

# Physics 500A – Quantum Mechanics – Homework assignment 4

Due November 6, 2019, in class

To be picked up by November 20, 2019

## Time-dependent perturbation theory [2pt]

Consider the Hamiltonian  $H(t) = H_0 + V(t)$ , where  $H_0$  does not depend on time and has eigenstates  $|n\rangle$  such that  $H_0|n\rangle = E_n|n\rangle$ . In class, we derived the basic differential equations for the coefficients  $c_n(t)$  in the expansion  $|\psi, t\rangle = \sum_n c_n(t)e^{-iE_n t}|n\rangle$ , see Sakurai eq.(5.5.15). Derive the expansion of  $c_n(t)$  in powers of  $V$ , i.e. derive Sakurai (5.6.17) [1-st edition] or (5.7.17) [2-nd edition]. This expansion is called the time-dependent perturbation theory. Sakurai's book uses the so-called "interaction picture", but you don't have to use it.

## Spin magnetic resonance [10pt]

As promised some time ago, here is the problem to work out the two-state problem with a periodic time-dependent potential (see Sakurai, section 5.5). The Hamiltonian is  $H = H_0 + V(t)$ , where  $H_0 = E_1|1\rangle\langle 1| + E_2|2\rangle\langle 2|$ , and  $V(t) = \gamma e^{i\omega t}|1\rangle\langle 2| + \gamma e^{-i\omega t}|2\rangle\langle 1|$ . The parameters  $\gamma$  and  $\omega$  are real and positive, and we can take  $E_2 > E_1$ .

- a) [3pt] Recall that transition amplitudes satisfy  $i\hbar\dot{c}_n(t) = \sum_m V_{nm}(t)\exp(i\omega_{nm}t)c_m(t)$ . Take the system to be in the lower-energy state  $|1\rangle$  at  $t=0$ . Find the probability  $P(1\rightarrow 1)$  that the system is still in state  $|1\rangle$  at  $t > 0$ . Find the probability  $P(1\rightarrow 2)$  that the system is in state  $|2\rangle$  at  $t > 0$ . (I.e., derive the answer that I wrote down in class.) Solve the equations by hand, do not use any computer programs. Check that your answers for  $c_1$  and  $c_2$  satisfy  $|c_1|^2 + |c_2|^2 = 1$ .
- b) [3pt] Now use time-dependent perturbation theory to lowest non-vanishing order to find the same probabilities  $P(1\rightarrow 1)$  and  $P(1\rightarrow 2)$ . [You will have to go to 1-st order for  $c_2$  and to 2-nd order for  $c_1$ .] Compare your answer with the exact result of part a) for small  $\gamma$ .
- c) [2pt] Take the frequency of the external perturbation to be far from the resonance,  $\hbar|\omega - \omega_{21}| \gg \gamma$ . When is the perturbation theory answer of part b) reliable? Now take the frequency of the external perturbation to be close to the resonance,  $\hbar|\omega - \omega_{21}| \ll \gamma$ . When is the perturbation theory answer of part b) reliable?
- d) [2pt] Write a short (one page) essay on how you understand the phenomenon and applications of the nuclear magnetic resonance, and what this problem has to do with NMR.

## Neutron interference in magnetic field [4pt]

Read the papers by Werner *et al.* in *Phys. Rev. Lett.* **35**, 1053 (1975), and by Rauch *et al.* in *Phys. Lett. A* **54**, 425 (1975), and the comments on them in Sakurai. Write a 1–2 page summary of how you understand the papers, and what the results are. [In

particular, explain how the phase difference comes about, and what the experiment has to do with  $2\pi$  or  $4\pi$  rotations.]

### Orbital angular momentum [8pt]

The orbital angular momentum operator in three dimensions is  $\mathbf{L} = \mathbf{x} \times \mathbf{p}$ , or in components  $L_i = \epsilon_{ijk} x_j p_k$ , where small Latin indices label  $x$ ,  $y$ , and  $z$ . It is easy to write down how the  $L_i$  operator acts on position-space wave functions in Cartesian coordinates  $\psi(x, y, z)$ , because you know how  $p_i$  acts on  $\psi(x, y, z)$ . Study Sections 3.5 and 3.6 in Sakurai if you haven't yet. For the calculations below, you can do them by hand, but also feel free to use a computer algebra program such as **Mathematica**.

- a) [2pt] Transform from Cartesian to spherical coordinates, and derive how the orbital angular momentum operator  $L_i$  acts on the position-space wave function  $\psi(r, \theta, \varphi)$  in spherical coordinates. In other words, derive equations (3.6.9), (3.6.11), (3.6.12) in Sakurai.
- b) [2pt] Find how the operator  $\mathbf{L}^2$  acts on the position-space wave function  $\psi(r, \theta, \varphi)$  in spherical coordinates. In other words, derive equation (3.6.15) in Sakurai.
- c) [4pt] At the beginning of Section 3.6 in Sakurai, there is an argument that  $\mathbf{L}$  is a generator of rotations, based on the action of  $\mathbf{L}$  on position eigenstates. Extend this argument to finite rotations for both position and momentum eigenstates. In other words, show that

$$\begin{aligned} e^{-\frac{i\phi L_z}{\hbar}} |\mathbf{x}\rangle &= |\mathbf{x}'\rangle, \\ e^{-\frac{i\phi L_z}{\hbar}} |\mathbf{p}\rangle &= |\mathbf{p}'\rangle, \end{aligned}$$

where  $x'_i = R_{ij} x_j$ ,  $p'_i = R_{ij} p_j$ , and  $R$  is the rotation matrix, see Eq.(3.1.3) in Sakurai. You may find it helpful to show first that the position operator is a generator of momentum translations.