

## Experimental Physics Vb (WS 2023/2024)

### Exercise 6

Tutorial: 1

Deadline: 16.12.2024

#### Task 1: Local gauge invariance of the Schrödinger equation

Applying the transformation to the right hand side yields:

$$\begin{aligned}
 iD'_0\Psi' &= i\left(\frac{\partial}{\partial t} + iQV'\right)e^{iQ\alpha(t,x)}\Psi \\
 &= i\left(\frac{\partial}{\partial t} + iQ(V - \partial_t\alpha)\right)e^{iQ\alpha(t,x)}\Psi \\
 &= i(\psi\partial_t + \partial_t\psi + iQ(V - \partial_t\alpha)\psi)e^{iQ\alpha(t,x)} \\
 &= i(iQ(\partial_t\alpha)\psi + \partial_t\psi + iQ(V - \partial_t\alpha)\psi)e^{iQ\alpha(t,x)} \\
 &= e^{iQ\alpha(t,x)}i(\partial_t + iQV)\psi
 \end{aligned}$$

$$\boxed{iD'_0\Psi' = e^{iQ\alpha(t,x)}iD_0\psi}$$

And to the left hand side:

$$\begin{aligned}
 D'_x\psi' &= \left(-\frac{\partial}{\partial x} + iQA'\right)\psi' \\
 &= \left(-\frac{\partial}{\partial x} + iQ(A + \partial_x\alpha)\right)e^{iQ\alpha(t,x)}\psi \\
 &= (-\partial_x\psi - \psi\partial_x + iQ(A + \partial_x\alpha)\psi)e^{iQ\alpha(t,x)} \\
 &= (-\partial_x\psi - iQ(\partial_x\alpha)\psi + iQ(A + \partial_x\alpha)\psi)e^{iQ\alpha(t,x)} \\
 &= e^{iQ\alpha(t,x)}(-\partial_x + iQA)\psi \\
 &= e^{iQ\alpha(t,x)}D_x\psi
 \end{aligned}$$

$$\boxed{D'^2_x\psi' = D'_x(e^{iQ\alpha(t,x)}D_x\psi) = e^{iQ\alpha(t,x)}D^2_x\psi}$$

All in all this implies that the Schrödinger equation is locally gauge invariant:

$$\begin{aligned}
 \frac{1}{2m}(iD'_x)^2\psi' &= iD'_0\Psi' \\
 e^{iQ\alpha(t,x)}\frac{1}{2m}(iD_x)^2\psi &= e^{iQ\alpha(t,x)}iD'_0\Psi' \\
 (e^{iQ\alpha(t,x)} \neq 0) \quad \frac{1}{2m}(iD_x)^2\psi &= iD_0\Psi
 \end{aligned}$$

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**Task 2: Gauge invariance and photon mass**

(a)

$$\begin{aligned} j^\mu &= \partial_\nu \partial^\nu \Lambda'^\mu - \partial^\mu \partial^\nu \Lambda'_\nu \\ &= \partial_\nu \partial^\nu (\Lambda^\mu - \partial^\mu \alpha) - \partial^\mu \partial^\nu (\Lambda_\nu - \partial_\nu \alpha) \\ &= \partial_\nu \partial^\nu \Lambda^\mu - \partial^\mu \partial^\nu \Lambda_\nu - \partial_\nu \partial^\nu \partial^\mu \alpha + \partial^\mu \partial^\nu \partial_\nu \alpha \\ &= \partial_\nu \partial^\nu \Lambda^\mu - \partial^\mu \partial^\nu \Lambda_\nu - \partial^\mu \partial_\nu \partial^\nu \alpha + \partial^\mu \partial_\nu \partial^\nu \alpha \\ &= \partial_\nu \partial^\nu \Lambda^\mu - \partial^\mu \partial^\nu \Lambda_\nu \end{aligned}$$

i.e the wave equation is invariant under the gauge transformation  $A^\mu \rightarrow A^\mu - \partial^\mu \alpha$ .

(b)

$$\begin{aligned} j^\mu &= (\partial_\nu \partial^\nu + m^2) \Lambda'^\mu - \partial^\mu \partial^\nu \Lambda'_\nu \\ &= \partial_\nu \partial^\nu \Lambda'^\mu - \partial^\mu \partial^\nu \Lambda'_\nu + m^2 \Lambda'^\mu \\ &= \partial_\nu \partial^\nu \Lambda^\mu - \partial^\mu \partial^\nu \Lambda_\nu + m^2 \Lambda'^\mu \\ 0 &= m^2 \underbrace{\Lambda'^\mu}_{\neq 0} \end{aligned}$$

i.e. wave equation for a massive vector field is only invariant under the same gauge transformation for  $m = 0$ .

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**Task 3: Higgs factory**

(a)

Electrons have no inner structure, which means that the initial state of a  $e^+e^-$  collision is known, where as with a  $p\bar{p}$  the quarks constituents have different energy levels resulting in a variety of possible initial states. Another advantage is that  $e^+e^-$  colliders produce fewer hadronic backgrounds.

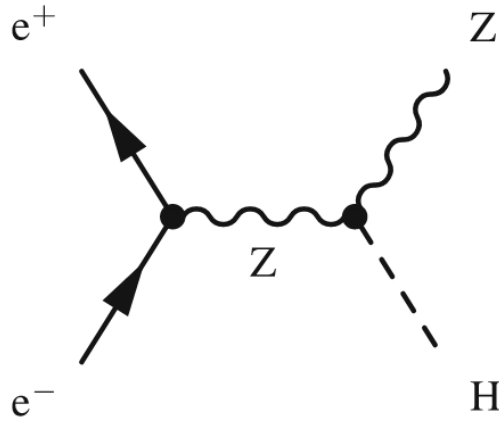
(b)

The minimum center of mass energy has to suffice to produce a non virtual  $Z$  and higgs boson:

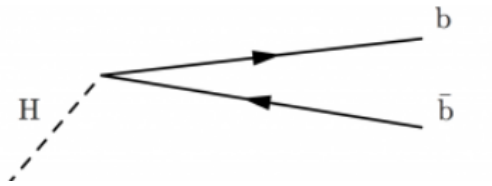
$$\sqrt{s_{\min}} = m_Z + m_H \approx 216 \text{ GeV}$$

(c)

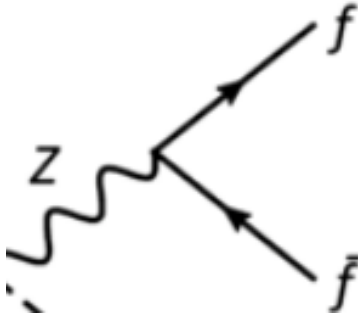
The  $e^+e^-$  collision produces a  $Z$  which radiates a higgs boson. The Feynman-diagram of this production is:



The heaviest fermion decay that can happen non virtually is  $H \rightarrow b\bar{b}$ , with the following feynman diagram:



While the  $Z$  decays either into two leptons  $Z \rightarrow l\bar{l}$  or two quarks  $Z \rightarrow q\bar{q}$ :



The experimental signature is therefore:

1. a pair of  $b$ -jets with invariant mass of  $m_H$ , corresponding to the Higgs
2. a pair of leptons or quark jets with invariant masses matching the  $Z$ -Bosons mass  $m_Z$ .

(d)

It follows from conservation of four-momentum that:

$$\begin{aligned}
 p' &= p \\
 p_{e^+} + p_{e^-} &= p_Z + p_H \\
 \boxed{p_H} &= p_{e^+} + p_{e^-} - p_Z
 \end{aligned}$$

The mass is then given by:

$$m_H^2 = p_H^2$$

$$\begin{aligned}
&= (p_{e^+} + p_{e^-} - p_Z)^2 \\
&= (p_{e^+} + p_{e^-})^2 - 2(p_{e^+} + p_{e^-})p_Z + p_Z^2 \\
&= s - 2(p_{e^+} + p_{e^-})p_Z + m_Z^2 \\
(\text{CoM frame}) \quad &= s - 2(\sqrt{s}, \vec{0}) \cdot (E_Z, \vec{p}_Z) + m_Z^2
\end{aligned}$$

$$m_H = \sqrt{s - 2E_Z\sqrt{s} + m_Z^2}$$

This relation can be used to calculate the Higgs mass without reconstructing the decay product of the Higgs boson. One just has to measure the center-of-mass energy  $\sqrt{s}$  along with the energy of the  $E_Z$  (assuming  $m_Z$  is known with sufficient accuracy).

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#### Task 4: Neutrino oscillations

(a)

$$\begin{pmatrix} |\nu_\mu\rangle \\ |\nu_\tau\rangle \end{pmatrix} = \underbrace{\begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix}}_R \begin{pmatrix} |\nu_2\rangle \\ |\nu_3\rangle \end{pmatrix}$$

$$\begin{aligned}
P_{\nu_\mu \rightarrow \nu_\mu}(t) &= |\langle \nu_\mu | \psi(t) \rangle|^2 \\
&= |\langle \nu_\mu | e^{-iHt/\hbar} | \nu_\mu \rangle|^2 \\
(E := \text{diag}(E_2, E_3)) \quad &= |\langle \nu_\mu | R e^{-iEt/\hbar} R^\dagger | \nu_\mu \rangle|^2 \\
(\text{sympy}) \quad &= \left( e^{\frac{iE_2 t}{\hbar}} \cos^2(\theta) + e^{\frac{iE_3 t}{\hbar}} \sin^2(\theta) \right) \left( e^{-\frac{iE_3 t}{\hbar}} \sin^2(\theta) + e^{-\frac{iE_2 t}{\hbar}} \cos^2(\theta) \right) \\
&= (\cos^2\theta + \sin^2\theta e^{i(E_2-E_3)t/\hbar}) (\cos^2\theta + \sin^2\theta e^{-i(E_2-E_3)t/\hbar}) \\
&= \cos^4\theta + \sin^4\theta + \sin^2\theta \cos^2\theta (e^{i(E_2-E_3)t/\hbar} + e^{-i(E_2-E_3)t/\hbar}) \\
&= \cos^4\theta + \sin^4\theta + 2\sin^2\theta \cos^2\theta \cos((E_2 - E_3)t/\hbar) \\
&= \frac{\cos(4\theta) + 3}{4} + \frac{1 - \cos(4\theta)}{4} \cos((E_2 - E_3)t/\hbar) \\
&= \frac{\cos(4\theta) + 3}{4} + \frac{1 - \cos(4\theta)}{4} \left( 1 - 2\sin^2\left(\frac{E_2 - E_3}{2}t/\hbar\right) \right) \\
&= 1 - \frac{1 - \cos(4\theta)}{2} \sin^2\left(\frac{E_2 - E_3}{2}t/\hbar\right) \\
&= 1 - \sin^2(2\theta) \sin^2\left(\frac{E_2 - E_3}{2}t/\hbar\right)
\end{aligned}$$

(b)

The code is appended on the last pages.

$$P_{\nu_\mu \rightarrow \nu_\mu}(t) = 1 - \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2 L}{4E}\right)$$

The conversion factor from km to 1/eV can be derived as:

$$\hbar c = 197 \text{ MeV} \cdot \text{fm} = 197 \cdot 10^6 \text{ eV} \cdot 10^{-18} \text{ km}$$

$$\text{km} = \frac{\hbar c}{197 \cdot 10^{-12} \text{ eV}} \frac{1}{\text{eV}} \approx 5.076 \cdot 10^9 \frac{1}{\text{eV}}$$

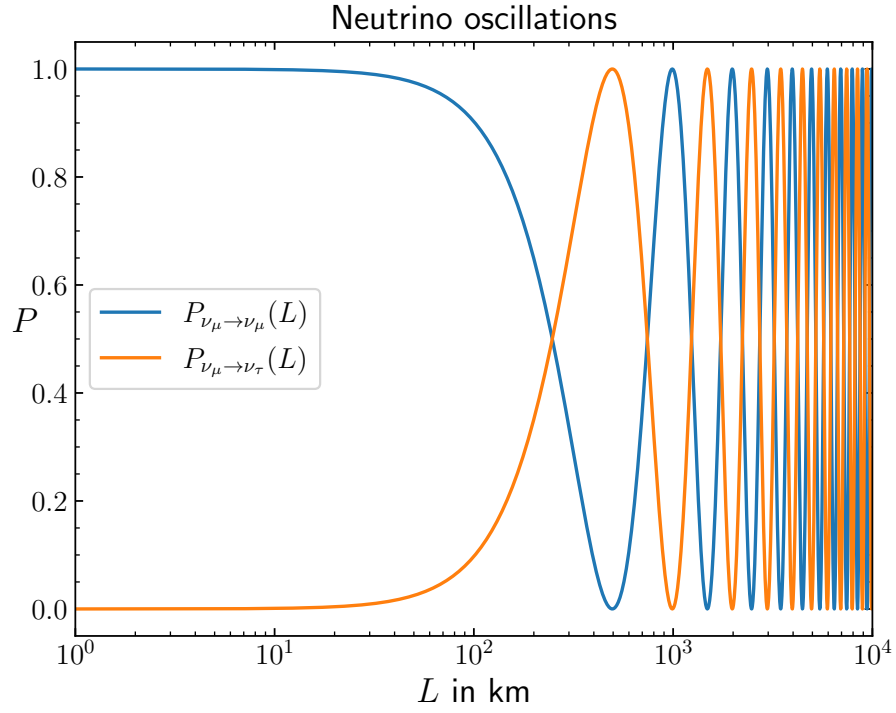


Figure 1: resulting plot

# jupyter

December 31, 2024

```
[35]: import numpy as np
from sympy import *
init_printing(use_latex="mathjax")
```

```
[10]: from sympy import *
theta, E2, E3, t, hbar = symbols("theta E2 E3 t hbar", real=True)

# Nr. 4
# (a)
psi = Matrix([1, 0])
R = Matrix([[cos(theta), sin(theta)], [-sin(theta), cos(theta)]])
Rinv = R.adjoint()
E = Matrix([[E2, 0], [0, E3]])

A = psi.T * R * exp(-I * E * t / hbar) * Rinv * psi
P = conjugate(A) * A
simplify(P[0])
```

```
[10]:  $\left( e^{\frac{iE_2 t}{\hbar}} \cos^2(\theta) + e^{\frac{iE_3 t}{\hbar}} \sin^2(\theta) \right) \left( e^{-\frac{iE_3 t}{\hbar}} \sin^2(\theta) + e^{-\frac{iE_2 t}{\hbar}} \cos^2(\theta) \right)$ 
```

```
[73]: # (b)
import scipy.constants as c
import matplotlib.pyplot as plt
plt.rcParams.update({"xtick.top": True, "ytick.right": True,
                    "xtick.minor.visible": True, "ytick.minor.visible": True,
                    "xtick.direction": "in", "ytick.direction": "in",
                    "axes.labelsize": "large", "text.usetex": True, "font.
↪size": 13
                    })

theta = np.pi/4
Delta_m_sq = 0.0025 / c.c**2 * c.e # J
Delta_E = c.c * Delta_m_sq
E = 1e9 # eV
L = np.logspace(0, 4, 1000) # km
t = L / c.c

P_mu = 1 - np.sin(2*theta)**2 * np.sin(Delta_E / 2 * t / c.hbar)**2
```

```

P_tau = 1-P_mu

plt.plot(L, P_mu, label="$P_{\nu_\mu \to \nu_\mu}(L)$")
plt.plot(L, P_tau, label="$P_{\nu_\mu \to \nu_\tau}(L)$")
plt.xscale("log")
plt.xlim(min(L), max(L))
plt.xlabel("$L$ in km")
plt.ylabel("$P$",rotation=0)
plt.legend()
plt.title("Neutrino oscillations")
plt.savefig("neutrino_oscillations.pdf")

```

