### ADVANCED REVIEW



# Recent advances in quantum-mechanical molecular dynamics simulations of proton transfer mechanism in various water-based environments

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#### **Abstract**

Proton transfer in water-based environments occurs because of hydrogen-bond interaction. There are many interesting physicochemical phenomena in this field, causing fast structural diffusion of hydronium and hydroxide ions. During the last few decades, to support experimental observations and measurements, quantum-mechanical molecular dynamics (QMMD) simulations with reasonable accuracy and efficiency have significantly unraveled structural, energetic, and dynamical properties of excess proton in aqueous environments. This review summarizes the state-of-the-art QMMD studies of proton transfer processes in aqueous solutions and complex systems including bulk liquid water, ice phases, and confined water in nanochannel/nanoporous materials as well as reports on CO<sub>2</sub> scrubbing by amine-based chemical absorption.

This article is categorized under:

Structure and Mechanism > Reaction Mechanisms and Catalysis

Molecular and Statistical Mechanics > Molecular Dynamics and Monte-Carlo

Methods

Electronic Structure Theory > Semiempirical Electronic Structure Methods Theoretical and Physical Chemistry > Reaction Dynamics and Kinetics

#### KEYWORDS

density functional tight binding method, divide and conquer method, hydrogen bond, proton transfer, quantum-mechanical molecular dynamics

# 1 | INTRODUCTION

Proton transfer (PT) is the crucial trigger for functionalization of various materials and biological systems.<sup>1–5</sup> The PT process has received considerable attention at the fundamental level in order to interpret the remarkably fast proton diffusivity in the hydrogen-bond (H-bond) network. Since the pioneering description of proton hopping was proposed more than 200 years ago,<sup>6</sup> which is referred to as structural diffusion or the Grotthuss mechanism, the contribution of water's ionic ingredients—hydronium (H<sub>3</sub>O<sup>+</sup>) and hydroxide (OH<sup>-</sup>) ions—to PT has been widely investigated.<sup>7–12</sup> Owing to a wide variety of chemical environments and the insufficient and controversial information about molecular mechanisms, gaining deeper insights on PT events is required.

Molecular dynamics (MD) simulation is an effective computational tool that enables tracking of atomic/molecular motions and analyzing the associated microscopic properties with qualitative and quantitative accuracy. With respect to PT, theoretical

information has been utilized for complementing experimental research, in which the issues related to measurement methods are difficult to resolve. Simulation of PT events requires MD techniques with several capabilities including appropriate treatment of covalent bond formation/cleavage for structural diffusion, reasonable computational efficiency for sufficient sampling, and concurrency of multiple proton hopping in a physically rational model system.

The parameterized reactive force field (RFF)<sup>13–17</sup> is one of the practical approaches for describing chemical reactions in large molecular systems. The RFF potentials specifically adjusted to study hydrated protons have been developed and applied to MD simulation of an excess proton in a water system.<sup>18–24</sup> Among the RFF approaches, the most popular method for revealing mechanistic details and elucidating structural and dynamical properties of the PT process is the multiple empirical valence bond (MS-EVB).<sup>25–27</sup> The application of MS-EVB to examine proton dynamics has been investigated in diverse model systems ranging from  $H_3O^+$  <sup>28–45</sup> and  $OH^-$  <sup>41,46</sup> diffusion in bulk water, water–ice interface,<sup>47</sup> and confined water environment task before starting MD simulations, thereby limiting the applicability of MS-EVB over a broad range of systems.

In this review, we focus on the recent progress in computational studies related to the PT process by ab initio molecular dynamics (AIMD).<sup>53–55</sup> The so-called Born–Oppenheimer molecular dynamics (BOMD) and Car–Parrinello MD<sup>56</sup> are two common approaches in AIMD, in which density functional theory (DFT)<sup>57–59</sup> is typically adopted as a solver of the electronic structure problem. Although the simulations of large systems are limited because of expensive computational requirements, AIMD is highly appreciated in terms of general applicability across the periodic table and explicit incorporation of polarization and charge transfer effects.

The numerous applications of AIMD have been widely conducted and reported in order to clarify the diffusion of H<sub>3</sub>O<sup>+</sup> and OH<sup>-</sup> ions in aqueous solutions. <sup>9-12,53-55</sup> Instead of repeating reported literature, this review emphases on some advances and challenges in the simulation of water's ionic species in bulk systems utilizing either sophisticated or approximated DFT. To capture the dynamics of ubiquitous PT processes, a full QM method is essential, thus MD simulations with a hybrid quantum mechanical/molecular mechanics (QM/MM) approach are not discussed in the present review. Subsequently, we review theoretical results of PT events in ice for emphasizing the difference in diffusion behavior of liquid and solid phases. Following the discussion about structural and dynamical characteristics of the proton in a system constituting only water molecules, the importance of PT in more complex systems is presented. As representative examples, we briefly account for the features of proton dynamics of water confined in nanosized hydrophobic channels and the key roles of the proton in the chemistry of CO<sub>2</sub> capture.

# 2 | HYDRONIUM AND HYDROXIDE ION DIFFUSIONS IN BULK WATER

After the initial pilot studies by Tuckerman et al., $^{60-62}$  many fundamental aspects of  $H_3O^+$  and  $OH^-$  diffusion in the liquid phase of water have been accumulated by DFT-based AIMD simulations. The following are notable insights about the molecular mechanism involving the  $H_3O^+$  ion: solvation structure of Zundel  $(H_5O_2^+)^{63}$  and Eigen  $(H_9O_4^+)^{64}$  cation species, $^{65-67}$  validation of the PT process through interconversion of the two above-mentioned protonated water complexes, $^{7.68-70}$  concept of proton rattling $^{65,71-73}$  often termed as special pair dance, $^{70}$  and coupling of multiple PT events in the concerted manner. $^{74,75}$  The corresponding mechanisms for  $OH^-$  species include coordination patterns of the hydrated structure  $^{71,76-80}$  and plausibility of the mechanism with a dynamical interplay between three- and four-fold structure motifs $^{10,11,71,76-82}$  in comparison to other mechanisms supposing mirror image or static solvation complexes individually. $^{83-85}$  Moreover, the influence of nuclear quantum effects on diffusion properties have been addressed. $^{68,76,86-90}$ 

One important direction of the study on PT in aqueous solution is to describe accurate trajectories with high-level electronic structure methods. Several benchmarks of pure liquid water for DFT functionals beyond the commonly employed generalized gradient approximation (GGA) and for wave-function-based approaches<sup>91–99</sup> can be explained by their performance on describing a quantitative illustration of PT events because the H-bond structure is closely linked with the structural diffusion. Currently, the highest quality proton dynamics is obtained at the PBE0-TS level,<sup>100</sup> namely the hybrid PBE0 functional combined with the Tkatchenko–Scheffler (TS) dispersion correction scheme. Compensating deficiencies of GGA functionals for relatively poor molecular polarizability, H-bond strength, and underestimation tendency of liquid density attributed to the lack of long-range van der Waals forces<sup>91,104,105</sup> leads to a satisfactory agreement of calculated diffusion coefficients with experimental ones, as shown in Table 1. Furthermore, the simulation has provided new perspectives that approximately halved OH<sup>-</sup> diffusion speed with respect to that of H<sub>3</sub>O<sup>+</sup> is caused by unfavorable correlation of the fourfold

**TABLE 1**  $H_3O^+$  and  $OH^-$  diffusion constants (in  $\mathring{A}^2/ps$ ) and their ratios estimated from various AIMD simulations

Method	Basis set/parameter	$N_{ m water}$	$D(\mathrm{H_3O}^+)$	$D(\mathrm{OH}^-)$	$D(\mathrm{H_3O^+})/D(\mathrm{OH^-})$
PW91 <sup>106</sup>	Plane wave	32	0.32 <sup>a</sup>	1.85 <sup>a</sup>	0.18
BLYP <sup>107,108</sup>	Plane wave	32	0.28 <sup>a</sup>	$0.19^{a}$	1.47
HCTH <sup>109</sup>	Plane wave	32	$0.33^{a}$	$0.04^{a}$	7.39
PBE <sup>101</sup>	Plane wave	64	1.08 <sup>b</sup>	1.82 <sup>b</sup>	0.59
PBE-TS <sup>101,103</sup>	Plane wave	64	1.28 <sup>b</sup>	$0.83^{b}$	1.54
PBE0-TS <sup>101-103</sup>	Plane wave	64	0.83 <sup>b</sup>	0.37 <sup>b</sup>	2.24
DFTB <sup>110,111</sup>	mio, c mio-IBI-BLYPd	523	0.91°	0.67 <sup>d</sup>	1.36
Experimental (H <sub>2</sub> O)			0.94 <sup>e</sup>	0.52 <sup>e</sup>	1.80 <sup>e</sup>
Experimental (D <sub>2</sub> O)			0.67 <sup>e</sup>	0.31 <sup>e</sup>	2.15 <sup>e</sup>

Abbreviation: BLYP, Becke-Lee-Yang-Parr.

Experimental data are shown for comparison.

hypercoordinated solvation structure with simultaneous proton hopping. Note that encouragement of the concerted PT process in an aqueous environment is recently revisited. 116

Despite a good agreement of the estimated diffusion constants, one should also note that the use of plane wave basis may also lead to various discrepancies, namely, overstructured radial distribution functions (RDFs) that leads to a more rigid H-bond network and slow water diffusivity. Earlier study also pointed out the needs of a larger basis set TZV2P when the meta-GGA functionals were employed for improving the RDFs and self-diffusion of water. 94,118

The other intriguing progress is to resolve the issue of system size by using computationally cost-effective quantum chemical methods. The implementation of a scalable algorithm for DFT-based AIMD simulations has been proposed as a solution. <sup>94,117,118</sup> Embracing computational models that have the performance in between the DFT and RFF can be an elective choice. The widely tested method for QMMD simulations of the condensed phase system is the density functional tight-binding (DFTB) method, <sup>119–121</sup> which is derived from DFT with careful approximations. The DFTB method satisfies both accuracy and efficiency with elaborated parameterizations. Previous studies have reported the imperfect water interactions in the bulk condition, <sup>122–126</sup> theoretical extensions for the improved description of the H-bond, <sup>110,111,127–134</sup> and the specific parameter optimization, <sup>113,135,136</sup> which may improve the performance of the DFTB method for conducting extensive sampling of PT events in aqueous solutions. To facilitate the large-scale BOMD simulations, the DFTB method may be combined with linear-scaling techniques such as the fragment molecular orbital method <sup>137–139</sup> and the divide-and-conquer (DC) method <sup>140–145</sup> or with some sparse algebraic schemes. <sup>146,147</sup>

The DC-type DFTB-based MD simulations have been lately applied to proton dynamics in bulk water for a system containing a large number of atoms. <sup>112,113</sup> An interesting observation for H<sub>3</sub>O<sup>+</sup> diffusion is the consistent description of both vehicular and structural diffusion constants relative to experimental ones <sup>112</sup> (see Table 1 for overall diffusion constants defined as the sum of two terms) as opposed to the constants calculated by estimation from smaller unit cell sizes. <sup>123,125</sup> This is likely a manifestation of the existence of an artifact in the small model system, namely the strengthened H-bond network introduced by the contributions from periodic images. In Reference 112, not only dynamical properties but also the proton diffusion barrier have agreed with the experimental values, exhibiting less than 0.5 kJ/mol error. Regarding OH<sup>-</sup> diffusion, the DFTB potential fitted to structural properties at the Becke-Lee-Yang-Parr level has also resulted in qualitative production of dynamical and energetic properties. <sup>113</sup> Moreover, the simulation has confirmed the decrease of the proton diffusion barrier with respect to the less-coordinated solvation structures. This observation supports the widely accepted PT mechanism. <sup>10,11,71,76–82</sup> Moreover, these simulation results indicate the importance of treating large and realistic systems. A larger system is beneficial to minimize unphysical mirror image interactions. In addition, it was shown that adopting a large system size will reduce the temperature fluctuations during the equilibration and production run. <sup>112</sup>

<sup>&</sup>lt;sup>a</sup>Reference 10 (production runs: 20–50 ps).

<sup>&</sup>lt;sup>b</sup>Reference 100 (production runs: 14–55 ps).

<sup>&</sup>lt;sup>c</sup>Reference 112 (production runs: 20 ps).

<sup>&</sup>lt;sup>d</sup>Reference 113 (production runs: 40 ps).

eReference 114.

# 3 | PROTON TRANSFER IN ICE PHASES

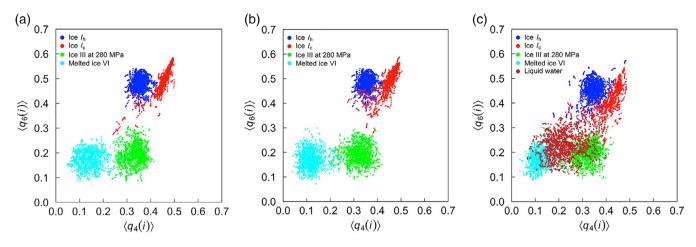
The application of AIMD to the PT process in ice, either crystalline or amorphous forms of water,  $^{148,149}$  can be straightforwardly motivated once its reasonable performance has been achieved in bulk water. In such solid systems, the presence of an excess proton or a proton hole is represented as an ionic defect, which becomes the origin of pure structural diffusion under the condition that the oxygen atom positions rarely changes. For example, experimental measurements have shown proton diffusivity in an ordinary hexagonal (ice  $I_h$ ) or metastable cubic (ice  $I_c$ ) form of ice approximately 10 times faster than that in the liquid phase at the same temperature.  $^{64,150-156}$  Despite several explanations have been proposed,  $^{157-163}$  the investigation of the anomalously fast PT event in the ice phases remains an active research field.

Similar to the case of bulk water, the description of H-bond as accurate as possible in ice phases is critical. The static properties of various polymorph ice structures, including equilibrium crystal structure, binding energy, and bulk modulus, have been assessed for different levels of quantum chemical methods. 164–169 One crucial conclusion from the comparison among DFT functionals is that long-range dispersion correction benefits the improved performance. 164,166,168,169 Popular PBE functional was found to underestimate the sublimation energy of high-pressure ice by 15% compared to the experiment. The Sufficiently large basis sets ("tier2" for H and "tier3" for O) are required since it was shown to significantly affect the calculated lattice energies. Perdew and coworkers show that the meta-GGA functionals, namely revised Tao-Perdew-Staroverov-Scuseria and meta-generalized gradient approximation made simple (MGGA-MS) yield a better description on the sublimation energy and equilibrium volume of low-pressure phase of ice, even without employing the dispersion correction. The dispersion correction is necessary for the GGA functionals to improve their accuracy in describing these properties. The Despite the MGGA-MS underestimates the sublimation energy of high-pressure phases of ice, the dispersion correction was confirmed to improve the description for ambient-pressure phases of ice. The extensive benchmark has confirmed the adequate accuracy of dispersion-corrected DFTB for low-density phases. These results support the application of AIMD to proton dynamics.

Beside the numerical assessments, a number of first-principles studies have been performed for the ice system in the scope of vibrational frequency analyses  $^{172-182}$  and the PT phenomena.  $^{183-186}$  As for the latter, Ohmine and co-workers have determined the potential energy surfaces of PT in proton-ordered/disordered ice structures using the QM/MM approach. An important highlight is that the distances between oxygen atoms in the excess proton moiety are shortened by two types of forces—an attractive force due to the presence of the excess proton in the middle of two oxygen atoms (O—H<sup>+</sup>—O) and a repulsive force from the fourth-coordinating water molecule on the protonated water molecule. The existing local structure of water molecules around the  $H_3O^+$  ion strictly differs from that of the liquid water, in which the term of Eigen or Zundel no longer exists. Due to the environment constraining the oxygen distances, the formation of the Zundel species is prevented. The absence of Zundel–Eigen interchange during the transfer leads to concerted PT events, which then enhance the PT rate. Moreover, the Zundel form may exist in the liquid water due to the absence of the repulsive force from the fourth-coordinating water molecule.

For elucidating the relationship between the confined structure in ice phases and proton diffusivity, the BOMD simulations have been recently carried out for ice  $I_h$  and ice  $I_c$  with low density as well as two high-density structures (ice III and melted ice VI) at the DFTB level. <sup>183</sup> The averaged local bond order parameters, <sup>187</sup> which quantify the degree of cubic ( $\langle q_4 \rangle$ ) and hexagonal ( $\langle q_6 \rangle$ ) symmetries in the crystal structure, have been employed to extract the regularity of water molecule configuration. Figure 1 shows the correlation between  $\langle q_4 \rangle$  and  $\langle q_6 \rangle$  values, in which the clustered points clearly distinguish the different ice crystal structures. Importantly, the low- and high-density variants are distinguishable by hexagonality. The calculated structural diffusion constants at several temperatures are summarized in Table 2. The proton diffusivity in ice  $I_h$  and ice  $I_c$  is found to be faster than that in ice III and melted ice VI. The evaluation of structural and dynamical properties indicates that a higher hexagonal symmetry inducing delocalization of a positive charge is likely to be an essential factor for the fast PT event. In addition, the estimated proton diffusion barrier was only 0.5 kJ/mol higher than the experimental value, establishing the validity of DFTB simulations. <sup>183,188</sup>

On the other side of theoretical analysis using relatively low quantum chemical methods, the calibration of low-temperature AIMD simulations by nuclear quantum effects has been examined. The recent study by Drechsel-Grau and Marx has demonstrated collective proton tunneling in a six-membered ring structure of ice  $I_h$ , which is schematically represented in Figure 2. According to the AIMD approach combined with the path integral technique, seeing the first PT of H<sup>(1)</sup> to the adjacent oxygen atom O<sup>(2)</sup> is observed followed by the concerted movements of five other protons H<sup>(2)</sup>–H<sup>(6)</sup>. The energy barrier for the collective PT event is independent of the local structure in terms of the free energy profile (Figure 3) and the similarity of the calculated properties along the PT coordinates between chair and boat conformations. Note that neither many-



**FIGURE 1** Correlation between the cubic ( $\langle q_4(i) \rangle$ ) and hexagonal ( $\langle q_6(i) \rangle$ ) symmetries for ice structures and liquid water estimated at (a) 230, (b) 250, and (c) 270 K. (Reprinted with permission from Reference 183. Copyright 2018 American Chemical Society)

**TABLE 2** Estimated structural diffusion constants (in  $\text{Å}^2/\text{ps}$ ) for ice structures based on MD simulations at the DFTB level. (Reprinted with permission from Reference 183. Copyright 2018 American Chemical Society)

T(K)	Ice I <sub>h</sub>	Ice I <sub>c</sub>	Ice III	Melted ice VI
50	0.10	0.00		
70	0.36	0.32		
120	0.94	0.72		
230	1.84	2.04	0.25	0.36
250	1.98	1.14	0.48	1.00
270	1.90	2.28	0.54	1.21

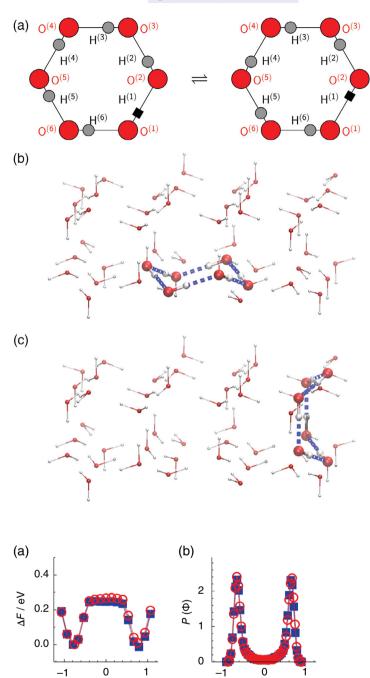
body proton tunneling nor multiple PT processes have been achieved at temperatures less than  $100~\mathrm{K}$  in the aforementioned DFTB simulations.  $^{178}$ 

In addition to the fundamental aspects of the PT process in the ice phase, recent AIMD studies successfully address other related issues such as chemical reactions on top of the ice surface,  $^{192-195}$  ionization dynamics of the water dimer on the ice surface,  $^{196-198}$  and ionization of acids on the quasi-liquid layer of ice.  $^{199}$  Moreover, the PT event under the presence of an electric field has been investigated, in which the enhanced coupling between dipole moments and internal electric field leads to a faster diffusion in a more ferroelectric phase than that in ice  $I_h$ . Elucidation of the PT process in ice from manifold theoretical viewpoints will enrich the information about H-bond systems.

# 4 | PROTON TRANSFER IN CONFINED WATER ENVIRONMENT

A confined water environment is a specific system created by the surrounding material in which water molecules are placed in a nanometer-sized vessel. The confined water is found in many materials including cells, <sup>201–203</sup> macromolecules, <sup>204–206</sup> supramolecule structures, <sup>207,208</sup> and gels. <sup>209–212</sup> Such a confined environment may lead to the phase transformation that is not present in the bulk condition under the same thermodynamical condition. <sup>213–215</sup> Moreover, the confined environment causes both slow and fast dynamics of water molecules, thus garnering significant attention from many research disciplines. <sup>216,217</sup> The slow dynamics of water in the region close to the hydrophilic chains of protein has been comprehensively reviewed by Bagchi<sup>218</sup> and theoretical research on MD simulations performed for water at the protein–solvent interface has been conducted. <sup>219</sup> It is noteworthy that the confinement results change with several properties, such as highly anisotropic orientational dynamics in carbon nanostructures, <sup>220</sup> dipole orientation in a narrow nanotube, <sup>221</sup> and liquid–liquid transition. <sup>222</sup>

The PT process in the confined water environment is one of the fascinating topics in AIMD. For example, Dellago et al. have reported fast proton mobility in a carbon nanotube (CNT) via both AIMD and RFF simulations. The reported proton diffusion constant of  $17 \text{ Å}^2/\text{ps}$  is 18.1 times higher than the experimentally estimated diffusion constant. Another AIMD



(d)

 $C(\tau, \phi_{_1}) / \mathring{A}$ 

0

0.0

Фс / Å

0.4

Δ/Å

0.6

0.2

(c)

*P*(∆)

8

6

4 2

0

0.0

**FIGURE 2** (a) Schematic representation of collective PT events in a six-membered ring of ice  $I_h$  structure. Red and gray circles represent the oxygen and hydrogen atoms, respectively, whereas the black square represents the deuteron in the case of the partially deuterated systems. The numbers in parentheses describe the label of the H-bonds. (b) and (c) display the chair and boat conformations existing in the proton-ordered ice  $I_h$  structure. (Reprinted with permission from Reference 160. Copyright 2017 Royal Society of Chemistry)

**FIGURE 3** Comparison of the collective PT in chair (filled blue squares) and boat (empty red circles) conformations of the proton-ordered six-membered ring structure. (a) Free energy profile along the constraint PT coordinate, (b) probability distribution of the collective PT coordinate defined as  $\phi = \sum_{i=1}^{6} \phi_i/6$ , where  $\phi_i = d(O^{(i)}H^{(i)}) - d(H^{(i)}O^{(i+1)})$  in a six-membered ring, (c) probability distribution of

 $\Delta = \left(\sum_{i=1}^{6} |\phi(\gamma, s) - \phi_i(\gamma, s)|^2 / 6\right)^{1/2}, \text{ and (d) imaginary time correlation}$  function  $C(\tau, \phi_1) = \langle |\phi_1(\tau) - \phi_1(0)|^2 \rangle^{1/2}$ . (Reprinted with permission from Reference 160. Copyright 2017 Royal Society of Chemistry)

simulation has been performed to investigate the effect of the diameter and fluorination of the inner wall of CNT on the PT behavior.<sup>225</sup> The upper panel of Figure 4 shows the amount of H-bond formation between fluorine atoms and water molecules as a function of simulation time. The observed peaks clearly indicate that the smaller diameter particles prefer a greater

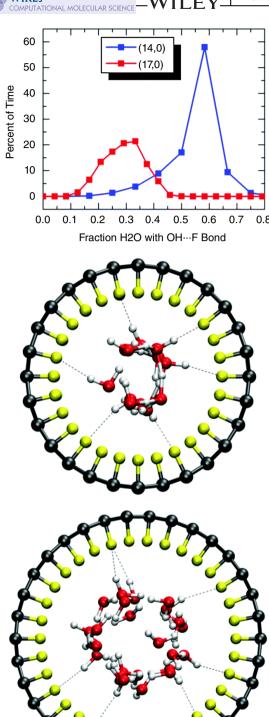
Φ/Å

0.5

 $\tau / \beta \hbar$ 

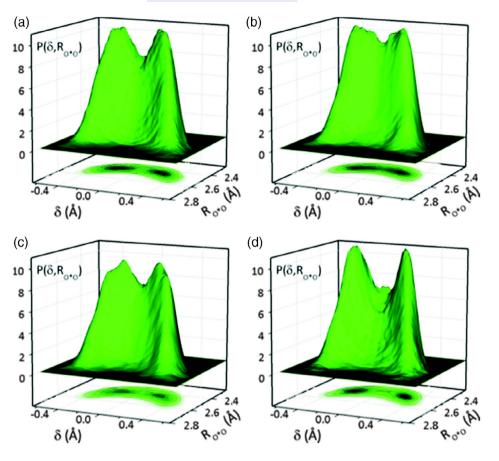
1.0

**FIGURE 4** Water molecule configurations in the presence of an excess proton in two different fluorinated carbon nanotube (CNT) sizes, namely, 13.3 Å (bottom) and 11.0 Å (middle). The CNT with a larger and smaller diameter exhibits (14,0) and (17,0) chirality, respectively. Percent of time throughout the trajectory  $OH \cdots F$  H-bonds exist with respect to the fraction of interacting water molecules is also shown (top). (Reprinted with permission from Reference 224. Copyright 2014 Royal Society of Chemistry)



H-bond formation, which has been confirmed by the typical snapshots in the bottom panel of Figure 4. In addition, a comparison with a system without an excess proton in the confined water environment exhibited the superior interaction of the charged endohedral water structure with the functionalized sidewall.

To establish a qualitative explanation of change in PT behavior due to the fluorination of CNT, probability distribution functions as a function of PT coordinates have been evaluated in Reference 224. In Figure 5, a local minimum at  $\delta = 0$  indicates a metastable Zundel form, that is, the proton is located between two water molecules. The Zundel form has been slightly dominant only in the case of fluorinated CNT with a small diameter. On the other hand, for fluorinated CNT with a large diameter, the opposite Zundel versus Eigen ratio has been obtained. The overall simulation results have implied that the



**FIGURE 5** Probability distribution functions,  $P(\delta,R_{\mathrm{O^*O}})$ , of the PT coordinate defined as  $\delta = R_{\mathrm{O^*H^*}} - R_{\mathrm{H^*O}}$  and the distance between the oxygen atoms,  $R_{\mathrm{O^*O}}$ , where "\*" indicates the hydronium ion. Non-fluorinated carbon nanotubes (CNTs) with chirality (14,0) and (17,0) are shown in (a) and (b), respectively, whereas the fluorinated CNTs with chirality (14,0) and (17,0) are shown in (c) and (d), respectively. (Reprinted with permission from Reference 224. Copyright 2014 Royal Society of Chemistry)

stronger interaction of tightly confined water with fluorinated groups induces a directional PT process perpendicular to the tube axis. This explanation has been further supported by projection maps of the most active proton in the AIMD trajectories.<sup>224</sup>

Despite excluding many factors such as nuclear quantum effects, density of water, and distribution of fluorine atoms covering the inner channel, AIMD simulations have provided the valuable molecular mechanism and a possible way to further design the channel space with improved proton conduction performance. However, currently, the estimation of the diffusion constant or proton conductivity for systems more complex than pure water structures is considered to be a challenging task. This may be due to the difficulty of conducting sufficient statistical sampling in the limited system size. At least, it may be necessary to perform the production runs for tens of picoseconds, for example, longer than 20 or 100 ps for a large and small system sizes, respectively. The simulations at the DFTB level may help to enhance the trajectory sampling as performed for the bulk liquid water and modeled channel.

In spite of the computational cost issue, structural and dynamical properties of  $H_3O^+$  and  $OH^-$  in CNT were successfully assessed at the DFT level. The simulations revealed that each  $H_3O^+$  or  $OH^-$  ion is two-coordinated to donate or accept a proton. Since no H-bond breaking is required, the PT rates become faster than the case of the bulk liquid water. It was also confirmed that the migration kinetics of  $OH^-$  is slightly faster than that of  $H_3O^+$  ion, which is opposite to the nature of their diffusivity in the bulk liquid water. Another theoretical study reported that the local structure around the  $H_3O^+$  ion is not significantly affected by the confinement. As a result, the dynamics of PT event is similar to that in the bulk phase, which depends on the material that confines the ionic defect. Note that for the latter case, the confinement was created by incorporating water in the mackinawite sheets FeS mineral.  $^{226,227}$ 

In contrast to the case of bulk liquid water and ice systems, the simulations of PT in confined water system can be performed using at the PBE level with ultrasoft pseudopotentials, without the dispersion correction. Adding the dispersion correction may deteriorate the system containing H-bond network. Such evidence may limit the applicability of the dispersion-corrected functional on the system containing a more complex material than the bulk liquid water system. Since the plane wave basis is usually employed, the size of basis set is defined by its cutoff, namely, 25 Ry (~340 eV)<sup>226</sup> and 400 eV.

# 5 | PROTON TRANSFER IN CARBON DIOXIDE CAPTURE AND STORAGE MATERIALS

Carbon dioxide capture and storage (CCS) is one of the industrial technologies aimed at solving the problem of excess  $CO_2$  present in the atmosphere. To urgently resolve such an environmental issue, the practical use of  $CO_2$  chemical absorption and regeneration processes is of paramount importance. While novel absorbents have been developed, there has been an increasing interest in developing a rational material design scheme on the basis of the elementary reactions of  $CO_2$  capture in order to further enhance the process performance. The computational studies have made a significant impact on the understanding of the mechanism and the evaluation of the microscopic properties, as observed in the recent comprehensive reviews. <sup>237–239</sup> In this section, we highlight the success of AIMD simulations in exploring the dynamic process of  $CO_2$  absorption in which PT is an inseparable event.

Among the CO<sub>2</sub> chemical absorbents, aqueous amine solution is the most widely used.<sup>240</sup> The application of AIMD to CO<sub>2</sub> absorption and regeneration processes has unveiled many mechanistic details that occur between CO<sub>2</sub> and amines, solvent water molecules, and water constituent ions.<sup>241–257</sup> The extensively investigated reaction is the formation of carbamate via a zwitterion intermediate, <sup>241–247,249,251–255,257</sup> possibly because it is experimentally difficult to probe the metastable chemical species. After the carbon atom of CO<sub>2</sub> approaches the nitrogen atom of amine molecules, the binding complex experiences deprotonation of the nitrogen atom by being attacked by a base with high proton affinity. The representative snapshots of CO<sub>2</sub> capture in a monoethanolamine solution obtained by Hwang and co-workers are shown in Figure 6, where the solvent water molecule serves as a base.<sup>246</sup> The PT event in Figure 6b proceeded rapidly because PT may be associated with the insignificant energy barrier of less than 0.1 eV, which is estimated from static quantum chemical calculations.<sup>238</sup> Note that other investigated elementary reactions in Figure 6c–e also involve the PT processes with structural diffusions or concerted migrations.

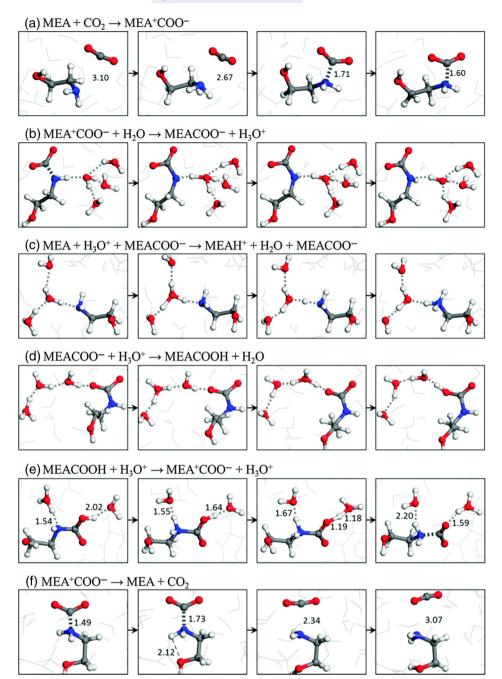
The recent large-scale DFTB-based MD simulations have investigated the participation of the OH<sup>-</sup> ion in the zwitterion mechanism.<sup>254</sup> The explicit inclusion of OH<sup>-</sup> in the system has led to the abstraction of the prevailing proton by the anion instead of the solvent water molecules. This observation arises from two factors: the stronger proton acceptor of negatively charged species and the increased probability of establishing a close contact with CO<sub>2</sub> in favor of the structural diffusivity of OH<sup>-</sup>. Sakti et al.<sup>254</sup> have further estimated rate constants of the zwitterion formation reaction by fitting time-course changes of chemical species to the kinetic model. Their analysis showed a reaction rate of zwitterion deprotonation slower than that of formation; this may be interpreted as the importance of OH<sup>-</sup> ion concentration for producing more CO<sub>2</sub> absorbed compound (carbamate ion).

The bicarbonate ion is another chemical product formed during the CO<sub>2</sub> capture process. <sup>258,259</sup> The AIMD simulations have not extensively targeted the bicarbonate formation reaction in an aqueous system. <sup>250,256,258,260,261</sup> Nakai et al. <sup>250</sup> have gained molecular insights from structural and electronic analyses that the Grotthuss-type diffusion of the OH<sup>-</sup> ion plays the key role in CO<sub>2</sub> interaction. They have also observed various chemical reactions involving the PT process during the sampling of 100 trajectories and further illustrated the direct PT mechanism in the CO<sub>2</sub> regeneration process. The PT process can be considered to dominate the major part of the chemical absorption in CCS.

The theoretical estimation of acidity (p $K_a$ ) of the protonated amine serve as an important research area related to CCS. The AIMD simulations have contributed to compute p $K_a$  values of acid molecules by considering the explicit solution environment and the temperature effect with the help of the metadynamics sampling technique. Recently, the excellent performance of the DFTB-based metadynamics approach has been verified for the calculation of p $K_a$  values of amines. The comparison with more than 30 experimental data are related to CCS. The area are related to CCS. and the simulation of acid to CCS. The sample of acid molecules by considering the explicit solution environment and the temperature effect with the help of the metadynamics sampling technique. The acid molecules by considering the explicit solution environment and the temperature effect with the help of the metadynamics sampling technique. The acid molecules by considering the explicit solution environment and the temperature effect with the help of the metadynamics sampling technique. The acid molecules by considering the explicit solution environment and the temperature effect with the help of the metadynamics sampling technique. The acid molecules by considering the explicit solution environment and the temperature effect with the help of the metadynamics sampling technique. The acid molecules by considering the explicit solution environment and the temperature effect with the help of the metadynamics sampling technique. The acid molecules by considering the explicit solution environment and the temperature effect with the help of the metadynamics sampling technique. The acid molecules by considering the explicit solution environment and the temperature effect with the help of the metadynamics and the explicit solution environment and the temperature effect with the help of the metadynamics and the explicit solution environment a

# 6 | CONCLUSION

Recent theoretical studies on PT in various water environments have highlighted the importance of the theoretical development for describing the water systems. Despite some theoretical issues and limitations, qualitative and quantitative properties associated with the PT process have been successfully evaluated and are found to be comparable with the experimental evidence. Along with more developed theoretical methods, the reason for OH<sup>-</sup> diffusion to be slower than that of H<sub>3</sub>O<sup>+</sup> in bulk liquid water has been revealed <sup>100</sup> and complemented the pre-solvation scenarios explained by AIMD simulations at the GGAs and DFTB levels. <sup>10,112,113</sup> While the accuracy of the DFT-based MD simulations depend on the choice of functional and the size of basis set, the DFTB methods relies on the performance of the parameter set. As previously examined by Choi et al., <sup>129</sup>



**FIGURE 6** Snapshots of the ab initio molecular dynamics simulations describing the elementary reactions during the CO<sub>2</sub> absorption in an aqueous monoethanolamine solution. (Reprinted with permission from Reference 246. Copyright 2015 Royal Society of Chemistry)

the extension of the second-order to the third-order DFTB did not significantly improve the structure and dynamics of PT in liquid water. In fact, it was also found that the former one has better description of PT dynamics. Incorporating the diffusion correction in the MD simulation has not been confirmed to improve the structure and dynamics of  $H_3O^+$  and  $OH^-$  ions in liquid water. Therefore, more thorough benchmark simulations at the DFTB level will be valuable assets for further improvement on the computational methodology.

At a low-temperature regime, in spite of high translational and rotational barriers that prevent vehicular diffusion, proton diffusivity is not negligible. The computer simulations have provided a theoretical evidence of faster diffusivity in the ice phase than that in bulk liquid water, <sup>176</sup> though occasionally the PT process is suppressed by the formation of a H<sub>3</sub>O<sup>+</sup> ion with long shelf life. Moreover, fast proton diffusivity can be observed in a confined water environment created inside the nanochannel or porous materials. In such a confined water structure, once a charge defect is created, the proton diffusivity can be accelerated up to 18 times faster than that in liquid water. <sup>221</sup> H<sub>3</sub>O<sup>+</sup> ion formation and water mobility are likely to be controlled by temperature, chemical environment related to hydrophobicity, and size of the nanoporous materials. <sup>222</sup> Furthermore, the PT process is essential in the chemical absorption and regeneration processes of the CCS technology. The in-depth



**TABLE 3** Experimental and estimated  $pK_a$  values calculated from density functional tight-binding-based metadynamics simulations <sup>269</sup>

Entry	Amine	Abbreviation	$pK_a$ (Calculated)	$pK_a$ (Experimental)
1	1-Piperidinepropanolamine	1-PPP	9.32	9.43 <sup>a</sup>
2	2-(2-Aminoethoxy)ethanol	2-AEE	9.49	9.42 <sup>b</sup>
3	2-(Diisopropylamino)ethanol	2-DIPA	9.27	9.11°
4	2-Piperidineethanol	2-PPE	11.08	10.90 <sup>a</sup>
5	2-Piperidinemethanol	2-PPM	10.47	10.60 <sup>a</sup>
6	3-Dimethylamino-1-propanol	3-DMAP	9.30	9.27 <sup>a</sup>
7	3-Morpholinopropylamine	3-MOPA	9.89	9.95 <sup>d</sup>
8	3-Piperidinemethanol	3-PPM	10.98	10.80 <sup>a</sup>
9	3-Piperidino-1,2-propanediol	3-PPPD	8.79	8.91 <sup>a</sup>
10	3-Quinoclidinol	3-QCD	10.22	10.10 <sup>a</sup>
11	4,2-Hydroxyethylmorpholine	4,2-HEMO	6.82	$6.90^{d}$
12	2-Amino-2-ethyl-1,3-propanediol	AEPD	8.78	8.82 <sup>b</sup>
13	2-Amino-2-methyl-1-propanol	AMP	9.16	9.23 <sup>b</sup>
14	Diethanolamine	DEA	8.80	8.88 <sup>a</sup>
15	Diisopropanolamine	DIPA	8.80	8.88 <sup>a</sup>
16	N,N-Dimethylisopropanolamine	DMIPA	9.45	9.47 <sup>d</sup>
17	2-(Ethylamino)ethanol	EAE	10.00	10.00 <sup>a</sup>
18	Ethyldiethanolamine	EDEA	8.66	$8.80^{\rm d}$
19	Methyldiethanolamine	MDEA	8.43	8.31°
20	Monoethanolamine	MEA	9.14	9.09 <sup>a</sup>
21	Methylmonoethanolamine	MMEA	9.94	9.85 <sup>b</sup>
22	3-Amino-1-propanol	MPA	9.86	9.96 <sup>b</sup>
23	Piperazine	PZ	9.30	9.38°
24	Serinol(2-aminopropane-1,3-diol)	SAPD	8.58	8.55 <sup>a</sup>
25	2-(tert-Butylamino)ethanol	TBAE	9.76	9.70 <sup>a</sup>
26	Triethanolamine	TEA	7.41	7.45 <sup>b</sup>
27	Tris(hydroxymethyl)aminomethane	THMAM	8.16	8.08 <sup>a</sup>
28	Tris[2-(2-methoxyethoxy)ethyl]amine	TMEEA	6.67	6.69 <sup>a</sup>
29	Tricine	TRC	8.00	8.10 <sup>a</sup>
30	Triethylamine	TREA	10.30	10.32 <sup>b</sup>
31	n-Cyclohexylethanolamine	n-CHEA	9.89	10.10 <sup>a</sup>
32	n-Cyclopentylethanolamine	n-CPEA	10.25	10.10 <sup>a</sup>
33	n-Cyclopropylethanolamine	n-CPREA	8.50	8.40 <sup>a</sup>
34	tert-Butyldiethanolamine	t-BDEA	8.92	9.03 <sup>d</sup>
Mean absolu	ate deviation		0.09	

<sup>&</sup>lt;sup>a</sup>Reference 270.

understanding of the reaction mechanisms and calculation of the  $pK_a$  values as a useful factor affecting  $CO_2$  absorption capability will boost future industrial applications of the PT process.

Classical MD simulations can hardly treat the PT event due to the inclusion of formation and cleavage of chemical bonds. On the other hand, recent QMMD simulations successfully rendered the mechanism, structure, dynamic, and energetic properties of the PT event, based on the careful treatments. First, the accuracy of the DFT method to describe the PT depends on the

<sup>&</sup>lt;sup>b</sup>Reference 271.

<sup>&</sup>lt;sup>c</sup>Reference 272.

dReference 273.

choice of functionals<sup>94</sup> and the size of basis set.<sup>118</sup> One can tune the accuracy by adding the dispersion correction with many possible choice of dispersion models.<sup>131,176</sup> Similarly, the accuracy of the DFTB methods rely on their inbuilt parameter set. A good DFTB water models were developed to describe the structural<sup>139</sup> and dynamical properties via a specific parameterization scheme that limits the applicability for complex applications. A more general parameterization scheme is necessary for improving the transferability of parameters.<sup>274</sup> Second, frontier PT applications require a computational method that able to treat a large system size. Despite the QM/MM approach is one of sophisticated methods to tackle a big system, the ubiquitous nature of PT requires one to carefully choose the size of QM layer that can correctly capture the proton movements. A full QM treatment is arguably necessary for handling such ubiquitous nature of PT. With a support from a massive parallelization, one could develop linear-scaling DFT<sup>119,275,276</sup> and DFTB methods.<sup>147</sup> The latter is more practical owe to a double benefit from cost efficiency of the DFTB method assisted by the DC method in a massive parallel computation scheme.<sup>143–147</sup>

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#### CONFLICT OF INTEREST

The authors have declared no conflicts of interest for this article.

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