

Mecânica Quântica Avançada

Lista 1

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Exercise 1 (4.1). Show that a necessary and sufficient condition for $|\psi\rangle$ to be an eigenvector of a Hermitian operator A is that the dispersion (4.8) $\Delta_\psi A = 0$.

Answer 1. Vamos iniciar mostrando que se $|\psi\rangle = 0$, então $\Delta_\psi A = 0$. Ora, por definição:

$$\Delta_\psi A = \langle \psi | A^2 | \psi \rangle - \langle \psi | A | \psi \rangle^2 \quad (1)$$

É fácil ver que

$$\langle \psi | A^2 | \psi \rangle = \langle \psi | A^\dagger | A | \psi \rangle = a^2 \quad (2)$$

E também

$$\langle \psi | A | \psi \rangle^2 = (a)^2 = a^2 \quad (3)$$

Logo, é trivial que $\Delta_\psi A = 0$.

Agora vamos assumir que a dispersão é nula, ou seja, $\Delta A = 0$. Então, por definição:

$$0 = \sqrt{\langle A^2 \rangle - \langle A \rangle^2} \quad (4)$$

$$\Rightarrow \langle A^2 \rangle = \langle A \rangle^2 \quad (5)$$

Exercise 2 (4.4.2 - 1). Let $|\psi\rangle$ be a vector (not normalized) in the Hilbert space of states and H be a Hamiltonian. The expectation value $\langle H \rangle_\psi$ is

$$\langle \psi \rangle_\psi = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} \quad (6)$$

Show that if the minimum of this expectation value is obtained for $|\psi\rangle = |\psi_m\rangle$ and the maximum for $|\psi\rangle = |\psi_M\rangle$, then

$$H |\psi_m\rangle = E_m |\psi_m\rangle, \quad H |\psi_M\rangle = E_M |\psi_M\rangle \quad (7)$$

where E_m and E_M are the smallest and largest eigenvalues.

Answer 2. É evidente que

$$\langle H \rangle_{\psi_m} = \frac{\langle \psi_m | H | \psi_m \rangle}{\langle \psi_m | \psi_m \rangle} = E_m \quad (8)$$

Portanto, é evidente que se $\langle H \rangle_{\psi_m}$ for mínimo, então E_m também é. Vale um raciocínio análogo para $E_M = \langle H \rangle_{\psi_M}$.

Exercise 3 (4.4.2 - 2). We assume that the vector $|\varphi\rangle$ depends on a parameter α : $|\varphi\rangle = |\varphi(\alpha)\rangle$. Show that if

$$\left. \frac{\partial \langle H \rangle_{\varphi(\alpha)}}{\partial \alpha} \right|_{\alpha=\alpha_0} = 0, \quad (9)$$

then $E_m \leq \langle H \rangle_{\varphi(\alpha_0)}$ if α_0 corresponds to a minimum of $\langle H \rangle_{\varphi(\alpha)}$, and $\langle H \rangle_{\varphi(\alpha_0)} \leq E_M$ if α_0 corresponds to a maximum. This result forms the basis of an approximation method called the variational method (Section 14.1.4).

Answer 3. Vamos abrir a derivada:

$$\frac{\partial \langle H \rangle_{\varphi(\alpha)}}{\partial \alpha} = \frac{1}{\langle \psi | \psi \rangle} (\langle \psi | H | \partial_\alpha | \psi \rangle + (\partial_\alpha \langle \psi |) | H | \psi \rangle) - \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle^2} ((\partial_\alpha \langle \psi |) | \psi \rangle + \langle \psi | \partial_\alpha | \psi \rangle) \quad (10)$$

Então, em α_0

$$\langle \psi | H | \partial_\alpha \psi \rangle + \langle \partial_\alpha \psi | H | \psi \rangle = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} (\langle \partial_\alpha \psi | \psi \rangle + \langle \psi | \partial_\alpha \psi \rangle) \quad (11)$$

Isolando a quantidade de interesse:

$$\langle H \rangle_{\varphi(\alpha_0)} = \frac{\langle \psi | H | \partial_\alpha \psi \rangle + \langle \partial_\alpha \psi | H | \psi \rangle}{\langle \partial_\alpha \psi | \psi \rangle + \langle \psi | \partial_\alpha \psi \rangle} \quad (12)$$

Podemos reescrever, usando que $H = H^\dagger$:

$$\langle H \rangle_{\varphi(\alpha_0)} = \frac{\langle \partial_\alpha \psi | H | \psi \rangle^\dagger + \langle \partial_\alpha \psi | H | \psi \rangle}{\langle \partial_\alpha \psi | \psi \rangle^\dagger + \langle \partial_\alpha \psi | \psi \rangle} \quad (13)$$

Podemos expandir qualquer estado usando os autokets do hamiltoniano, que são ortonormais:

$$|\psi\rangle = \sum c_j |\psi_j\rangle \quad (14)$$

De modo que

$$\partial_\alpha |\psi\rangle = \sum \frac{\partial c_j}{\partial \alpha} |\psi_j\rangle \quad (15)$$

Assim, vale que

$$\langle \partial_\alpha \psi | \psi \rangle = \sum c_j \frac{\partial c_j^*}{\partial \alpha} \quad (16)$$

E também

$$\langle \partial_\alpha \psi | H | \psi \rangle = \sum E_j c_j \frac{\partial c_j^*}{\partial \alpha} \quad (17)$$

Considere o denominador:

$$\begin{aligned} \langle \partial_\alpha \psi | \psi \rangle^\dagger + \langle \partial_\alpha \psi | \psi \rangle &= \sum c_j \frac{\partial c_j^*}{\partial \alpha} + c_j^* \frac{\partial c_j}{\partial \alpha} \\ &= \sum \frac{\partial}{\partial \alpha} (c_j^* c_j) \\ &= \sum \partial_\alpha |c_j|^2 \\ &= \partial_\alpha \sum |c_j|^2 \end{aligned}$$

Considerando, enfim, o numerador e fazendo os mesmos cálculos:

$$\begin{aligned}
\langle \partial_\alpha \psi | H | \psi \rangle^\dagger + \langle \partial_\alpha \psi | H | \psi \rangle &= \sum E_j c_j \frac{\partial c_j^*}{\partial \alpha} + E_j c_j^* \frac{\partial c_j}{\partial \alpha} \\
&= \sum E_j \partial_\alpha |c_j|^2 \\
&= \partial_\alpha \sum E_j |c_j|^2
\end{aligned}$$

É óbvio, então, que

$$E_m \partial_\alpha \sum |c_j|^2 \leq \partial_\alpha \sum E_j |c_j|^2 \leq E_M \partial_\alpha \sum |c_j|^2 \quad (18)$$

Portanto, concluímos que

$$E_m \leq \langle H \rangle_{\varphi(\alpha_0)} \leq E_M \quad (19)$$