, dt = 1$$

In the strict mathematical sense neither of these limits exists, but they provide useful conceptual ways towards an understanding of the delta function.

Letting f(t) = 1 in the definition of the delta function,

$$\int_{-\infty}^{+\infty} \delta(t-a) \, dt = 1$$

For this reason  $\delta(t-a)$  is called the unit impulse function.

In the strict mathematical sense neither of these limits exists, but they provide useful conceptual ways towards an understanding of the delta function.

Letting f(t) = 1 in the definition of the delta function,

$$\int_{-\infty}^{+\infty} \delta(t-a) \, dt = 1$$

For this reason  $\delta(t-a)$  is called the unit impulse function.

▶ The unit step function is related to the unit impulse function

$$H(t-a) = \int_{-\infty}^{t} \delta(u-a) du \Rightarrow \frac{d}{dt} H(t-a) = \delta(t-a)$$

In the strict mathematical sense neither of these limits exists, but they provide useful conceptual ways towards an understanding of the delta function.

Letting f(t) = 1 in the definition of the delta function,

$$\int_{-\infty}^{+\infty} \delta(t-a) \, dt = 1$$

For this reason  $\delta(t-a)$  is called the unit impulse function.

▶ The unit step function is related to the unit impulse function

$$H(t-a) = \int_{-\infty}^{t} \delta(u-a) du \Rightarrow \frac{d}{dt} H(t-a) = \delta(t-a)$$

$$L[\delta(t-a): t \rightarrow s] =$$

In the strict mathematical sense neither of these limits exists, but they provide useful conceptual ways towards an understanding of the delta function.

Letting f(t) = 1 in the definition of the delta function,

$$\int_{-\infty}^{+\infty} \delta(t-a) \, dt = 1$$

For this reason  $\delta(t-a)$  is called the unit impulse function.

The unit step function is related to the unit impulse function

$$H(t-a) = \int_{-\infty}^{t} \delta(u-a) du \Rightarrow \frac{d}{dt} H(t-a) = \delta(t-a)$$

$$L[\delta(t-a): t \to s] = \int_{0}^{\infty} \delta(t-a)e^{-st} dt$$

In the strict mathematical sense neither of these limits exists, but they provide useful conceptual ways towards an understanding of the delta function.

▶ Letting f(t) = 1 in the definition of the delta function,

$$\int_{-\infty}^{+\infty} \delta(t-a) \, dt = 1$$

For this reason  $\delta(t-a)$  is called the unit impulse function.

▶ The unit step function is related to the unit impulse function

$$H(t-a) = \int_{-\infty}^{t} \delta(u-a) du \Rightarrow \frac{d}{dt} H(t-a) = \delta(t-a)$$

$$L[\delta(t-a)\colon t\to s] = \int_0^\infty \delta(t-a)e^{-st} dt$$
$$= \int_0^\infty \delta(t-a)e^{-st} dt =$$

In the strict mathematical sense neither of these limits exists, but they provide useful conceptual ways towards an understanding of the delta function.

▶ Letting f(t) = 1 in the definition of the delta function,

$$\int_{-\infty}^{+\infty} \delta(t-a) \, dt = 1$$

For this reason  $\delta(t-a)$  is called the unit impulse function.

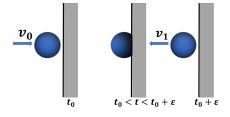
▶ The unit step function is related to the unit impulse function

$$H(t-a) = \int_{-\infty}^{t} \delta(u-a) du \Rightarrow \frac{d}{dt} H(t-a) = \delta(t-a)$$

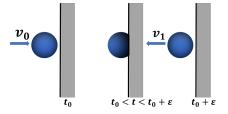
$$L[\delta(t-a): t \to s] = \int_0^\infty \delta(t-a)e^{-st} dt$$
$$= \int_0^\infty \delta(t-a)e^{-st} dt = e^{-as}, \ a > 0$$

Consider an elastic ball of mass m moving at velocity  $v_0$  toward a rigid wall.

Consider an elastic ball of mass m moving at velocity  $v_0$  toward a rigid wall.

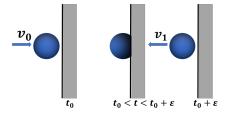


Consider an elastic ball of mass m moving at velocity  $v_0$  toward a rigid wall.



At time  $t_0$ , the ball collides with the wall, the wall exerts a force f(t) on the ball over a short period of time  $\varepsilon$ .

Consider an elastic ball of mass m moving at velocity  $v_0$  toward a rigid wall.

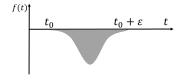


- At time  $t_0$ , the ball collides with the wall, the wall exerts a force f(t) on the ball over a short period of time  $\varepsilon$ .
- During this time, the ball is in contact with the wall and the velocity of the ball reduces from  $v_0$  to 0 and then changes its direction, finally leaving the wall with velocity  $-v_1$ .

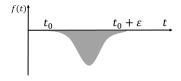
▶ The force f(t) depends on the contact between the elastic ball and the rigid wall.

▶ The force f(t) depends on the contact between the elastic ball and the rigid wall. It is opposite to the direction of the initial velocity  $v_0 \Rightarrow f(t)$  is negative.

▶ The force f(t) depends on the contact between the elastic ball and the rigid wall. It is opposite to the direction of the initial velocity  $v_0 \Rightarrow f(t)$  is negative.

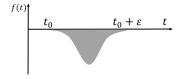


▶ The force f(t) depends on the contact between the elastic ball and the rigid wall. It is opposite to the direction of the initial velocity  $v_0 \Rightarrow f(t)$  is negative.



▶ The area under the force curve is called the impulse I

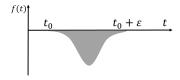
▶ The force f(t) depends on the contact between the elastic ball and the rigid wall. It is opposite to the direction of the initial velocity  $v_0 \Rightarrow f(t)$  is negative.



▶ The area under the force curve is called the impulse *I* 

$$I = \int_{t_0}^{t_0 + \varepsilon} f(t) \, dt$$

▶ The force f(t) depends on the contact between the elastic ball and the rigid wall. It is opposite to the direction of the initial velocity  $v_0 \Rightarrow f(t)$  is negative.



▶ The area under the force curve is called the impulse *I* 

$$I = \int_{t_0}^{t_0 + \varepsilon} f(t) \, dt$$

► The Impulse-Momentum Principle: the change in momentum of mass m is equal to the total impulse on m:  $m(-v_1) - mv_0 = I$ .

► A mathematical idealization:

A mathematical idealization: consider an impulse function f(t) over a time interval  $t_0 < t < t_0 + \varepsilon$  with constant amplitude  $I/\varepsilon$  such that the area under the function f(t) is I

▶ A mathematical idealization: consider an impulse function f(t) over a time interval  $t_0 < t < t_0 + \varepsilon$  with constant amplitude  $I/\varepsilon$  such that the area under the function f(t) is I

$$f(t) = \frac{I}{\varepsilon} (H(t-t_0) - H(t-(t_0+\varepsilon)))$$

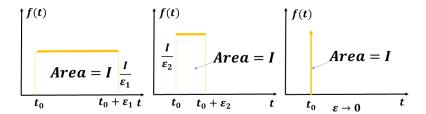
▶ A mathematical idealization: consider an impulse function f(t) over a time interval  $t_0 < t < t_0 + \varepsilon$  with constant amplitude  $I/\varepsilon$  such that the area under the function f(t) is I

$$f(t) = \frac{I}{\varepsilon} (H(t-t_0) - H(t-(t_0+\varepsilon)))$$

For I = 1, the limiting function

$$\lim_{\varepsilon \to 0} f(t) = \lim_{\varepsilon \to 0} \frac{H(t - t_0) - H(t - (t_0 + \varepsilon))}{\varepsilon}$$

is called the unit impulse function



Solve the IVP

$$\ddot{y} + 2\dot{y} + y = \delta(t-1), \ y(0) = 2, \ \dot{y}(0) = 3$$

that represents a damped oscillator with an external driving force of infinite magnitude acting for an infinitely short time such that unit impulse is imparted to the system at time t=1.

Solve the IVP

$$\ddot{y} + 2\dot{y} + y = \delta(t-1), \ y(0) = 2, \ \dot{y}(0) = 3$$

that represents a damped oscillator with an external driving force of infinite magnitude acting for an infinitely short time such that unit impulse is imparted to the system at time t=1. Let  $\bar{y}(s)=L[y(t)\colon t\to s]$  and take the Laplace transform of the differential equation:

Solve the IVP

$$\ddot{y} + 2\dot{y} + y = \delta(t-1), \ y(0) = 2, \ \dot{y}(0) = 3$$

that represents a damped oscillator with an external driving force of infinite magnitude acting for an infinitely short time such that unit impulse is imparted to the system at time t=1. Let  $\bar{y}(s)=L[y(t)\colon t\to s]$  and take the Laplace transform of the differential equation:

$$s^2\bar{y}(s) - 2s - 3 + 2(s\bar{y}(s) - 2) + \bar{y}(s) = e^{-s}$$
.

Solve the IVP

$$\ddot{y} + 2\dot{y} + y = \delta(t-1), \ y(0) = 2, \ \dot{y}(0) = 3$$

that represents a damped oscillator with an external driving force of infinite magnitude acting for an infinitely short time such that unit impulse is imparted to the system at time t = 1. Let  $\bar{v}(s) = I[v(t): t \rightarrow s]$  and take the Laplace transform of the

Let  $\bar{y}(s) = L[y(t) \colon t \to s]$  and take the Laplace transform of the differential equation:

$$s^2\bar{y}(s) - 2s - 3 + 2(s\bar{y}(s) - 2) + \bar{y}(s) = e^{-s}.$$

$$\Rightarrow \bar{y}(s) = \frac{2s + 7 + e^{-s}}{s^2 + 2s + 1} =$$

Solve the IVP

$$\ddot{y} + 2\dot{y} + y = \delta(t-1), \ y(0) = 2, \ \dot{y}(0) = 3$$

that represents a damped oscillator with an external driving force of infinite magnitude acting for an infinitely short time such that unit impulse is imparted to the system at time t=1.

Let  $\bar{y}(s) = L[y(t): t \to s]$  and take the Laplace transform of the differential equation:

$$s^2\bar{y}(s) - 2s - 3 + 2(s\bar{y}(s) - 2) + \bar{y}(s) = e^{-s}$$
.

$$\Rightarrow \bar{y}(s) = \frac{2s+7+e^{-s}}{s^2+2s+1} = \frac{2}{s+1} + \frac{5}{(s+1)^2} + \frac{e^{-s}}{(s+1)^2}$$

after taking partial fractions.

Solve the IVP

differential equation:

$$\ddot{y} + 2\dot{y} + y = \delta(t-1), \ y(0) = 2, \ \dot{y}(0) = 3$$

that represents a damped oscillator with an external driving force of infinite magnitude acting for an infinitely short time such that unit impulse is imparted to the system at time t=1. Let  $\bar{y}(s)=L[y(t)\colon t\to s]$  and take the Laplace transform of the

$$s^2\bar{y}(s) - 2s - 3 + 2(s\bar{y}(s) - 2) + \bar{y}(s) = e^{-s}$$
.

$$\Rightarrow \bar{y}(s) = \frac{2s + 7 + e^{-s}}{s^2 + 2s + 1} = \frac{2}{s + 1} + \frac{5}{(s + 1)^2} + \frac{e^{-s}}{(s + 1)^2}$$

after taking partial fractions. Apply the first and the second shifting properties to obtain

$$y(t) = 2e^{-t} + 5te^{-t} + (t-1)e^{-(t-1)}H(t-1)$$

# Laplace Transforms of Selected Functions

$$f(t) = L^{-1} \left[ \bar{f}(s) \colon s \to t \right] \qquad \bar{f}(s) = L \left[ f(t) \colon t \to s \right]$$

$$H(t-a) \qquad \qquad \frac{e^{-as}}{s} \qquad Re(s) > 0$$

$$H(t-a)f(t-a) \qquad \qquad e^{-as}\bar{f}(s)$$

$$e^{at}f(t) \qquad \qquad \bar{f}(s-a)$$

$$f(at) \qquad \qquad \frac{1}{a}\bar{f}\left(\frac{s}{a}\right) \qquad a > 0$$

$$\int_0^t f(t-w)g(w) \, dw \qquad \qquad \bar{f}(s)\bar{g}(s)$$

$$\delta(t-a) \qquad \qquad e^{-as}$$

$$(-t)^n f(t) \qquad \qquad \bar{f}^{(n)}(s)$$

$$f^{(n)}(t) \qquad \qquad s^n \bar{f}(s) - \sum_{j=0}^{n-1} s^{n-j-1} f^{(j)}(0)$$

Let f(t) be a real-valued function defined for  $0 \le t < \infty$  with periodicity T:

$$f(t) = f(t + kT), \quad k = 1, 2, ...$$

Let f(t) be a real-valued function defined for  $0 \le t < \infty$  with periodicity T:

$$f(t) = f(t + kT), \quad k = 1, 2, ...$$

The Laplace transform of the periodic function f is

$$L[f(t): t \to s] = \bar{f}(s) = \int_0^\infty f(t)e^{-st}dt = \sum_{k=0}^\infty \int_{kT}^{(k+1)T} f(t)e^{-st}dt$$

Use the substitution t = u + kT

Let f(t) be a real-valued function defined for  $0 \le t < \infty$  with periodicity T:

$$f(t) = f(t + kT), \quad k = 1, 2, ...$$

The Laplace transform of the periodic function f is

$$L[f(t): t \to s] = \bar{f}(s) = \int_0^\infty f(t)e^{-st}dt = \sum_{k=0}^\infty \int_{kT}^{(k+1)T} f(t)e^{-st}dt$$

Use the substitution t = u + kT and exploit the periodicity of f:

$$\bar{f}(s) = \sum_{k=0}^{\infty} \int_0^T f(u+kT)e^{-s(u+kT)}du =$$

Let f(t) be a real-valued function defined for  $0 \le t < \infty$  with periodicity T:

$$f(t) = f(t + kT), \quad k = 1, 2, ...$$

The Laplace transform of the periodic function f is

$$L[f(t)\colon t o s]=ar{f}(s)=\int_0^\infty f(t)e^{-st}dt=\sum_{k=0}^\infty \int_{kT}^{(k+1)T} f(t)e^{-st}dt$$

Use the substitution t = u + kT and exploit the periodicity of f:

$$\bar{f}(s) = \sum_{k=0}^{\infty} \int_{0}^{T} f(u+kT)e^{-s(u+kT)}du = \sum_{k=0}^{\infty} e^{-kTu} \int_{0}^{T} f(u)e^{-su}du$$

Let f(t) be a real-valued function defined for  $0 \le t < \infty$  with periodicity T:

$$f(t) = f(t + kT), \quad k = 1, 2, ...$$

The Laplace transform of the periodic function f is

$$L[f(t): t \to s] = \bar{f}(s) = \int_0^\infty f(t)e^{-st}dt = \sum_{k=0}^\infty \int_{kT}^{(k+1)T} f(t)e^{-st}dt$$

Use the substitution t = u + kT and exploit the periodicity of f:

$$\bar{f}(s) = \sum_{k=0}^{\infty} \int_{0}^{T} f(u+kT)e^{-s(u+kT)}du = \sum_{k=0}^{\infty} e^{-kTu} \int_{0}^{T} f(u)e^{-su}du$$

$$\bar{f}(s) = \frac{1}{1 - e^{-sT}} \int_0^T f(t)e^{-st} dt$$

To demonstrate the validity of the obtained formula, we shall use it to calculate the Laplace transform of  $\sin at$ .

To demonstrate the validity of the obtained formula, we shall use it to calculate the Laplace transform of sin *at*. We have already calculated

$$L[\sin at \colon t \to s] = \frac{a}{a^2 + s^2}$$

To demonstrate the validity of the obtained formula, we shall use it to calculate the Laplace transform of sin *at*. We have already calculated

$$L[\sin at \colon t \to s] = \frac{a}{a^2 + s^2}$$

To demonstrate the validity of the obtained formula, we shall use it to calculate the Laplace transform of sin *at*. We have already calculated

$$L[\sin at \colon t \to s] = \frac{a}{a^2 + s^2}$$

$$\int_0^T e^{-st} \sin at dt =$$

To demonstrate the validity of the obtained formula, we shall use it to calculate the Laplace transform of sin *at*. We have already calculated

$$L[\sin at \colon t \to s] = \frac{a}{a^2 + s^2}$$

$$\int_0^T e^{-st} \sin at dt = \left. \frac{-e^{-st}}{s} \sin at \right|_0^T + \frac{a}{s} \int_0^T e^{-st} \cos at \, dt$$

To demonstrate the validity of the obtained formula, we shall use it to calculate the Laplace transform of sin *at*. We have already calculated

$$L[\sin at: t \to s] = \frac{a}{a^2 + s^2}$$

$$\int_0^T e^{-st} \sin at dt = \frac{-e^{-st}}{s} \sin at \Big|_0^T + \frac{a}{s} \int_0^T e^{-st} \cos at dt$$
$$= \frac{a}{s} \frac{-e^{-st}}{s} \cos at \Big|_0^T - \frac{a^2}{s^2} \int_0^T e^{-st} \sin at dt$$

To demonstrate the validity of the obtained formula, we shall use it to calculate the Laplace transform of sin at. We have already calculated

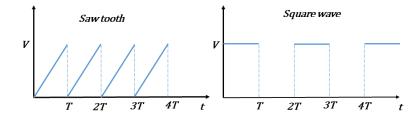
$$L[\sin at: t \to s] = \frac{a}{a^2 + s^2}$$

$$\int_0^T e^{-st} \sin at dt = \frac{-e^{-st}}{s} \sin at \Big|_0^T + \frac{a}{s} \int_0^T e^{-st} \cos at dt$$

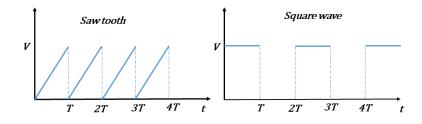
$$= \frac{a}{s} \frac{-e^{-st}}{s} \cos at \Big|_0^T - \frac{a^2}{s^2} \int_0^T e^{-st} \sin at dt$$

$$= \frac{a}{s^2} (1 - e^{-sT}) - \frac{a^2}{s^2} \int_0^T e^{-st} \sin at dt$$

The periodic waveforms saw tooth and square wave of amplitude V and periods T and 2T are shown below



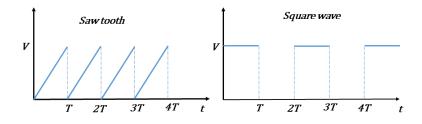
The periodic waveforms saw tooth and square wave of amplitude V and periods T and 2T are shown below



Since

$$\int_0^T f(t)e^{-st}dt =$$

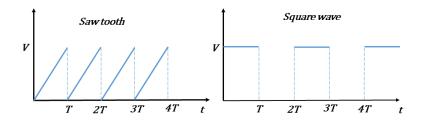
The periodic waveforms saw tooth and square wave of amplitude V and periods T and 2T are shown below



Since

$$\int_0^T f(t)e^{-st}dt = \frac{V}{T}\int_0^T te^{-st}dt =$$

The periodic waveforms saw tooth and square wave of amplitude V and periods  $\mathcal{T}$  and  $2\mathcal{T}$  are shown below



Since

$$\int_{0}^{T} f(t)e^{-st}dt = \frac{V}{T}\int_{0}^{T} te^{-st}dt = \frac{V}{Ts^{2}}(1 - e^{-sT} - sTe^{-sT}),$$

so the Laplace transform of the saw tooth waveform of amplitude  $\ensuremath{\textit{V}}$  and period  $\ensuremath{\textit{T}}$  is

$$\bar{f}(s) = \frac{V}{Ts^2} \left( 1 - \frac{sTe^{-sT}}{1 - e^{-sT}} \right)$$

so the Laplace transform of the saw tooth waveform of amplitude  $\ensuremath{V}$  and period  $\ensuremath{\mathcal{T}}$  is

$$\bar{f}(s) = \frac{V}{Ts^2} \left( 1 - \frac{sTe^{-sT}}{1 - e^{-sT}} \right)$$

The Laplace transform of the square wave of amplitude V and period 2T is

$$\bar{f}(s) = \frac{V}{1 - e^{-2sT}} \int_0^T t e^{-st} dt = \frac{V}{s} \frac{1}{1 + e^{-sT}}$$

so the Laplace transform of the saw tooth waveform of amplitude V and period T is

$$ar{f}(s) = rac{V}{Ts^2} \left( 1 - rac{sTe^{-sT}}{1 - e^{-sT}} 
ight)$$

The Laplace transform of the square wave of amplitude V and period 2T is

$$\bar{f}(s) = \frac{V}{1 - e^{-2sT}} \int_{0}^{T} te^{-st} dt = \frac{V}{s} \frac{1}{1 + e^{-sT}}$$

#### Other periodic pulses.

Meander function Triangular wave 
$$f(t) = \begin{cases} V & t \in (0,T) \\ -V & t \in (T,2T) \end{cases} \qquad f(t) = \begin{cases} tV/T & t \in (0,T) \\ \frac{(2T-t)V}{T} & t \in (T,2T) \end{cases}$$
Full wave rectification Half wave rectification 
$$f(t) = \begin{cases} V \sin(\frac{\pi t}{T}) & t \in (0,T) \\ V \sin(\pi t/T) & t \in (0,T) \end{cases}$$

Full wave rectification
$$f(t) = \left\{ \begin{array}{ll} V \sin(\frac{\pi t}{T}) & t \in (T, 2T) \\ \text{Full wave rectification} \end{array} \right.$$

$$f(t) = \left\{ \begin{array}{ll} V \sin(\frac{\pi t}{T}) & t \in (0, T) \\ 0 & t \in (T, 2T) \end{array} \right.$$

Determine the current in a series circuit containing an inductor of L Henry and a resistor of R Ohm when the circuit is driven by a saw tooth voltage of period T and amplitude V. Assume that no current is flowing initially.

Determine the current in a series circuit containing an inductor of L Henry and a resistor of R Ohm when the circuit is driven by a saw tooth voltage of period T and amplitude V. Assume that no current is flowing initially.

The current I(t) in the circuit is governed by the differential equation

$$L\dot{I} + RI = U_{saw}(T), \quad I(0) = 0.$$

Determine the current in a series circuit containing an inductor of L Henry and a resistor of R Ohm when the circuit is driven by a saw tooth voltage of period T and amplitude V. Assume that no current is flowing initially.

The current I(t) in the circuit is governed by the differential equation

$$L\dot{I} + RI = U_{saw}(T), \quad I(0) = 0.$$

Therefore, the Laplace transform  $\bar{I}(s) = L[I(t)\colon t \to s]$  satisfies the algebraic equation

Determine the current in a series circuit containing an inductor of L Henry and a resistor of R Ohm when the circuit is driven by a saw tooth voltage of period T and amplitude V. Assume that no current is flowing initially.

The current I(t) in the circuit is governed by the differential equation

$$L\dot{I} + RI = U_{saw}(T), \quad I(0) = 0.$$

Therefore, the Laplace transform  $\bar{I}(s) = L[I(t): t \to s]$  satisfies the algebraic equation

$$(Ls+R)\overline{I}=rac{V}{T}rac{1}{s^2}\left(1-rac{sTe^{-sT}}{1-e^{-sT}}
ight).$$

Determine the current in a series circuit containing an inductor of L Henry and a resistor of R Ohm when the circuit is driven by a saw tooth voltage of period T and amplitude V. Assume that no current is flowing initially.

The current I(t) in the circuit is governed by the differential equation

$$L\dot{I} + RI = U_{saw}(T), \quad I(0) = 0.$$

Therefore, the Laplace transform  $\bar{I}(s) = L[I(t)\colon t \to s]$  satisfies the algebraic equation

$$(Ls+R)\overline{I}=\frac{V}{T}\frac{1}{s^2}\left(1-\frac{sTe^{-sT}}{1-e^{-sT}}\right).$$

Denote a = R/L. Then

Determine the current in a series circuit containing an inductor of L Henry and a resistor of R Ohm when the circuit is driven by a saw tooth voltage of period T and amplitude V. Assume that no current is flowing initially.

The current I(t) in the circuit is governed by the differential equation

$$L\dot{I} + RI = U_{saw}(T), \quad I(0) = 0.$$

Therefore, the Laplace transform  $\bar{I}(s) = L[I(t): t \to s]$  satisfies the algebraic equation

$$(Ls+R)\overline{I} = \frac{V}{T}\frac{1}{s^2}\left(1 - \frac{sTe^{-sT}}{1 - e^{-sT}}\right).$$

Denote a = R/L. Then

$$\bar{I} = \frac{V}{RT} \frac{a}{s^2(s+a)} \left( 1 - \frac{sTe^{-sT}}{1-e^{-sT}} \right) = \frac{V}{RT} \frac{1}{s^2} \left( 1 - \frac{sTe^{-sT}}{1-e^{-sT}} \right) -$$

$$-\frac{V}{RT}\frac{1}{s(s+a)}\left(1-\frac{sTe^{-sT}}{1-e^{-sT}}\right)$$

$$-\frac{V}{RT}\frac{1}{s(s+a)}\left(1-\frac{sTe^{-sT}}{1-e^{-sT}}\right)$$

Note, that the first term in the last expression is a multiple of  $L[U_{saw}(t):t \to s]$  and hence,

$$-\frac{V}{RT}\frac{1}{s(s+a)}\left(1-\frac{sTe^{-sT}}{1-e^{-sT}}\right)$$

Note, that the first term in the last expression is a multiple of  $L[U_{saw}(t):t\rightarrow s]$  and hence,

$$\bar{I} - \frac{1}{R}\bar{U}_{saw} = -\frac{VL}{R^2T}\frac{a}{s(s+a)} + \frac{V}{R}\frac{1}{s+a}\frac{e^{-sT}}{1-e^{-sT}}$$

$$-\frac{V}{RT}\frac{1}{s(s+a)}\left(1-\frac{sTe^{-sT}}{1-e^{-sT}}\right)$$

Note, that the first term in the last expression is a multiple of  $L[U_{saw}(t):t \to s]$  and hence,

$$ar{I} - rac{1}{R} ar{U}_{saw} = -rac{VL}{R^2 T} rac{a}{s(s+a)} + rac{V}{R} rac{1}{s+a} rac{e^{-sT}}{1 - e^{-sT}}$$

$$= -rac{VL}{R^2 T} rac{a}{s(s+a)} + rac{V}{R} \sum_{i=1}^{\infty} rac{e^{-skT}}{s+a}$$

Using the inverse Laplace transforms

$$L^{-1}\left[\frac{a}{s(s+a)}:s\to t\right]=1-e^{-at}$$

Using the inverse Laplace transforms

$$L^{-1}\left[\frac{a}{s(s+a)}:s\to t\right] = 1 - e^{-at}$$
 
$$L^{-1}\left[\frac{e^{-skT}}{(s+a)}:s\to t\right] = e^{-a(t-kT)}H(t-kT)$$

Using the inverse Laplace transforms

$$L^{-1}\left[\frac{a}{s(s+a)}:s\to t\right] = 1 - e^{-at}$$

$$L^{-1}\left[\frac{e^{-skT}}{(s+a)}:s\to t\right] = e^{-a(t-kT)}H(t-kT)$$

we obtain that

$$I(t) = rac{U_{saw}(t)}{R} - rac{VL}{R^2T}(1 - e^{-at}) + rac{V}{R} \sum_{k=1}^{\infty} e^{-a(t-kT)} H(t-kT)$$

### **Exercises**

- 1. Find the Laplace transform of  $f(t) = e^{-t-1/2}$  and  $f(t) = \cos(at + b)$ .
- 2. Find the inverse Laplace transform of  $\bar{f}(s) = (s-2)/(s^2-2)$ ,  $\bar{f}(s) = 3/(s^2+4s+9)$ .
- 3. Solve the convolution integral equations

A. 
$$f(t) = 1 + \int_0^t f(u) \cos(t - u) du$$
,

B. 
$$\sin t - t = \int_0^t (t - u)^2 f(u) du$$
.

4. Use Laplace transforms to solve the initial value problems

A. 
$$\ddot{y}(t) - 5\dot{y}(t) + 6y(t) = 0$$
,  $y(0) = 2$ ,  $\dot{y}(0) = 1$ ,  
B.  $\ddot{v}(t) - v(t) = te^{2t}$ ,  $v(0) = 0$ ,  $\dot{v}(0) = 1$ .

5. Using the definition of the Laplace transform, prove the second shifting property.

#### **Exercises**

6. Write f(t) as a step function and calculate its Laplace transform

$$f(t) = \begin{cases} 1 & t < 0 \\ 3 & 1 \le t < 7 \\ 5 & t \ge 7. \end{cases}$$

7. Use Laplace transforms to solve the initial value problem

$$4\ddot{y}(t)+4\dot{y}(t)+5y(t)=g(t)\,y(0)=\dot{y}(0)=0,$$
  $g(t)=\left\{egin{array}{ll} 4&0\leq t<\pi\ 0&t\geq\pi \end{array}
ight.$  and sketch the solution  $y(t)$ .

8. Solve the initial value problem

$$\ddot{y}(t) + 2\dot{y}(t) + 2y(t) = \delta(t - \pi)y(0) = \dot{y}(0) = 0,$$

by using Laplace transforms and sketch the solution y(t).

9. Determine the current in a series circuit containing an inductor of L Henry and a resistor of R Ohm when the circuit is driven by a square wave voltage of period 2T and amplitude V. No current is flowing initially.

How to solve an IVP using Laplace transforms in Matlab?

How to solve an IVP using Laplace transforms in Matlab? Here we are:

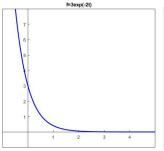
$$y''(t) + 6y'(t) + 5y(t) = f(t), \quad y(0) = 1, \ y'(0) = -1$$
  
Let  $f(t) = 3e^{-2t}$ 

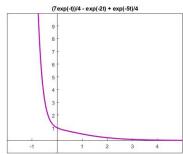
How to solve an IVP using Laplace transforms in Matlab? Here we are:

$$y''(t) + 6y'(t) + 5y(t) = f(t), \quad y(0) = 1, \, y'(0) = -1$$
 Let  $f(t) = 3e^{-2t}$ 

```
>> syms s t Y f = \exp(-2*t)*3 F = laplace(f) % define the Laplace transform of the RHS part f(t) \\ y2=Y*s.^2-1*s-(-1) % define the Laplace transform of y'' <math display="block">y1=Y*s-1 % define the Laplace transform of y' \\ Yimage = solve(y2 + 6*y1 + 5*Y-F,Y) % express the image of y ilaplace(Yimage,s,t) % find the original, i.e. the solution of the IVP
```

```
Yimage =  (s + 3/(s + 2) + 5)/(s^2 + 6*s + 5)  ans =  (7*exp(-t))/4 - exp(-2*t) + exp(-5*t)/4
```





$$f(t) = \begin{cases} \sin t & 0 < t < \pi \\ -1 & \pi < t < 2\pi \\ 0 & t > 2\pi \end{cases}$$

$$f(t) = \begin{cases} \sin t & 0 < t < \pi \\ -1 & \pi < t < 2\pi \\ 0 & t > 2\pi \end{cases}$$

$$f(t) = \sin t + H(t - \pi)(-1 - \sin t) + H(t - 2\pi)(0 - (-1))$$

$$f(t) = \begin{cases} \sin t & 0 < t < \pi \\ -1 & \pi < t < 2\pi \\ 0 & t > 2\pi \end{cases}$$

$$f(t) = \sin t + H(t - \pi)(-1 - \sin t) + H(t - 2\pi)(0 - (-1))$$

$$\Rightarrow f(t) = \sin(t) + (-1-\sin(t)) * \text{heaviside}(t-pi) + \text{heaviside}(t-2*pi)$$

$$f(t) = \begin{cases} \sin t & 0 < t < \pi \\ -1 & \pi < t < 2\pi \\ 0 & t > 2\pi \end{cases}$$
 
$$f(t) = \sin t + H(t - \pi)(-1 - \sin t) + H(t - 2\pi)(0 - (-1))$$
 
$$\Rightarrow f = \sin(t) + (-1 - \sin(t)) * \text{heaviside}(t - \text{pi}) + \text{heaviside}(t - 2*\text{pi})$$
 
$$f(t) = 1 + 2\delta(t - 1)$$

$$f(t) = \begin{cases} \sin t & 0 < t < \pi \\ -1 & \pi < t < 2\pi \\ 0 & t > 2\pi \end{cases}$$

$$f(t) = \sin t + H(t - \pi)(-1 - \sin t) + H(t - 2\pi)(0 - (-1))$$

$$>> f = \sin(t) + (-1 - \sin(t)) * \text{heaviside}(t - \text{pi}) + \text{heaviside}(t - 2*\text{pi})$$

$$f(t) = 1 + 2\delta(t - 1)$$

$$>> f = 1 + 2* \text{dirac}(t - 1)$$

$$f = 2* \text{dirac}(t - 1) + 1$$

$$>> F = \text{laplace}(f)$$

$$F = 2* \exp(-s) + 1/s$$