shallow_water01.f90 – semi-implicit one-dimensional shallow water equations

The model uses Arakawa C-grid staggering. Linearized equation of motion. Bottom friction.

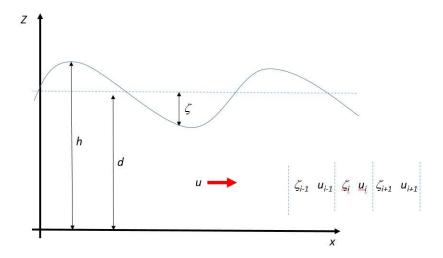


Fig 1: Geometry of the 1-D shallow water problem. Disturbed water depth h, un-disturbed water depth d, sea surface elevation ζ , horizontal velocity u. The lower right corner shows the C-grid staggering.

The momentum equation is formulated on the old (n) and new time level (n+1). The weight between both levels is w_{imp} .

$$u^{n+1} = u^n - \Delta t \, c_D \, \frac{1}{h} \, u^n |u^n| - \Delta t \, g \, \left(1 - w_{imp}\right) \frac{\partial \zeta^n}{\partial x} - \Delta t \, g \, w_{imp} \, \frac{\partial \zeta^{n+1}}{\partial x} \tag{1}$$

The same is valid for the equation of continuity:

$$\zeta^{n+1} = \zeta^n - \Delta t \left(1 - w_{imp} \right) \frac{\partial (h \, u)^n}{\partial x} - \Delta t \, w_{imp} \frac{\partial (h \, u)^{n+1}}{\partial x} \tag{2}$$

with

$$h = d + \zeta \tag{3}$$

Our strategy is to form a system of equations for all elements of ζ^{n+1} , solve this system iteratively and finally compute the velocity at the new time step u^{n+1} .

Inserting equation (1) into (2) yields:

$$\zeta^{n+1} = \zeta^n - \Delta t \left(1 - w_{imp} \right) \frac{\partial (h \, u)^n}{\partial x} - \Delta t \, w_{imp} \frac{\partial (h^n \, u^*)}{\partial x} + \Delta t^2 \, w_{imp}^2 \, g \, h \, \frac{\partial^2 \zeta^{n+1}}{\partial x^2} \tag{4}$$

with the interim solution

$$u^* = u^n - \Delta t \ c_D \ \frac{1}{h} \ u^n |u^n| - \Delta t \ g \left(1 - w_{imp}\right) \frac{\partial \zeta^n}{\partial x} \tag{5}$$

If we define

$$div = -\Delta t \left(1 - w_{imp}\right) \frac{\partial (h \, u)^n}{\partial x} - \Delta t \, w_{imp} \frac{\partial (h^n \, u^*)}{\partial x} \tag{6}$$

then it follows

$$\zeta_i^{n+1} = \frac{1}{1 + c_E + c_W} \left[\zeta^n + div + c_E \zeta_{i+1}^{n+1} + c_W \zeta_{i-1}^{n+1} \right] \tag{7}$$

with

$$c_E = \begin{cases} \frac{\Delta t^2 w_{imp}^2 g h_E}{\Delta x^2} & \text{if } h_{i+1} > 0\\ 0. & \text{otherwise} \end{cases}$$
 (8)

$$c_W = \begin{cases} \frac{\Delta t^2 w_{imp}^2 g h_W}{\Delta x^2} & \text{if } h_{i-1} > 0\\ 0. & \text{otherwise} \end{cases}$$
 (9)

with

$$h_E = 0.5 (h_i + h_{i+1}) (10)$$

$$h_W = 0.5 (h_i + h_{i-1}) (11)$$

i.e. the water depth at the velocity point u_i and u_{i-1} .

After the system of equations (6) is solved iteratively with the Gauss-Seidel-Method, the velocity at the new time level is computed with

$$u^{n+1} = u^* - \Delta t \ g \ w_{imp} \frac{\partial \zeta^{n+1}}{\partial x}$$
 (12)

Reference:

Backhaus JO (1983) A semi-implicit scheme for the shallow water equations for application to shelf sea modelling. Continental Shelf Research, Vol. 2, No. 4, pp. 243 - 254. https://doi.org/10.1016/0278-4343(82)90020-6