

Radiative Processes: Homework 1

1) Derive stefan-Boltzmann law from Planck fxn,

$$F = \pi \int_0^{\infty} B_{\nu} d\nu = \sigma T^4, \text{ stefan-Boltzmann law}$$

generally Planck's law,

$$B_{\nu}(T) = \frac{8\pi h \nu^3 / c^3}{e^{h\nu/kT} - 1}$$

$$\begin{aligned} F &= \pi \int_0^{\infty} B_{\nu} d\nu = \pi \int_0^{\infty} \frac{8\pi h \nu^3 / c^3}{e^{h\nu/kT} - 1} d\nu \\ &= \frac{8\pi h}{c^3} \int_0^{\infty} \frac{\nu^3}{e^{h\nu/kT} - 1} d\nu \end{aligned}$$

$$\begin{aligned} \text{Let } x &= h\nu/kT \rightarrow dx = h/kT d\nu \\ &\rightarrow d\nu = kT/h dx \end{aligned}$$

applying change of variables,

$$\begin{aligned} F &= \frac{8\pi h}{c^3} \int_0^{\infty} \left(\frac{x kT}{h} \right)^3 dx \cdot \frac{kT}{h} \\ &= \frac{8\pi h}{c^3} \left(\frac{kT}{h} \right)^4 \int_0^{\infty} \frac{x^3}{e^x - 1} dx \\ &\quad \quad \quad = \pi^4/15 \text{ (from hint)} \end{aligned}$$

$$\begin{aligned} \Rightarrow F &= \frac{8\pi h}{c^3} \left(\frac{kT}{h} \right)^4 \cdot \frac{\pi^4}{15} = \frac{8\pi^5 k (kT)^4}{15 c^3 h^4} \\ &= \frac{8\pi^5 (kT)^4}{15 c^3} \end{aligned} \rightarrow$$

and the Stephan-Boltzmann constant is,

$$\sigma_{SB} \equiv \frac{8\pi^5 k^4}{15 c^3}$$

$$\Rightarrow \boxed{F = \sigma_{SB} T^4}$$

2) Plot the Planck f(x)

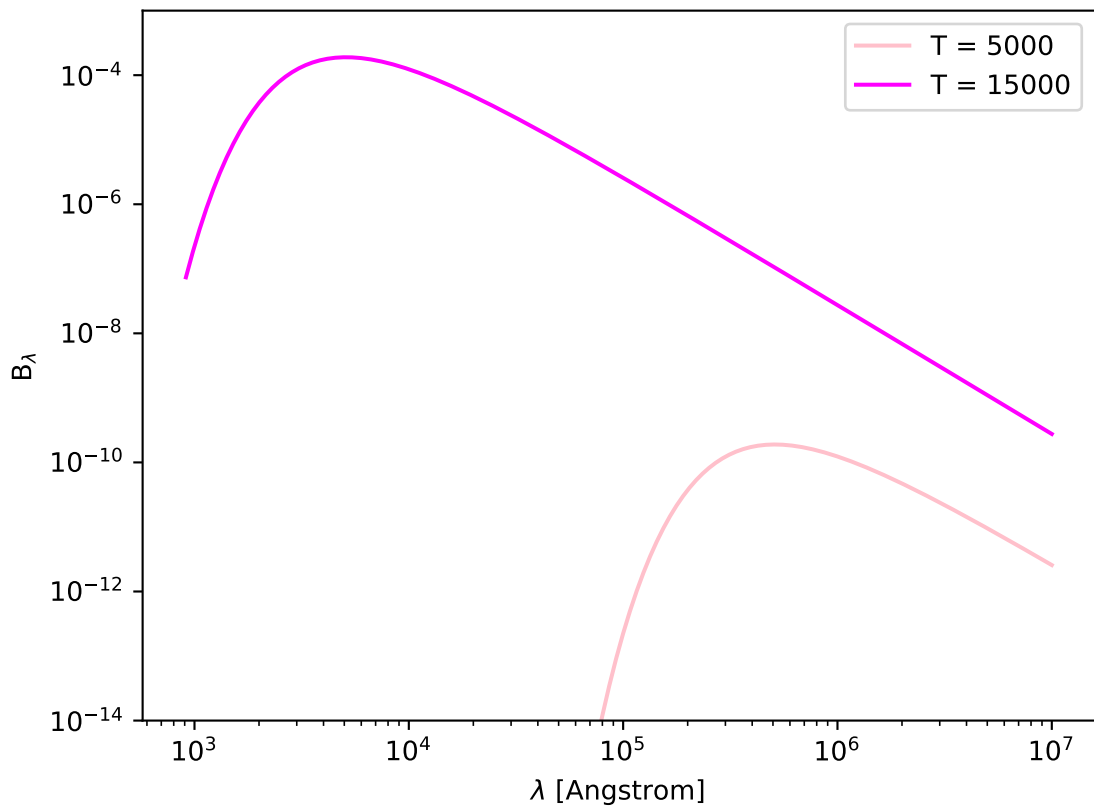
a. $T = 100 \text{ K}$, warm interstellar dust

b. $T = 10000 \text{ K}$, hot star

For the hot star, B_ν peaks around 10^{-4} . But for the warm interstellar dust B_ν peaks much lower $\sim 10^{-10}$.

(see plots on next page)

#2



3) $\lambda_{\min} = ?$ (series limits)

a. Lyman ($n=1$)

b. Balmer ($n=2$)

c. Paschen ($n=3$)

$$\Delta E = -13.6 \left(\frac{1}{n_2^2} - \frac{1}{n_1^2} \right)$$

$$c = \lambda \nu \quad \text{and} \quad E = h \nu \quad \text{and} \quad E = \frac{hc}{\lambda}$$

It can be helpful to write this equation in terms of Rydberg constant,

$$R_H = \frac{13.6}{hc \cdot Z^2} \sim 109678 \text{ cm}^{-1} = 1.09 \times 10^5 \text{ cm}^{-1}$$

$$\Rightarrow E = \frac{1}{\lambda} = R \left(\frac{1}{n_2^2} - \frac{1}{n_1^2} \right), \text{ Rydberg formula}$$

a. Lyman series

$$n_1 = 2, \quad n_2 = 1$$

$$\frac{1}{\lambda} = R \left(\frac{1}{1} - \frac{1}{4} \right) = \frac{3}{4} R$$

$$\Rightarrow \lambda = \frac{4}{3R} = 1.2 \times 10^{-5} \text{ cm}$$

$$\lambda = 120 \text{ nm}$$

→ Vacuum UV
(EUV)



b. Balmer series

$$n_1 = 3, n_2 = 2$$

$$\frac{1}{\lambda} = R \left(\frac{1}{4} - \frac{1}{9} \right) = \frac{5R}{36}$$

$$\lambda = \frac{36}{5R} = 6.61 \times 10^{-5} \text{ cm}$$

$$\lambda = 661 \text{ nm} \quad \text{Optical}$$

c. Paschen series

$$n_1 = 4, n_2 = 3$$

$$\frac{1}{\lambda} = R \left(\frac{1}{9} - \frac{1}{16} \right) = \frac{7R}{144}$$

$$\lambda = \frac{144}{7R} = 1.89 \times 10^{-4} \text{ cm}$$

$$\lambda = 1890 \text{ nm} \quad \text{near infrared}$$

$= 0.00189 \text{ mm}$

4) pure He stellar atmosphere

$T = ?$, where half of the He I atoms have been ionized

$P_e = 200 \text{ dyne} \cdot \text{cm}^{-2}$, constant electron pressure

$$X_i = 24.6 \text{ eV} = 3.9 \times 10^{-11} \text{ erg}$$

$$X_{ii} = 54.4 \text{ eV} = 8.7 \times 10^{-11} \text{ erg}$$

$$U_i = 1$$

$$U_{ii} = 2$$

$$U_{iii} = 1$$

} partition fns

a. $\frac{N_{ii}}{N_i} = ?$ and $\frac{N_{iii}}{N_{ii}} = ?$

a. $T = 5000 \text{ K}$

b. $T = 15000 \text{ K}$

c. $T = 25000 \text{ K}$

The Saha eqn,

$$\frac{N_+}{N_i} = \frac{2kT}{P_e} \frac{g_+}{g_i} \left(\frac{2\pi m_e kT}{h^2} \right)^{3/2} e^{-X_i/kT}$$

So generally,

$$\frac{N_{ii}}{N_i} = \frac{2kT}{P_e} \frac{U_{ii}}{U_i} \left(\frac{2\pi m_e kT}{h^2} \right)^{3/2} e^{-X_i/kT}$$

$$\frac{N_{iii}}{N_{ii}} = \frac{2kT}{P_e} \frac{U_{iii}}{U_{ii}} \left(\frac{2\pi m_e kT}{h^2} \right)^{3/2} e^{-X_{ii}/kT} \rightarrow$$

now just plug in numbers,
for $T = 5000 \text{ K}$,

$$\frac{N_{11}}{N_1} = \frac{2k(5000)}{Pe} \frac{1}{2} \left(\frac{2\pi m_e(5000)}{h^2} \right)^{3/2} e^{-X_{11}/k(5000)}$$

$$\Rightarrow \frac{N_{11}}{N_1} = 3.62 \times 10^{-18}, \quad T = 5000 \text{ K}$$

$$\frac{N_{111}}{N_{11}} = \frac{2k(5000)}{Pe} \frac{1}{2} \left(\frac{2\pi m_e(5000)}{h^2} \right)^{3/2} e^{-X_{111}/k(5000)}$$

$$\Rightarrow \frac{N_{111}}{N_{11}} = 5.74 \times 10^{-49}, \quad T = 5000 \text{ K}$$

for $T = 15000 \text{ K}$,

I followed a similar process as above to get,

$$\frac{N_{11}}{N_1} = 1.22$$

$$\text{for } T = 15000 \text{ K}$$

$$\frac{N_{111}}{N_{11}} = 2.62 \times 10^{-11}$$

for $T = 25000 \text{ K}$,

$$\frac{N_{11}}{N_1} = 8160.38$$

$$\text{for } T = 25000 \text{ K}$$

$$\frac{N_{111}}{N_{11}} = 1.86 \times 10^{-3}$$

$$b. \frac{N_{II}}{N_{TOT}} = \frac{N_{II}}{(N_I + N_{II} + N_{III})} \quad \text{total}$$

express in terms of $\frac{N_{II}}{N_I}$ and $\frac{N_{III}}{N_{II}}$

Total number density,

$$N_{TOT} = N_I + N_{II} + N_{III}$$

$$\text{let's say, } N_{II} = \left(\frac{N_{II}}{N_I} \right) \cdot N_I$$

$$N_{III} = \left(\frac{N_{III}}{N_{II}} \right) \cdot N_{II}$$

Substitute into N_{TOT} ,

$$N_{TOT} = N_I + \left(\frac{N_{II}}{N_I} \right) \cdot N_I + \left(\frac{N_{III}}{N_{II}} \right) N_{II}$$

$$N_{TOT} = N_I + \left(\frac{N_{II}}{N_I} \right) N_I + \left(\frac{N_{III}}{N_{II}} \right) \left(\frac{N_{II}}{N_I} \right) N_I$$

$$N_{TOT} = N_I \left[1 + \frac{N_{II}}{N_I} + \left(\frac{N_{III}}{N_{II}} \right) \left(\frac{N_{II}}{N_I} \right) \right]$$

$$\Rightarrow \frac{N_{II}}{N_{TOT}} = \frac{N_{II}}{N_I} \cdot \frac{N_I}{N_I + \frac{N_{II}}{N_I} N_I + \left(\frac{N_{III}}{N_{II}} \right) \left(\frac{N_{II}}{N_I} \right) N_I} = \frac{N_{II}}{N_I} \cdot \frac{1}{\left[1 + \frac{N_{II}}{N_I} + \left(\frac{N_{III}}{N_{II}} \right) \left(\frac{N_{II}}{N_I} \right) \right]}$$

c. Plot $\frac{N_{II}}{N_{TOT}} (T)$

→ See plot on next page.

He I partial ionization zone
when $\frac{N_{II}}{N_{TOT}} = 0.5$

$\sim 14850 \text{ K}$

⇒ a white dwarf's
atmosphere is at
least partially ionized.

#4 part c

