

Experimental physik Vb (WS 2023/2024)

Exercise 4

Tutorial: 1 Deadline: 25.11.2024

Aufgabe 1: SU(3) Quark model of the hadrons

(a)

Let $A \in SU(3)$ and $B \in \mathfrak{su}(3)$. The defining property of SU(3) is that the determinant is one. This reduced the amount of free parameters, from nine in a 3×3 matrix to 8. SU(3) is thereby an 8-dimensional space, where as $\mathfrak{su}(3)$ must be at least 8-dimensional to be able to generate SU(3).

Starting with the definition of SU(3), one can derive a restraint for matrices of $\mathfrak{su}(3)$:

$$1 \stackrel{!}{=} \det A$$

$$= \det(\exp iB)$$

$$= \exp(i \operatorname{tr} B)$$

$$\implies \operatorname{tr} B = 2\pi n , n \in \mathbb{Z}$$

Arbitrarily choosing n = 0, since it is the simplest, one finds $3 \cdot 3 - 1 = 8$ traceless and linearly independent matrices, which will generate eight linearly independent members of SU(3). Considering that SU(3) is an 8-dimensional space, one can say that they generate SU(3) in its entirety. Taking a look at λ_1 , one an easily verify, that they indeed span the space of traceless 3×3 matrices, and therefore generate SU(3).

Note: Allowing for $n \neq 0$ will introduce more possible linearly independant generators of SU(3), however the space they span does not generate members of SU(3) in general, for example:

$$A = \begin{pmatrix} 2\pi & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \text{ tr } A = 2\pi$$

$$e^{iA} = \mathbb{1} \in SU(3)$$

$$e^{i\frac{A}{2\pi}} = \begin{pmatrix} e^{i\pi^2} & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \not\in SU(3)$$

The eight traceless matrices are the only ones that do span a space of generators, since $\operatorname{tr}(\lambda A) = \lambda \operatorname{tr} A = 0$ for a traceless matrix A. Choosing them to be hermitian as well is mathematically not necessary but probably convenient in the context of physics.

(b)

The eigenvectors of \hat{I}_3 are:

$$\hat{I}_3 = \frac{1}{2} \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$
$$|I_3, 1/2\rangle = (1, 0, 0)$$
$$|I_3, -1/2\rangle = (0, 1, 0)$$
$$|I_3, 0\rangle = (0, 0, 1)$$

The eigenvectors of \hat{Y} are:

$$\hat{Y} = \frac{1}{3} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$$
$$|Y, 1/3\rangle = (1, 0, 0)$$
$$|Y, 1/3\rangle = (0, 1, 0)$$
$$|Y, -2/3\rangle = (0, 0, 1)$$

 \hat{Y} and \hat{I}_3 evidently share the same three eigenstates.

(c)

For the operator V_{\pm} :

$$\begin{aligned} [\hat{I}_{3}, \hat{V}_{+}] &= \frac{1}{2} \hat{V}_{+} \\ [\hat{I}_{3}, \hat{V}_{-}] &= -\frac{1}{2} \hat{V}_{-} \\ \Longrightarrow [\hat{I}_{3}, \hat{V}_{\pm}] &= \pm \frac{1}{2} \hat{V}_{\pm} \end{aligned}$$

$$\begin{split} \hat{I}_{3}\hat{V}_{\pm}\left|I_{3},Y\right\rangle &=\left(\hat{V}_{\pm}\hat{I}_{3}+\left[\hat{I}_{3},\hat{V}_{\pm}\right]\right)\left|I_{3},Y\right\rangle \\ &=\left(\hat{V}_{\pm}\hat{I}_{3}\pm\frac{1}{2}\hat{V}_{\pm}\right)\left|I_{3},Y\right\rangle \\ &=\hat{V}_{\pm}\left(\hat{I}_{3}\pm\frac{1}{2}\mathbb{1}\right)\left|I_{3},Y\right\rangle \\ &=\left(I_{3}\pm\frac{1}{2}\right)\hat{V}_{\pm}\left|I_{3},Y\right\rangle \end{split}$$

$$\implies \hat{V}_{\pm} \left| I_3 \right\rangle \propto \left| I_3 \pm \frac{1}{2} \right\rangle$$

$$[\hat{Y}, \hat{V}_{+}] = \hat{V}_{+}$$

 $[\hat{Y}, \hat{V}_{-}] = -\hat{V}_{-}$

$$\implies [\hat{Y}, \hat{V}_{\pm}] = \pm \hat{V}_{\pm}$$

$$\hat{Y}\hat{V}_{\pm}|Y\rangle = (\hat{V}_{\pm}\hat{Y} + [\hat{Y}, \hat{V}_{\pm}])|Y\rangle$$

$$= (\hat{V}_{\pm}\hat{I}_{3} \pm \hat{V}_{\pm})|Y\rangle$$

$$= \hat{V}_{\pm}(\hat{I}_{3} \pm \mathbb{1})|Y\rangle$$

$$= (I_{3} \pm 1)\hat{V}_{\pm}|Y\rangle$$

$$\implies \hat{V}_{\pm}|Y\rangle = |Y \pm 1\rangle$$

$$\implies \hat{V}_{\pm}|I_{3}, Y\rangle \propto \left|I_{3} \pm \frac{1}{2}, Y \pm 1\right\rangle$$

For the operator U_{\pm} :

$$[\hat{I}_{3}, \hat{U}_{+}] = -\frac{1}{2}\hat{U}_{+}$$

$$[\hat{I}_{3}, \hat{U}_{-}] = \frac{1}{2}\hat{U}_{-}$$

$$\implies [\hat{I}_{3}, \hat{U}_{\pm}] = \mp \frac{1}{2}\hat{U}_{\pm}$$

$$\hat{I}_{3}\hat{U}_{\pm} |I_{3}\rangle = (\hat{U}_{\pm}\hat{I}_{3} + [\hat{I}_{3}, \hat{U}_{\pm}]) |I_{3}\rangle$$

$$= (\hat{U}_{\pm}\hat{I}_{3} \mp \frac{1}{2}\hat{U}_{\pm}) |I_{3}\rangle$$

$$= \hat{U}_{\pm} (\hat{I}_{3} \mp \frac{1}{2}\mathbb{1}) |I_{3}\rangle$$

$$= (I_{3} \mp \frac{1}{2}) \hat{U}_{\pm} |I_{3}\rangle$$

$$\implies \hat{U}_{\pm} |I_{3}\rangle \propto |I_{3} \mp \frac{1}{2}\rangle$$

$$\begin{split} [\hat{Y}, \hat{U}_{+}] &= \hat{U}_{+} \\ [\hat{Y}, \hat{U}_{-}] &= -\hat{U}_{-} \\ \Longrightarrow [\hat{Y}, \hat{U}_{\pm}] &= \pm \hat{U}_{\pm} \\ \\ \hat{Y}\hat{U}_{\pm} |Y\rangle &= (\hat{U}_{\pm}\hat{Y} + [\hat{Y}, \hat{U}_{\pm}]) |Y\rangle \\ &= \left(\hat{U}_{\pm}\hat{Y} \pm \hat{U}_{\pm}\right) |Y\rangle \\ &= \hat{U}_{\pm} \left(\hat{Y} \pm \mathbb{1}\right) |Y\rangle \\ &= (Y \pm 1) \hat{U}_{+} |Y\rangle \end{split}$$

$$\Longrightarrow \widehat{\left|\hat{U}_{\pm}\left|I_{3},Y\right\rangle \propto \left|I_{3}\mp\frac{1}{2},Y\pm1\right\rangle\right|}$$

For the operator I_{\pm} :

$$\begin{aligned} [\hat{I}_3, \hat{I}_+] &= \hat{I}_+ \\ [\hat{I}_3, \hat{I}_-] &= -\hat{I}_- \\ \Longrightarrow [\hat{I}_3, \hat{I}_\pm] &= \pm \hat{I}_\pm \end{aligned}$$

$$\hat{I}_3 \hat{I}_\pm |I_3\rangle = (\hat{I}_\pm \hat{I}_3 + [\hat{I}_3, \hat{I}_\pm]) |I_3\rangle$$

$$= (\hat{I}_\pm \hat{I}_3 \pm \hat{I}_\pm) |I_3\rangle$$

$$= \hat{I}_\pm (\hat{I}_3 \pm \hat{1}_\pm) |I_3\rangle$$

$$= \hat{I}_\pm (\hat{I}_3 \pm \hat{1}_\pm) |I_3\rangle$$

$$= (I_3 \pm 1) \hat{I}_\pm |I_3\rangle$$

$$\Longrightarrow \hat{I}_\pm |I_3\rangle \propto |I_3 \pm 1\rangle$$

$$\begin{aligned} [\hat{Y}, \hat{I}_{+}] &= 0 \\ [\hat{Y}, \hat{I}_{-}] &= 0 \\ \Longrightarrow [\hat{Y}, \hat{I}_{\pm}] &= 0 \\ \\ \hat{Y}\hat{I}_{\pm} |Y\rangle &= (\hat{I}_{\pm}\hat{Y} + [\hat{Y}, \hat{I}_{\pm}]) |Y\rangle \\ &= \hat{I}_{\pm}\hat{Y} |Y\rangle \end{aligned}$$

$$\Longrightarrow \widehat{I}_{\pm} |I_3, Y\rangle \propto |I_3 \pm 1, Y\rangle$$

 $=Y\hat{I}_{\pm}|Y\rangle$

(d)

Isospin and hypercharge of the three basis vectors where shown to be (task (b)):

$$\hat{I}_3 |u\rangle = \frac{1}{2} |u\rangle$$

$$\hat{I}_3 |d\rangle = -\frac{1}{2} |d\rangle$$

$$\hat{I}_3 |s\rangle = 0 |s\rangle$$

$$\begin{split} \hat{Y} & |u\rangle = \frac{1}{3} |u\rangle \\ \hat{Y} & |d\rangle = \frac{1}{3} |d\rangle \\ \hat{Y} & |s\rangle = -\frac{2}{3} |s\rangle \end{split}$$

Now for the strangeness S:

$$Y = B + S$$

$$S = Y - B = Y - \frac{1}{3}$$

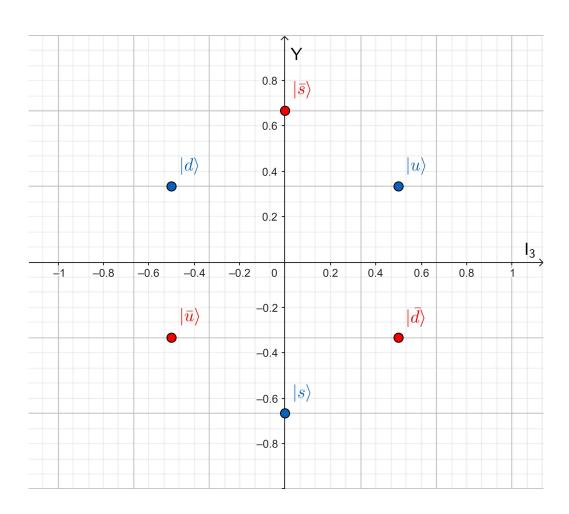
$$\implies \hat{S} = \hat{Y} - \frac{1}{3}\mathbb{1}$$

$$\hat{S} \left| u \right\rangle = 0 \left| u \right\rangle$$

$$\hat{S} | d \rangle = 0 | d \rangle$$

$$\hat{S}\left|s\right\rangle = -1\left|s\right\rangle$$

(e)



(f)

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Python-Code 1:
import numpy as np
from sympy import *
from sympy.abc import y
import matplotlib.pyplot as plt
plt.rcParams.update({"xtick.top": True , "ytick.right": True,
                     "xtick.minor.visible": True, "ytick.minor.visible": True,
                     "xtick.direction": "in", "ytick.direction": "in",
                     "axes.labelsize": "large", "text.usetex": True, "font.size": 13
                     })
11 = Matrix([[0,1,0],[1,0,0],[0,0,0]])
12 = Matrix([[0,-1j,0],[1j,0,0],[0,0,0]])
14 = Matrix([[0,0,1],[0,0,0],[1,0,0]])
15 = Matrix([[0,0,-1j],[0,0,0],[1j,0,0]])
16 = Matrix([[0,0,0],[0,0,1],[0,1,0]])
17 = Matrix([[0,0,0],[0,0,-1j],[0,1j,0]])
13 = Matrix([[1,0,0],[0,-1,0],[0,0,0]])
18 = 1/sqrt(3) * Matrix([[1,0,0],[0,1,0],[0,0,-2]])
Ip = (11 + 1j*12)/2
Im = (11 - 1j*12)/2
Vp = (14 + 1j*15)/2
Vm = (14 - 1j*15)/2
Up = (16 + 1j*17)/2
Um = (16 - 1j*17)/2
operators = [1,Ip,Im,Vp,Vm,Up,Um]
operator_names = ["", "$I_+$","$I_-$","$V_+$","$V_-$","$U_+$","$U_-$"]
I3 = 13/2
Y = 1/sqrt(3) * 18
u = Matrix([1,0,0])
d = Matrix([0,1,0])
s = Matrix([0,0,1])
quarks = np.array([u,d,s])
quark_names = [r"$|u\rangle$",r"$|d\rangle$",r"$|s\rangle$"]
quark_colors = ["b","g","r"]
com = lambda A,B: A*B - B*A
eigenval = lambda A,x: solve(A*x - y *x,y)[y]
getY = lambda x: eigenval(Y,x)
getI3 = lambda x: eigenval(I3,x)
getX = lambda x: np.array([getI3(x), getY(x)])
fig,ax = plt.subplots()
for quark,name,color in zip(quarks,quark_names,quark_colors):
```

```
t = 0.1 * getX(quark)
for op,op_name in zip(operators,operator_names):
    rot = op*quark
    if rot == 0*u: continue
    ax.scatter(*(getX(rot)+t), c=color)
    ax.text(*(getX(rot)+2.5*t-[0.05,0]), op_name+name)
plt.axhline(0,linestyle="--",c="gray",linewidth=0.5)
plt.axvline(0,linestyle="--",c="gray",linewidth=0.5)
ax.set(xlabel="$I_3$",ylabel="$Y$",xlim=(-0.8,0.8),ylim=(-0.9,0.6),title="Effects of shiftoperators")
plt.savefig("effects_of_shiftoperators.svg")
```

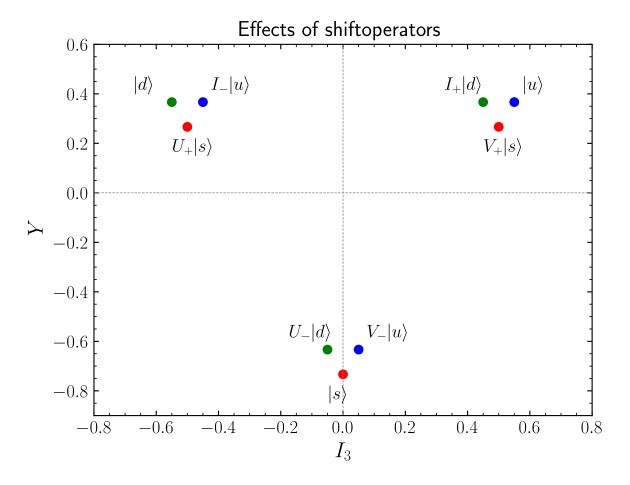


Figure 1: In the graphic, the triplet from the previous task is displayed, along 7.1 th all states affected by the shift operators. Most states get mapped to (0,0,0) and are therefore omitted. Additionally, states that are mapped to the same point have been slightly separated for better clarity.

(g)

i.

Using the graphics on the problem sheet, it is obvious that $[3] \otimes [\bar{3}]$ decomposes to $[1] \oplus [8]$.



ii.

To decompose $[3] \otimes [3] \otimes [3]$, it is easiest to decompose $[3] \otimes [3]$ first. Again taking a look at the graphical decomposition $[3] \otimes [3]$ decomposes to $[\bar{3}] \oplus [6]$.

$$[3] \otimes [3] \otimes [3] = [3] \otimes ([\overline{3}] \oplus [6])$$
$$= [3] \otimes [\overline{3}] \oplus [3] \otimes [6]$$
$$= [1] \oplus [8] \oplus [3] \otimes [6]$$

The remaining decomposition is $[3] \otimes [6]$ and it decomposes to $[10] \oplus [8]$. The composition of $[3] \otimes [3] \otimes [3]$ therefore is:

$$[3] \otimes [3] \otimes [3] = [1] \oplus [8] \oplus [3] \otimes [6]$$
$$= [1] \oplus [8] \oplus [8] \oplus [10]$$

(h)

```
Python-Code 2:
def add_particle(ax, x, y, label, offset=(0, 0), color='black', size=10):
    ax.scatter(x, y, c=color, s=size)
    ax.text(x + offset[0], y + offset[1], label, ha='center', fontsize=9, color=color)
shifts = {
    'ru': (0.1, 0.15), 'rd': (0.1, -0.2),
    'lu': (-0.1, 0.15), 'ld': (-0.2, -0.1),
    'l': (-0.15, 0), 'r': (0.15, 0),
    'u': (0, 0.15), 'd': (0, -0.15)
}
particles = [
    # Mesons - Octet
    (-1, 0.5, r"$K^+$", 'green', 'u'), (-1, -0.5, r"$K^0$", 'green', 'l'), # Kaons (Y = 1)
    (0, 1, r"$\pi^+$", 'green', 'd'), (0, 0, r"$\pi^0$", 'green', 'd'), (0, -1, r"$\pi^-$",
    \rightarrow 'green', 'l'), # Pions (Y = 0)
    (1, 0.5, r"$\overline{K}^0$", 'green', 'l'), (1, -0.5, r"$K^-$", 'green', 'd'), # Kaons
    \hookrightarrow (Y = -1)
    (0, 0, r"$\eta$", 'green', 'ru'),
    # Baryons - Octet
    (1, 0.5, r"$p$", 'blue', 'u'), (1, -0.5, r"$n$", 'blue', 'r'), # Proton & Neutron (Y = 1)
    (0, 1, r"$\Sigma^+$", 'blue', 'u'), (0, 0, r"$\Sigma^0$", 'blue', 'r'), (0, -1,
    \rightarrow r"$\Sigma^-$", 'blue', 'd'), # Sigmas (Y = 0)
    (-1, 0.5, r"$\Xi^0$", 'blue', 'l'), (-1, -0.5, r"$\Xi^-$", 'blue', 'u'), # Xis (Y = -1)
    (0, 0, r"$\Lambda$", 'blue', 'lu'),
```

```
# Baryons - Decuplet
    (1, 1.5, r"$\Delta^{++}$", 'red', 'u'), (1, 0.5, r"$\Delta^+$", 'red', 'r'),
    (1, -0.5, r"$\Delta^0$", 'red', 'l'), (1, -1.5, r"$\Delta^-$", 'red', 'd'), # Deltas (Y =
    → 1)
    (0, 1, r"$\Sigma^{*+}$", 'red', 'l'), (0, 0, r"$\Sigma^{*0}$", 'red', 'l'),
    (0, -1, r"$\Sigma^{*-}$", 'red', 'r'), # Sigma stars (Y = 0)
    (-1, 0.5, r"$\Xi^{*0}$", 'red', 'r'), (-1, -0.5, r"$\Xi^{*-}$", 'red', 'd'), # Xi stars
    \hookrightarrow (Y = -1)
    (-2, 0, r"\$\Omega^-, 'red', 'l') # Omega (Y = -2)
]
fig, ax = plt.subplots(figsize=(8, 8))
ax.axhline(0, color='black', linewidth=0.5)
ax.axvline(0, color='black', linewidth=0.5)
ax.set_xlim(-3, 3)
ax.set_ylim(-2, 2)
ax.set_xlabel(r"$Y$ (Hypercharge)", fontsize=12)
ax.set_ylabel(r"$I_3$ (Isospin projection)", fontsize=12)
ax.set_title(r"$SU(3)$ Multiplets for Lightest Mesons and Baryons", fontsize=14)
for x, y, label, color, shift_key in particles:
    offset = shifts.get(shift_key, (0, 0))
    add_particle(ax, x, y, label, offset=offset, color=color)
ax.grid(color='gray', linestyle='--', linewidth=0.5)
plt.tight layout()
plt.show()
fig.savefig("tmp.pdf")
```

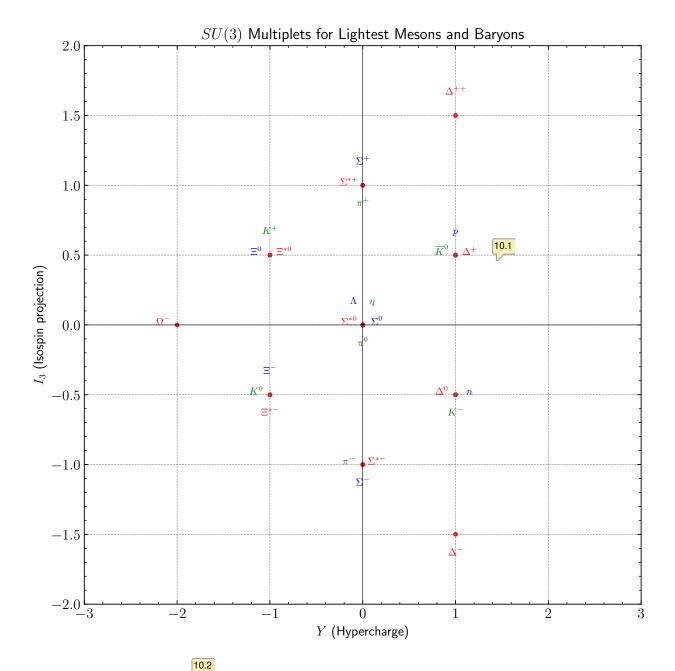


Figure 2: What a tedious task. Mesons, including the pions, kaons, and the η meson, are shown in green, while baryons from the SU(3) octet, such as protons, neutrons, and Λ baryons, are represented in blue. Baryons from the SU(3) decuplet, including the Δ , Σ^P , Ξ^* resonances, and the Ω^- baryon, are displayed in red.

(i)

The breaking of the SU(3) symmetrie is caused by the fact that the different quarks dont have the same mass and interact differently under the weak interaction. The symmetrie would be exact if only the strong force is looked at.

(j)

The relevant particles, their constituents and their masses are:

- Proton/Neutron (N): und or udd $\implies m_N = 3m_u + W_B$
- Xi (Ξ): uss or $dss \implies m_{\Xi} = m_u + 2m_s + W_B$

- Lambda (Λ): $uds \implies m_{\Lambda} = 2m_u + m_s + W_B$
- Sigma (Σ): uus, uds, or $dds \implies m_{\Sigma} = 2m_u + m_s + W_B$

Using this the mass relation can be shown:

$$\frac{m_N + m_{\Xi}}{2} = \frac{3m_{\Lambda} + m_{\Sigma}}{4}$$

$$\frac{(3m_u + W_B) + (m_u + 2m_s + W_B)}{2} = \frac{3(2m_u + m_s + W_B) + (2m_u + m_s + W_B)}{4}$$

$$2m_u + m_s + W_B = 2m_u + m_s + W_B$$

The experimental masses are approximatly:

$$m_N \approx 938 \,\mathrm{MeV}$$
 $m_\Xi \approx 1315 \,\mathrm{MeV}$
 $m_\Lambda \approx 1115 \,\mathrm{MeV}$
 $m_\Sigma \approx 1192 \,\mathrm{MeV}$

Plugging these values into the mass relation yields:

$$\frac{938 + 1315}{2} = \frac{3 \cdot 1115 + 1192}{4}$$
$$1126.5 = 1134.25$$

The relation holds quite well, with a small discrepancy, which can be attributed to the simplified assumptions made.



Index der Kommentare

- 7.1 Not sure what you mean by this. The shift operators change between the members of [3] and the origin is not one of them.
- 7.2 You still have to calculate it, or at least show pictorially.
- 10.1 The baryons are correct but the mesons are flipped with respect to where they should be. e.g. K+ should be at Y=1 and I3=1/2. A good way to check for this is that all particles on a point should have the same electromagnetic charge.
- 10.2 It's far quicker to do this by hand, if that's what you mean
- 11.1 1+2+4+1+3+2+1+2+1+2