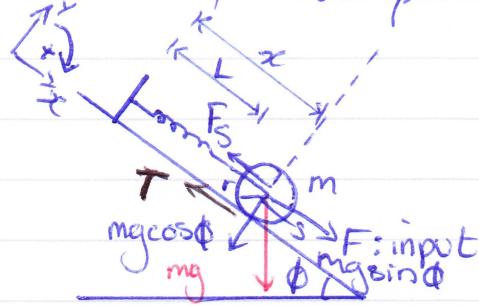


RLE 2024 Complete example: modelling to simulations.



Nonlinear spring

$$F_s = -k_1(x-L) - k_2(x-L)^3$$

$$k_1, k_2 > 0$$

No-slip condition, T causes rotation

First sketch all forces applied

x-axis forces:

$$F_s, F, mgsin\phi, T$$

→ Total force along the x-axis:

$$F + mgsin\phi - \underbrace{k_1(x-L) - k_2(x-L)^3}_{\text{unknown.}} = m\ddot{x} \quad (1)$$

→ Total force along the y-axis:

$$N - mgcos\phi = 0$$

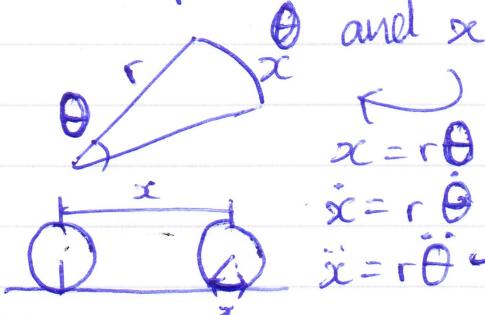


$$\underline{M_T^{(0)}} = +T \cdot r$$

$$\Leftrightarrow I\ddot{\Theta} = Tr \Leftrightarrow \frac{1}{2}mr^2\ddot{\Theta} = T$$

$$\Rightarrow \boxed{\frac{1}{2}mr\ddot{\Theta} = T}$$

relationship between Θ and x



$$\Rightarrow \frac{1}{2}m\ddot{x} = T$$

State: (x_1, x_2)
output: x
input: F

$$F + mgsin\phi - \frac{1}{2}m\ddot{x} - k_1(x-L) - k_2(x-L)^3 = m\ddot{x}$$

$$\Rightarrow \boxed{F + mgsin\phi - \cancel{k_1(x-L) - k_2(x-L)^3} = \frac{3}{2}m\ddot{x}}$$

$$F = mgsin\phi - k_1(x-L) - k_2(x-L)^3 = \frac{2}{3}m\ddot{x} \quad (3)$$

EQUILIBRIUM

multiply (3) by $\frac{2}{3m}$ $x_1 = x, x_2 = \dot{x}_1 = \dot{x}$

$$\Rightarrow \dot{x}_2 = \frac{2}{3m}F + \frac{2}{3}gsin\phi - \frac{2k_1}{3m}(x_1 - L) - \frac{2k_2}{3m}(x_1 - L)^3$$

$$\begin{cases} \dot{x}_1 = x_2 \\ \dot{x}_2 = \frac{2}{3m}F - \frac{2}{3}gsin\phi - \frac{2k_1}{3m}(x_1 - L) - \frac{2k_2}{3m}(x_1 - L)^3 \end{cases}$$

$$x = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}$$

SSR

$$\dot{x} = \begin{bmatrix} x_2 \\ \frac{2}{3m}F - \frac{2}{3}gsin\phi - \frac{2k_1}{3m}(x_1 - L) - \frac{2k_2}{3m}(x_1 - L)^3 \end{bmatrix}$$

$$\dot{x} = f(x, F)$$

→ A pair (x^{eq}, F^{eq}) is an EQUILIBRIUM POINT of the system if and only if:

$$f(x^e, F^e) = 0$$

so evaluate at equilibrium point:

x_2 is near velocity
So this makes sense

$$\begin{cases} x_2^{eq} = 0 \\ \frac{2}{3m}F^{eq} - \frac{2}{3}gsin\phi - \frac{2k_1}{3m}(x_1^{eq} - L) - \frac{2k_2}{3m}(x_1^{eq} - L)^3 = 0 \end{cases}$$

under these conditions, $(x_1^{eq}, x_2^{eq}, F^{eq})$ is an equilibrium point

$$\ddot{x}_1 = x_2$$

$$\ddot{x}_2 = \frac{2}{3m} F - \boxed{\frac{2}{3} g \sin \phi} - \frac{2k_1}{3m} (x_1 - L) - \frac{2k_2}{3m} \underline{(x_1 - L)^3}$$

constant

\downarrow
non-linear

LINEARISATION

$$g(x_i) = (x_i - L)^3$$

LINEARISE ↴

$$g'(x_1) = 3(x_1 - L)^2$$

$$g'(x, \overset{\text{ea}}{c}) = 3(x, \overset{\text{ea}}{c} - L)^2$$

$$g(x_1^{\text{eq}}) = (x_1^{\text{eq}} - L)^3$$

Taylor's approximation theorem on g at x_0

$$g(x_i) \approx g(x_i^{eq}) + g'(x_i^{eq}).(x_i - x_i^{eq})$$

for x_1 close to x_1^{eq}

$$(x_1 - L)^3 = \underbrace{(x_1^{\text{eq}} - L)^3}_{\text{constant}} + \underbrace{3(x_1^{\text{eq}} - L)^2 (x_1 - x_1^{\text{eq}})}_{\text{constant}}$$

$$\Leftrightarrow (x_1 - L)^3 - (x_1^{\text{eq}} - L)^3 = 3(x_1^{\text{eq}} - L)^2(x_1 - x_1^{\text{eq}}) \quad \text{OBSERVE THAT THIS}$$

IS A DEMONSTRATION VARIABLE

$$(-) \quad \ddot{x}_2 = \frac{2}{3m}F - \frac{2}{3}g\sin\phi - \frac{2k_1}{3m}(x_1 - L) - \frac{2k_2}{3m}(x_1 - L)^3$$

$$0 = \frac{2}{3m}F^{\text{eq}} - \frac{2}{3}g\sin\phi - \frac{2k_1}{3m}(x_1^{\text{eq}} - L) - \frac{2k_2}{3m}(x_1^{\text{eq}} - L)^3$$

$$\dot{\bar{x}}_2 = \frac{2}{3m} (\bar{F} - \bar{F}_{eq}) - \frac{2k_1}{3m} (\bar{x}_1 - \bar{x}_{eq}) - \frac{2k_2}{3m} [(\bar{x}_1 - \bar{L})^3 - (\bar{x}_{eq} - \bar{L})^3]$$

DEVIATION VARIABLE ANOTHER DEVIATION VARIABLE:
 \bar{F} \bar{x}_1

$$\Rightarrow \ddot{x}_2 = \frac{2}{3m} \bar{F} - \frac{2k_1}{3m} \bar{x}_1 - \frac{2k_2}{3m} \cdot 3(x_1^{e4} - L)^3 \cdot \cancel{(x_1 - x_2)} \cdot \bar{x}_1$$

$$\Rightarrow \ddot{\bar{x}}_2 = \frac{2}{3m} \bar{F} - \frac{2}{3m} \left(k_1 - k_2 3(x_1^{eq} - L)^3 \right) \cdot \bar{x}_1$$

$$\ddot{\bar{x}}_2 = \frac{2}{3m} \bar{F} = \frac{2}{3m} (k_1 + 3k_2 (\bar{x}_1^{\text{eq}} - 1)^3) \cdot \bar{x}_1$$

$$a \quad \boxed{\ddot{\bar{x}}_2 = a\bar{F} - b\bar{x}_1}$$

$$b \quad \ddot{\bar{x}}_1 = ?$$

(to make it simpler)

LINEARISED \Rightarrow

$$\begin{aligned} \dot{\bar{x}}_1 &= \bar{x}_2 \\ \dot{\bar{x}}_2 &= a\bar{F} - b\bar{x}_1 \end{aligned}$$

$$\begin{aligned} \dot{\bar{x}}_1 &= \bar{x}_2 \\ 0 &= \bar{x}_1 \end{aligned} \quad (-)$$

$$\begin{aligned} \dot{\bar{x}}_1 &= \bar{x}_2 - \bar{x}_1^{\text{eq}} \\ \Rightarrow \dot{\bar{x}}_1 &= \bar{x}_2 \end{aligned}$$

$$\bar{F} \rightarrow \boxed{G(s)} \rightarrow \bar{x} = \bar{x}_1$$

$$G(s) = \frac{\bar{x}(s)}{\bar{F}(s)}$$

Apply Laplace transform:

$$\begin{aligned} s\bar{x}_1 &= \bar{x}_2 \\ s\bar{x}_2 &= a\bar{F} - b\bar{x}_1 \end{aligned}$$

* plug in $s\bar{x}_1$ in $s\bar{x}_2 \rightarrow s^2\bar{x}_1 = a\bar{F} - b\bar{x}_1 \Leftrightarrow (s^2 + b)\bar{x}_1 = a\bar{F}$

in $G(s) = \frac{\bar{x}(s)}{\bar{F}(s)}$ form:

$$\left| \frac{\bar{x}_1}{\bar{F}} = \frac{a}{s^2 + b} \right| \quad \leftarrow \text{TRANSFER FUNCTION.}$$

(impulse) kick (when $\bar{F}(t) = \delta(t)$)

(step) push (when $\bar{F}(t) = A$)

(frequency) shake (when $\bar{F}(t) = A \sin(\omega t)$)

$$t \geq 0$$

or if a little more general, include 'A'

(Applications of Inverse Laplace transform)

• Impulse response

$$\text{If } \bar{F}(t) = \delta(t) \Rightarrow \bar{F}(s) = 1 \Rightarrow \bar{x}_1(s) = G(s) \cdot \bar{F}(s)$$

SIMULATIONS

$$\Rightarrow \bar{x}_1(t) = \mathcal{L}^{-1}\left\{\frac{a}{s^2+b}\right\} = \mathcal{L}^{-1}\left\{\frac{a}{s^2+\sqrt{b}^2}\right\}$$

$$= \frac{a}{\sqrt{b}} \cdot \mathcal{L}^{-1}\left\{\frac{\sqrt{b}}{s^2+\sqrt{b}^2}\right\}$$

$$= \frac{a}{\sqrt{b}} \sin(\sqrt{b} \cdot t), \quad t \geq 0$$

$$\bar{x}_1(t) = \frac{a}{\sqrt{b}} \cdot \sin(\sqrt{b} \cdot t)$$

$$\| \\ x_1(t) - x_1^{eq} = \frac{a}{\sqrt{b}} \sin(\sqrt{b} \cdot t)$$

$$\Rightarrow x_1(t) = x_1^{eq} + \frac{a}{\sqrt{b}} \cdot \sin(\sqrt{b} \cdot t).$$

• Step response

$$G(s) = \frac{a}{s^2+b}, \quad \bar{F}(t) = 1 \Rightarrow \bar{F}(s) = \frac{1}{s}$$

$$\bar{x}_1(s) = G(s) \cdot \bar{F}(s) = \frac{a}{(s^2+b)s}$$

$$\bar{x}_1(t) = \mathcal{L}^{-1}\left\{\frac{a}{(s^2+b)s}\right\}$$

↙ ↑
poles: 0, $\pm\sqrt{b} \cdot j$

PARTIAL FRACTIONS!!

$$\frac{a}{(s^2+b)s} = \frac{A_1}{s} + \frac{A_2 s + A_3}{s^2+b}$$

$$\therefore A_1 = \frac{a}{b}, \quad A_2 = -\frac{a}{b}, \quad A_3 = 0$$

$$\Rightarrow \frac{a/b}{s} + \frac{-a/b \cdot s}{s^2+b}$$

$$\mathcal{L}^{-1}\left\{\frac{a/b}{s}\right\} = \frac{a}{b}, \quad \mathcal{L}^{-1}\left\{\frac{-\frac{a}{b}s}{s^2+b}\right\} = -\frac{a}{b} \mathcal{L}^{-1}\left\{\frac{s}{s^2+b}\right\}$$

$$= -\frac{a}{b} \cdot \cos(\sqrt{b} \cdot t)$$

$$\bar{x}_1(t) = \mathcal{L}^{-1}\left\{\bar{x}_1(s)\right\}$$

$$= \frac{a}{b} - \frac{a}{b} \cdot \cos(\sqrt{b} \cdot t)$$

$$= \frac{a}{b} [1 - \cos(\sqrt{b} \cdot t)], \quad t \geq 0$$

- Frequency Response

$$\bar{F}(t) = \sin(\omega t), \quad \omega \in \mathbb{R}$$

$$\begin{aligned} \bar{F}(s) &= \frac{\omega}{s^2 + \omega^2} \\ C_1(s) &= \frac{a}{s^2 + b} \end{aligned}$$

$$\bar{x}_1(t) = \dots$$

approximation
for large t

$$\bar{x}_1(s) = \frac{\omega}{s^2 + \omega^2} \cdot \frac{a}{s^2 + b}$$

$$\downarrow \mathcal{L}^{-1}$$

- Partial fractions expansion (PFE)
- Inverse Laplace transform (ILT)

Careful!!

We have to treat the case

$\omega^2 = b$
 $w^2 = b$: because if

$$\bar{x}_1(s) = \frac{ca}{(s^2 + b)^2}$$

↑ Double pair of
imaginary poles.